



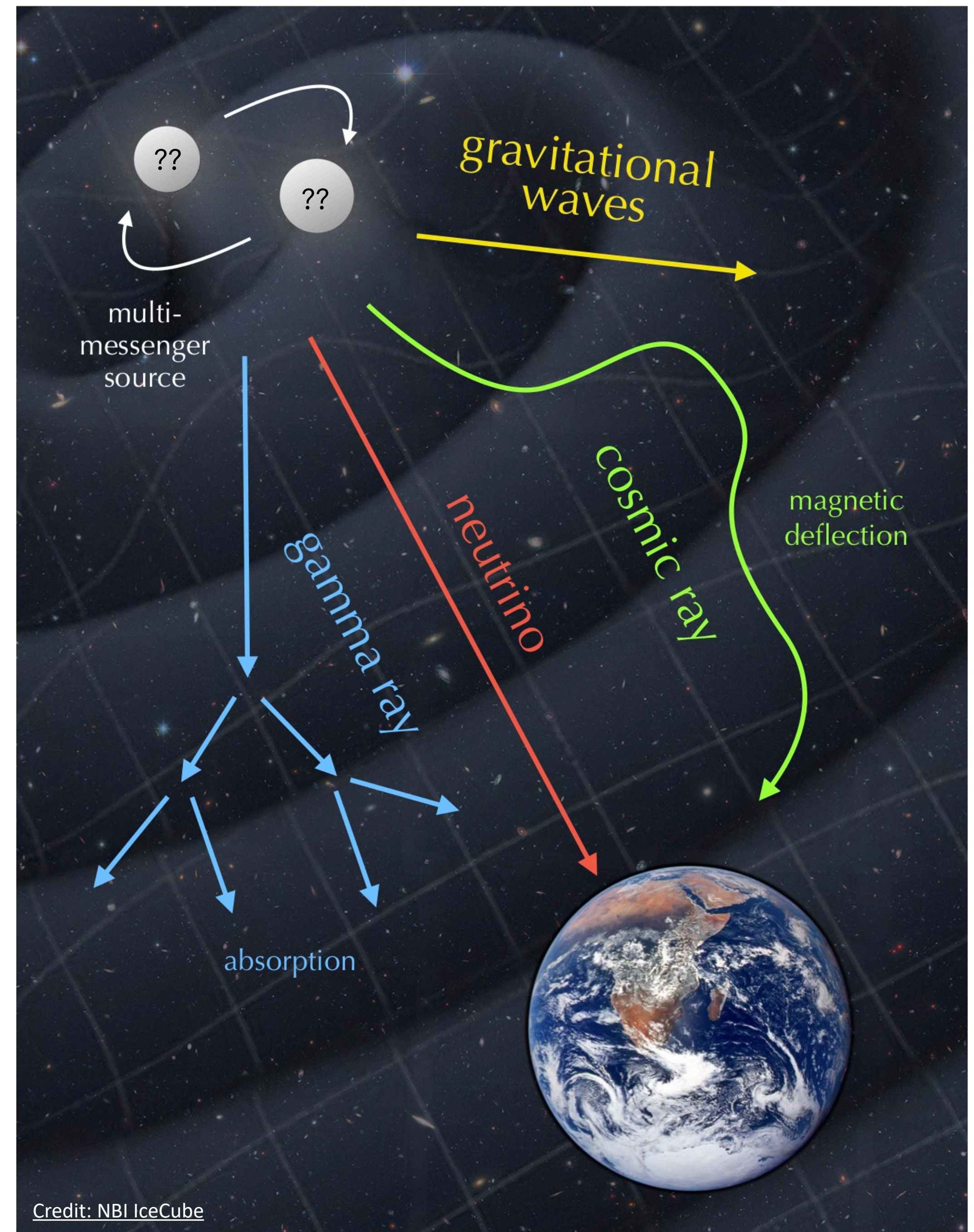
Indirect probes of dark matter

Francesca Calore, CNRS, LAPTh
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*ISAPP 2024: Particle candidates for dark matter
2024 Jul 1 — 4, Padova (IT)*

Plan of the lectures

- **Lecture 1:** Particles from the sky: **charged cosmic rays**, **gamma rays** (and **neutrinos**)
 - Basics and jargon
 - Production mechanisms
- **Lecture 2:** Sources of cosmic particles
 - Standard astrophysics
 - Generalities of dark matter searches
- **Lecture 3:** Probing **dark matter** with astroparticle observations
 - Particle** dark matter
 - Primordial black holes**
 - Anomalies and **excesses**



Some reading material

M. Longair, [*High energy astrophysics*](#), Cambridge Univ. Press (2012)

V. S. Berezhinskii et al., [*Astrophysics of cosmic rays*](#), Amsterdam: North-Hollans (1990)

G. Sigl, [*Astroparticle Physics: Theory and Phenomenology*](#), Atlantis Press Paris (2017)

G. Ghisellini, [*Radiative processes in high-energy astrophysics*](#), Lect. Notes Phys. (2013)

Very good lectures notes: [*Foundations of cosmic-ray astrophysics*](#), Varenna (2022)

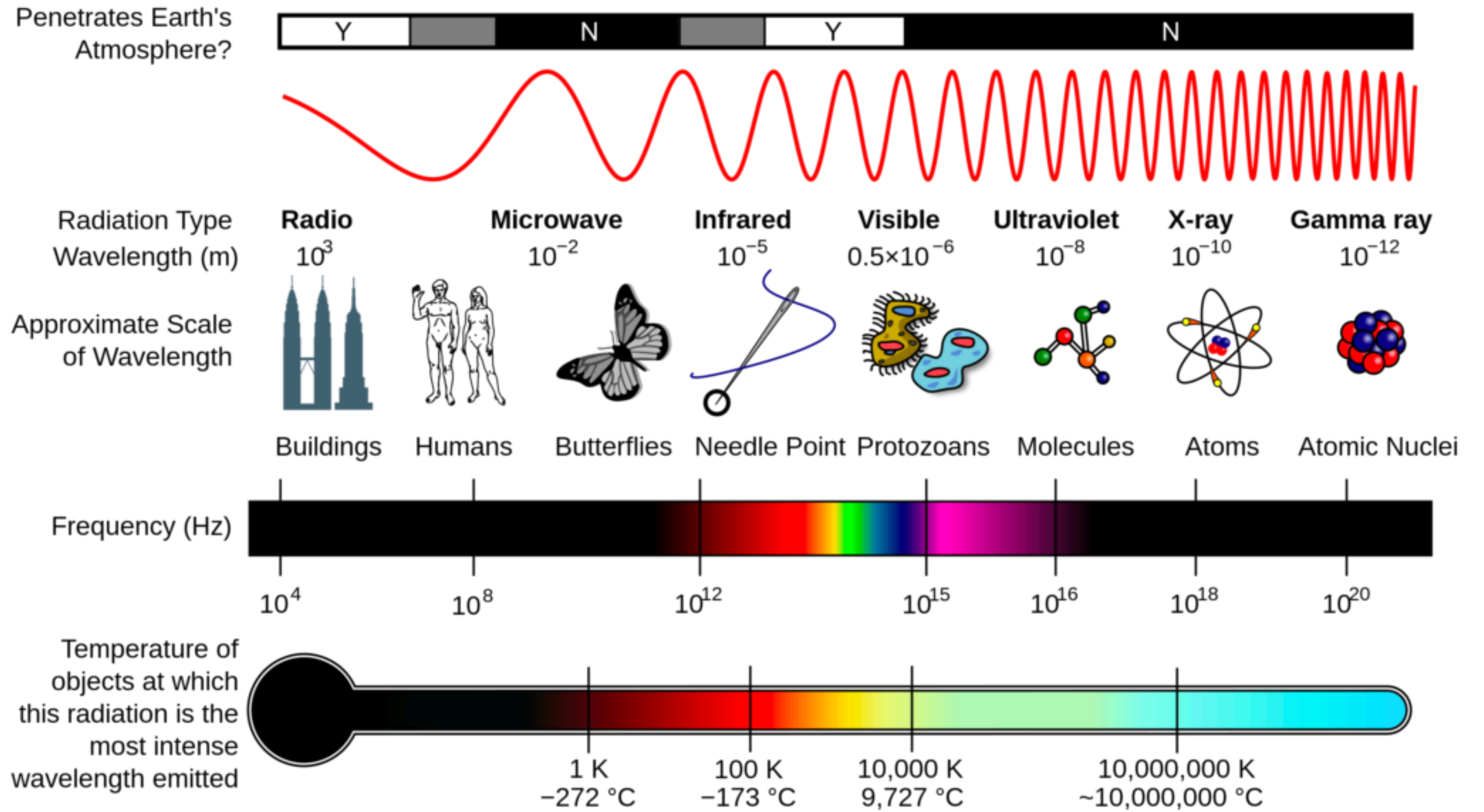
EuCAPT White Paper, [*Opportunities and Challenges for Theoretical Astroparticle Physics in the Next Decade*](#), *arXiv:2110.10074*



Feel free to email me at calore@lapth.cnrs.fr!

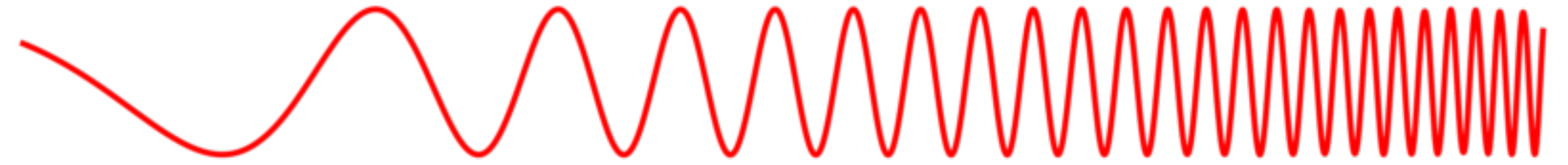
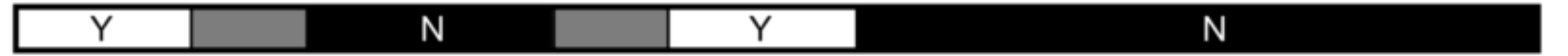
Prolegomena

The MW spectrum



The MW spectrum

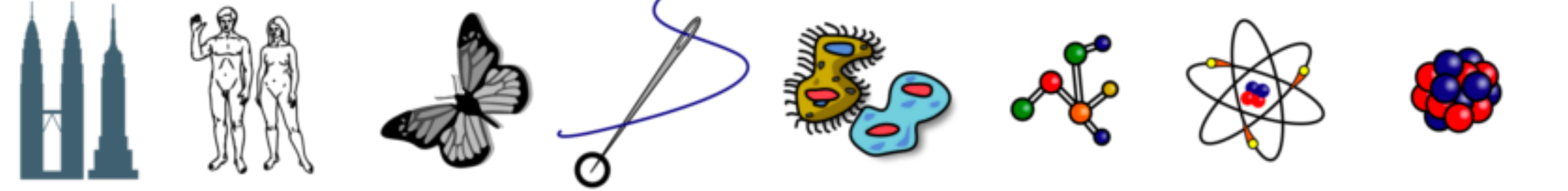
Penetrates Earth's Atmosphere?



Radiation Type
Wavelength (m)

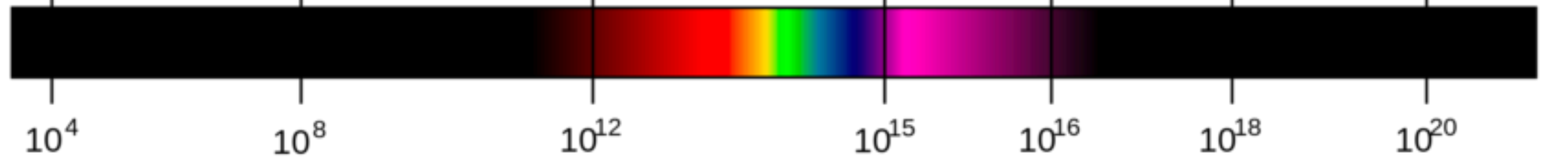
Radio	Microwave	Infrared	Visible	Ultraviolet	X-ray	Gamma ray
10^3	10^{-2}	10^{-5}	0.5×10^{-6}	10^{-8}	10^{-10}	10^{-12}

Approximate Scale of Wavelength

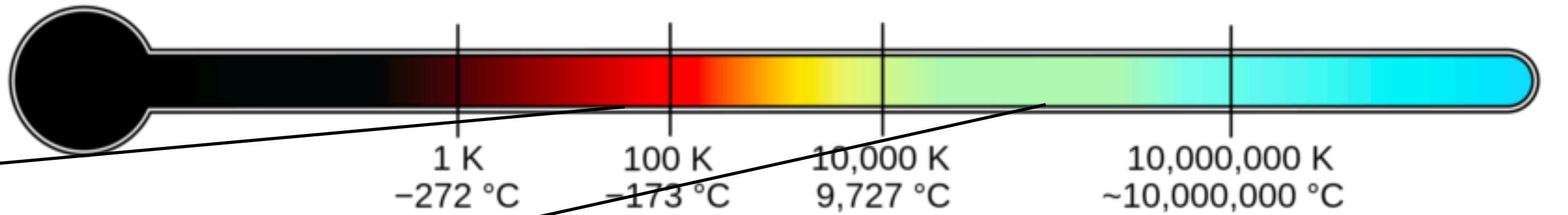


Buildings Humans Butterflies Needle Point Protozoans Molecules Atoms Atomic Nuclei

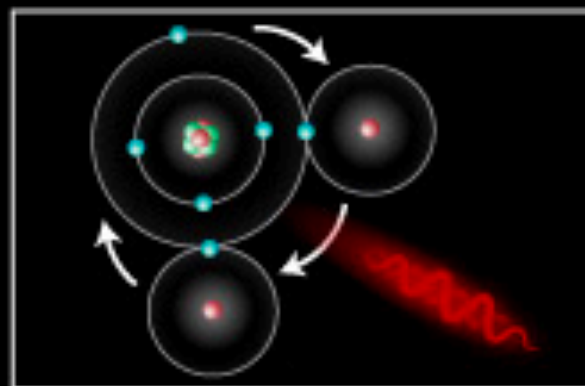
Frequency (Hz)



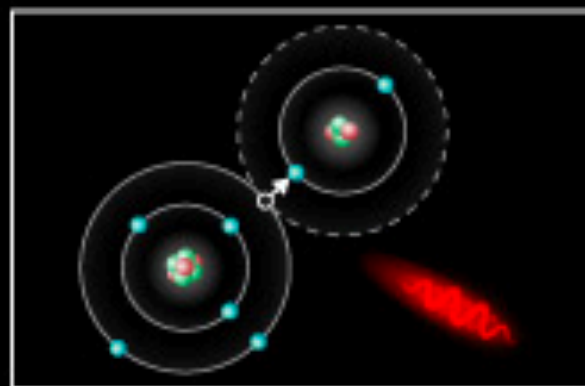
Temperature of objects at which this radiation is the most intense wavelength emitted



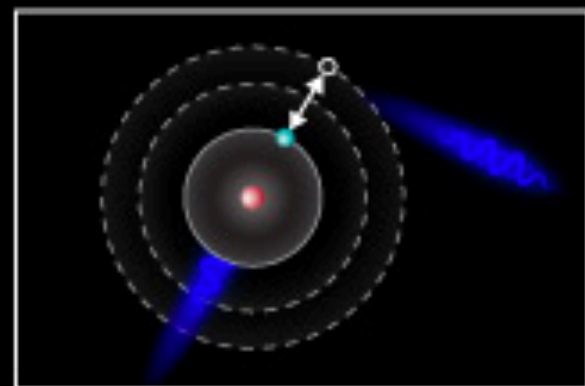
Thermal radiation



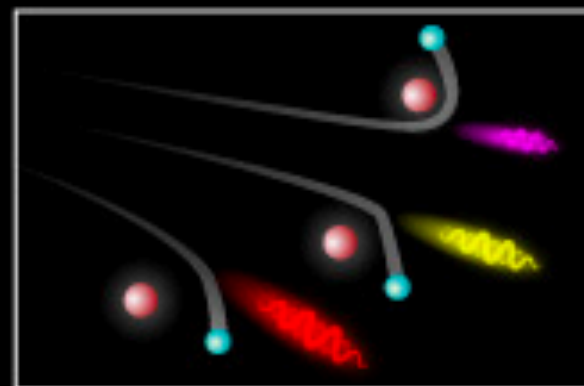
Molecular Motion
Low energy photons are created from molecules twisting and spinning.



Atoms Colliding
Low to high energy photons are created when atoms exchange electrons in collisions.



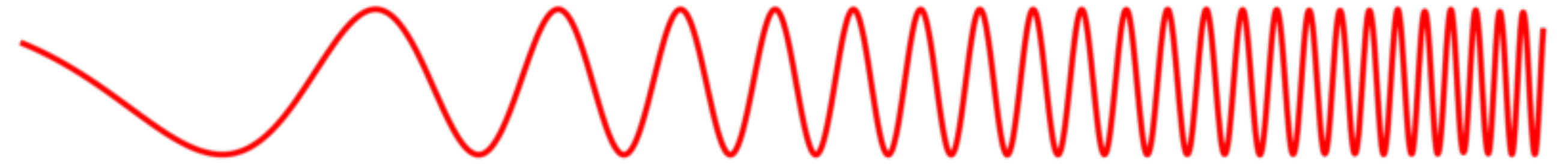
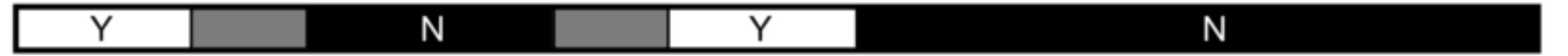
Photon Absorption
Atoms can absorb the energy of a photon, which excites an electron to release another photon.



Free Electrons
In an ionized gas, electrons fly past nuclei creating photons in a continuum of energy levels.

The MW spectrum

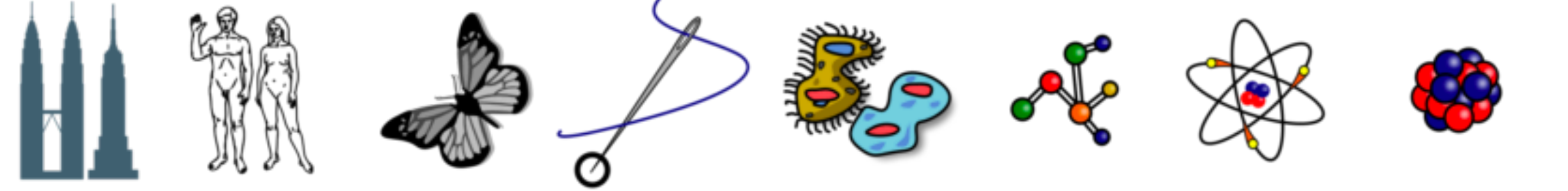
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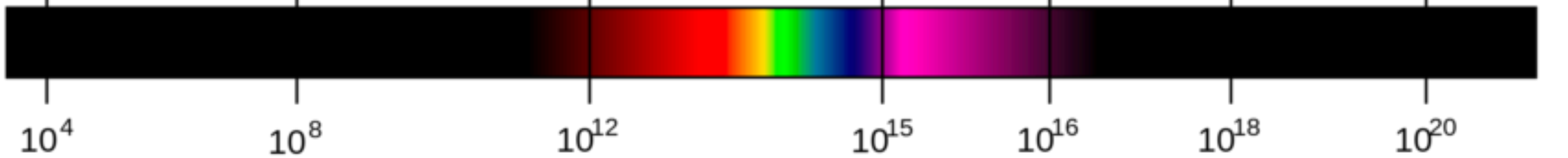
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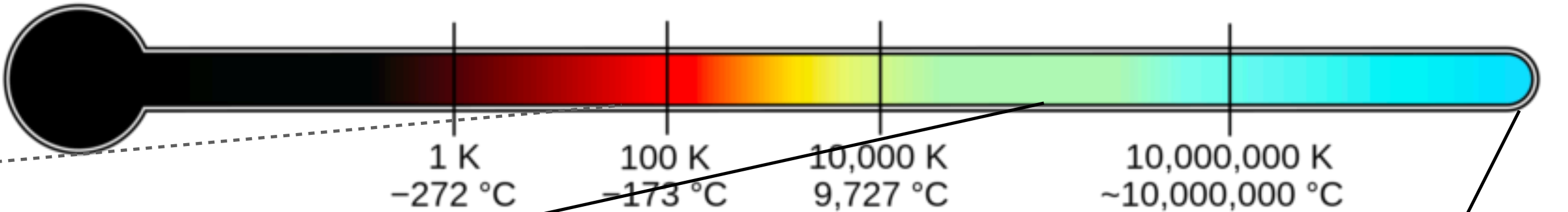


Buildings Humans Butterflies Needle Point Protozoans Molecules Atoms Atomic Nuclei

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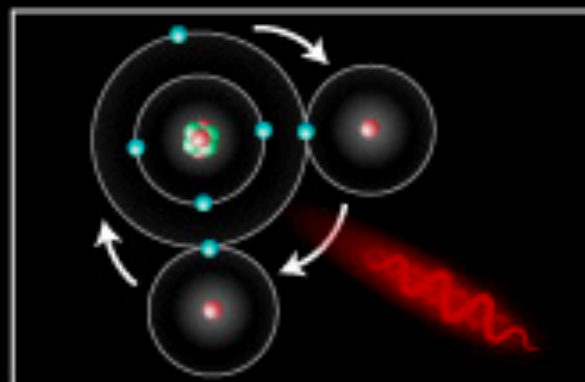


Temperature of objects at which this radiation is the most intense wavelength-emitted

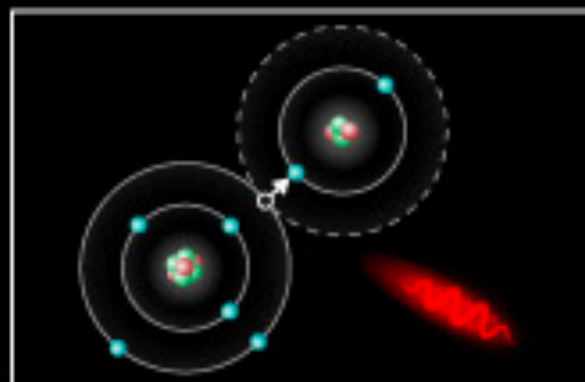


Thermal radiation

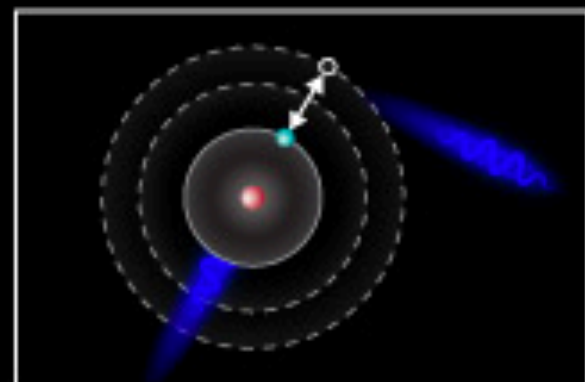
Non-thermal radiation



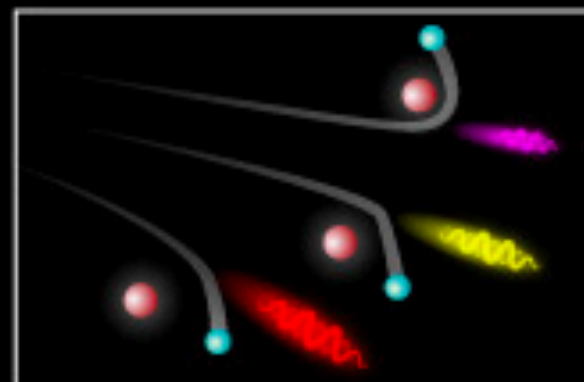
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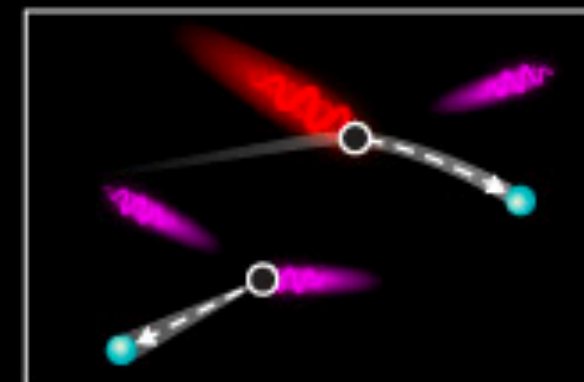
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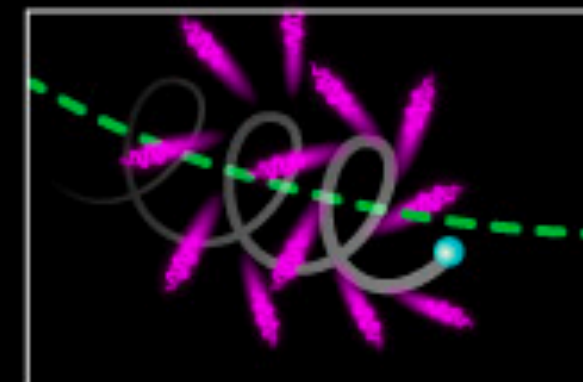
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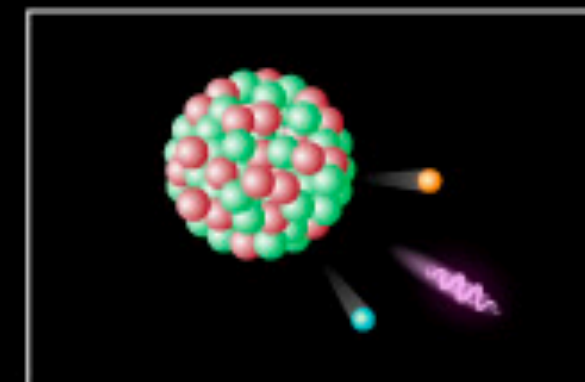
Free Electrons
In an ionized gas, electrons fly past nuclei creating photons in a continuum of energy levels.



Compton Scattering
Photons can collide with electrons, causing the electron to gain or lose energy releasing a photon.

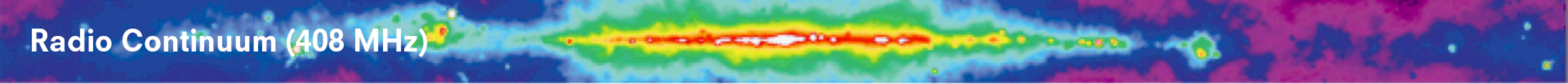


Synchrotron Radiation
Strong magnetic fields can accelerate electrons to release extremely high energy photons.

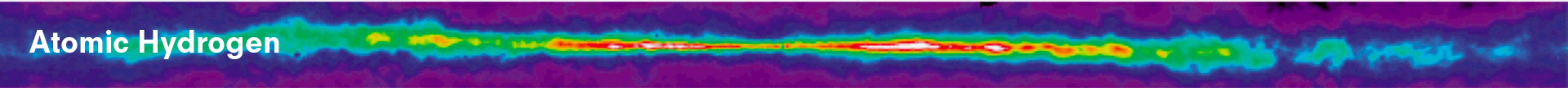


Nuclear Emission
High energy gamma rays and subatomic particles are emitted from the nuclei of radioactive atoms.

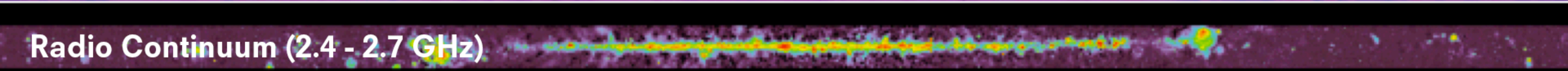
Radio Continuum (408 MHz)



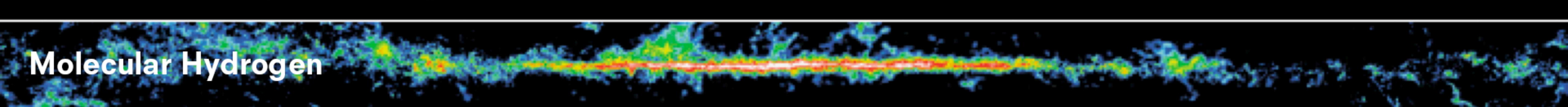
Atomic Hydrogen



Radio Continuum (2.4 - 2.7 GHz)



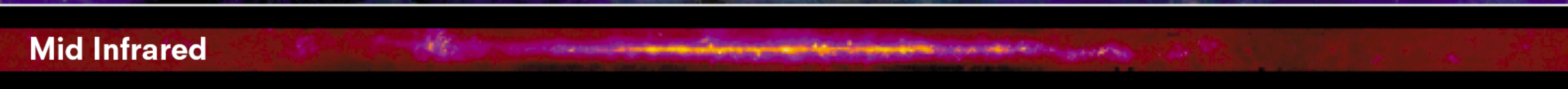
Molecular Hydrogen



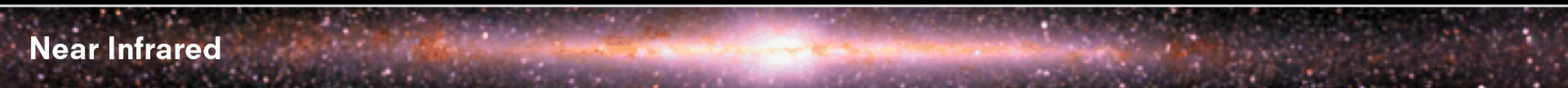
Mid Infrared



Mid Infrared



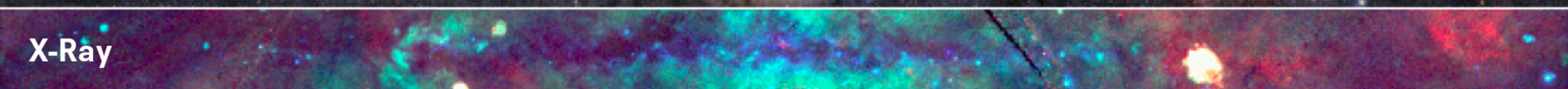
Near Infrared



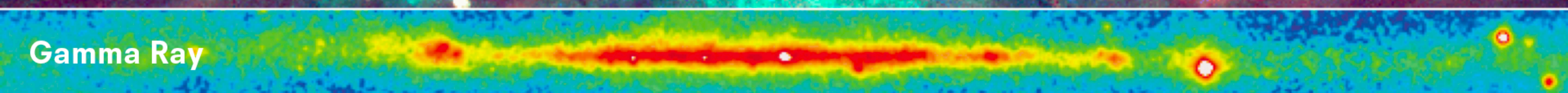
Visible



X-Ray



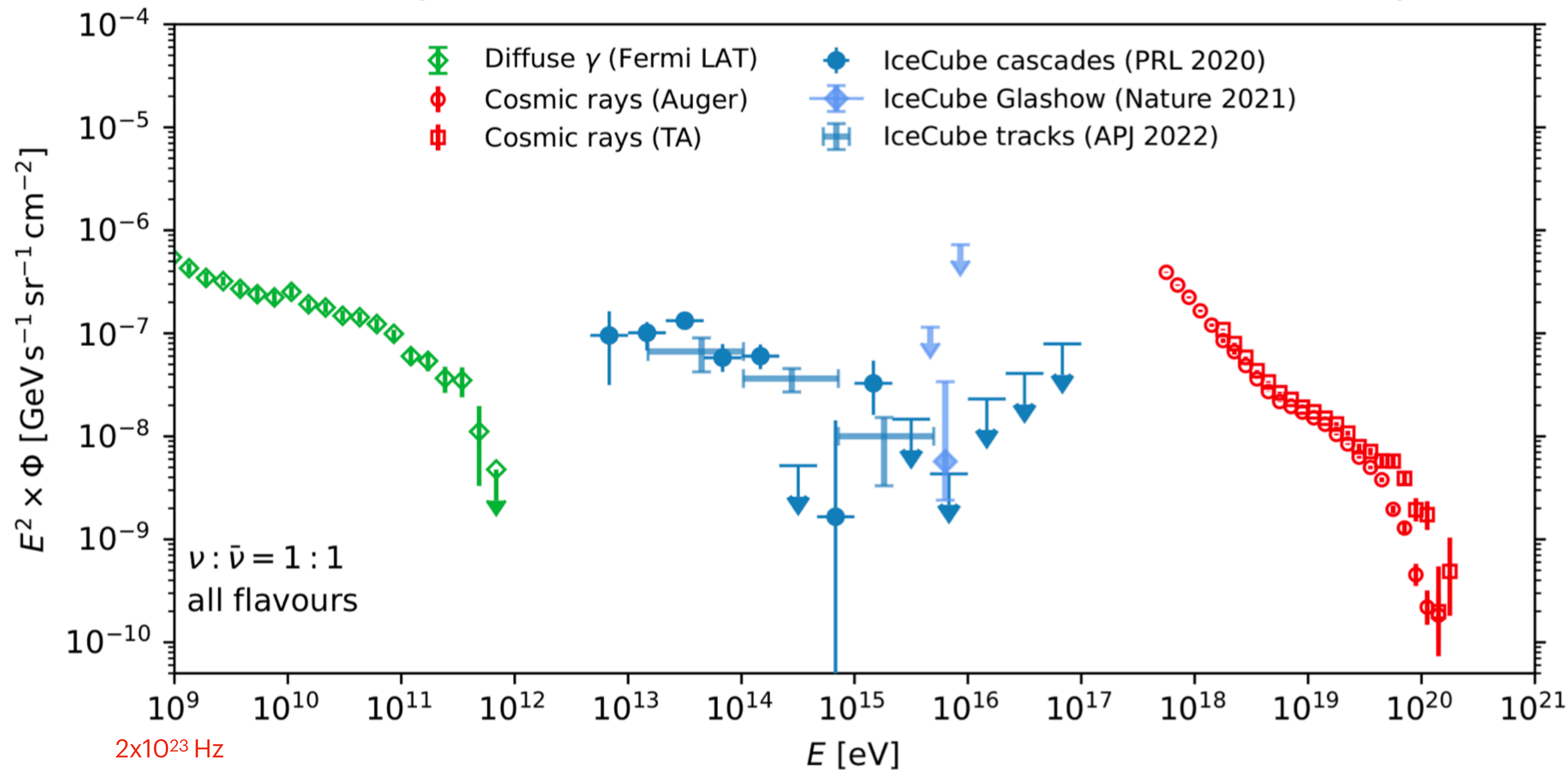
Gamma Ray



Gamma rays

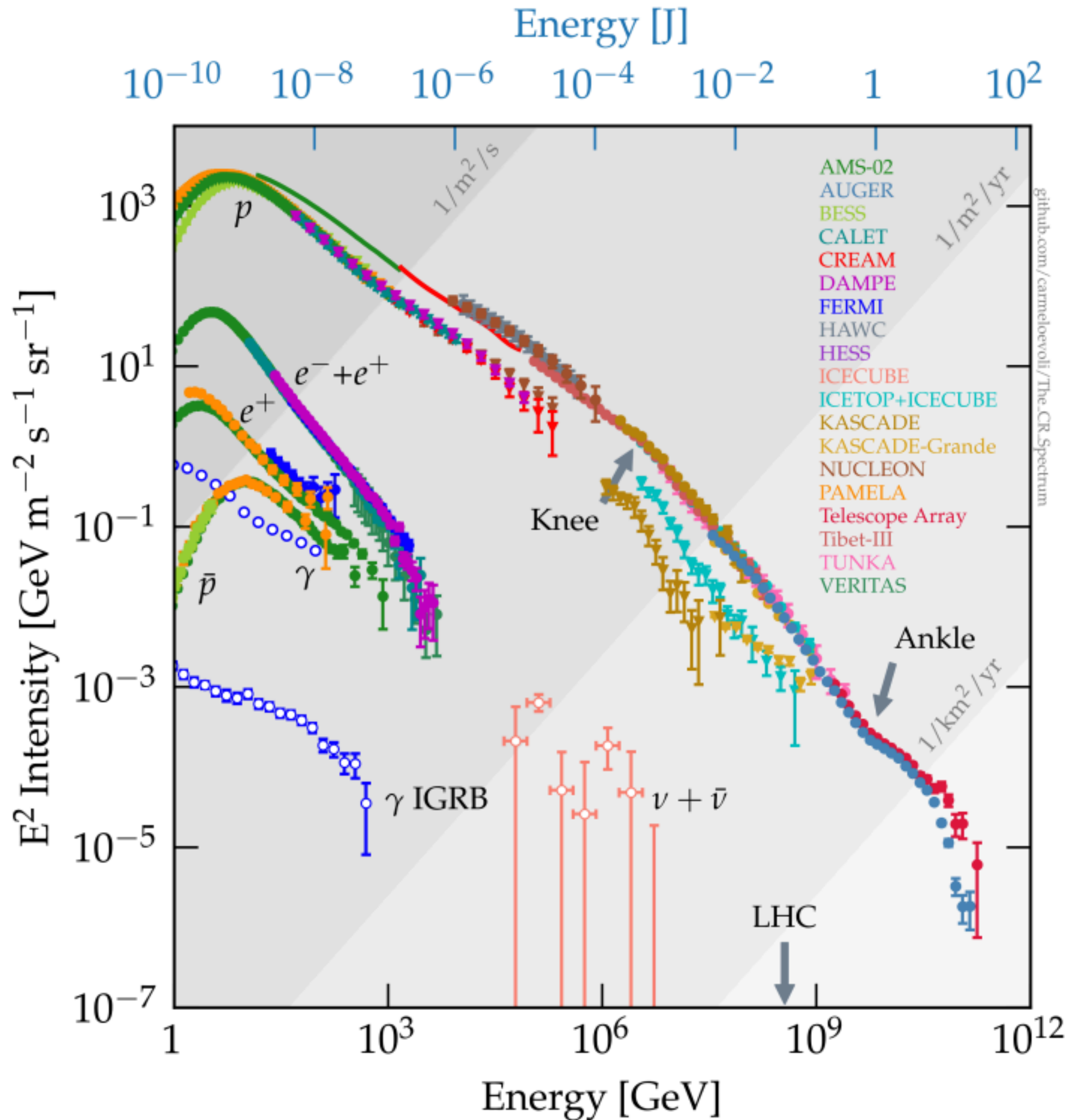
Neutrinos

Cosmic rays



Fundamental physics questions

- What is the **origin of charged particles** over this broad energy and flux range?
- How do **particles travel** in inter-stellar and inter-galactic media?
- What is the **nature of cosmic particle diffuse fluxes** (photons and neutrinos)?
- Are there **signatures of exotic physics** in astroparticle observations?



Galactic and extragalactic environments



$\lesssim 1 \text{ cm}^{-3}$
Several kpc
~230 Myr
one solar orbit
 $1 - 10 \mu\text{G}$

Extremely rarefied densities of matter
Very large distances and spatial scales
Very long timescales
Magnetised environments
Radiation fields

$$1 \text{ pc} \simeq 3.26 \text{ ly} \simeq 3.086 \times 10^{16} \text{ m}$$

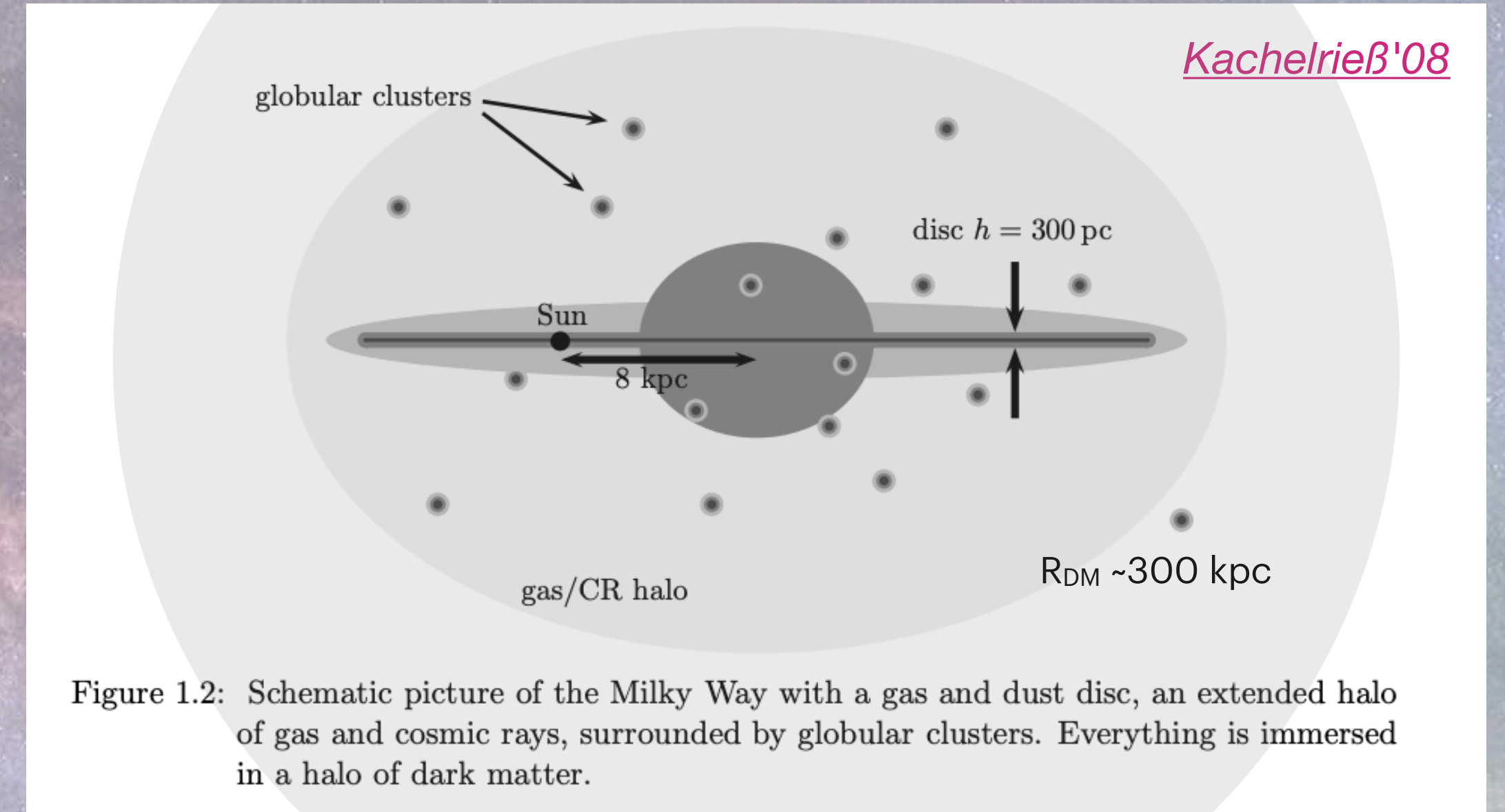
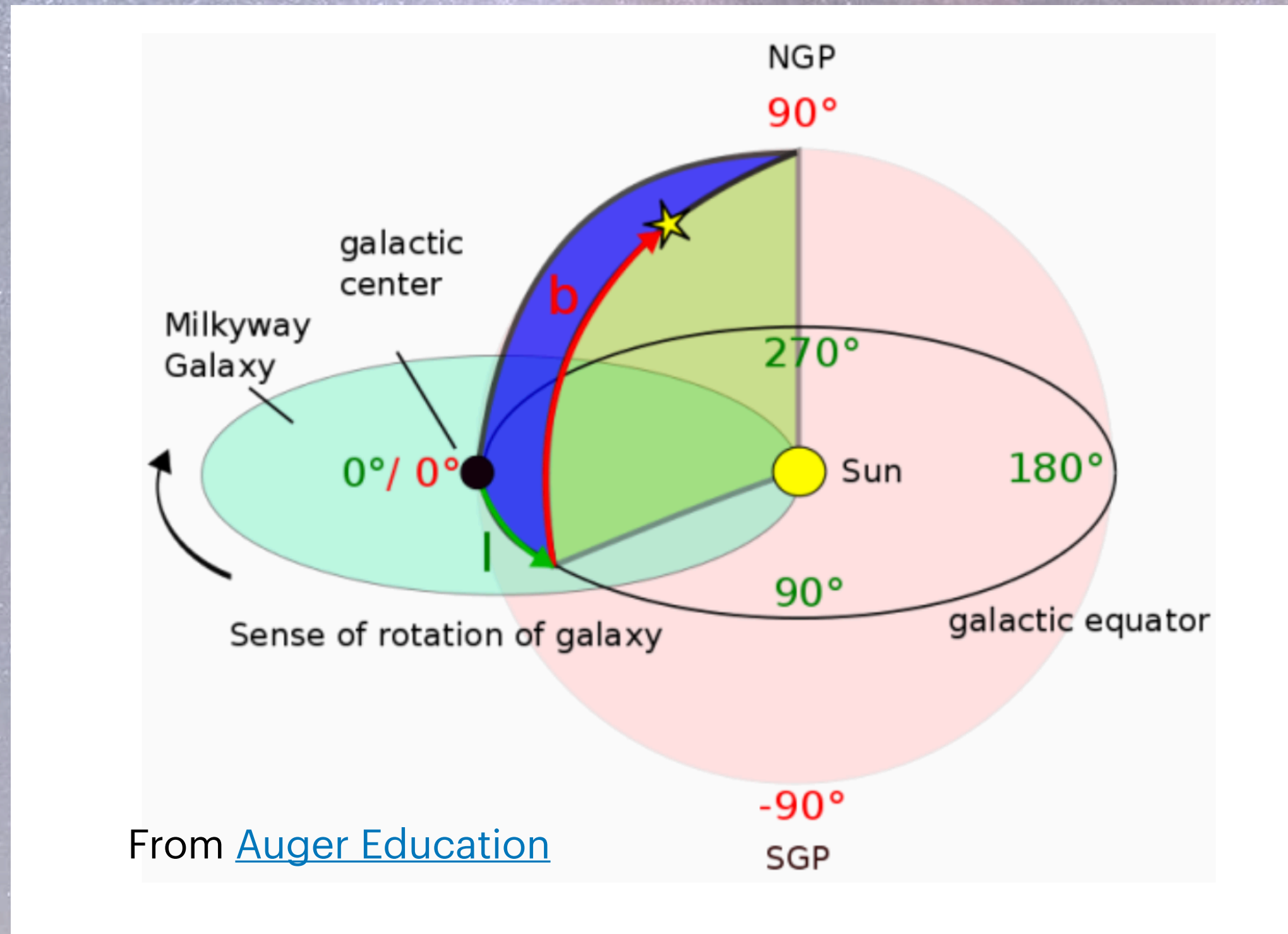
$$\lesssim 10^{-6} \text{ cm}^{-3}$$

100s Mpc or Gpc

~14 Gyr
age of the universe
 $2 - 4 \text{ nG}$



Galactic coordinate system



MILKY WAY ID:

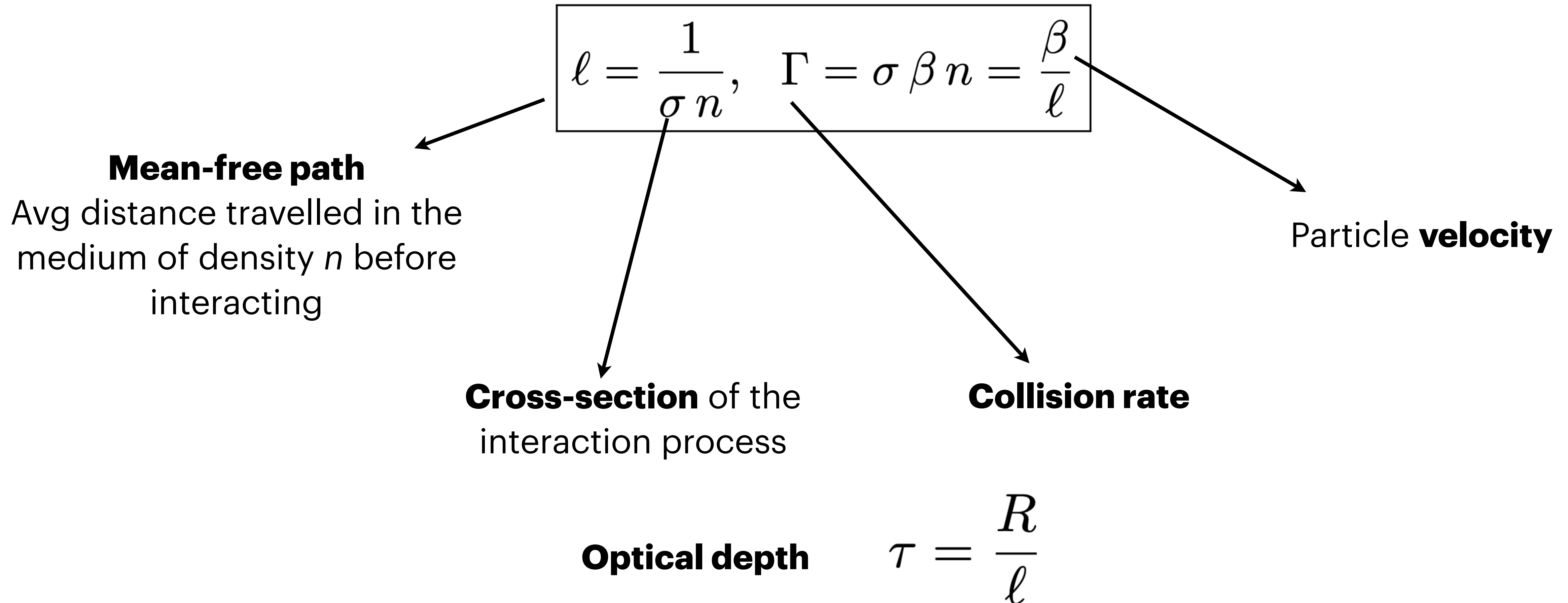
Stars $\sim 10^{11} \Rightarrow \sim 5 \times 10^{10} M_{\text{Sun}}$

Gas $\sim 10\% \Rightarrow \sim 5 \times 10^9 M_{\text{Sun}}$

Total Mass $\Rightarrow \sim 2 \times 10^{12} M_{\text{Sun}}$

Some basics: Mean free path

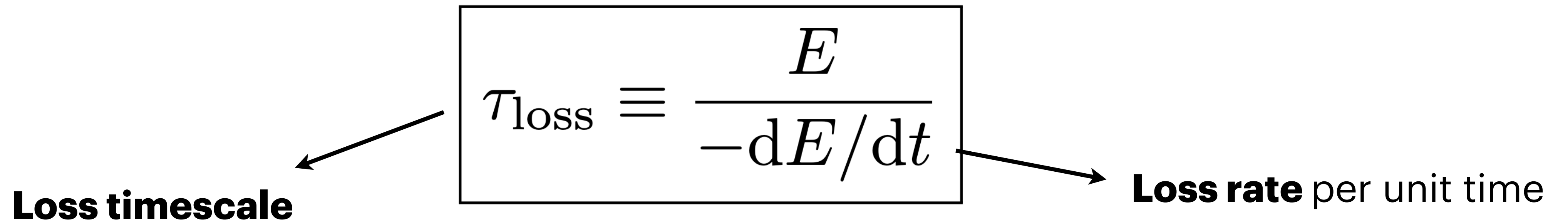
Particle interactions w/ associated **cross—section**



Some basics: Energy losses

$$\frac{d\sigma}{dE}$$

Energy transferred in a radiative process



Catastrophic losses: Particle disappears or energy transferred to other particle is seizable

$$\tau_{\text{loss}} \sim \Gamma^{-1} \quad d_{\text{loss}} \sim \text{mfp}$$

Continuous losses: Primary particle retains most of its energy and nature

$$\tau_{\text{loss}} \gg \Gamma^{-1} \quad d_{\text{loss}} \gg \text{mfp}$$

Some basics: Energy flux and emissivity

$$S = \frac{dE}{dA dt} \quad \left[\frac{\text{erg}}{\text{cm}^2 \text{s}} \right] \quad \text{Energy flux (all rays)}$$

$$I_\nu = \frac{dE}{dA dt d\Omega d\nu} \quad \left[\frac{\text{erg}}{\text{cm}^2 \text{s sr Hz}} \right] \quad \text{Specific intensity}$$

$$j_\nu = \frac{dE}{dV dt d\Omega d\nu} \quad \text{Spontaneous emission coefficient}$$

$$j = \frac{P}{4\pi}$$

For isotropic emitter
P = radiated power per unit time,
volume [erg/s/cm³]

$$\epsilon = \frac{dE}{\rho dV dt} \frac{4\pi}{d\Omega} \quad \text{Angle integrated emissivity}$$

Some basics: Energy and particle flux

$$S(\vec{r}, l, b) = \frac{1}{4\pi} \epsilon(\vec{r}, l, b) \rho(\vec{r}) ds d\Omega$$

Energy flux

Or

$$S(\vec{r}, l, b) = \frac{1}{4\pi} P(\vec{r}, l, b) ds d\Omega$$

Radiation \longrightarrow **Particle emission**

$$F(\vec{r}, l, b) = \frac{1}{4\pi} q(\vec{r}, l, b) ds d\Omega$$

Particle (local) emissivity q [# / cm³ / s] =
production rate / volume element

$$\frac{F(l, b)}{d\Omega} = \frac{1}{4\pi} \int_{l_{os}} q(\vec{r}(s), l, b) ds$$

$$\frac{F(E, l, b)}{d\Omega dE} = \frac{1}{4\pi} \int_{l_{os}} \tilde{q}(E, \vec{r}(s), l, b) ds$$

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Particles and radiation: production mechanisms

Collisional effects & radiative processes

- Photons interactions w/ radiation (**MW connection**) *Photons*
- Leptons interactions w/ radiation (**MW/MM connection**) *Photons*
 - ➔ Inverse Compton scattering
 - ➔ Synchrotron radiation
- Leptons interactions w/ matter (Ionisation/bremsstrahlung) *Photons*
- Hadrons interactions w/ radiation (**MW/MM connection**, extragalactic) *Photons, neutrinos*
- Hadrons interactions w/ matter (**MW/MM connection**, Galactic) *Photons, neutrinos*

Collisional effects & radiative processes

* Galactic CR

- prop. Timescale ~ 10 Myr
- $n_{\text{matter}} \sim 1/\text{cm}^3$

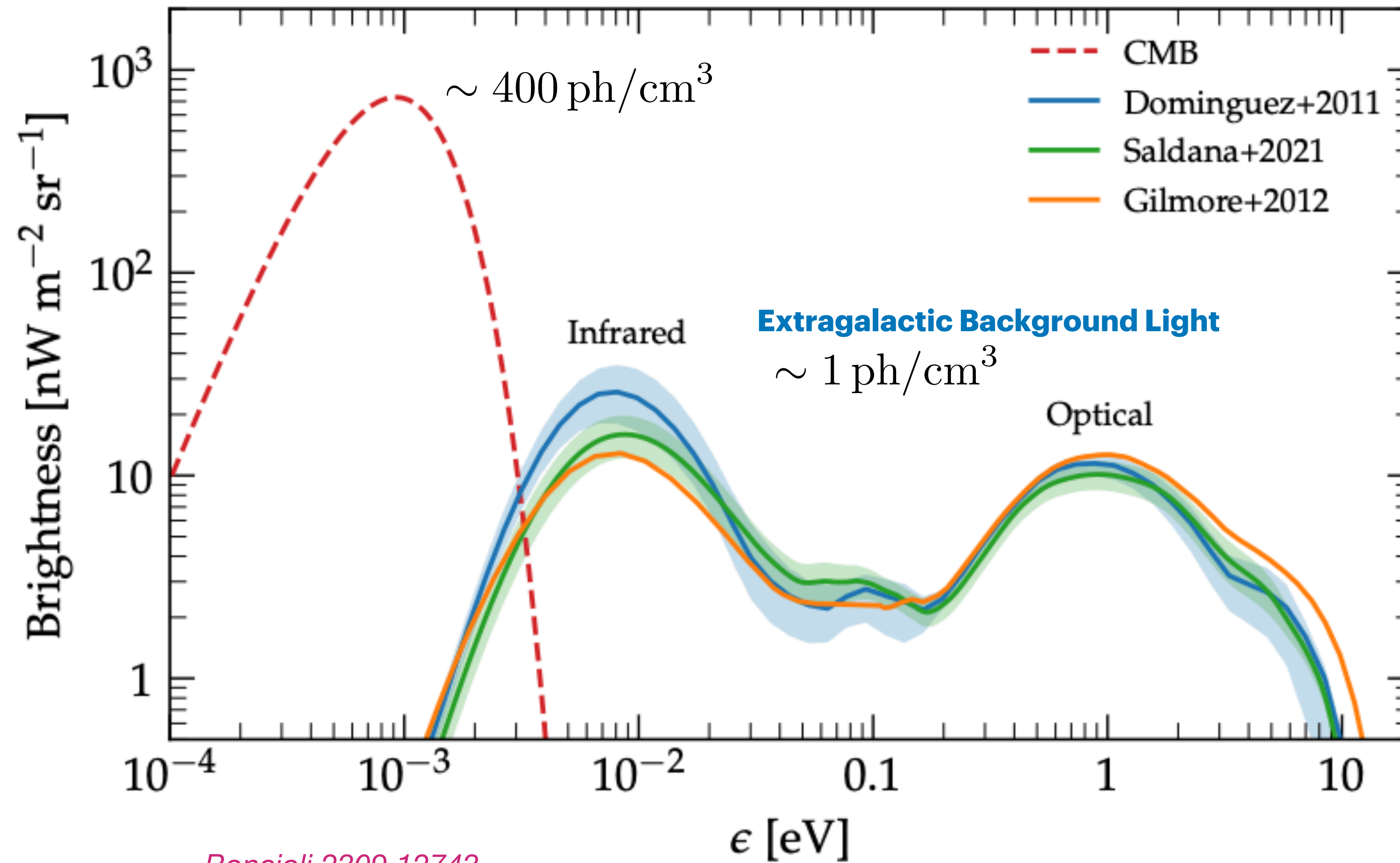
- CR proton $E = \gamma m_p$ sees
photons w/ $E_{\text{ph}} \sim \gamma E$ photon energy
in its rest frame \rightarrow preventing
inelastic processes

* Extragal. CR

- prop. Timescale ~ 10 Gyr
- $n_{\text{matter}} \sim 10^{-7} n_{\text{matter}}^{\text{gal}}$
- $\gamma \gtrsim 10^8 \rightarrow$ interesting kinemat.
coll thresholds open up
 \rightarrow inelastic processes

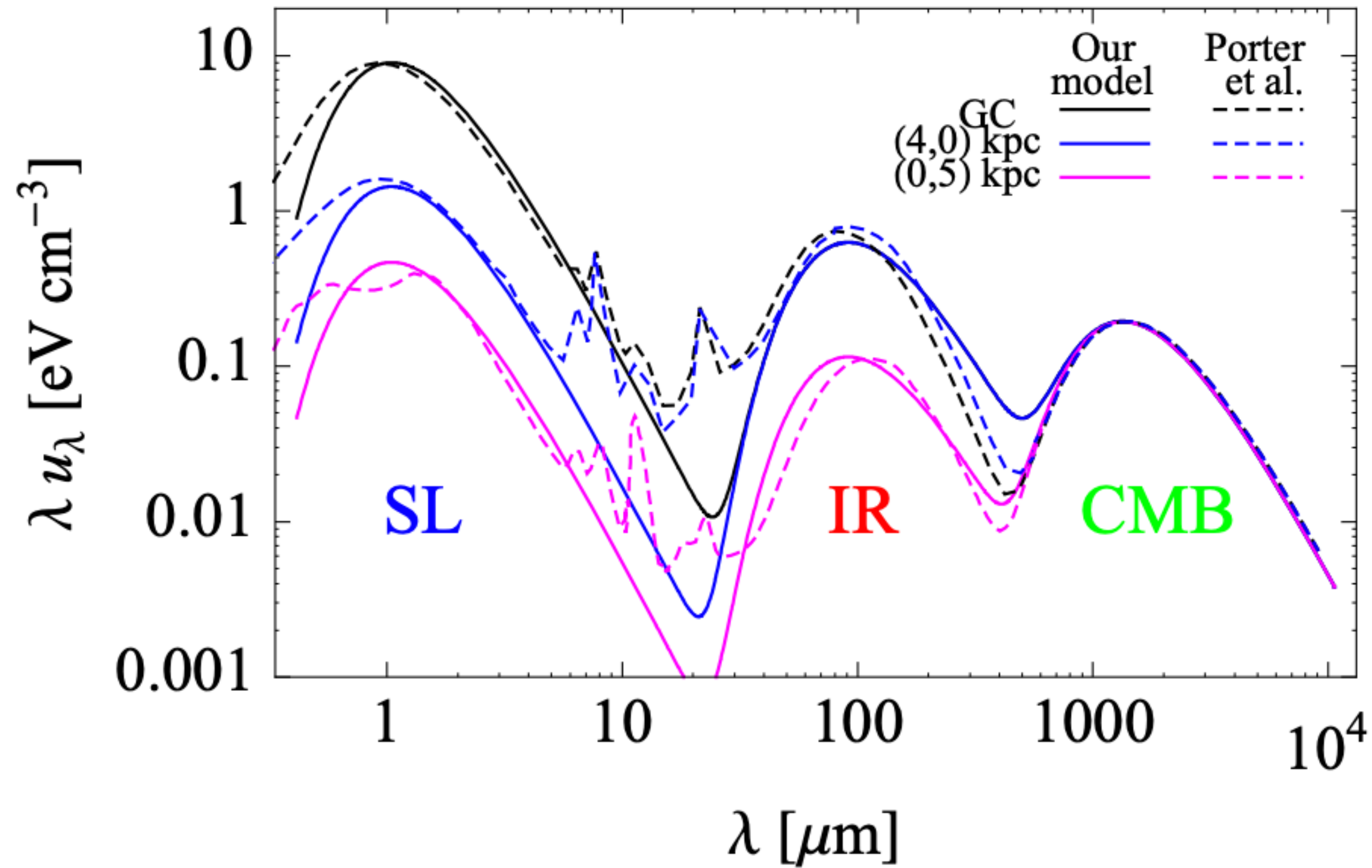
Extragalactic radiation backgrounds

$$\nu I_\nu \equiv \frac{c}{4\pi} \epsilon^2 \frac{dn}{d\epsilon}$$



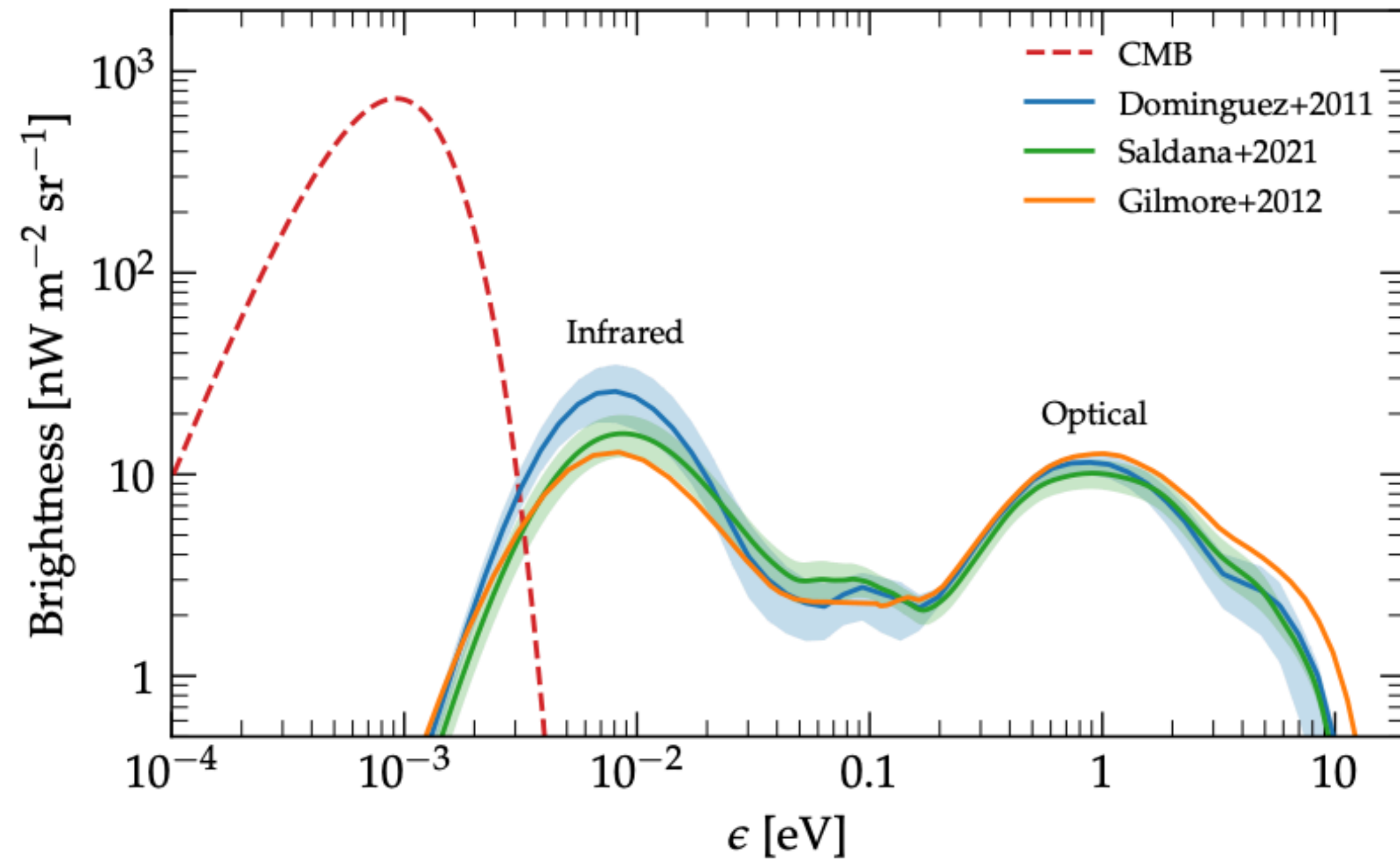
Boncioli 2309.12743

Galactic radiation backgrounds



Cirelli & Panci 0904.3830

Pair production



Pair-production



From invariance of s , COM energy

$$E'_\gamma \geq \frac{m_e^2}{\epsilon'}$$

in cosmological coving frame

EBL absorption of VHE photons

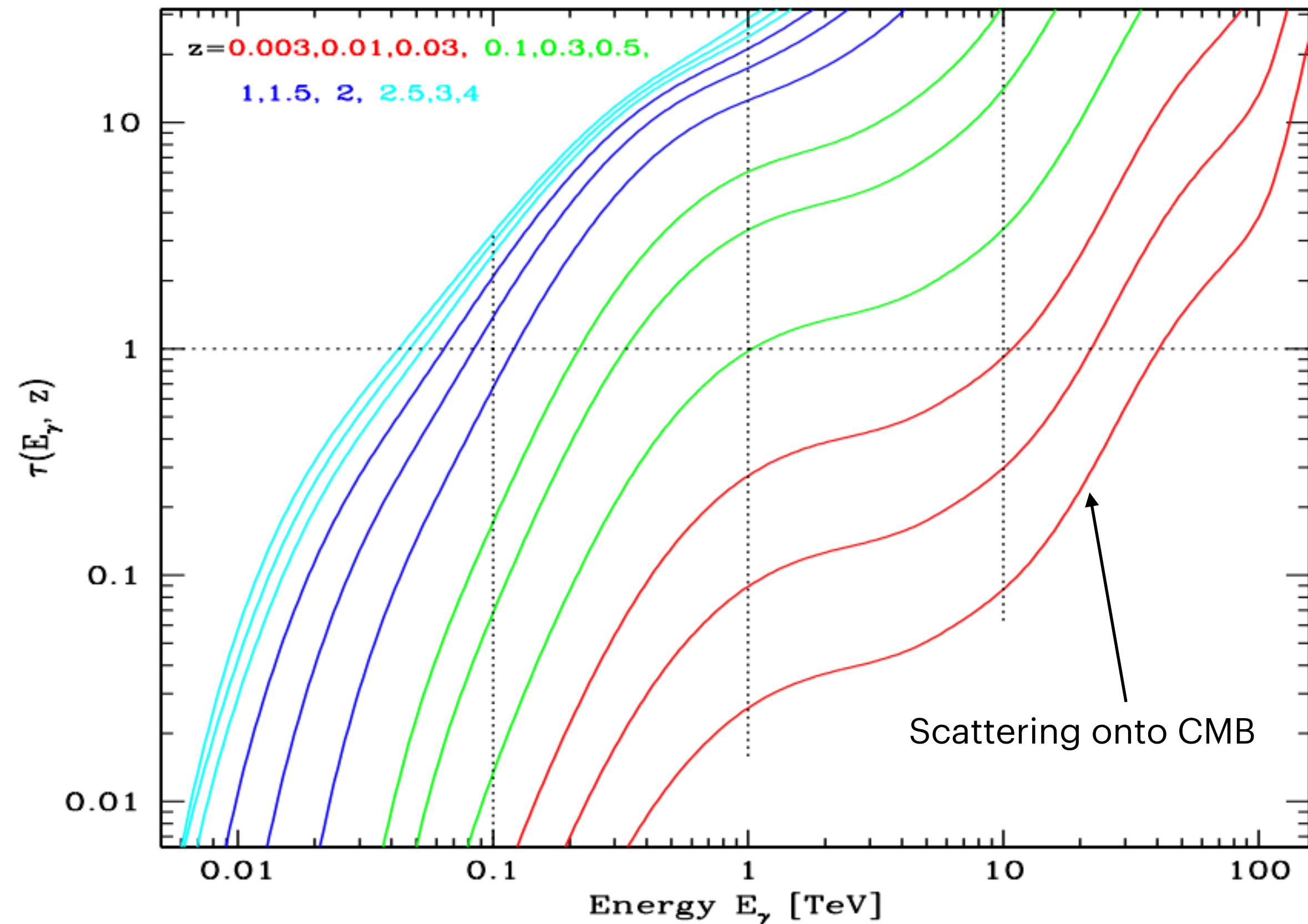
Optical depth \sim # interactions along the l.o.s

$$\tau(E_\gamma, z_s) = \mathcal{F}(n_\gamma(\epsilon', z), d_L(z), \sigma_{\gamma\gamma}(E_\gamma, \epsilon'))$$

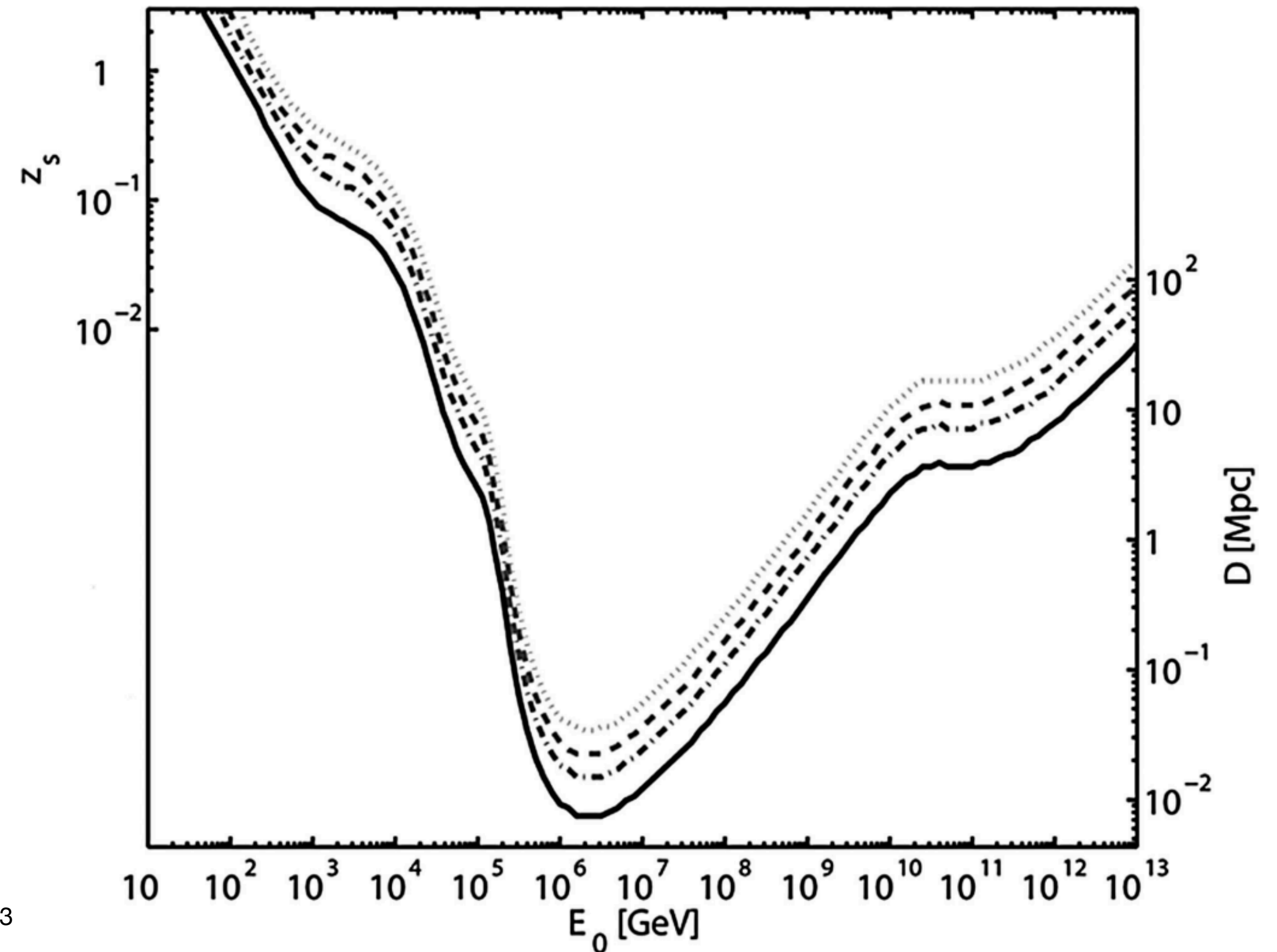
Gamma-ray cosmic horizon

$$\tau(E_\gamma, z_s) = 1$$

Franceschini Universe'21



De Angelis+MNRAS'13

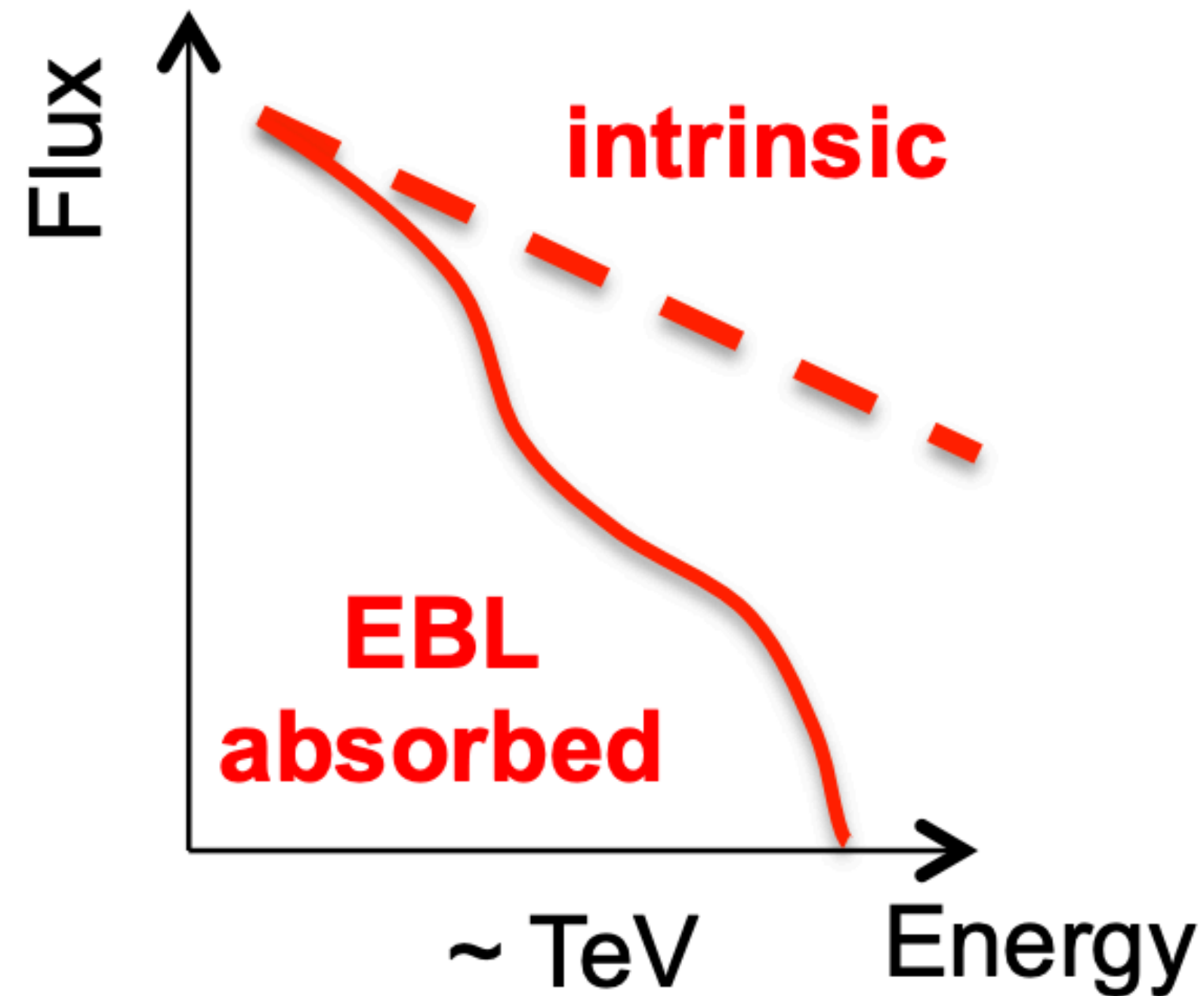


EBL absorption of VHE photons

Photon survival probability

$$P_{\gamma\gamma}(E_\gamma, z_s) \equiv e^{-\tau(E_\gamma, z_s)}$$

Also called transfer function $T_\gamma(E, L)$



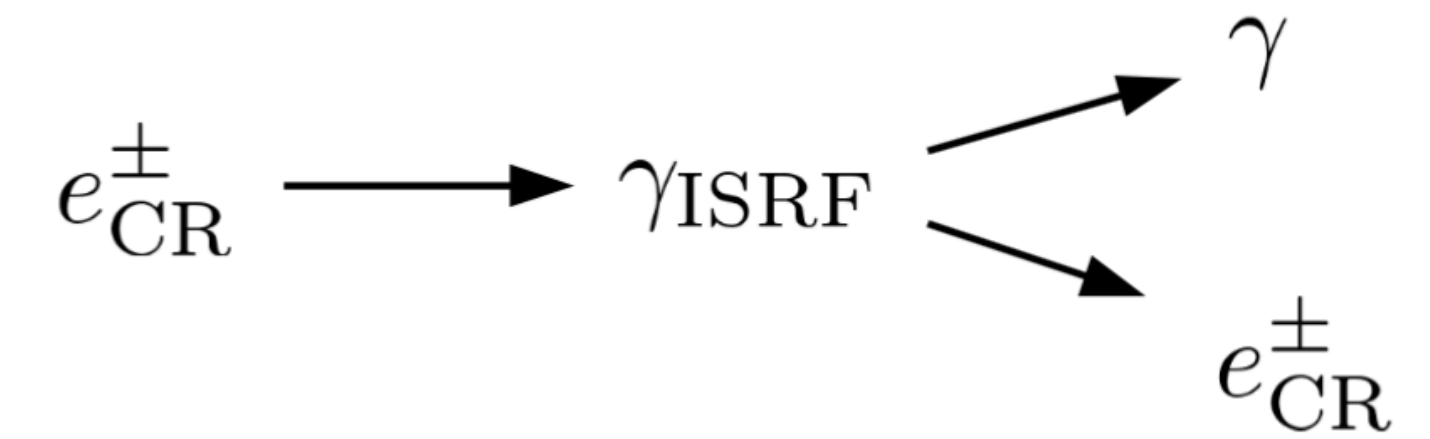
$$\Phi(E_\gamma, z_s) = \Phi(E_\gamma) \times e^{-\tau(E_\gamma, z_s)}$$



Differential (in energy) photon flux shows a characteristic EBL cutoff at about

$$E_{\gamma, \text{cutoff}}(z_s) \sim 800 (1 + z_s)^{-2.4} \text{ GeV}$$

Inverse Compton scattering

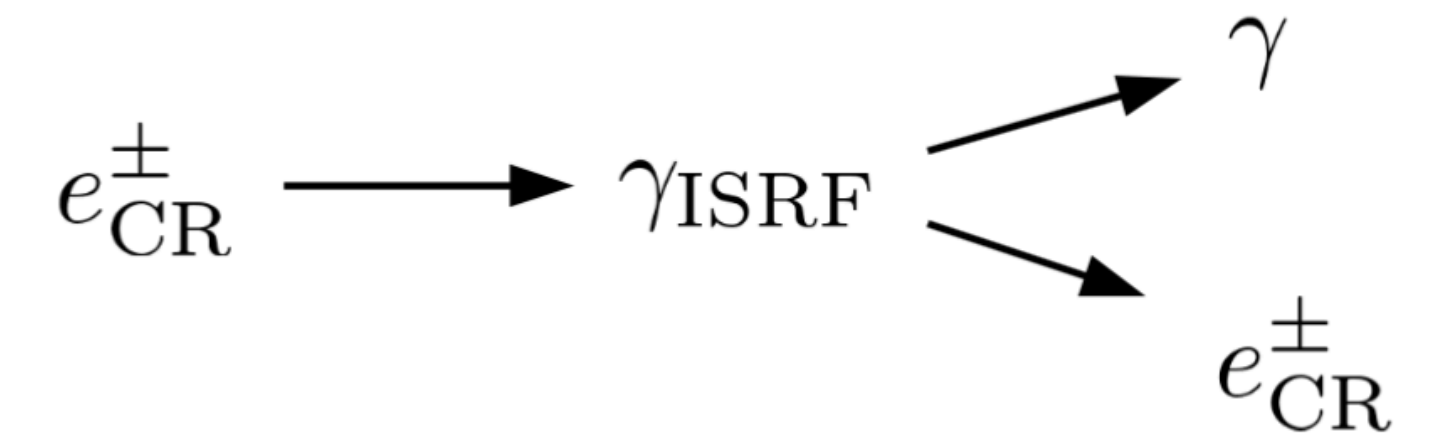


Classical treatment: Interaction of an e.m. wave w/ charged particles

- Electron subject to incoming e.m. wave is accelerated

$$\begin{aligned} \omega \ll 1 \\ \text{Lorentz force} = \vec{a} = \frac{q=e}{m} (\vec{E} + \vec{v} \times \vec{B}) \\ a = \frac{e}{m_e} E \quad \text{acc. along the direction of the field } \vec{E} \end{aligned}$$

Inverse Compton scattering



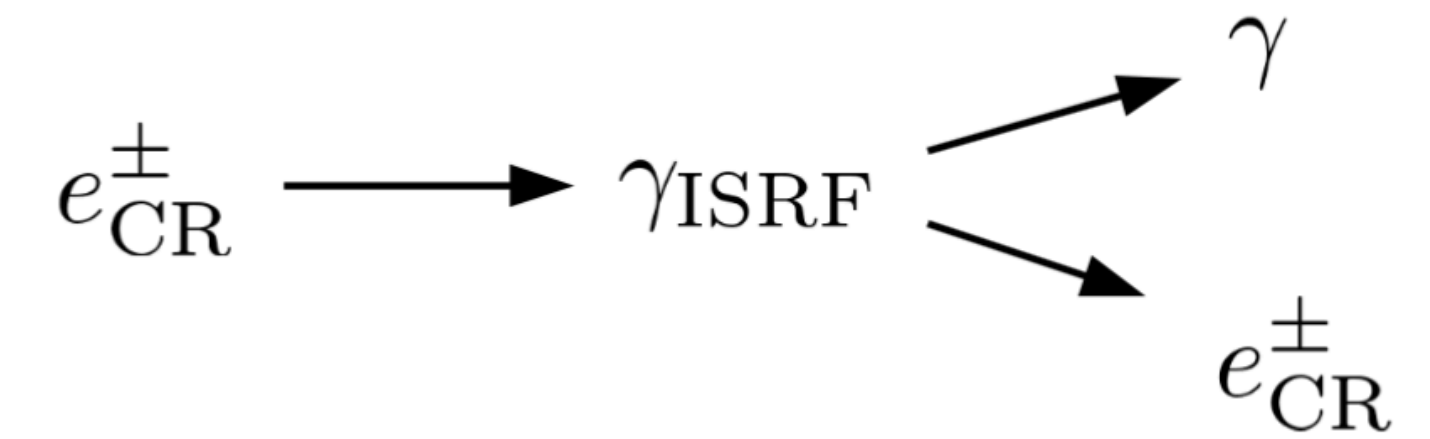
Classical treatment: Interaction of an e.m. wave w/ charged particles

- Power re-emitted by the electron is described by the **Larmor formula**

A handwritten equation on a grid background. The equation is $\langle P \rangle = \frac{2}{3} \frac{e^2}{m_e^2 c^3} |\dot{p}|^2$. Below the $\langle P \rangle$ term, there is a downward-pointing arrow and the text "Time averaged".

[NB: Inefficient for nuclei]

Inverse Compton scattering



Classical treatment: Interaction of an e.m. wave w/ charged particles

- **Cross-section** in the **Thomson limit**

$$\epsilon_{\text{ph}} \ll m_e$$

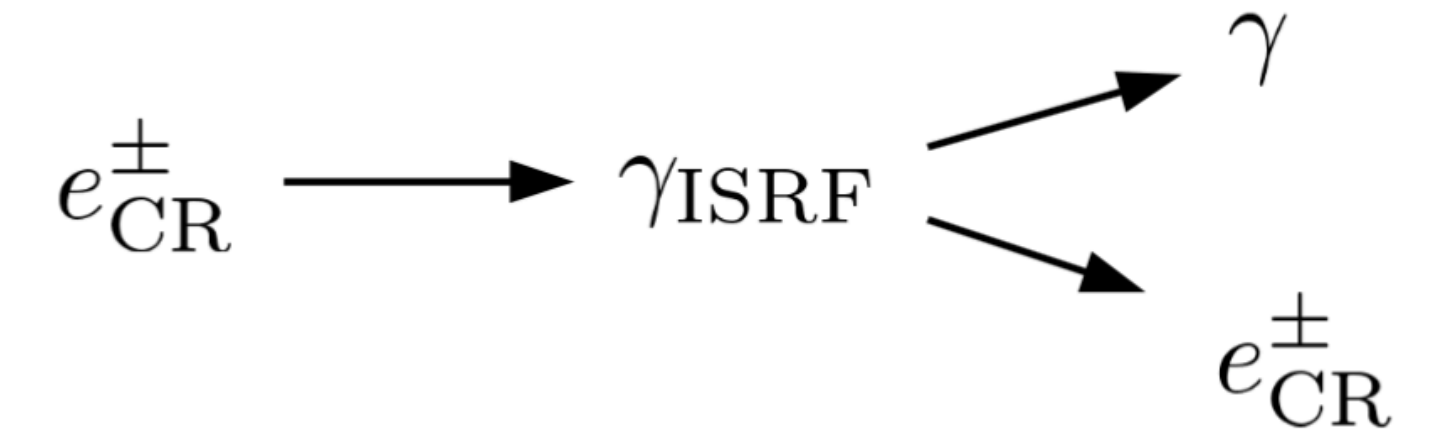
Handwritten text on a grid background: $\sigma \equiv \sigma_T$ Thomson cross-section $\sim 6.65 \cdot 10^{-25} \text{ cm}^2$

- **Energy-loss rate** (per particle) or **total IC power** is

Handwritten equations on a grid background: $P_{\text{IC}} = \frac{4}{3} \sigma_T c \beta^2 \gamma^2 u_{\text{rad}}$
 $[P_{\text{IC}}] = \text{erg/s}$

[NB: Same for synchrotron with magnetic energy density]

Inverse Compton scattering



Classical treatment: Interaction of an e.m. wave w/ charged particles

- **Energetics** of outgoing photons

$$\bar{E}_{\gamma} \approx \gamma^2 \varepsilon \quad ; \quad \varepsilon E_e \ll m_e^2$$

$$\bar{E}_{\gamma} \approx 30 \left(\frac{\varepsilon}{\text{eV}} \right) \left(\frac{E_e}{\text{GeV}} \right)^2 \text{ MeV}$$

typical photon fields in the galaxy

}	CMB	$\varepsilon \sim 2.35 \cdot 10^{-4} \text{ eV}$
	IR	$\varepsilon \sim 3.45 \cdot 10^{-3} \text{ eV}$
	OPT	$\varepsilon \sim 0.3 \text{ eV}$

$$\Rightarrow \bar{E}_{\gamma}^{\text{CMB}} \sim 7 \cdot 10^{-3} \text{ MeV} \sim \text{KeV}$$

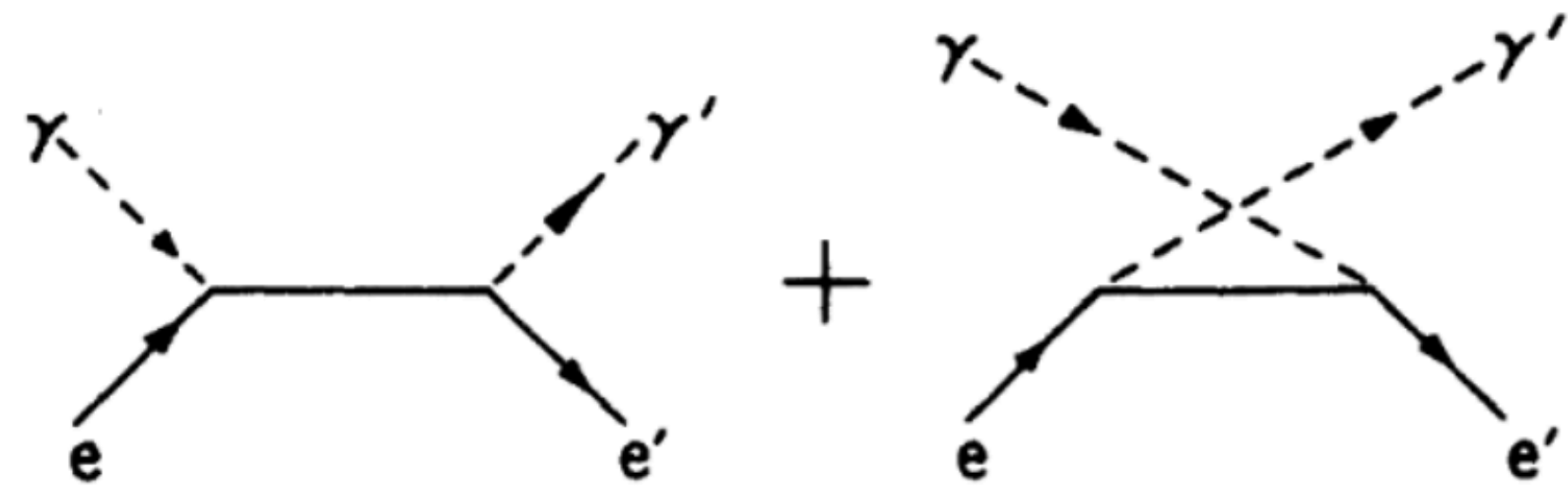
$$\bar{E}_{\gamma}^{\text{IR}} \sim 0.1 \text{ MeV}$$

$$\bar{E}_{\gamma}^{\text{OPT}} \sim 9 \text{ MeV}$$

scattering probability: $\sigma(\varepsilon)$

Inverse Compton scattering

Fully relativistic treatment

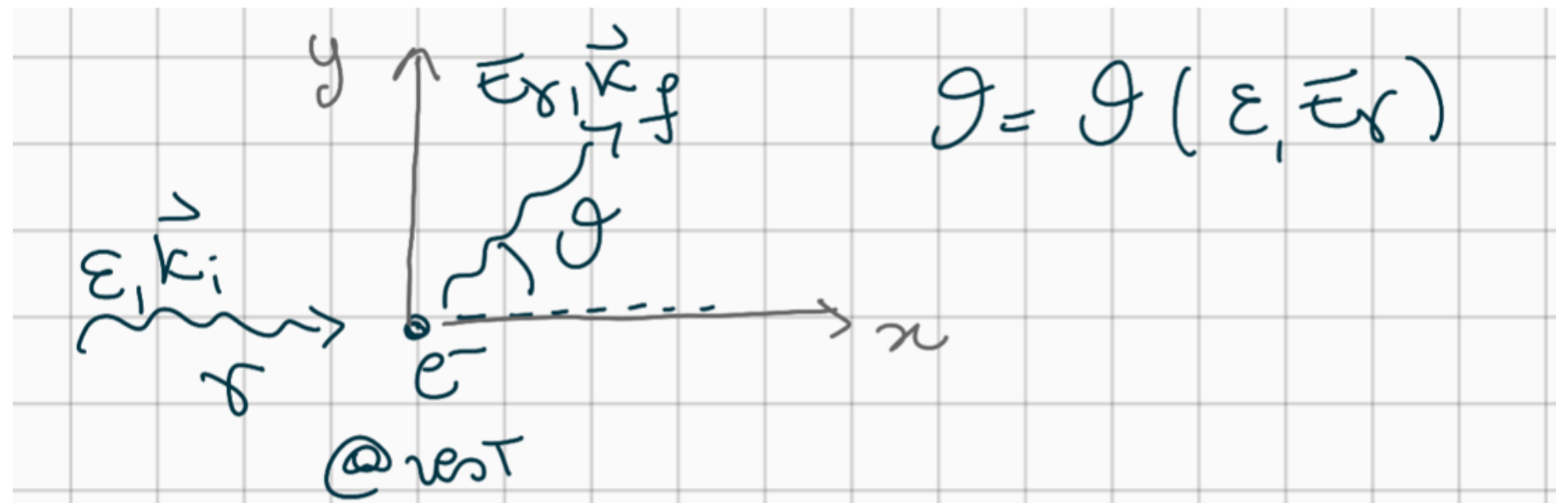


Klein-Nishina process: tree-level electron photon scattering in QED

$$\frac{d\sigma}{d\Omega} = \frac{3}{16\pi} \sigma_T \left(\frac{E_r}{\varepsilon}\right)^2 \left[\frac{\varepsilon}{E_r} + \frac{E_r}{\varepsilon} - 2m^2 \mathcal{G} \right]$$

$$\sigma = 2\pi \int_0^\pi \frac{d\sigma}{d\Omega} \sin\vartheta d\vartheta \equiv \sigma_{KN}(x)$$

$$x \equiv \frac{\varepsilon}{m_e}$$



Galactic Inverse Compton emission

Dominant production mechanism of photons ~MeV - GeV energies

$$\begin{aligned} \vec{q}(\vec{\tau}, \bar{\epsilon}_\gamma) &= 4\pi \int d\varepsilon \frac{dn}{d\varepsilon}(\varepsilon, \vec{\tau}) \times \\ &\quad \parallel \\ &\quad \text{local emissivity} \times \int d\bar{\epsilon}_e \frac{d\sigma_{ic}}{d\bar{\epsilon}_\gamma}(\bar{\epsilon}_e, \varepsilon, \bar{\epsilon}_\gamma) \times \\ &\quad \left[\frac{1}{\text{cm}^3 \text{ erg s}} \right] \times \Phi_e(\bar{\epsilon}_e, \vec{\tau}) \end{aligned}$$

Galactic Inverse Compton emission

Dominant production mechanism of photons ~MeV - GeV energies

$$\Phi(\bar{\epsilon}, \bar{\epsilon}_\gamma) = \frac{1}{\bar{\epsilon}_\gamma} \int \mathcal{P}_{ic}(\epsilon, \bar{\epsilon}_\gamma, \bar{\epsilon}_e, z) \Phi_e d\bar{\epsilon}_e$$
$$[\frac{1}{s}] = \mathcal{P}_{ic} = \frac{dP_{ic}}{d\bar{\epsilon}_\gamma} \text{ differential IC power}$$

Galactic Inverse Compton emission

Dundovic+ A&A'22

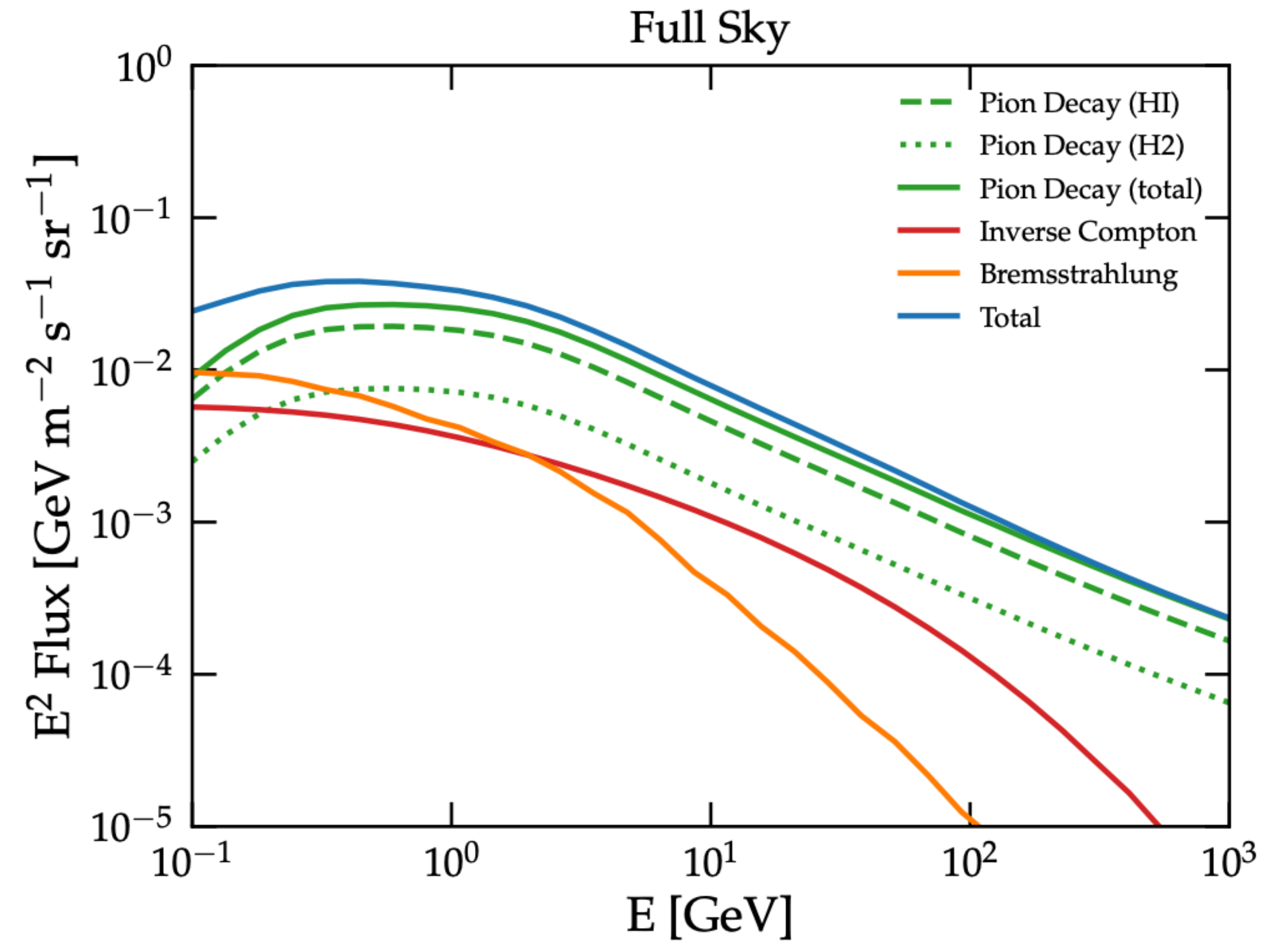
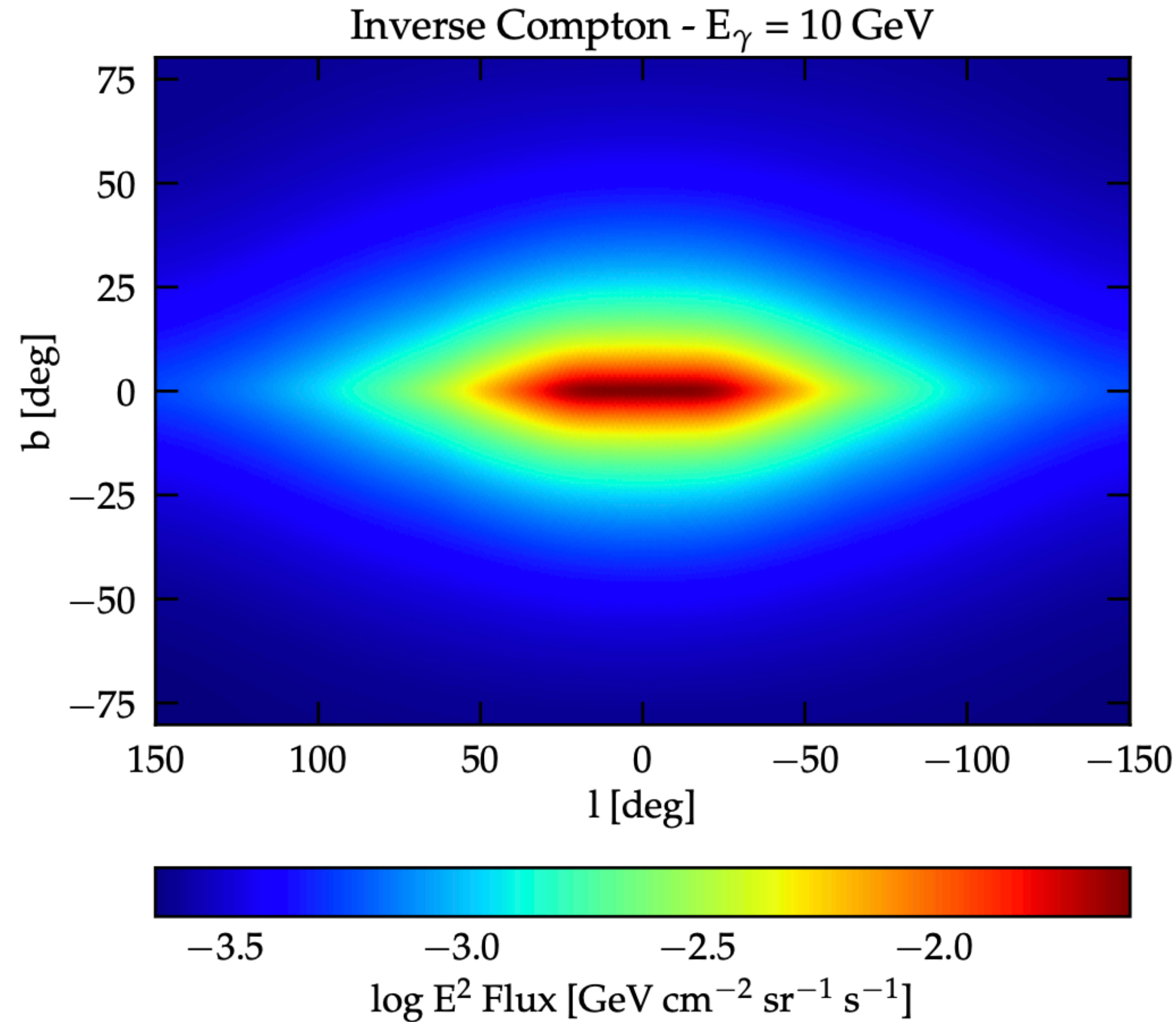
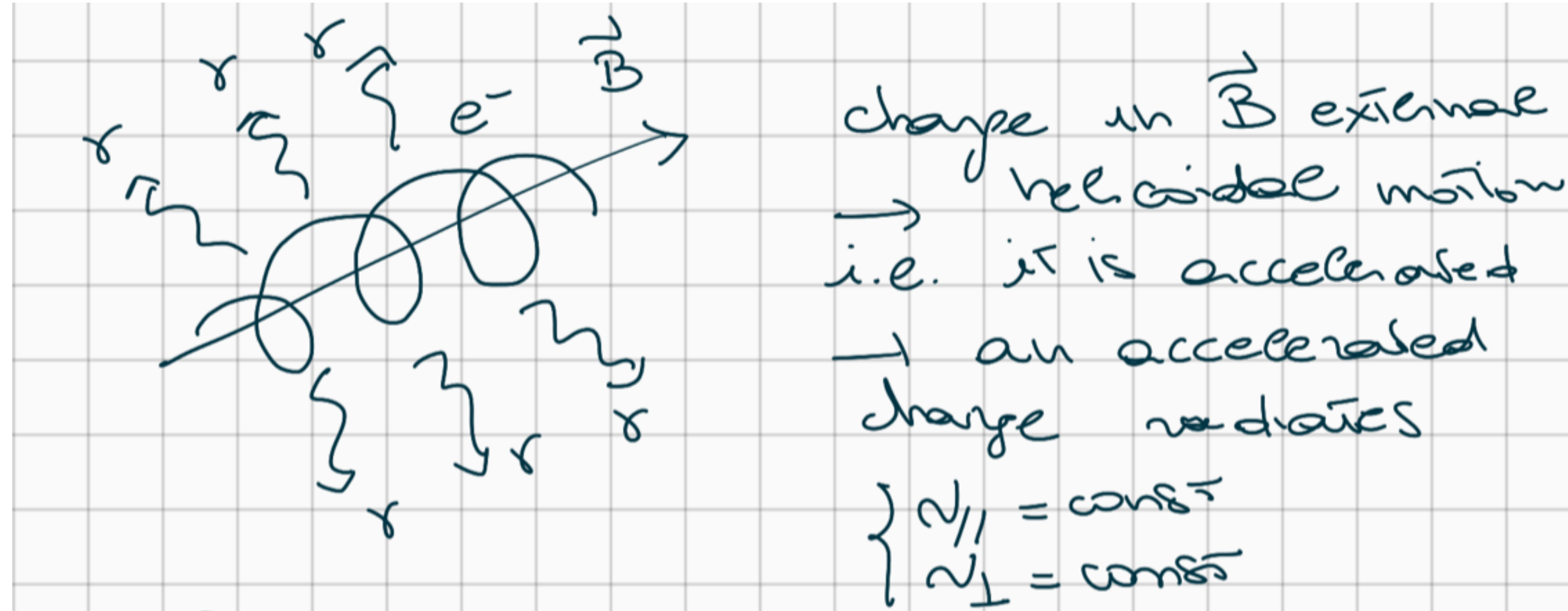


Fig. 7. Cartesian projection in Galactic coordinates of the IC gamma-ray flux at $E_\gamma = 10$ GeV.

Synchrotron radiation



$$r_L = \frac{m\gamma v_{\perp}}{qB}$$

Larmor radius [pc]

$$-\frac{dE}{dt} = \frac{4}{3}\sigma_T c u_B \gamma^2$$

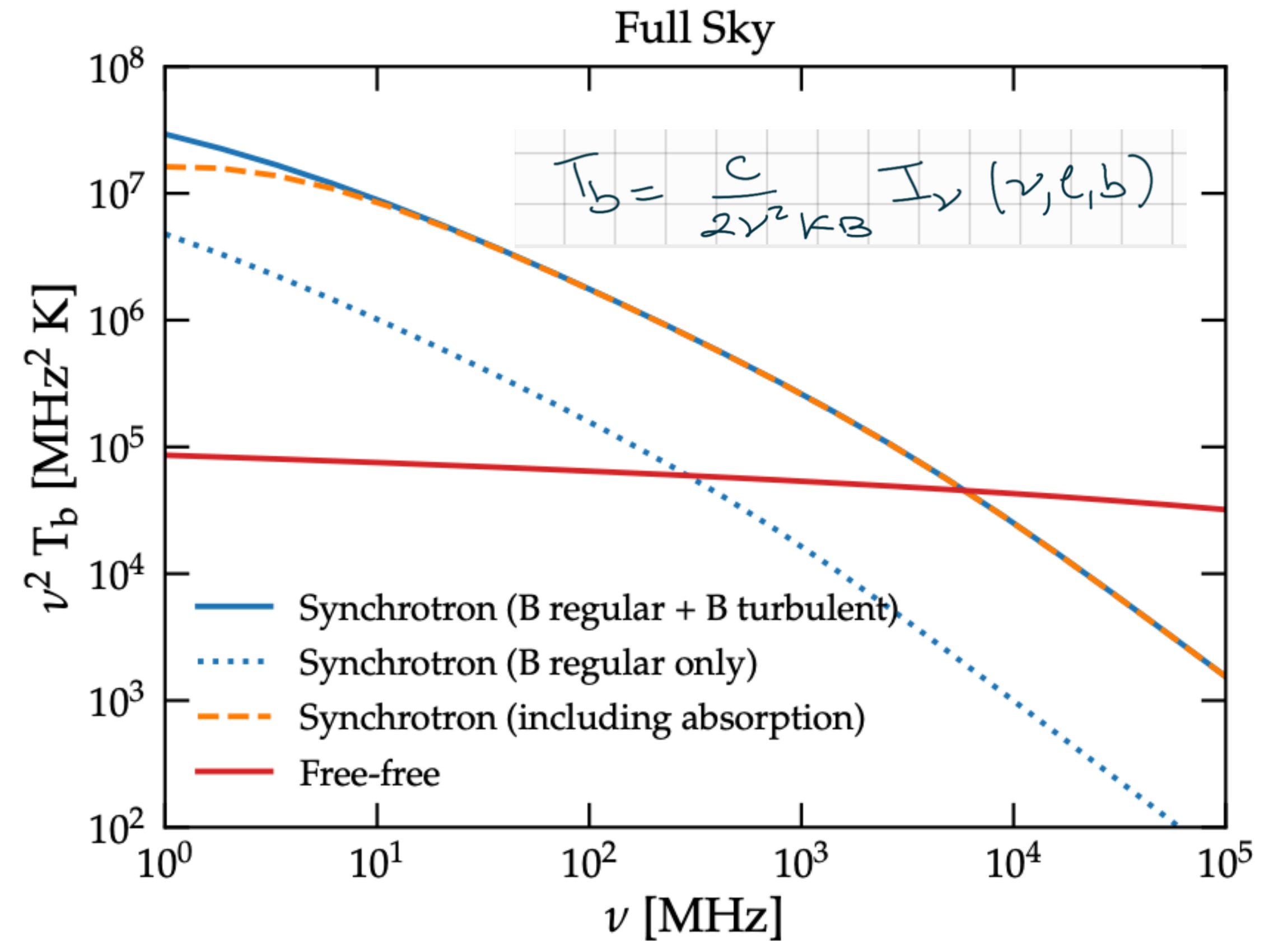
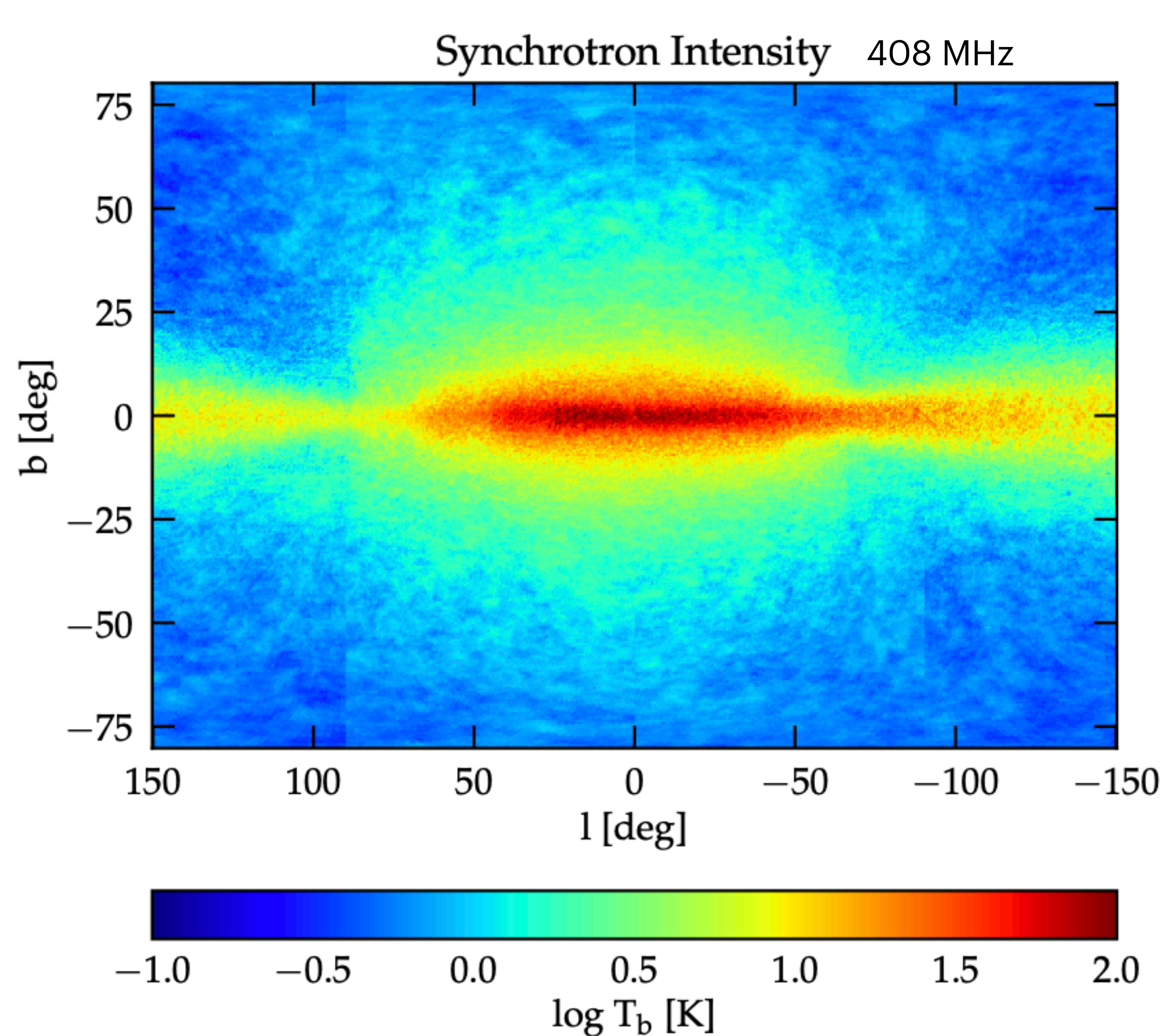
Average energy-loss rate

$$u_B = \frac{B^2}{2\mu_0}$$

$$E_s \simeq 500 \mu\text{eV} \frac{B}{\mu\text{G}} \left(\frac{E_e}{\text{GeV}} \right)^2$$

Energy of radiation for a CR electron

Galactic synchrotron radiation



Dundovic+ A&A'22

Peculiarity: The synchrotron radiation of ultra-relativistic electrons in a uniform magnetic field is expected to be **highly polarised**

Hadrons interactions with radiation

- **Elastic scattering** is suppressed

$$\sigma_{\gamma N} \propto \frac{Z^4}{A^2} \sigma_T$$

- **Inelastic scattering** is the most efficient channel of E-loss for UHE CRs

Collisional effect, catastrophic E loss mechanism

$$\text{Inelasticity: } k \sim \left\langle \frac{E_{\text{in}} - E_{\text{out}}}{E_{\text{in}}} \right\rangle$$

Pair production

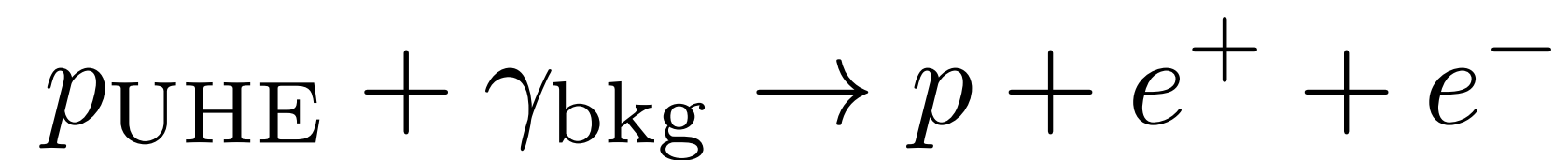
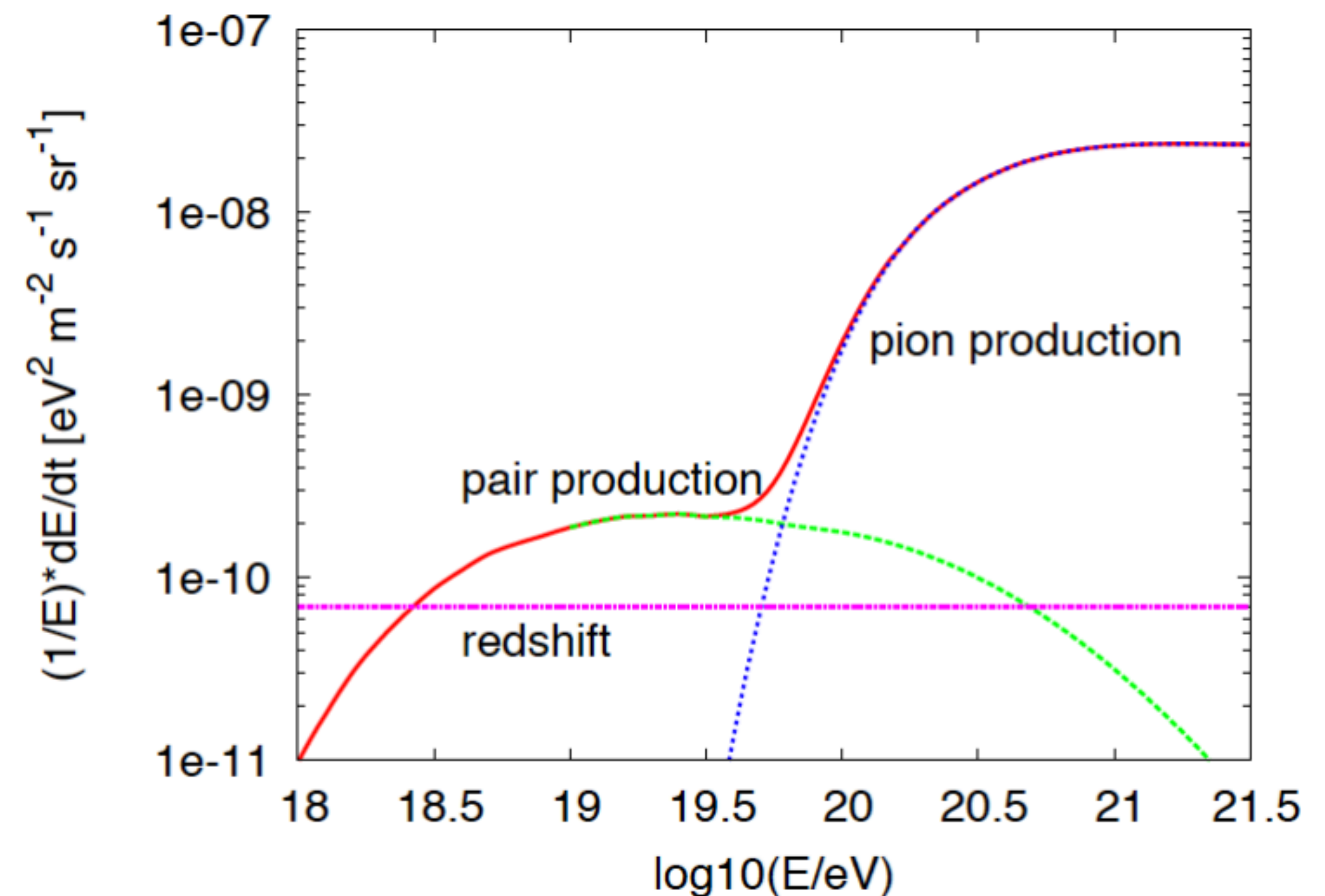
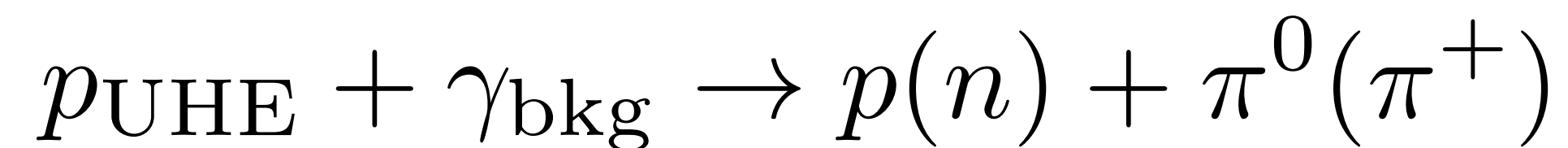


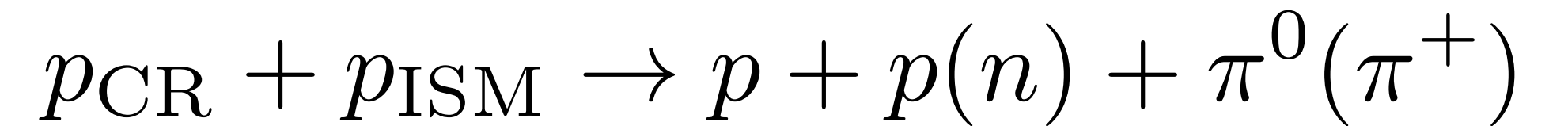
Photo-pion production



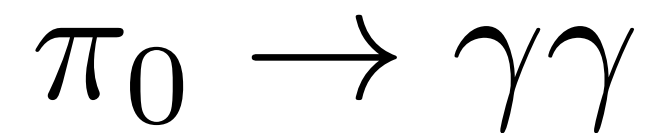
Hadrons interactions with matter

Inelastic proton-proton collisions

- Main E-loss mechanism for Galactic CR protons



- Main production mechanism of photons at GeV energies: **Neutral pion decay**

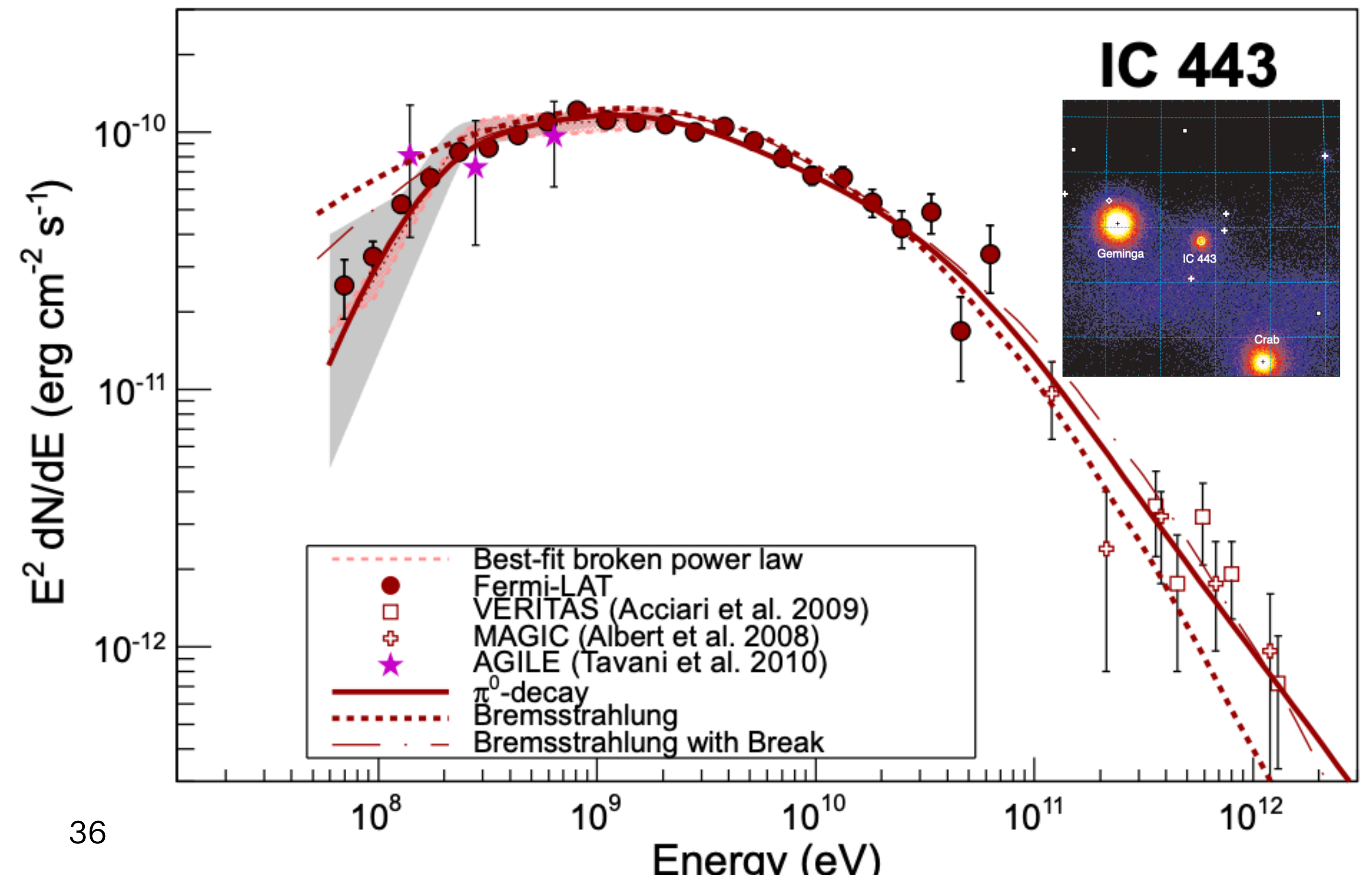


$$E_{\gamma} = \frac{m_{\pi^0}}{2} \approx 67.5 \text{ MeV}$$

Flat spectrum between E_{min} , E_{max} , always centred at the same energy regardless of pion energy

Pion bump

Ackermann+ Science'13



Galactic photons from neutral pion decay

$$\begin{aligned} \tilde{Q}_{\pi^0}(\bar{E}_{\pi}) &= 4\pi n_H \int d\bar{E}_p \delta(\bar{E}_{\pi} - k_{\pi} \bar{E}_p) \\ &\quad \times \frac{d\sigma_{pp \rightarrow \pi^0}}{d\bar{E}_{\pi}}(\bar{E}_p) \bar{\Phi}_p(\bar{E}_p) \\ &= 4\pi n_H \frac{d\sigma_{pp \rightarrow \pi^0}}{d\bar{E}_{\pi}}\left(\frac{\bar{E}_{\pi}}{k_{\pi}}\right) \cdot \bar{\Phi}_p\left(\frac{\bar{E}_{\pi}}{k_{\pi}}\right) \end{aligned}$$

$\tilde{Q}_{\pi^0}(\bar{E}_{\pi})$ is the **pion emissivity**, # pions / $dV d\bar{E} dt$.

Neutral pion spectrum has the same shape of the parent protons, shifted at lower energies

$$E_{\pi^0} \sim 0.17 E_p$$

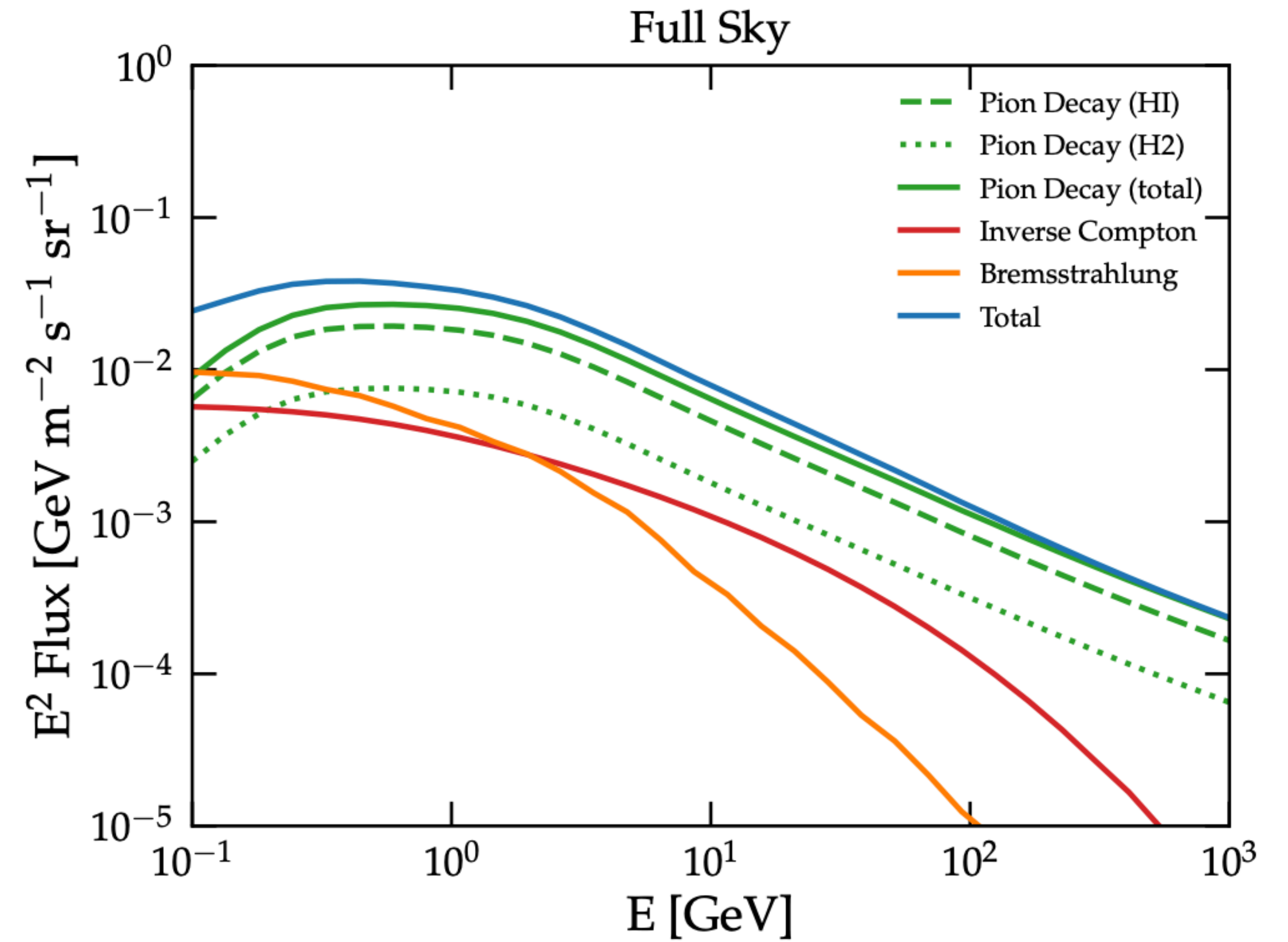
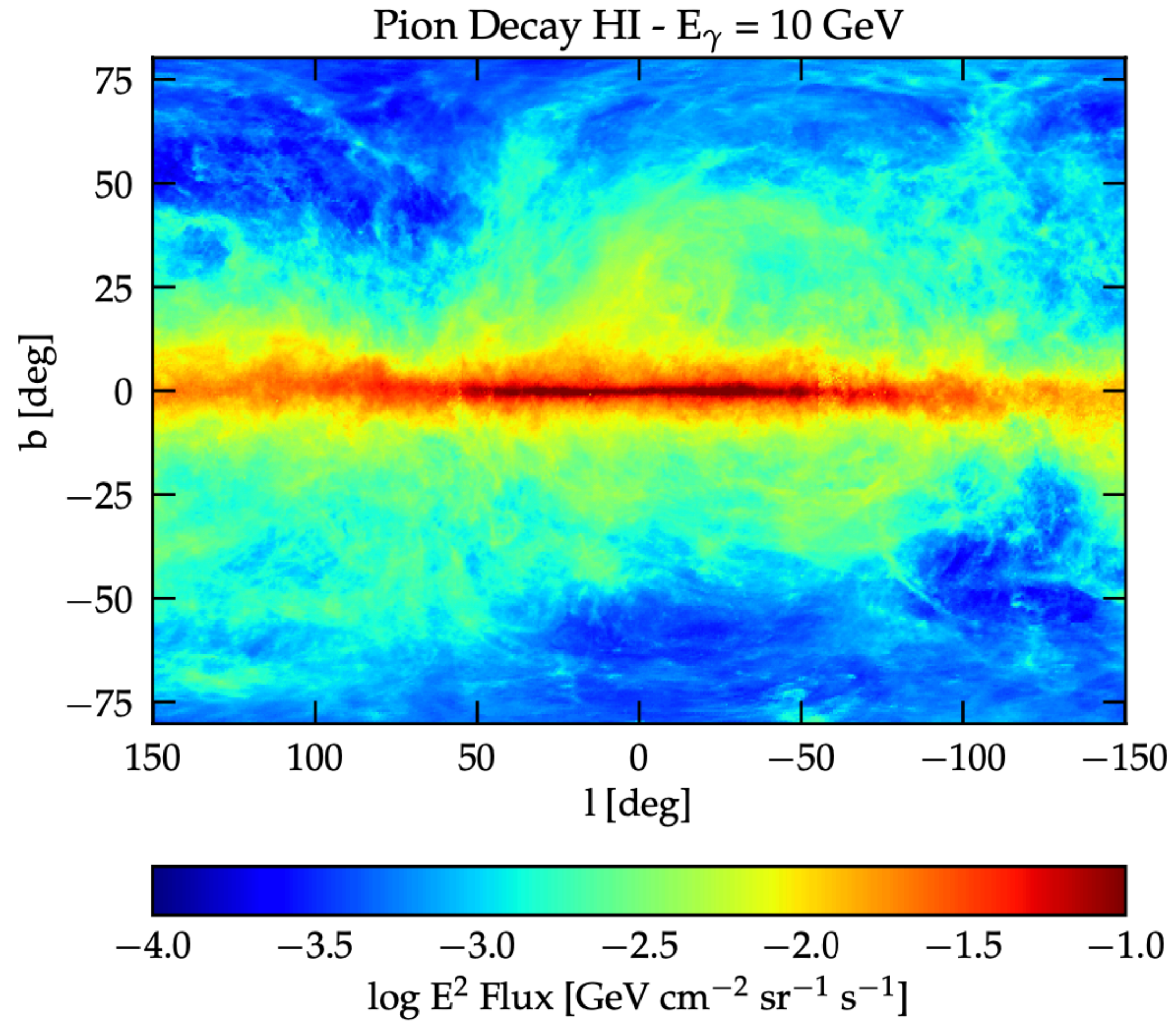
$$\tilde{Q}_{\gamma}(\bar{E}_{\gamma}) = \int_{\substack{\bar{E}_{\pi}^{\min} \\ 1}} d\bar{E}_{\pi} \frac{dn_{\gamma}}{d\bar{E}_{\gamma}} \tilde{Q}_{\pi^0}(\bar{E}_{\pi})$$

$\tilde{Q}_{\gamma}(\bar{E}_{\gamma})$ is the **gamma-ray emissivity**, [1/yr $\mu\text{m}^3 \text{s}$].

Gamma-ray spectrum at energies above the pion bump follows the parent proton distribution shifted by

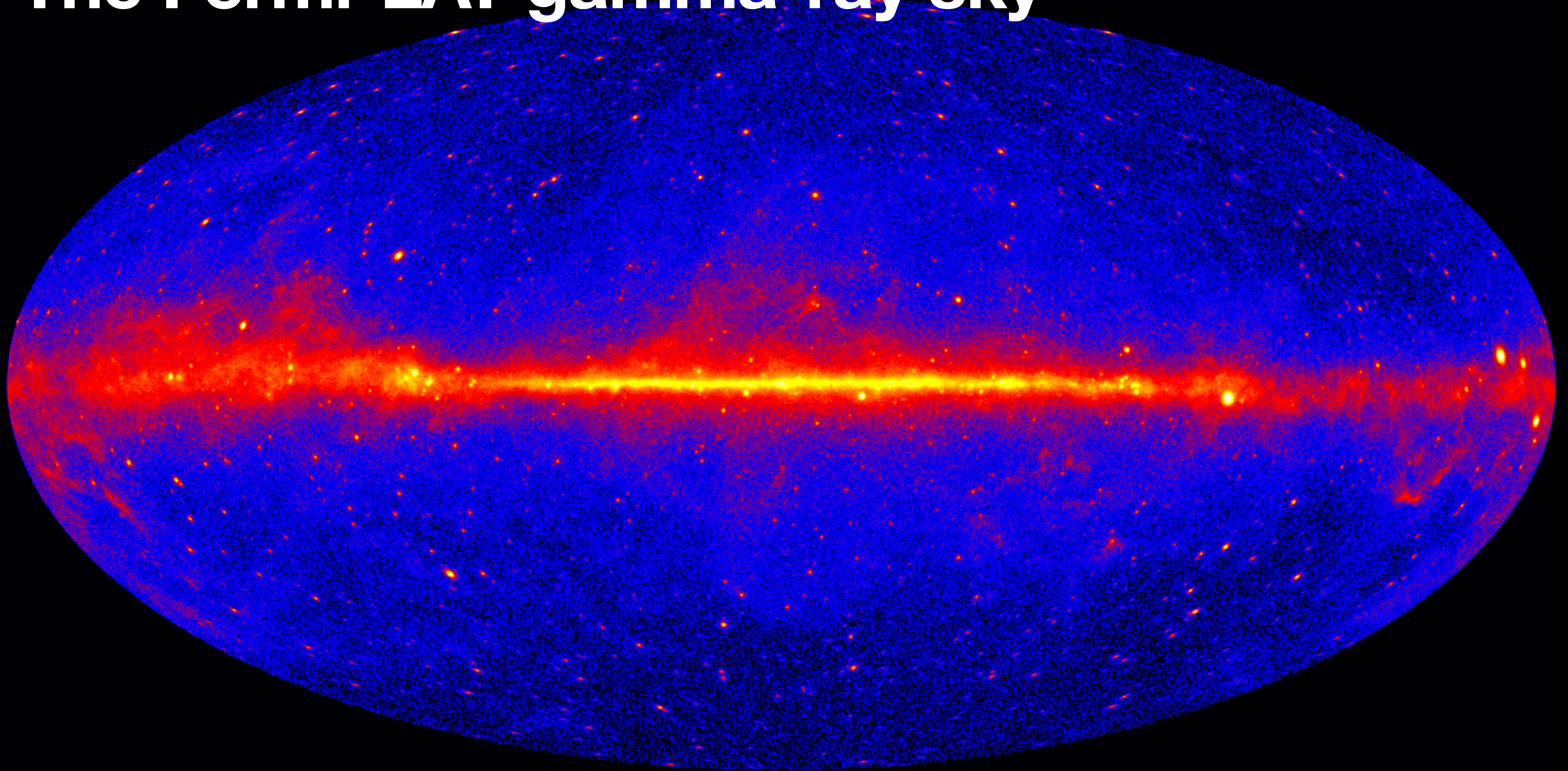
$$E_{\gamma} \sim \frac{1}{2} 0.17 E_p \sim \frac{E_p}{10}$$

Galactic photons from neutral pion decay



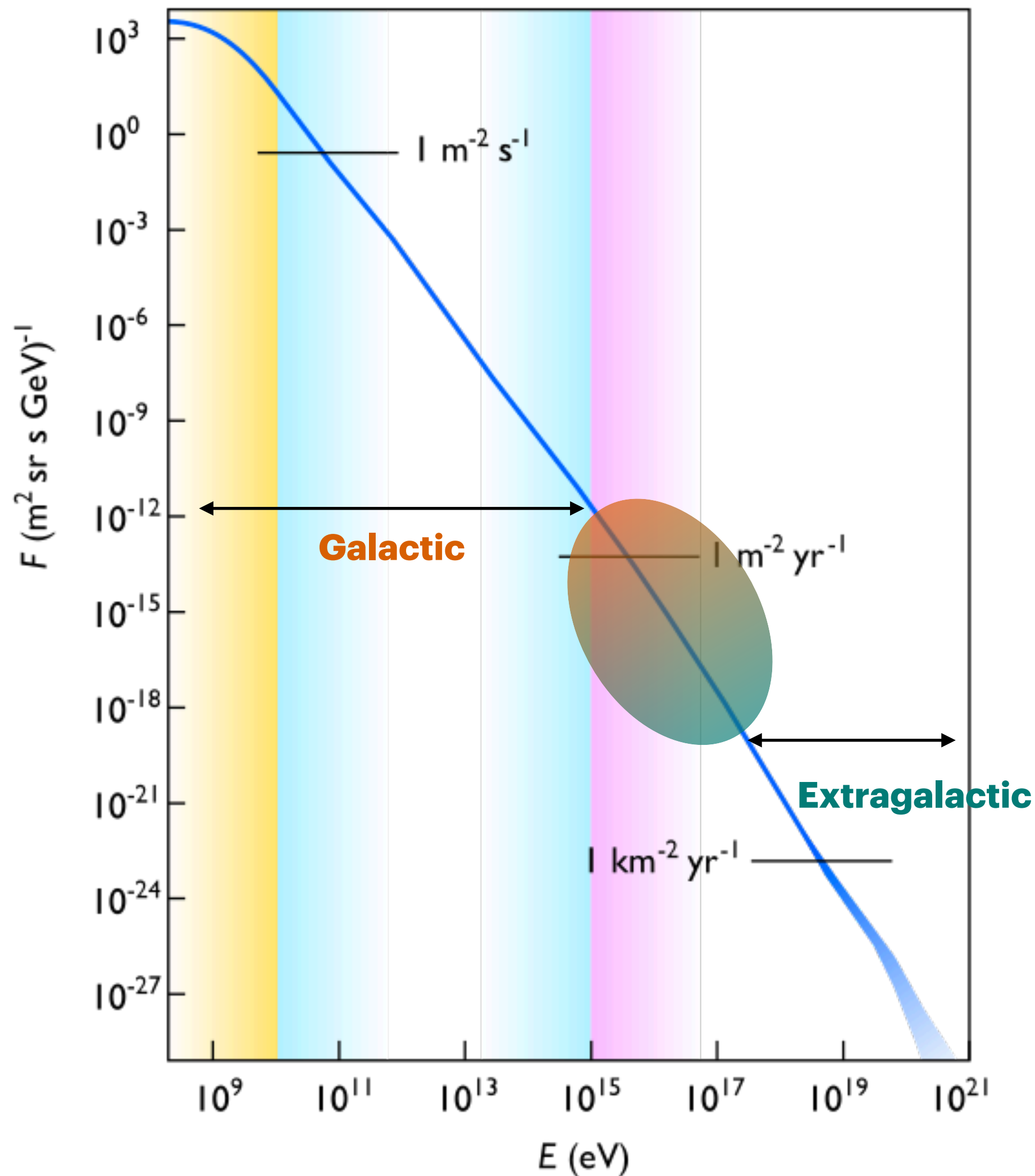
Dundovic+ A&A'22

The Fermi-LAT gamma-ray sky



Sources of cosmic particles: Standard astrophysics

Cosmic-ray sources: Galactic or extragalactic?



In the Galaxy we **observe** (in gamma rays)
CR factories up to
1 TeV = 10^{12} eV (HESS, VERITAS, MAGIC)
1 PeV = 10^{15} eV (LHAASO)

SNR, pulsars & neutron stars, binary, stellar clusters, PWN

AGN & jets, galaxy clusters, galaxies

Open questions:

1. What is the maximal energy CR are accelerated?
2. Where does the Gal-extragal transition occur?

How to accelerate CRs?

Some **requirements**:

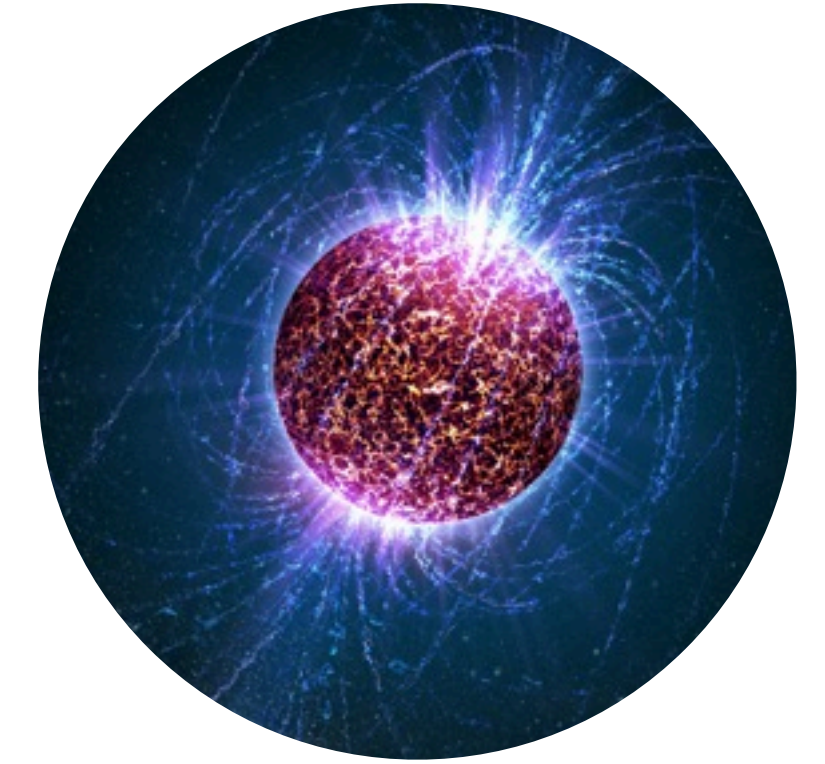
1. Energetics:

- Kinetic Energy (translational in SNRs, rotational in pulsars)
- Gravitational Energy (accretion disks)
- Magnetic (solar flares)

2. Mechanism for Energy Transfer: how to transfer energy from macroscopic objects into the (microscopic) acceleration of particles? (electromagnetic)

3. Confinement: particles must stay in the accelerator for the time needed to accelerate them

4. No significant E-losses



We need electric fields to accelerate particles!
They are generated by moving magnetic fields in the plasma
e.g. fast rotating B-field in pulsars, shock waves

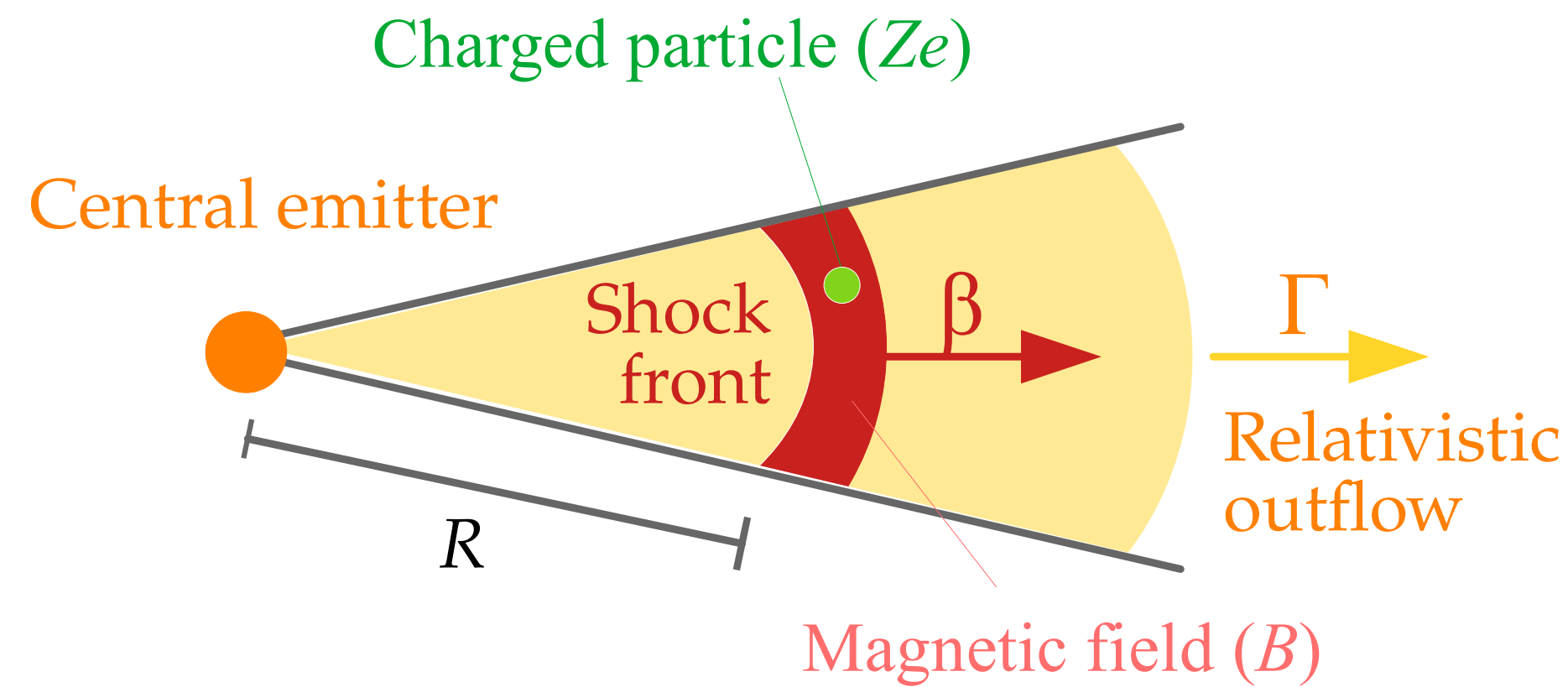
Confinement condition

Necessary condition for acceleration:

The system must be able to contain the particle

$$r_L \lesssim R$$

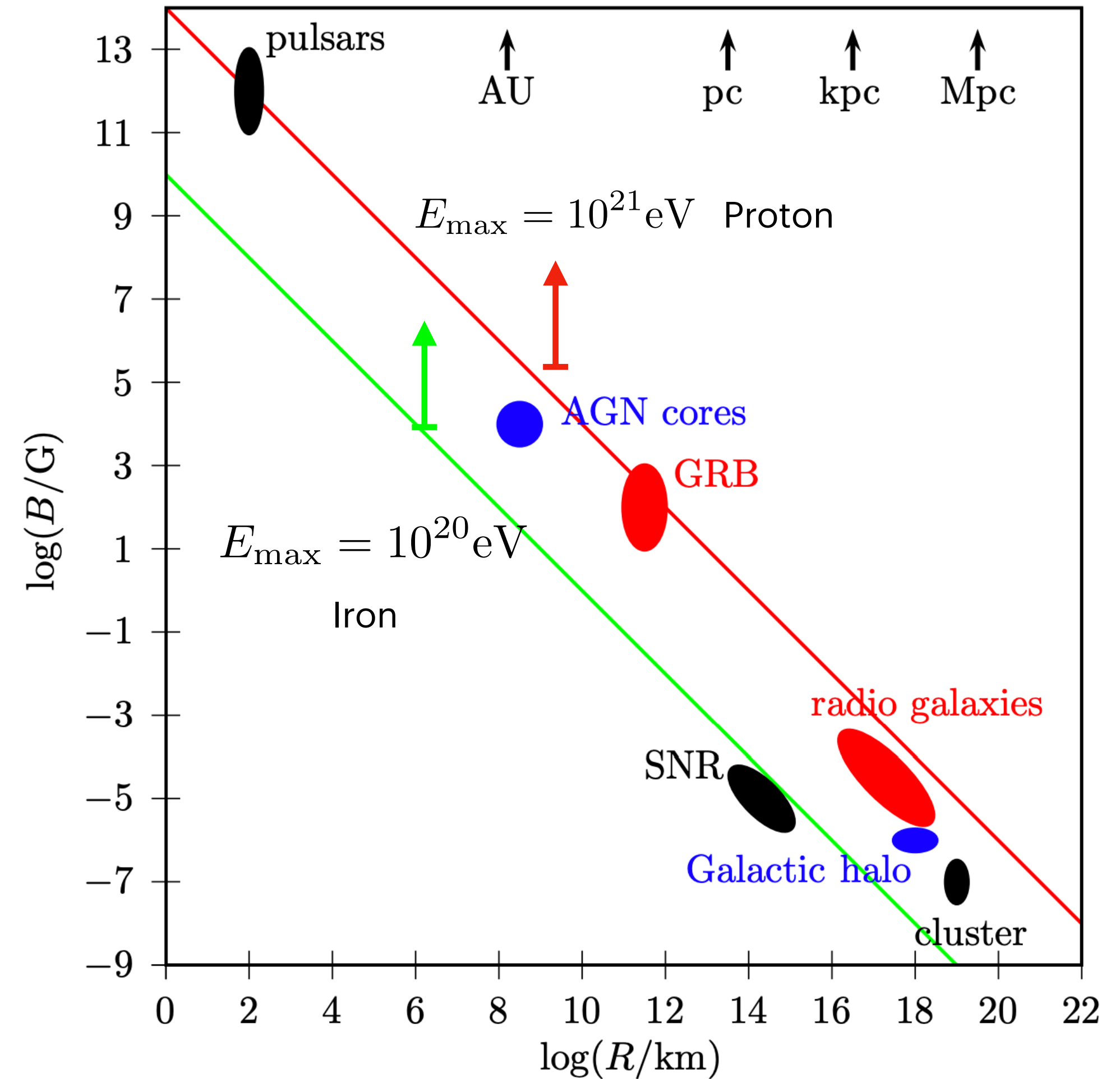
$$\beta = 1$$



Hillas Plot

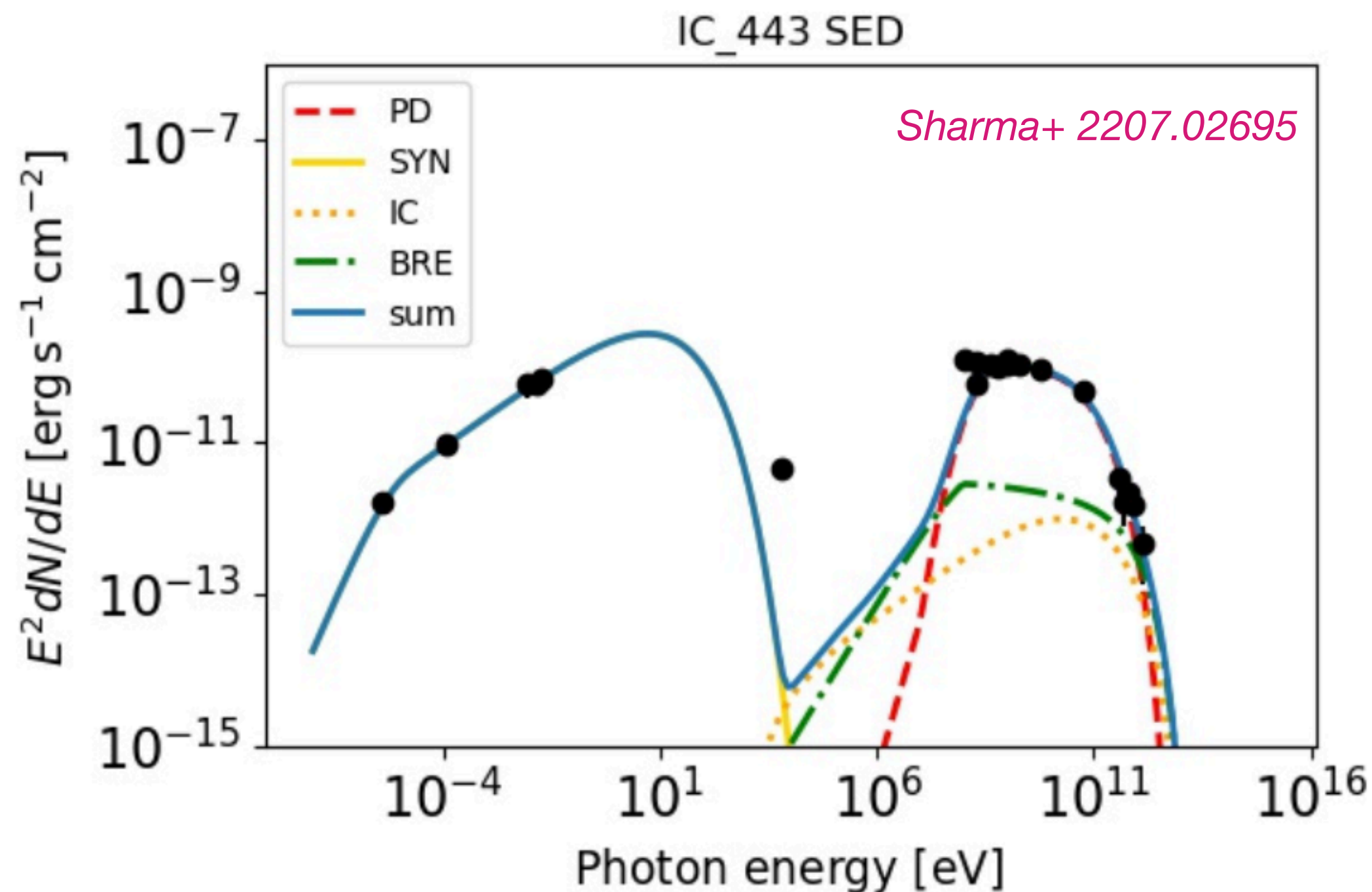
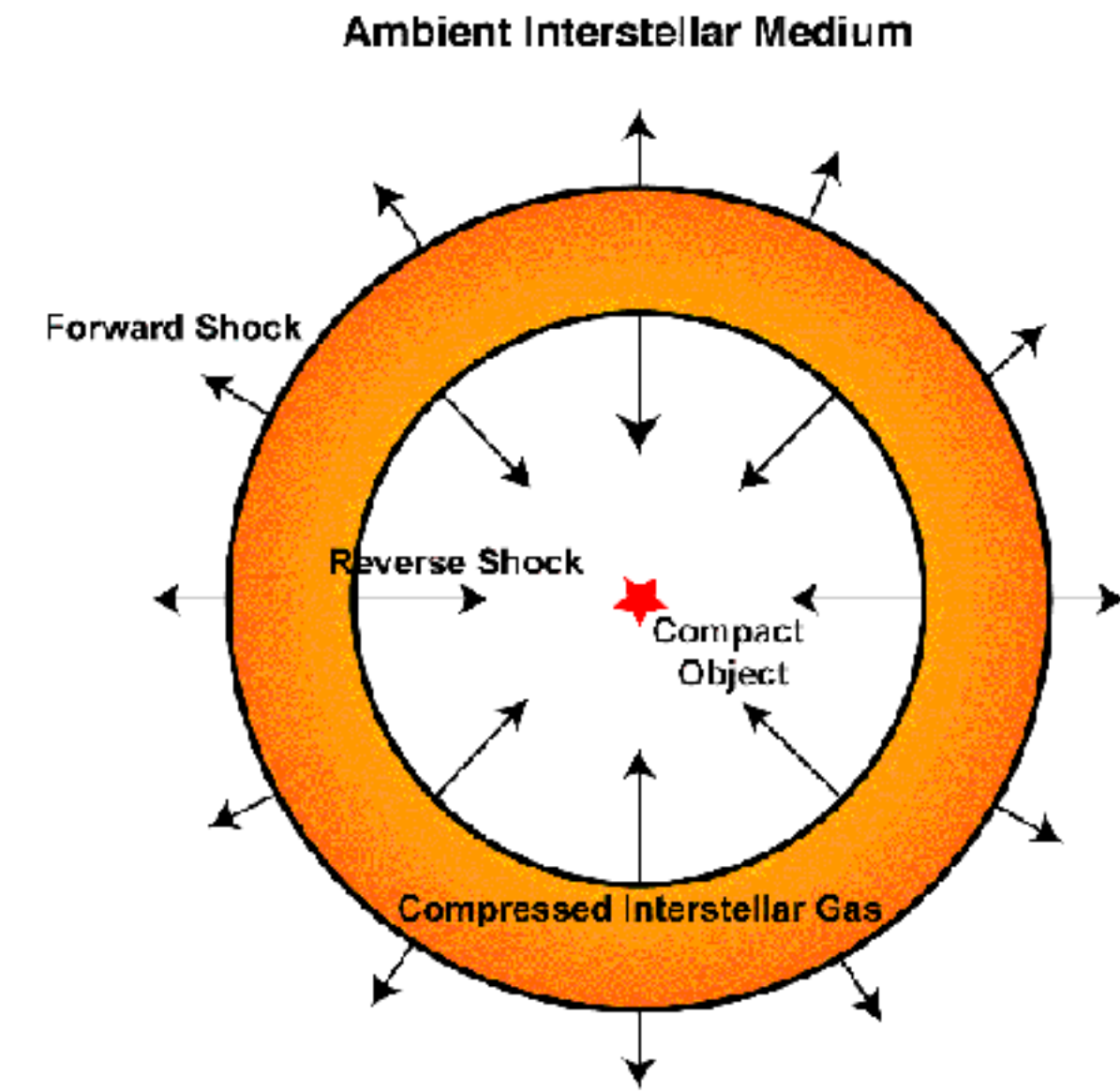
$$E_{\max} \propto RB$$

$$r_L = \gamma r_g = \sqrt{1 - \mu^2} \frac{\mathcal{R}}{B_0} \simeq 10^{-6} \sqrt{1 - \mu^2} \frac{\mathcal{R}}{\text{GV}} \frac{\mu\text{G}}{B_0} \text{pc}$$

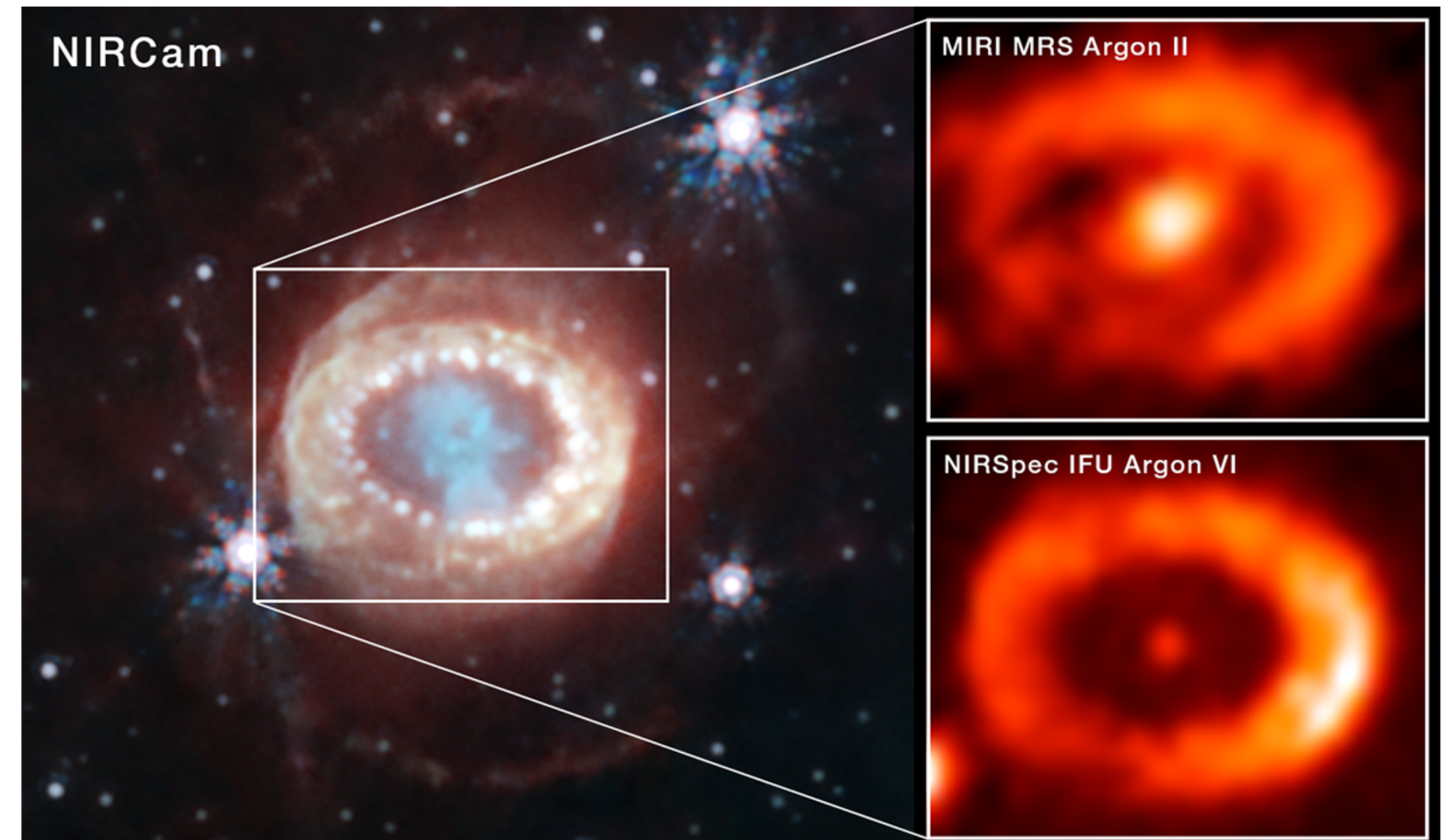


Supernova remnants

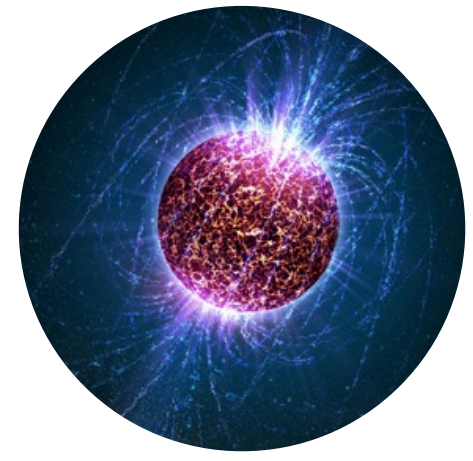
- Structure resulting from the explosion of a star in a SN of type I (white-dwarf accretion) or II (core-collapse)
- Expanding material ejected in the explosion and shocking the ISM along the way
- Strongly magnetised shocks 25 – 1000 μG



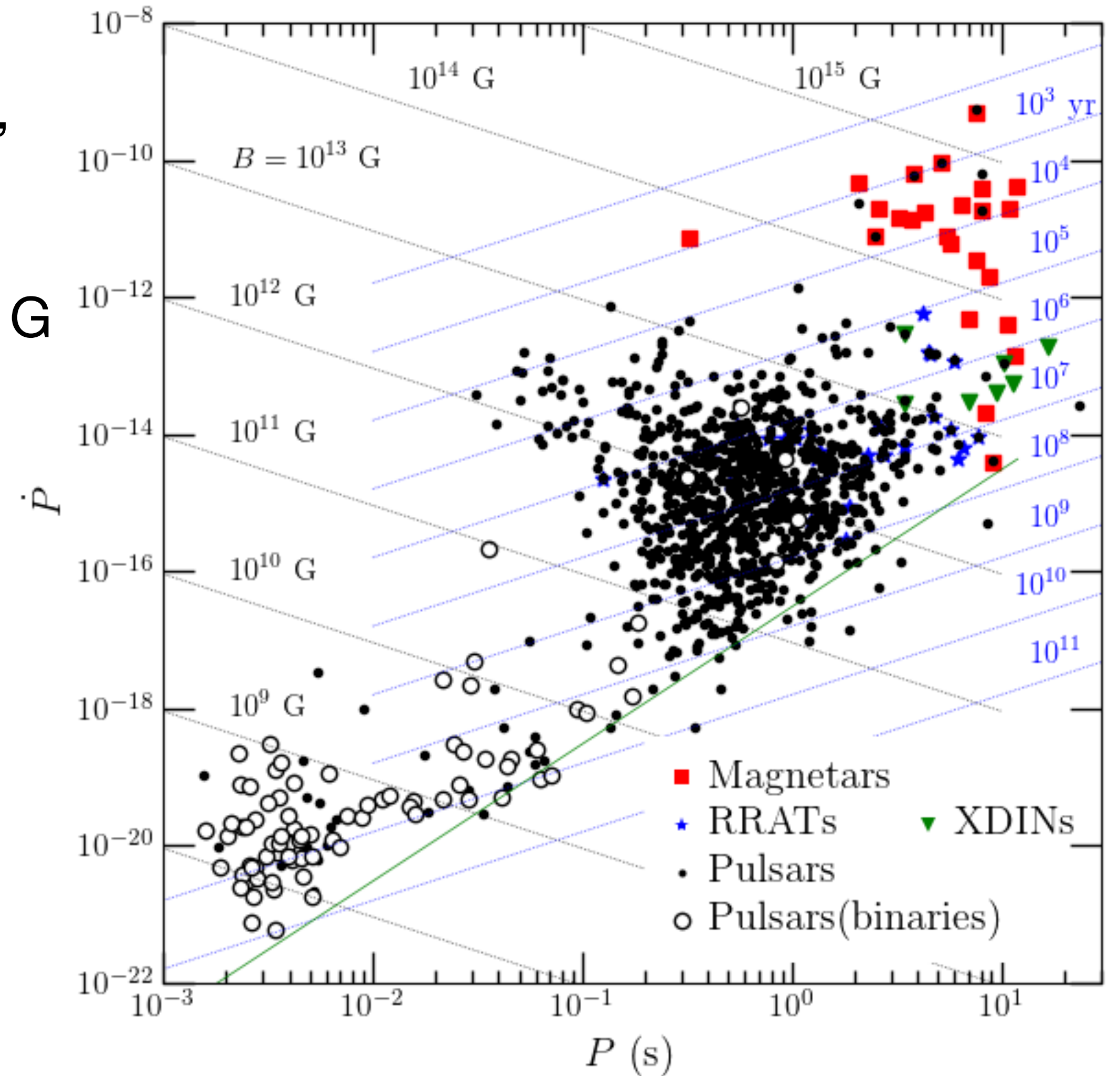
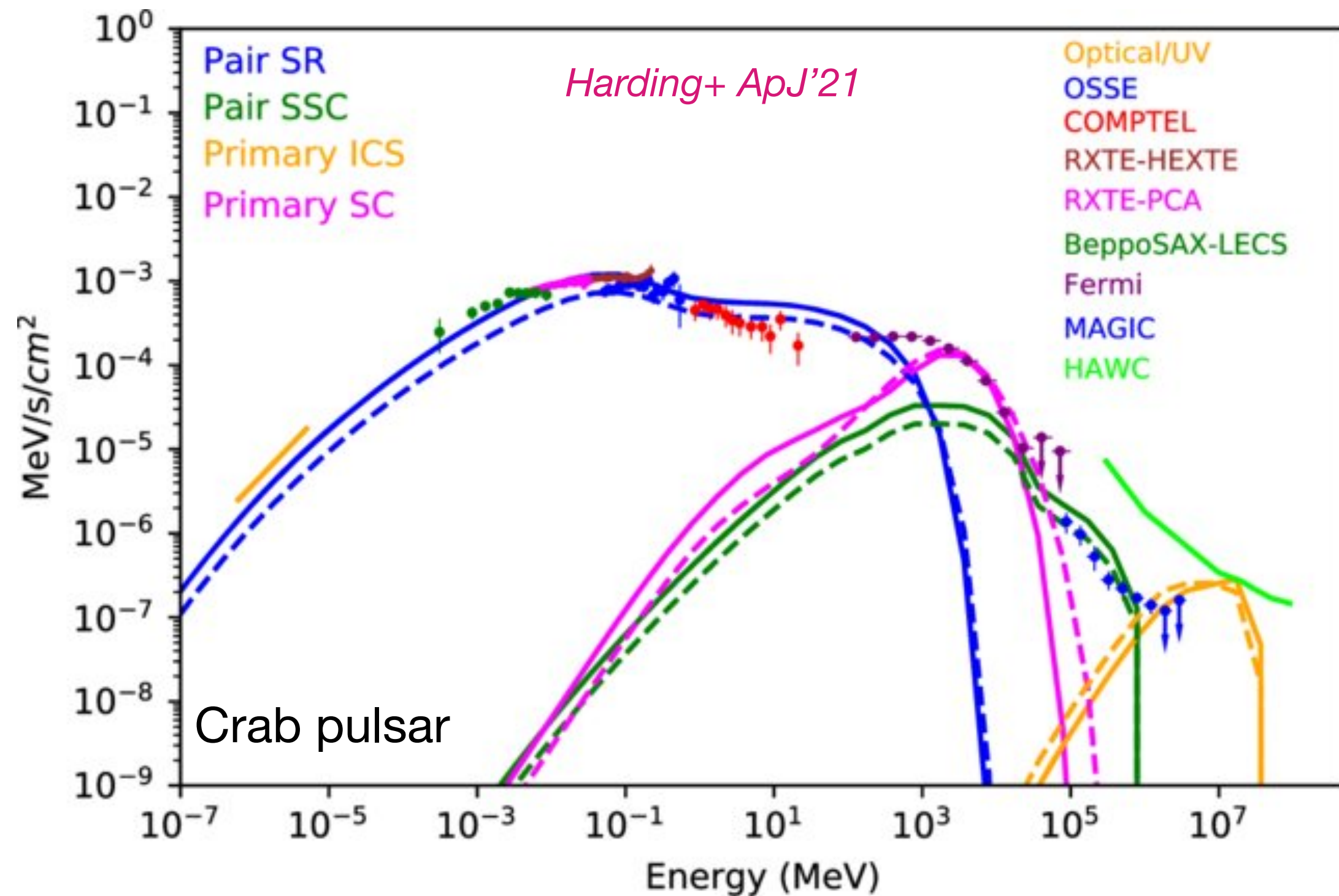
Webb's evidence of neutron star in SN1987A



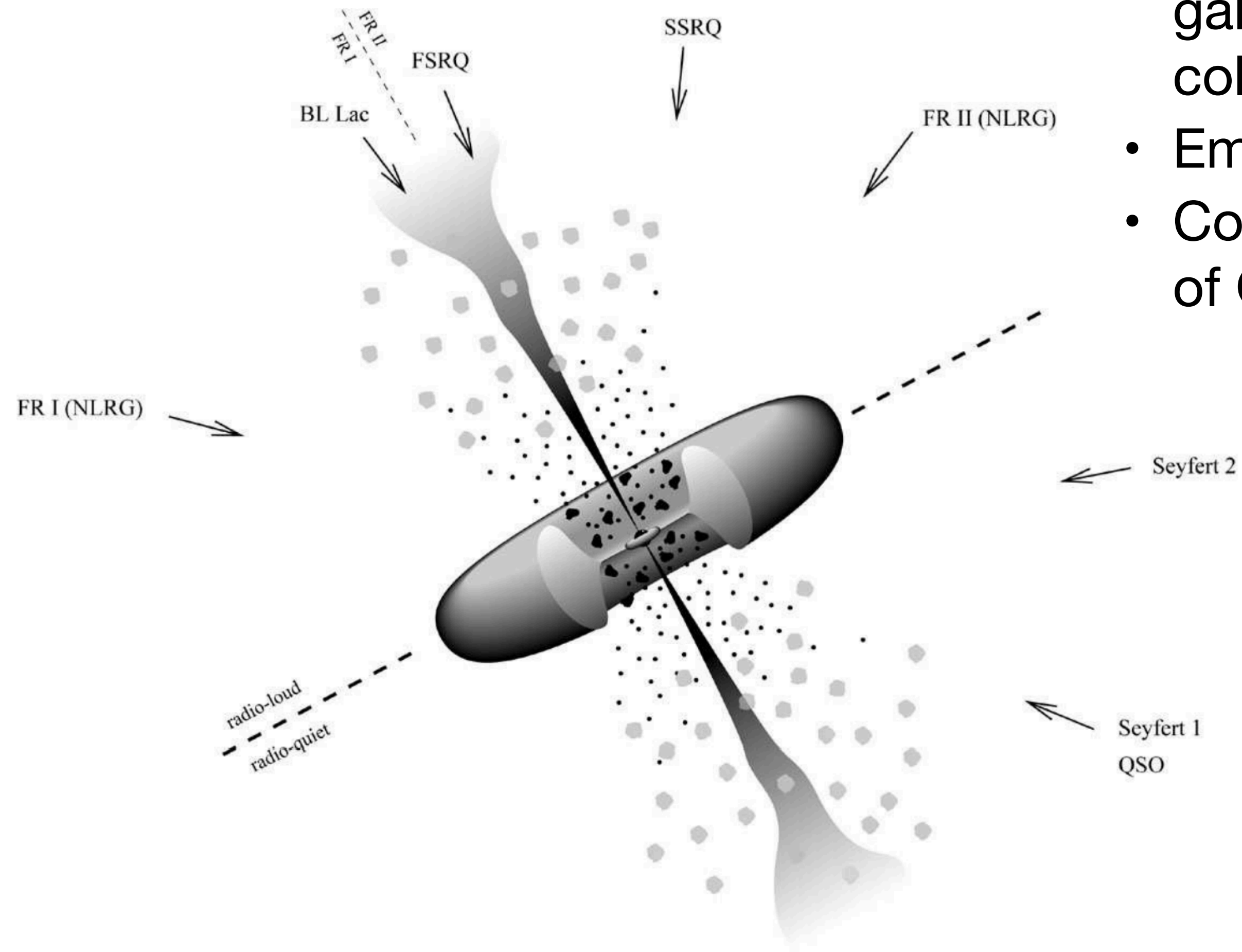
Pulsars and PSR wind nebulae



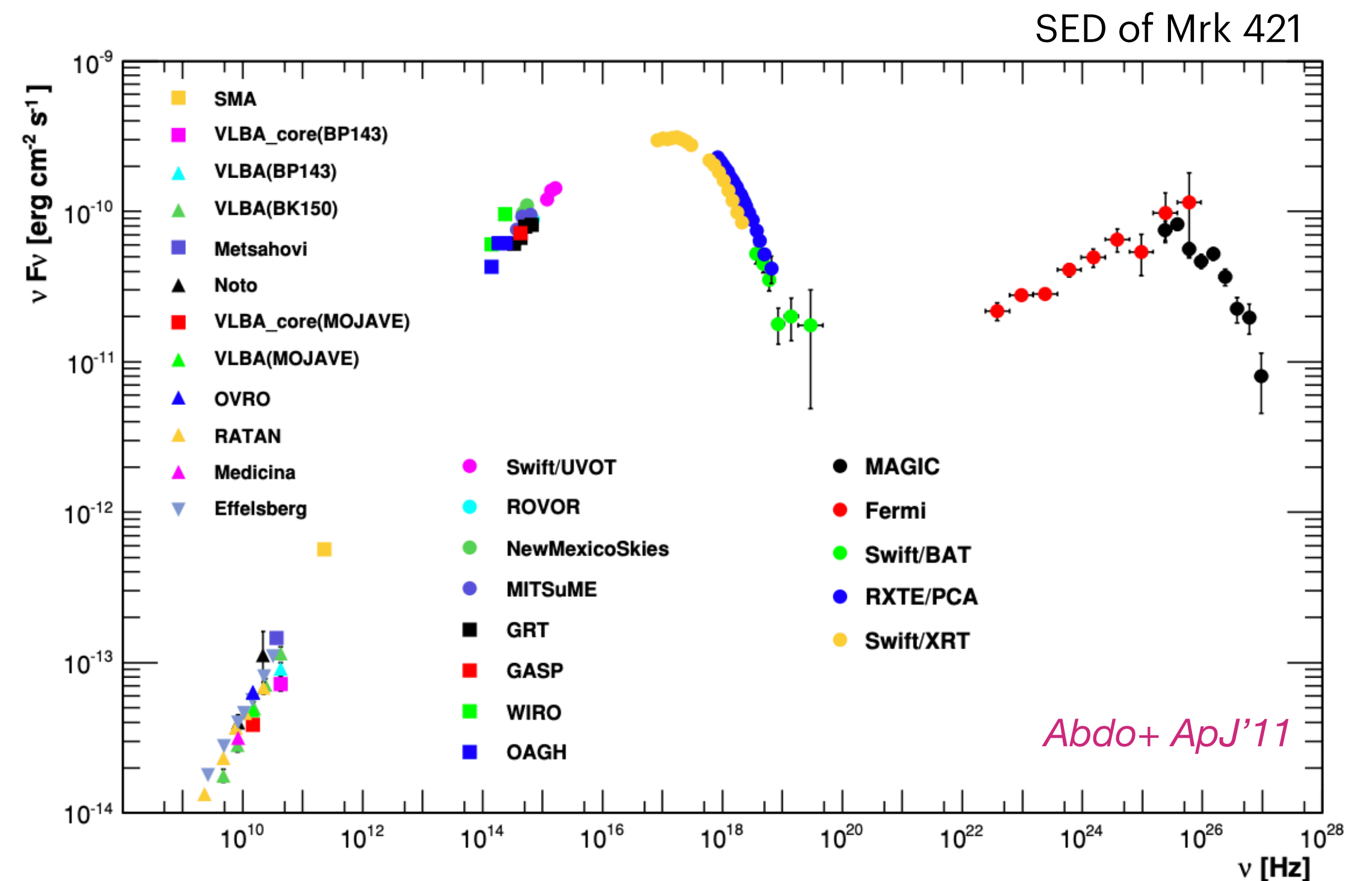
- Pulsars are rapidly rotating neutron stars with extremely high densities
- Emit beams of electromagnetic radiation (radio, X-ray, gamma-ray) from their magnetic poles through synchrotron and curvature radiation
- Dipolar magnetic field with strengths from 10^8 to 10^{15} G



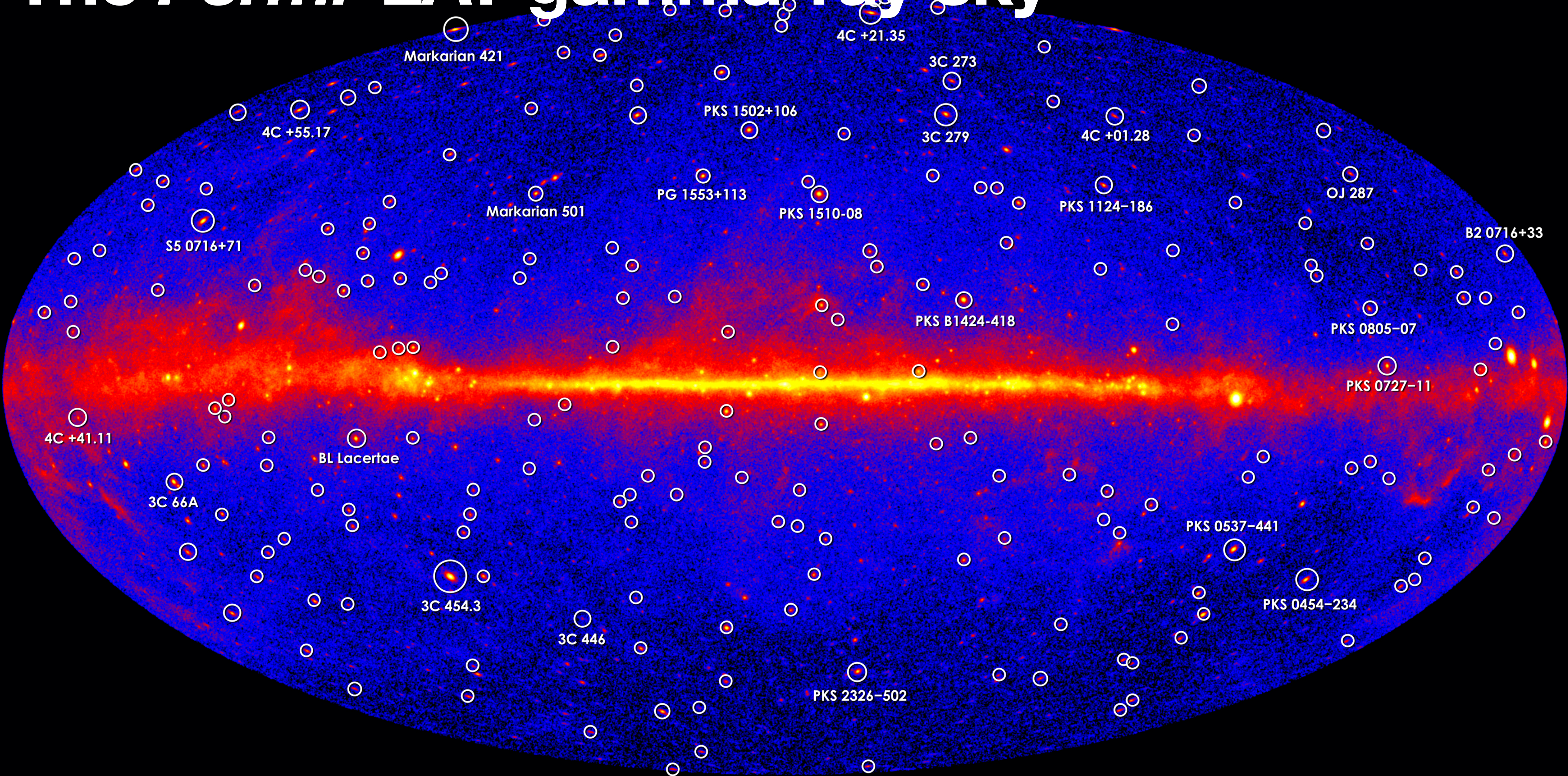
Active galactic nuclei



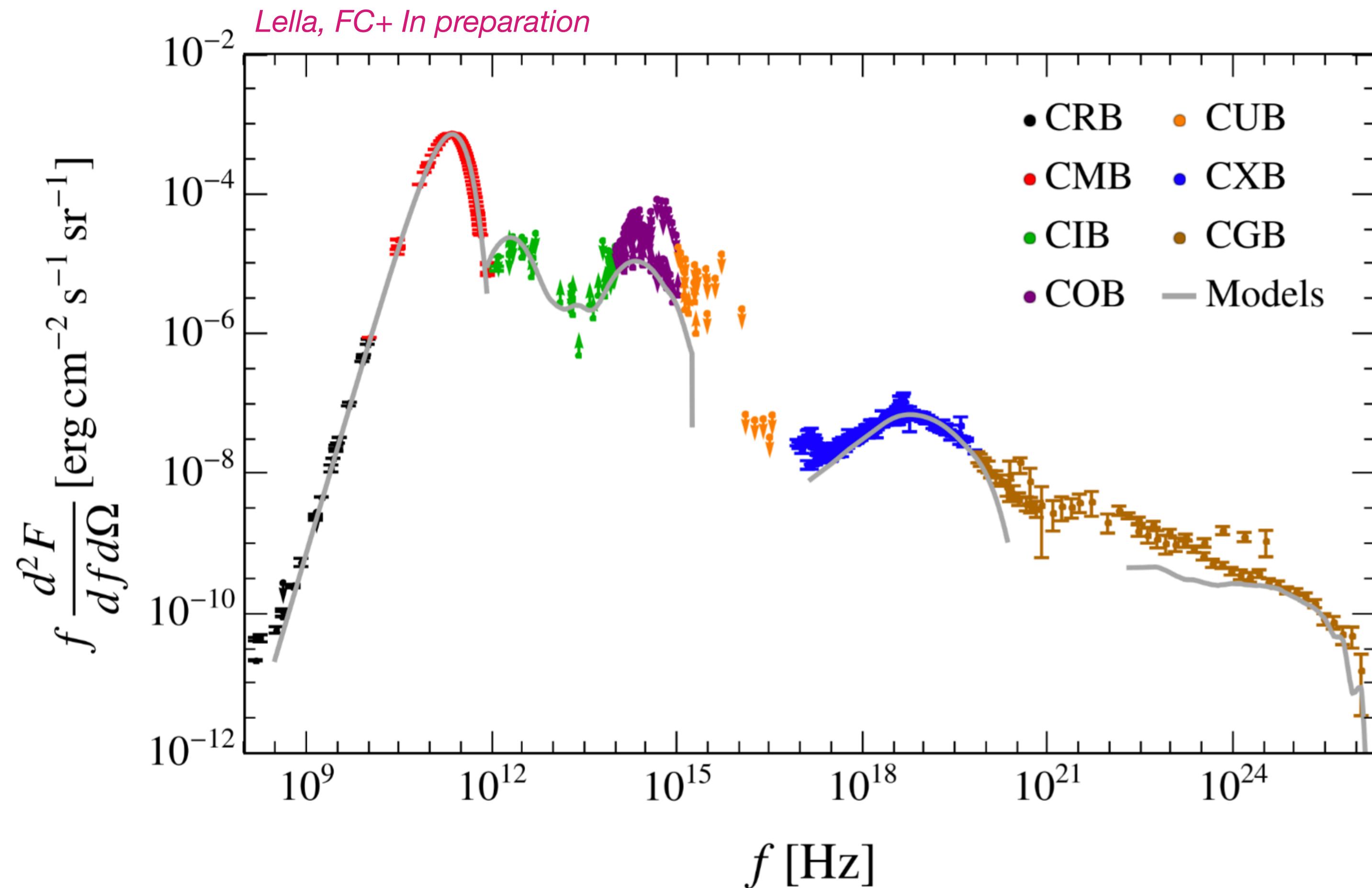
- Large fraction of the total luminosity of an active galaxy is non-thermal (accretion disk and collimated jet)
- Emitted by the nuclei of the galaxy
- Complex B fields structure (poloidal and toroidal) of $O(10^3 \text{ G})$



The *Fermi*-LAT gamma-ray sky



Cosmic photon background

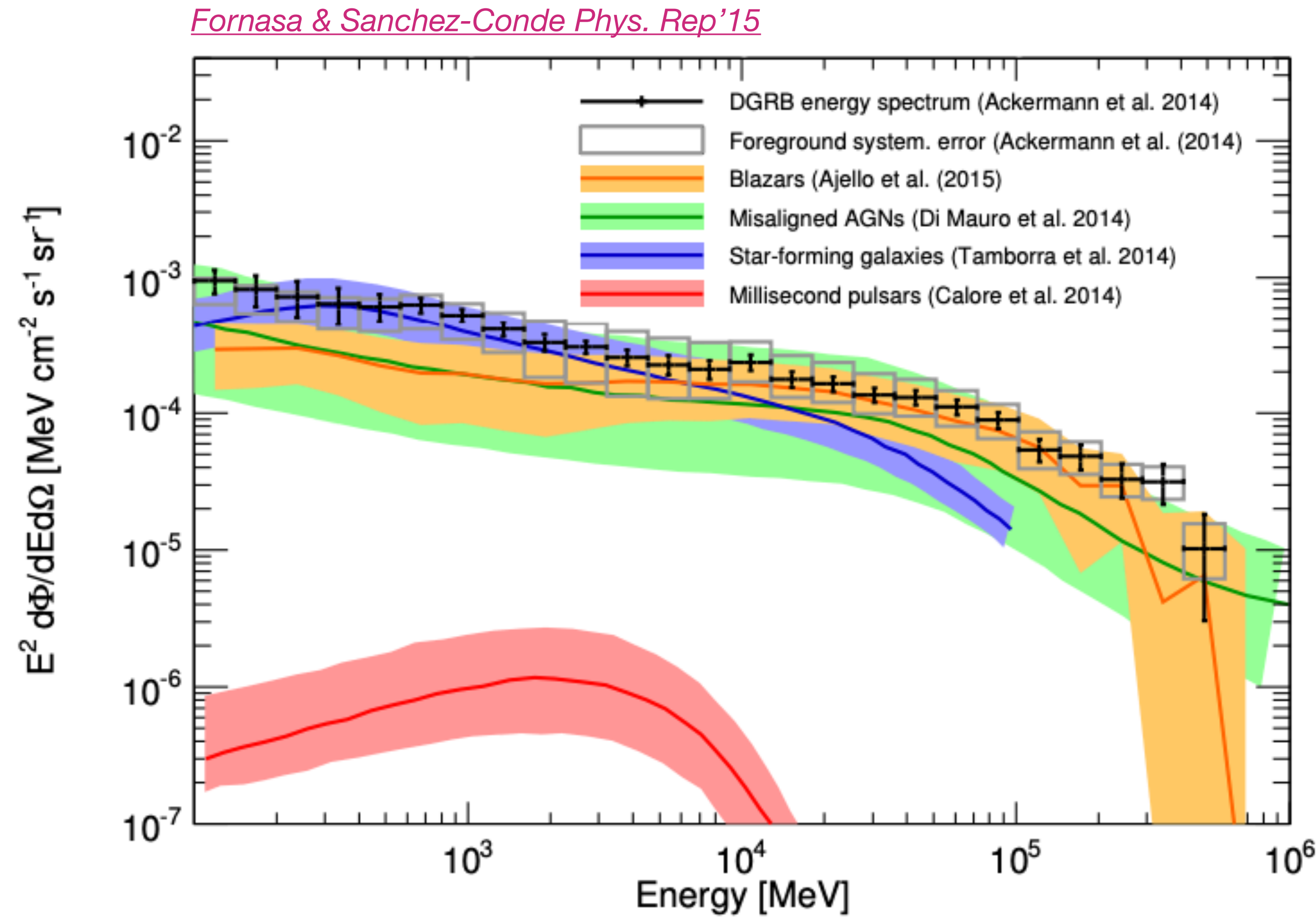
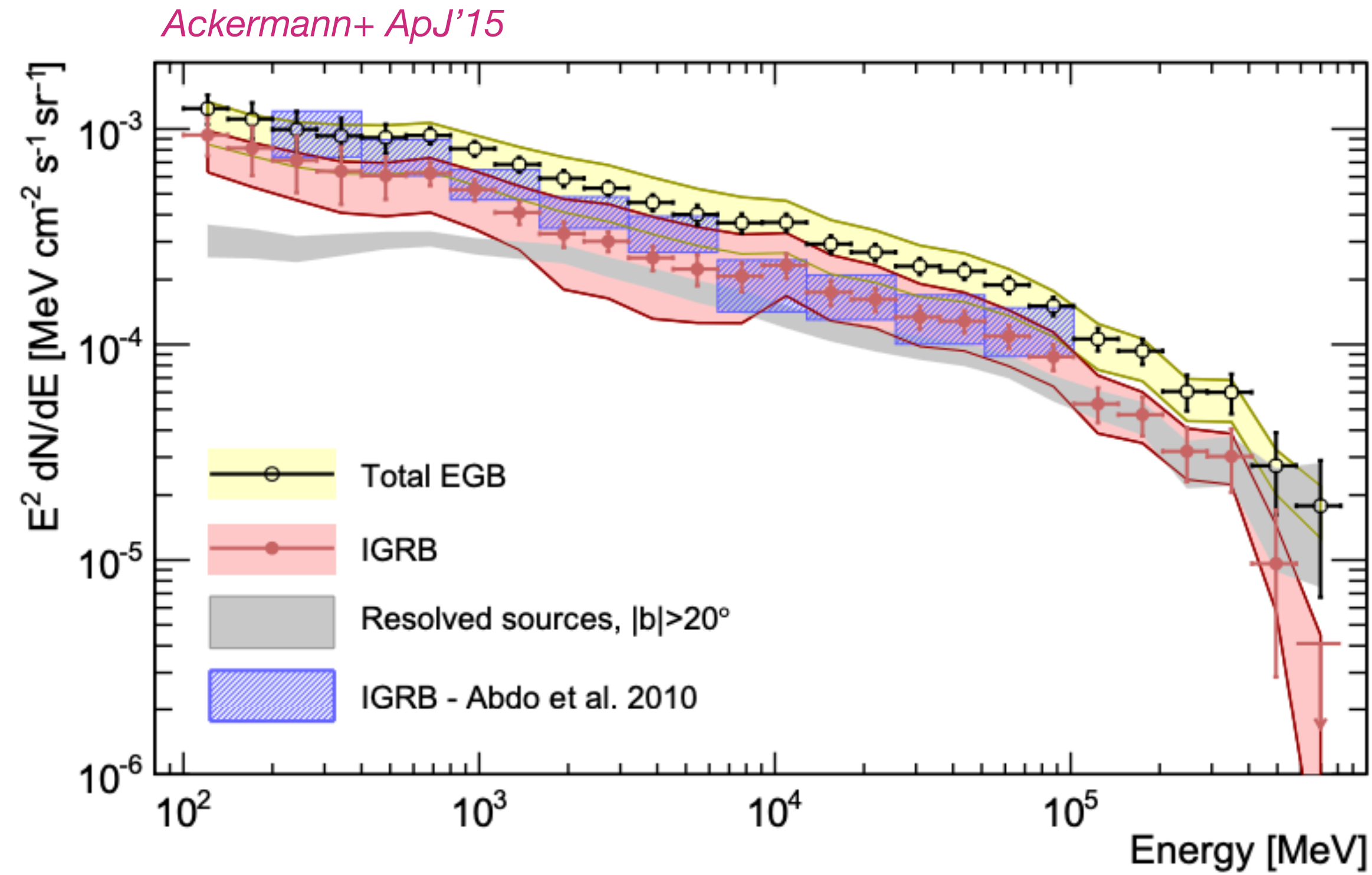


[For a review see: *Hill+ App. Spectrosc.'18*]

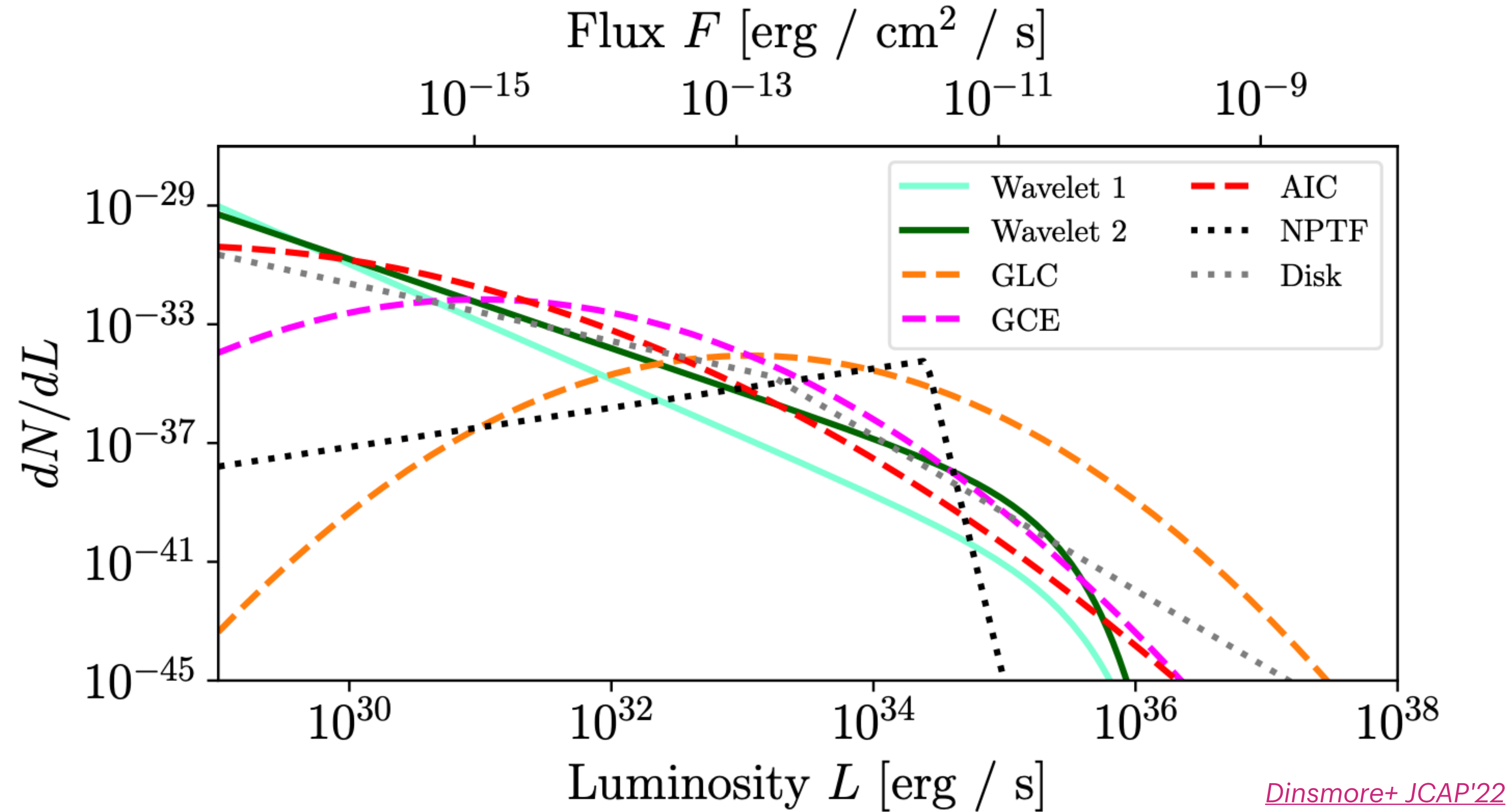
- **CRB & CMB**: Primordial background subtracted by Galactic foregrounds @ 2.7255 K
- **CIB**: Emission of dust heated by stars within unresolved galaxies, difficult to clean from Galactic foregrounds
- **COB**: Emission from stars, difficult to clean from Galactic foregrounds
- **CUB**: From all sources of ionising photons such as star-forming galaxies and quasars
- **CXB**: Dominate by bremsstrahlung in the hot accretion disks around AGNs
- **CGB**: Superposition of several source classes (AGNs, star-forming galaxies) and Galaxy emission

Common element: Mostly of extra-galactic origin, as superposition of faint photon emitters

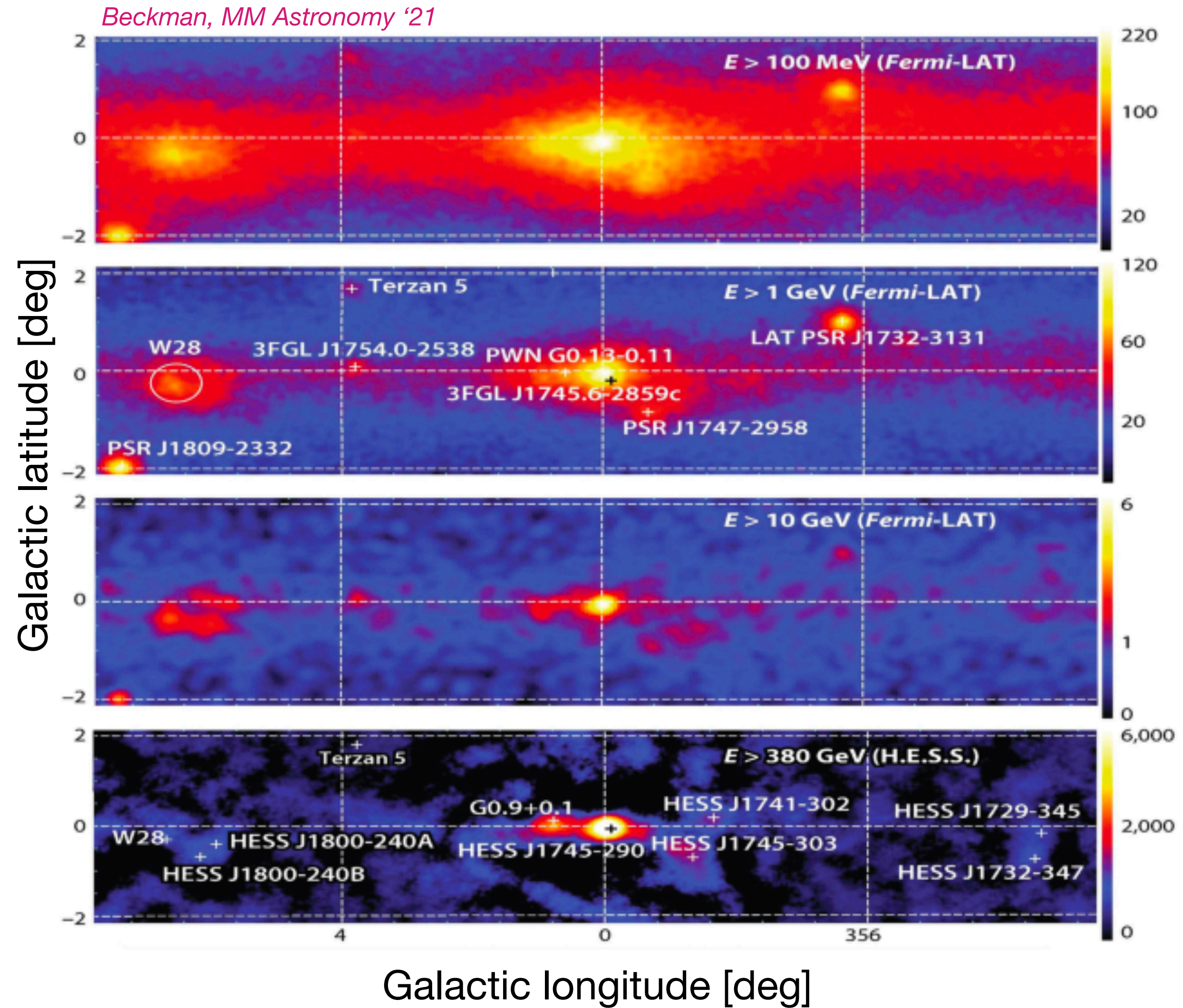
Unresolved sources and the Fermi IGRB



Galactic unresolved sources

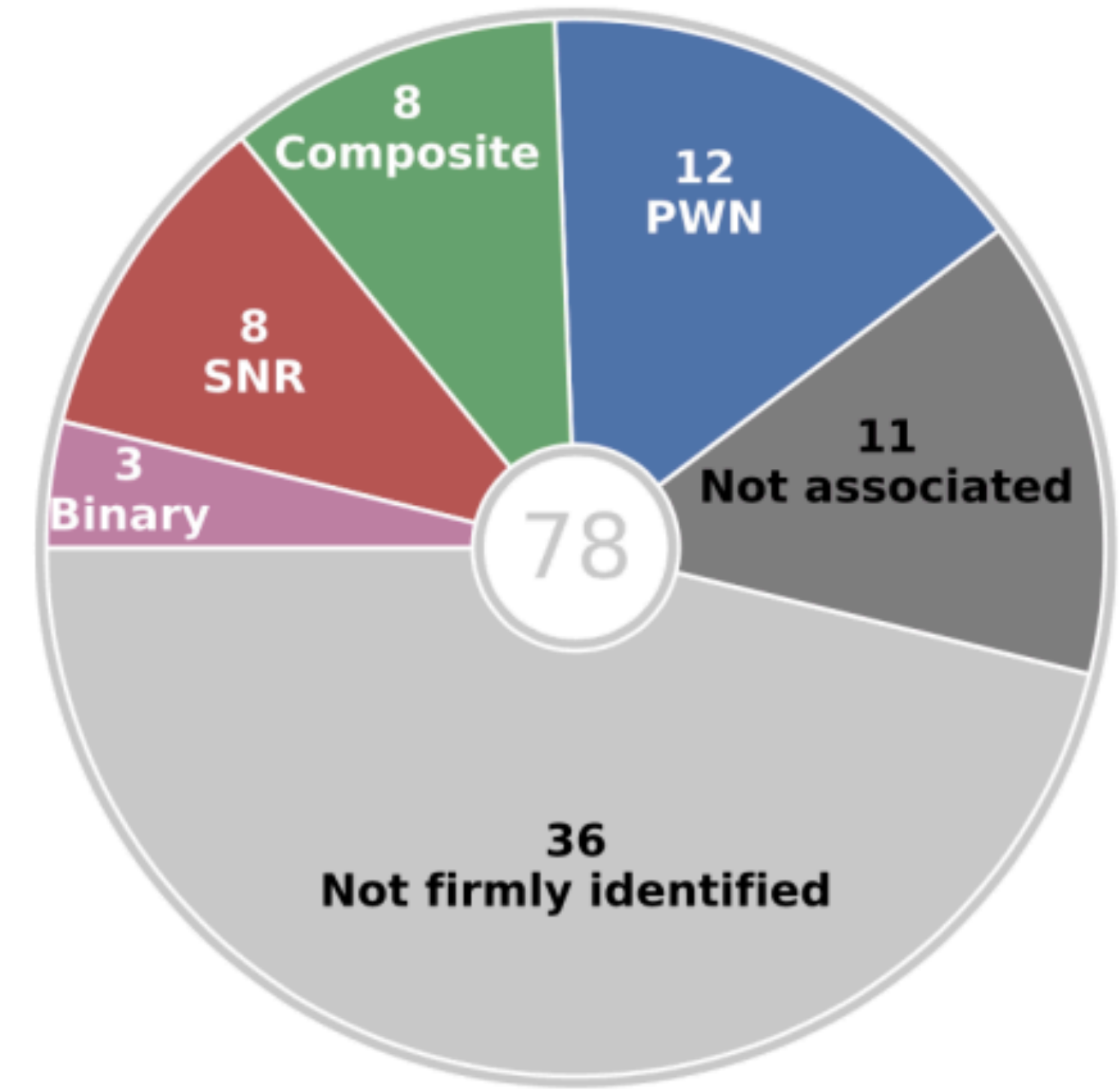
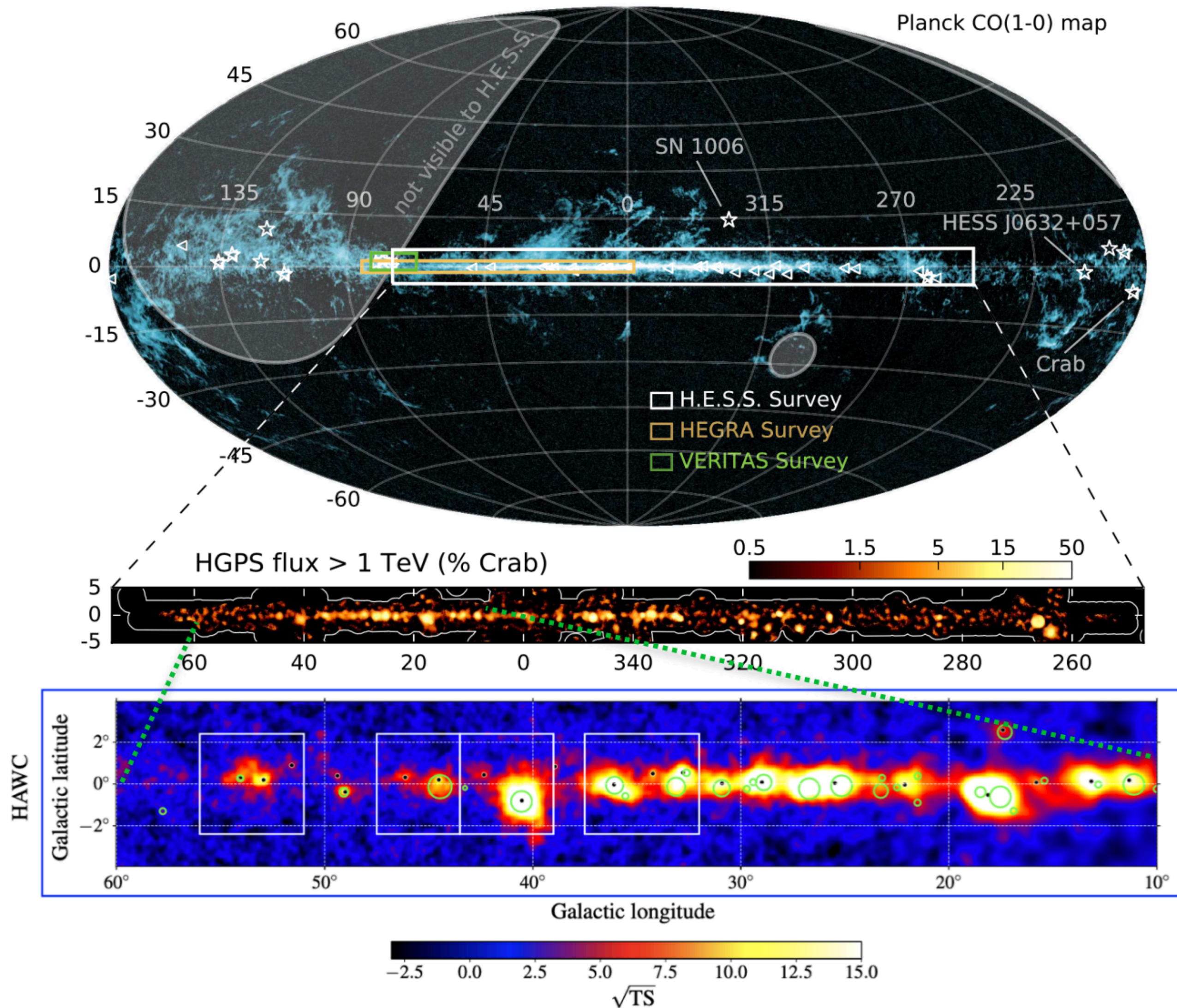


Diffuse emission from TeV to sub-PeV



Point sources at TeV energies

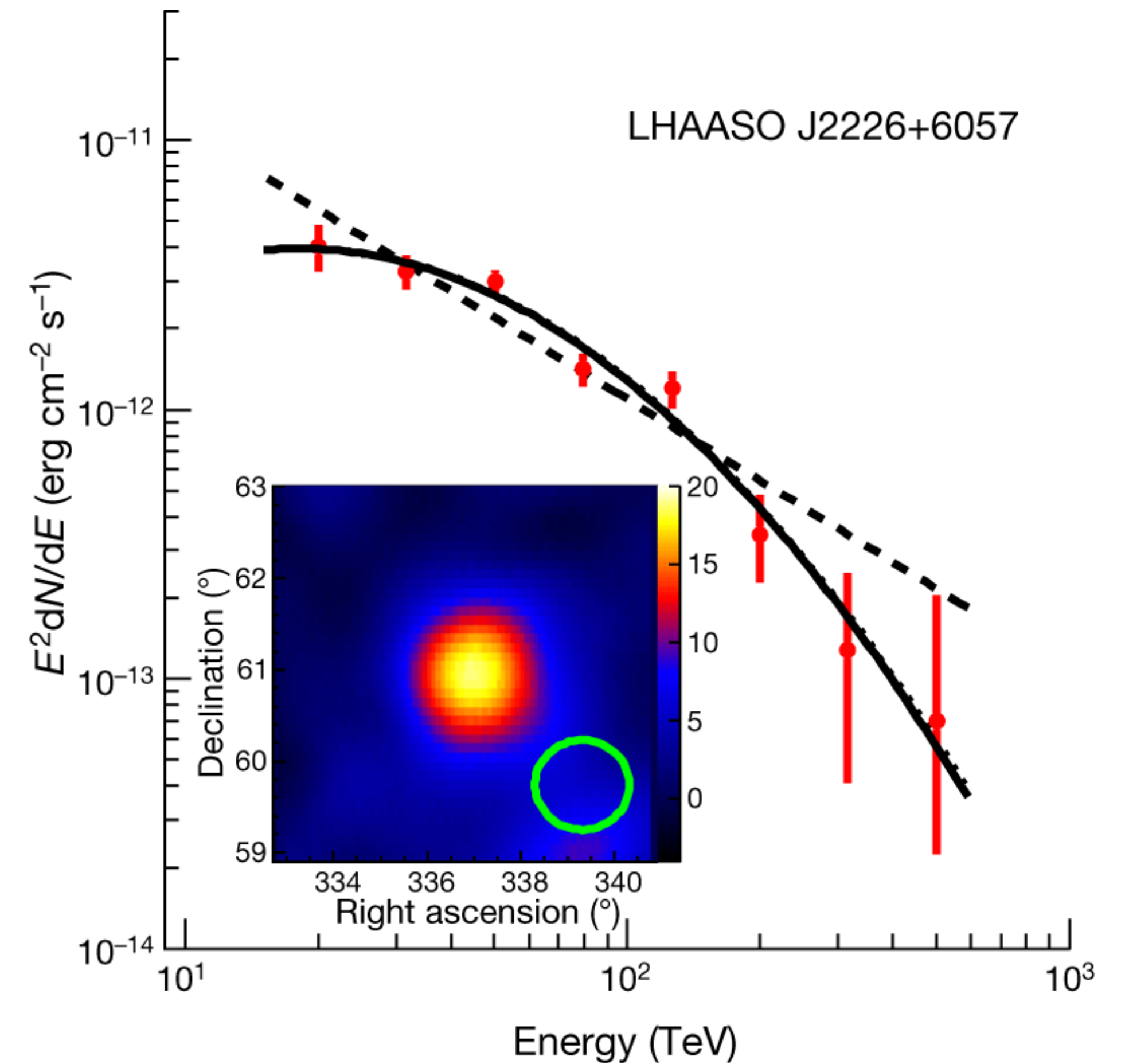
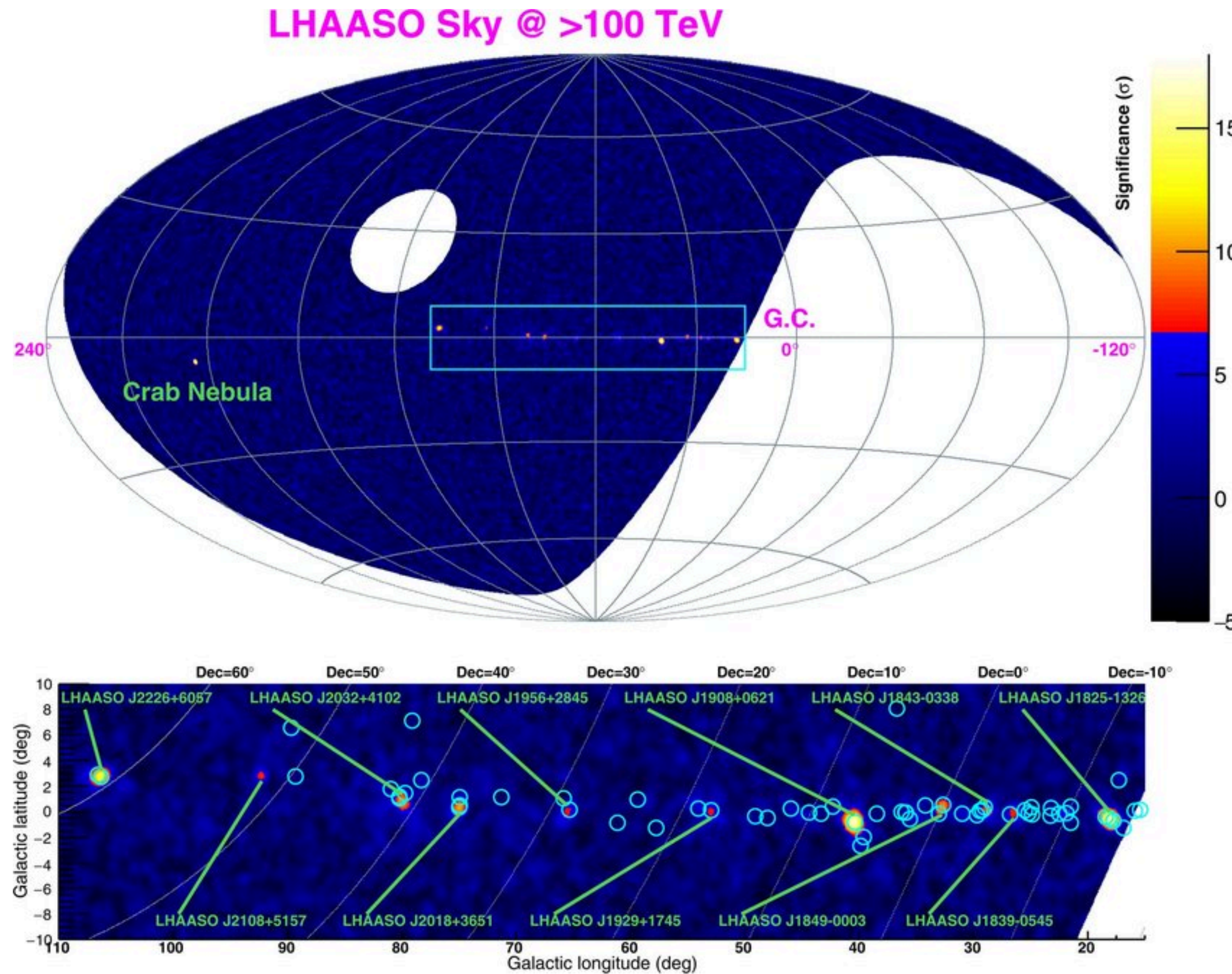
Despite small FoV IACTs can perform effective surveys



Galactic plane full of TeVatrons!

HAWC: Sources up to 100 TeV!

Point sources at sub-PeV energies

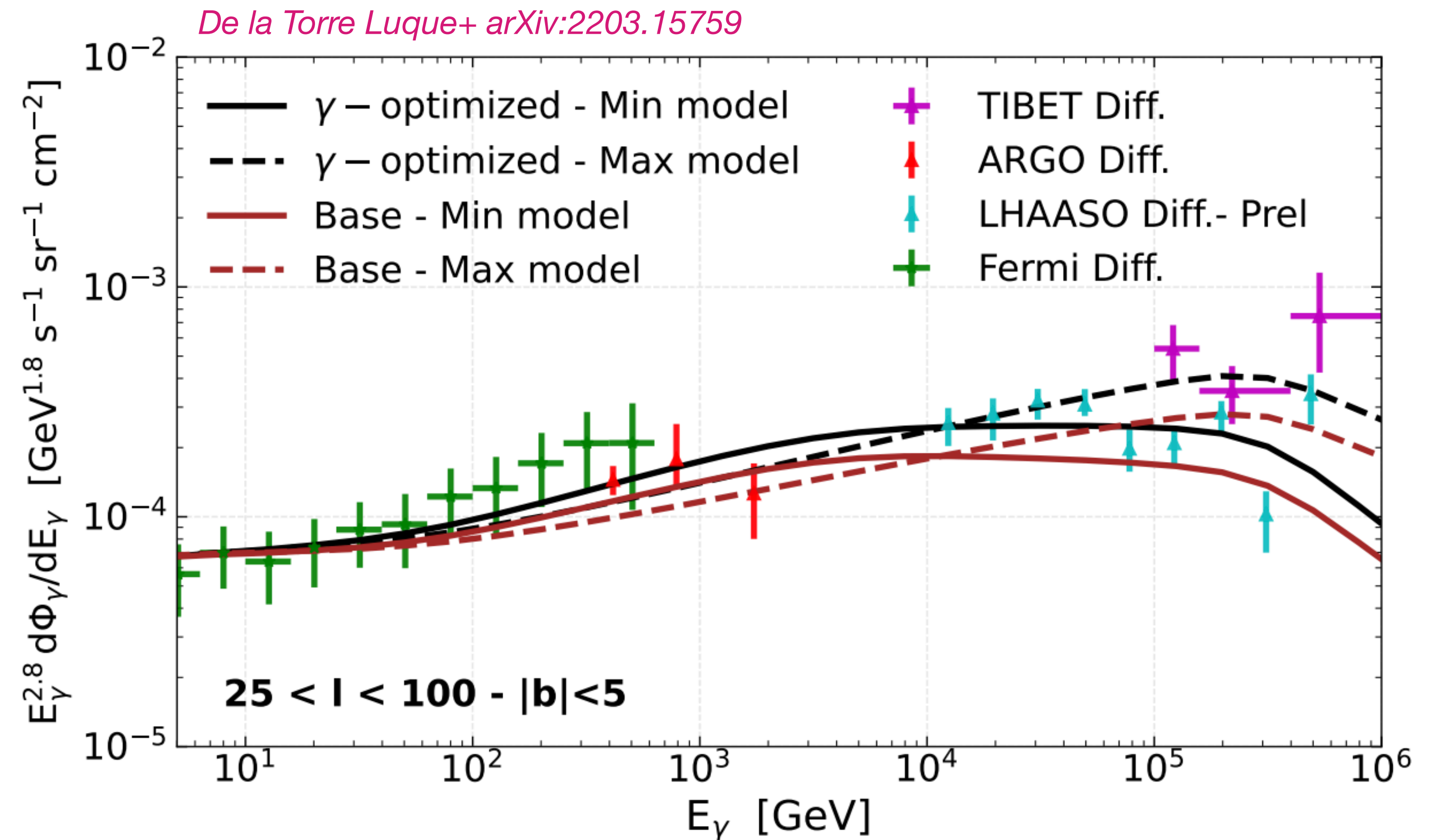
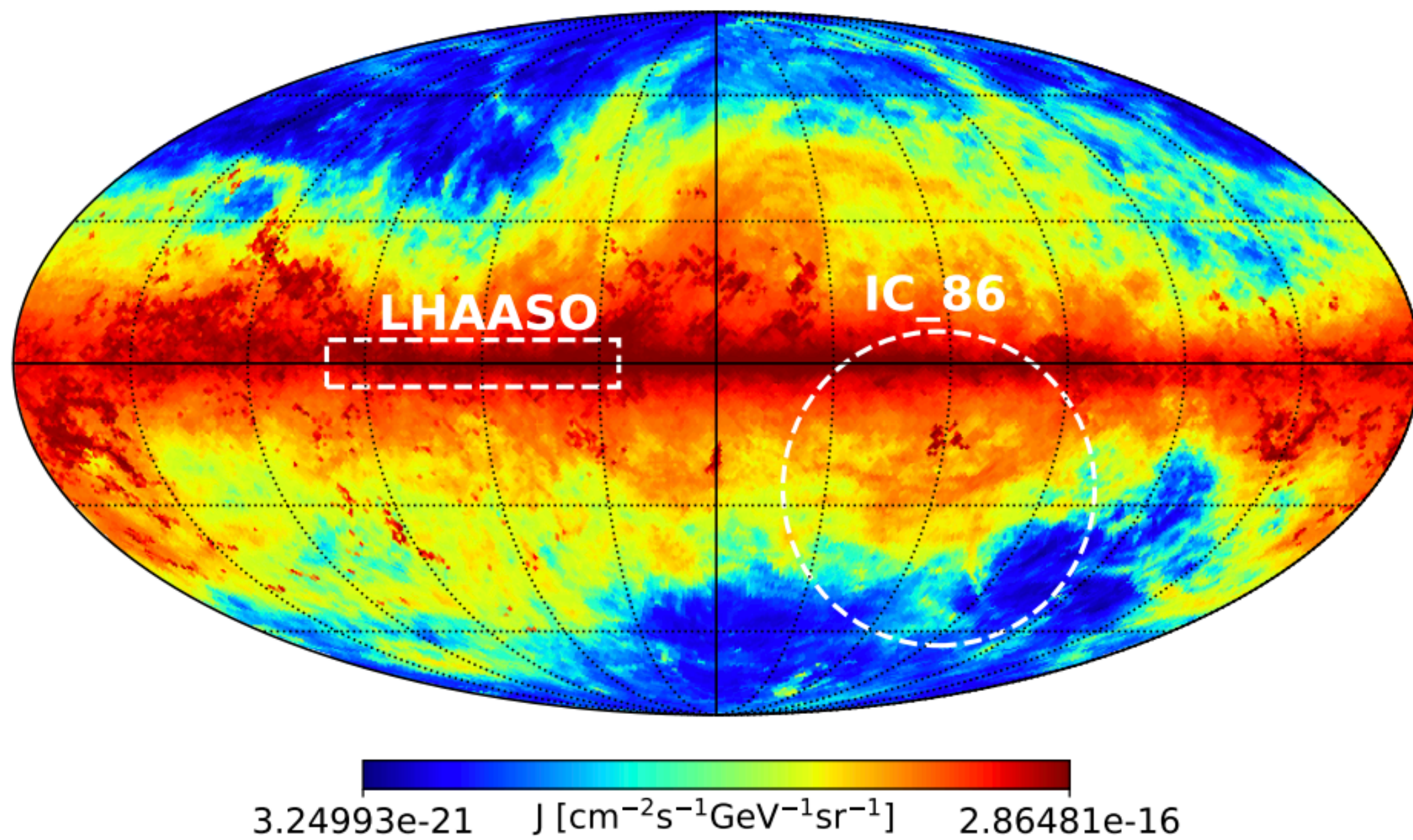
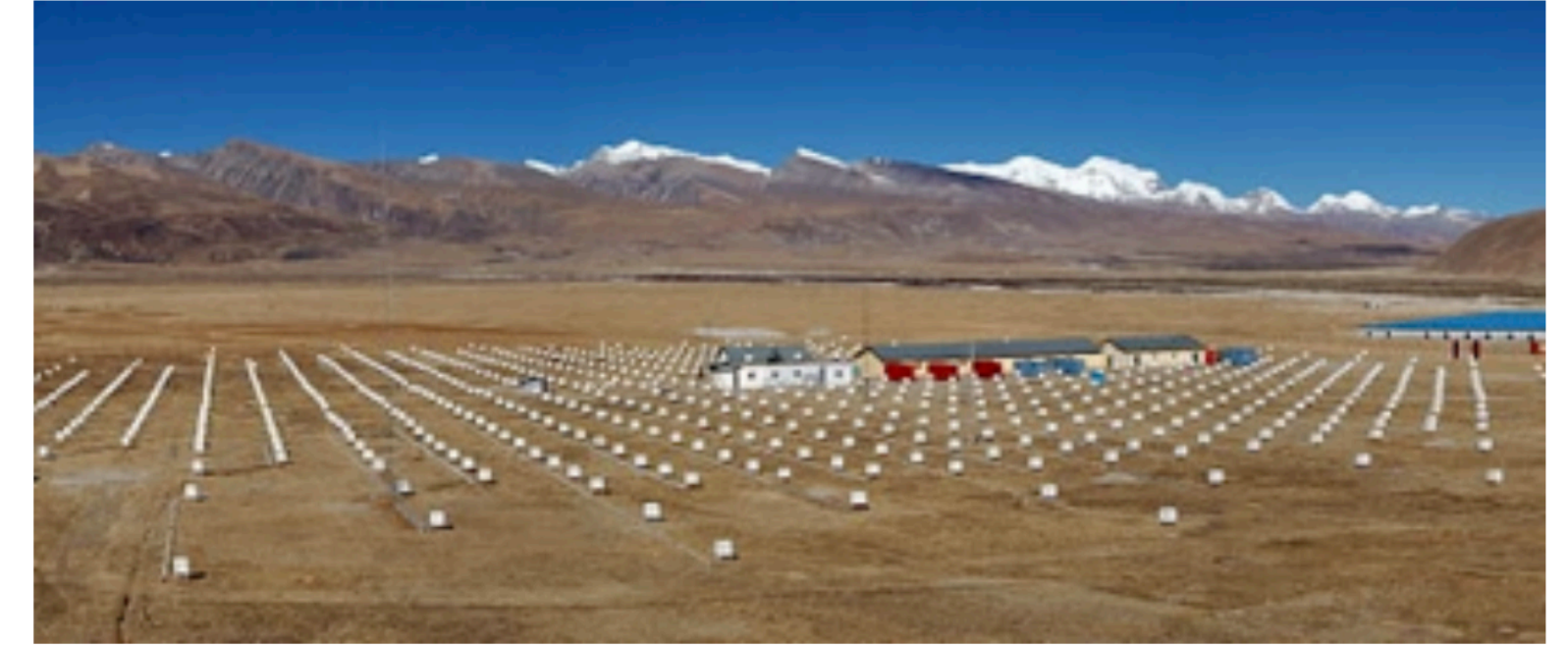


12 PeVatrons

Diffuse emission: from TeV to sub-PeV

First Detection of sub-PeV Diffuse Gamma Rays from the Galactic Disk: Evidence for Ubiquitous Galactic Cosmic Rays beyond PeV Energies

M. Amenomori *et al.* (Tibet AS $_{\gamma}$ Collaboration)
Phys. Rev. Lett. **126**, 141101 – Published 5 April 2021



Sources of cosmic particles: Exotic (astro)physics

Why do we want to look for new physics?

1. Physics BSM is strongly motivated
2. It already happened in the past to find some surprise in astro data which led to major discoveries (e.g. neutrino oscillations)
3. Anomalies in astroparticle observables do exist and BSM physics may be the answer
4. You expect astrophysical signatures of BSM physics. If you don't find them you can set constraints onto the relevant parameter space

Astroparticle observables can be either a **discovery** tool or a **constraint-setting** one

Challenges

Unusual scales of density, temperature, size, time, energy... if compared with what achievable in Earth laboratories!

Conceivable that some physics extrapolations may fail, highlighting new phenomena/regimes

We do not control the **environment**. This requires effort to understand astrophysics to devise 'robust' new physics **signatures** and validation tests

Main goal: Help you to schematise some research direction in **dark matter/BSM** searches

A poor phenomenologist perspective

- From a fundamental theory or more EFT approaches couplings with SM and new particles induce specific phenomenological signatures
- We are interested here in **signatures of dark matter particles in high-energy astrophysics**

What I consider as HE astro?

- Emission of **astrophysical sources** from X (\sim keV) to gamma rays (PeV)
- Astrophysical **diffuse backgrounds** at multiple wavelengths

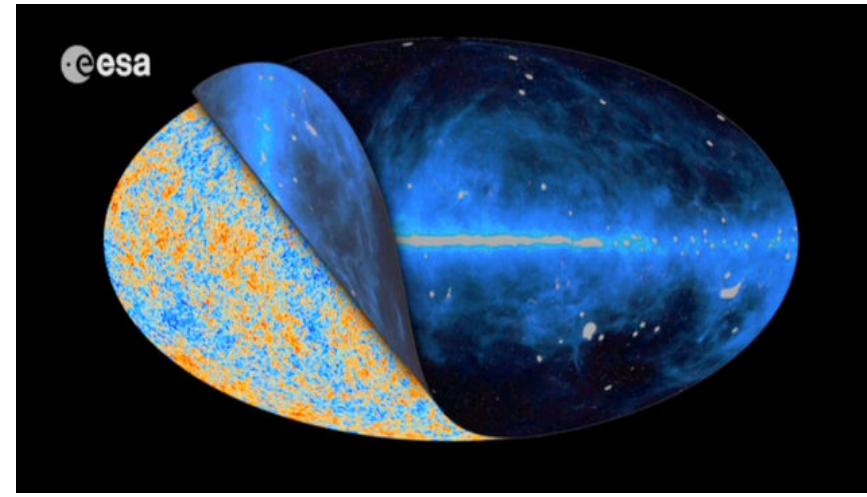
What is a signature?

- It is a **sizeable modification** of *standard* astrophysical signals
- Modification of spectral and spatial distributions of event counts, and/or polarisations

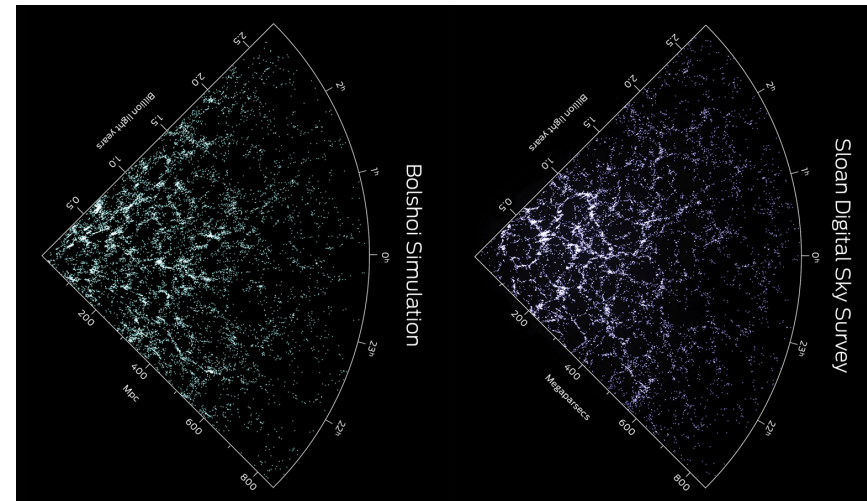
Detection vs constraints

Dark matter astro/cosmo evidence

Cosmic microwave background



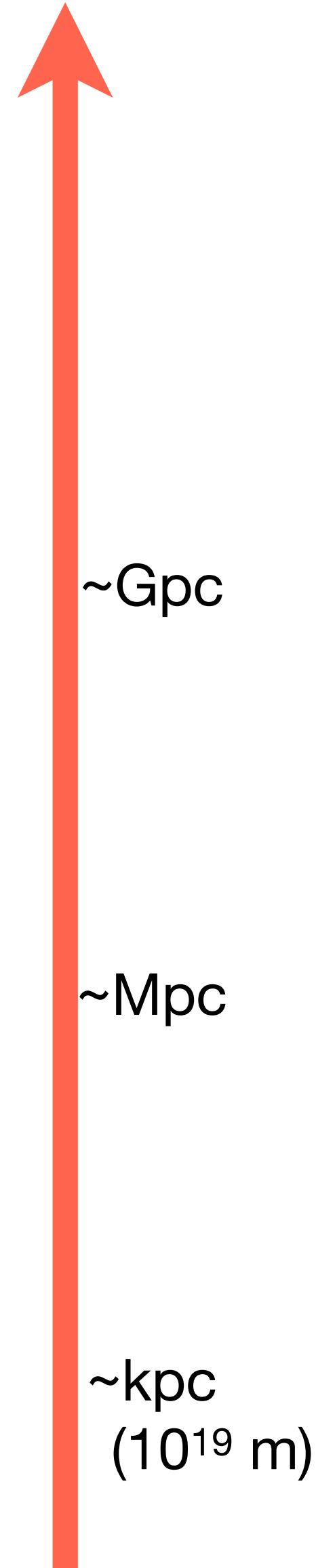
Large Scale Structures



Galaxy clusters



Galaxies



Astro/cosmo observations at multiple scales are instrumental in showing that an unknown ingredient of matter in the universe exists

Precise measurement of how much DM is present at cosmological scales (CMB)

$$\Omega_{\text{CDM}} \sim 0.26$$

Planck 2018, 68% CL

Inferred properties that DM should possess point to **new physics** beyond standard theories

The dark matter landscape

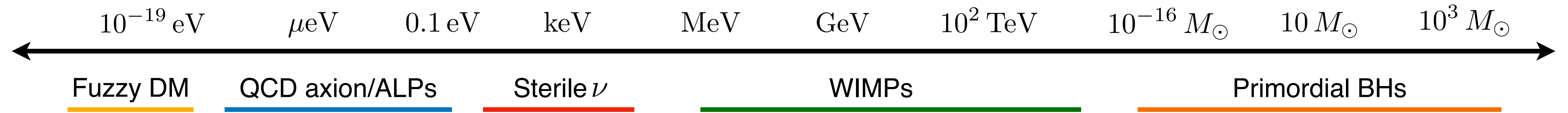


Vast parameter space in mass and interaction strength

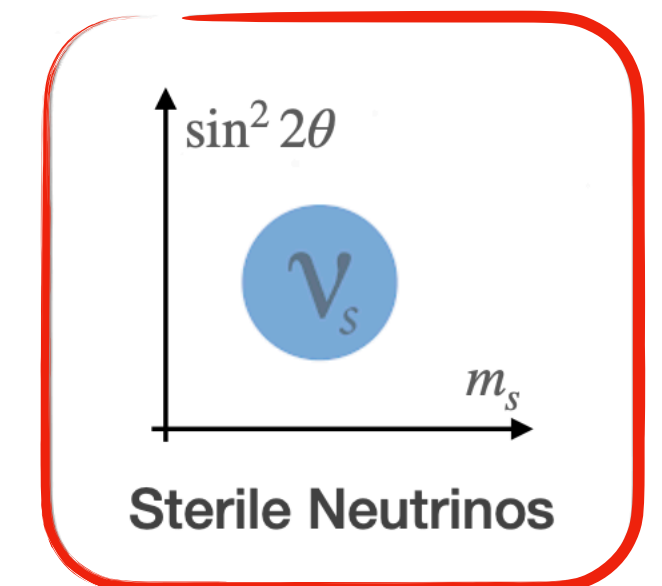
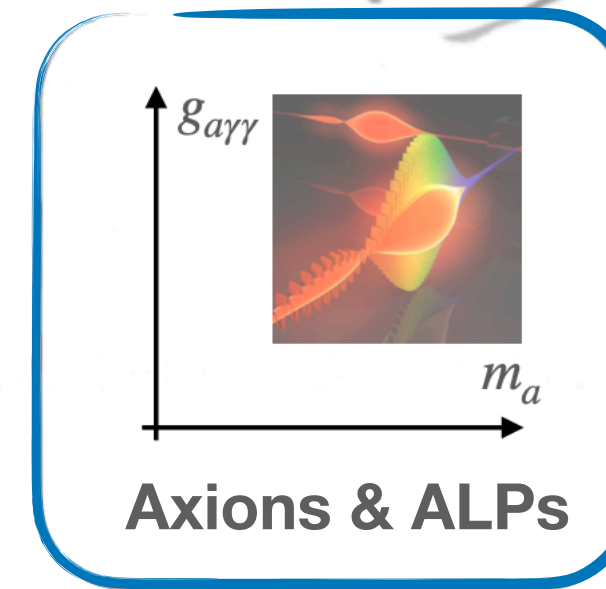
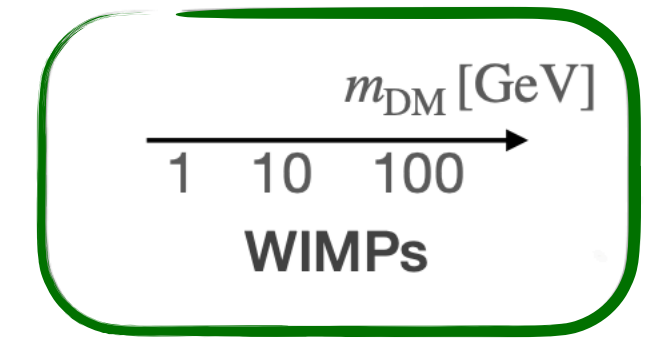
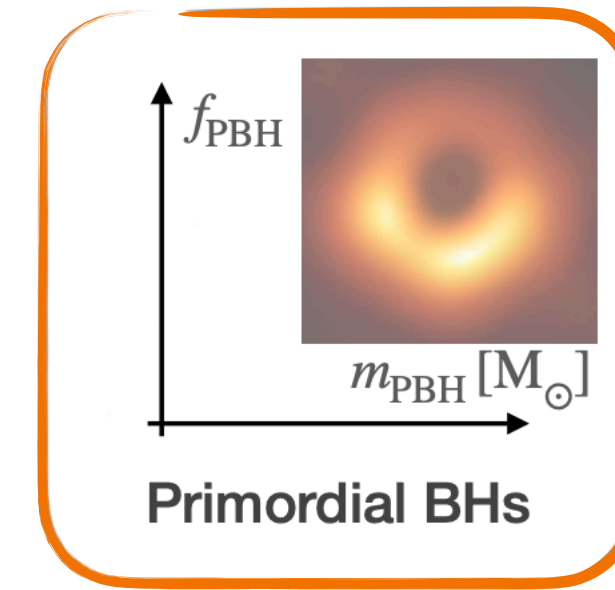
The dark matter landscape



The dark matter landscape



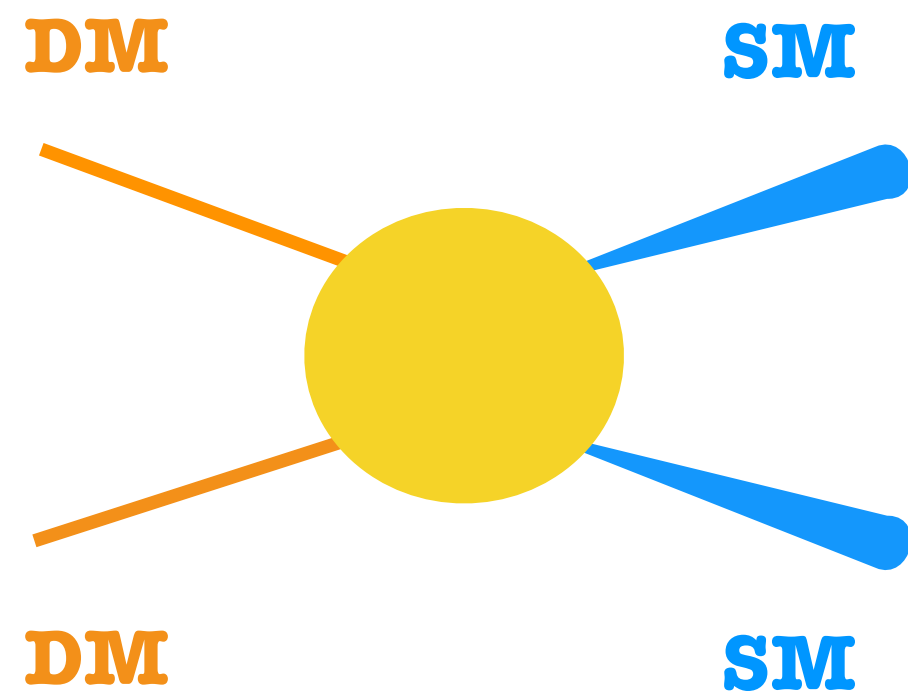
- Identification strategies are necessarily (more or less) **model dependent**
- The **theoretical prejudice** in dark matter searches is also set by what we can probe with available data
- You always need some sort of signature of your model!



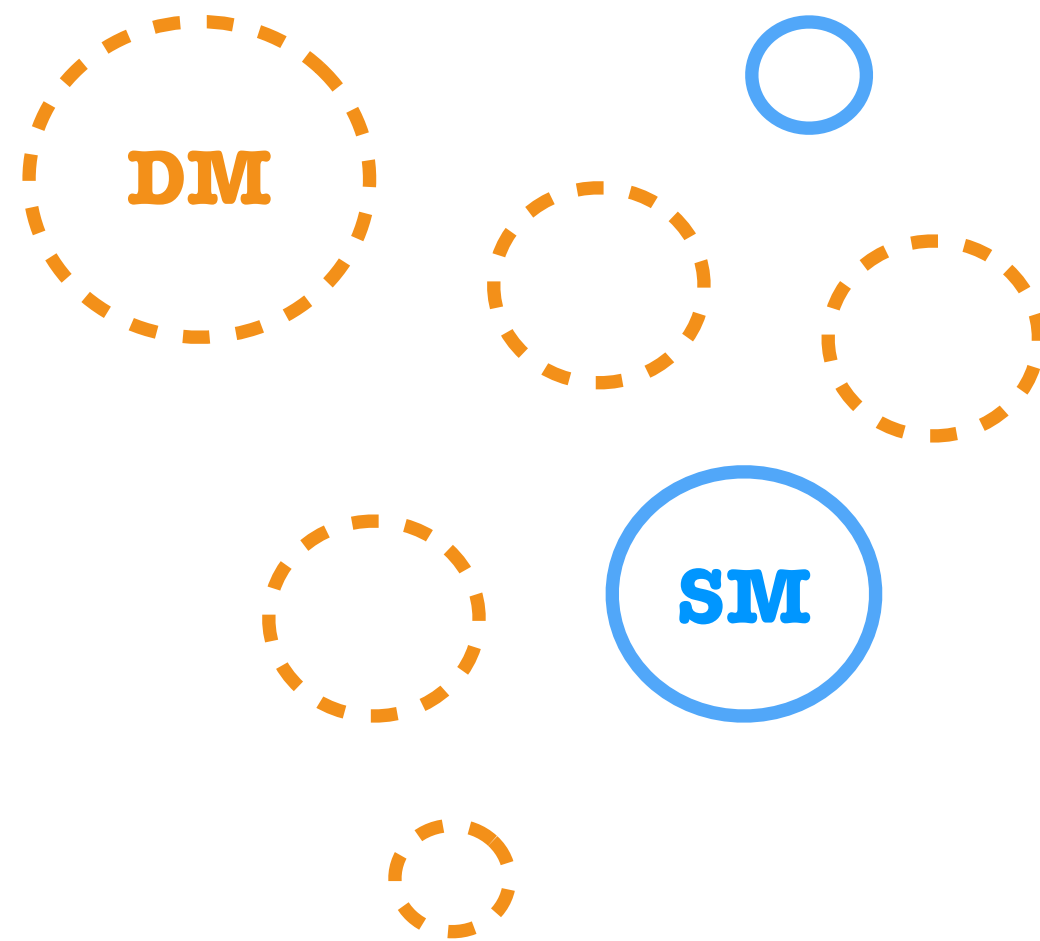
Credit: J. Alvey, EuCAPT Symposium 2021

Dark matter indirect detection

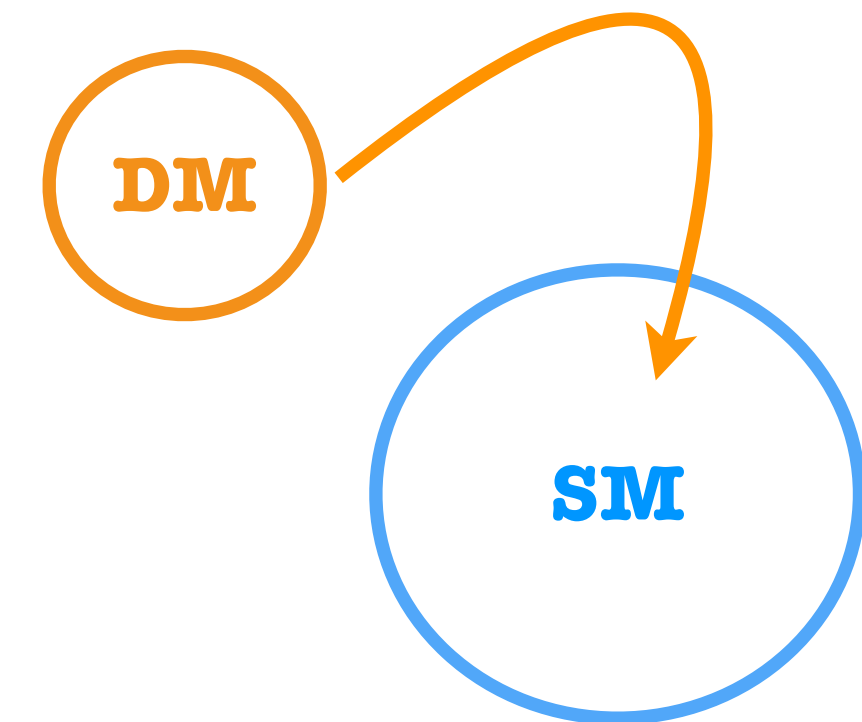
What dark matter does



Energy/particle injection



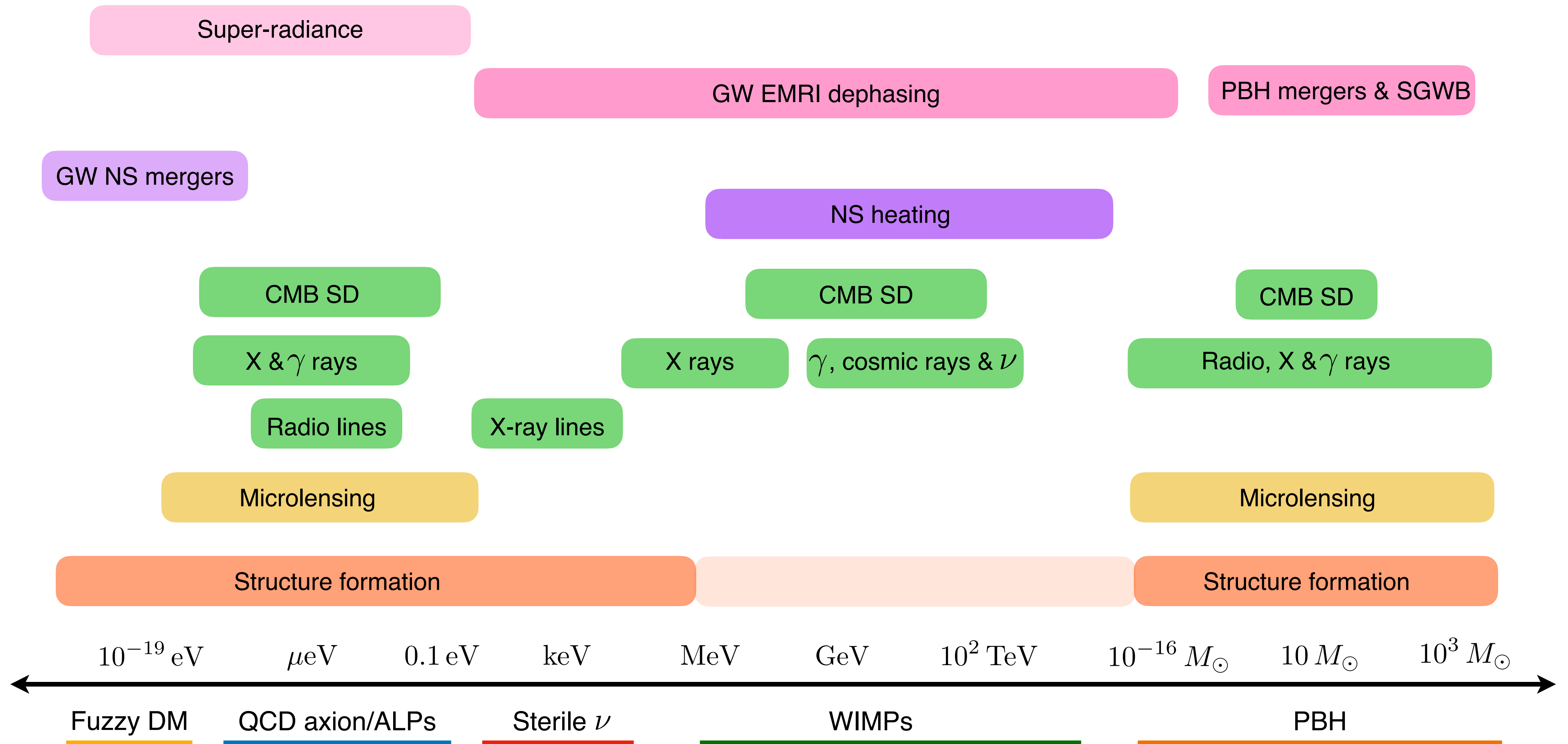
Gravitational interaction



*Capture/scattering/accretion
in/onto astrophysical objects*

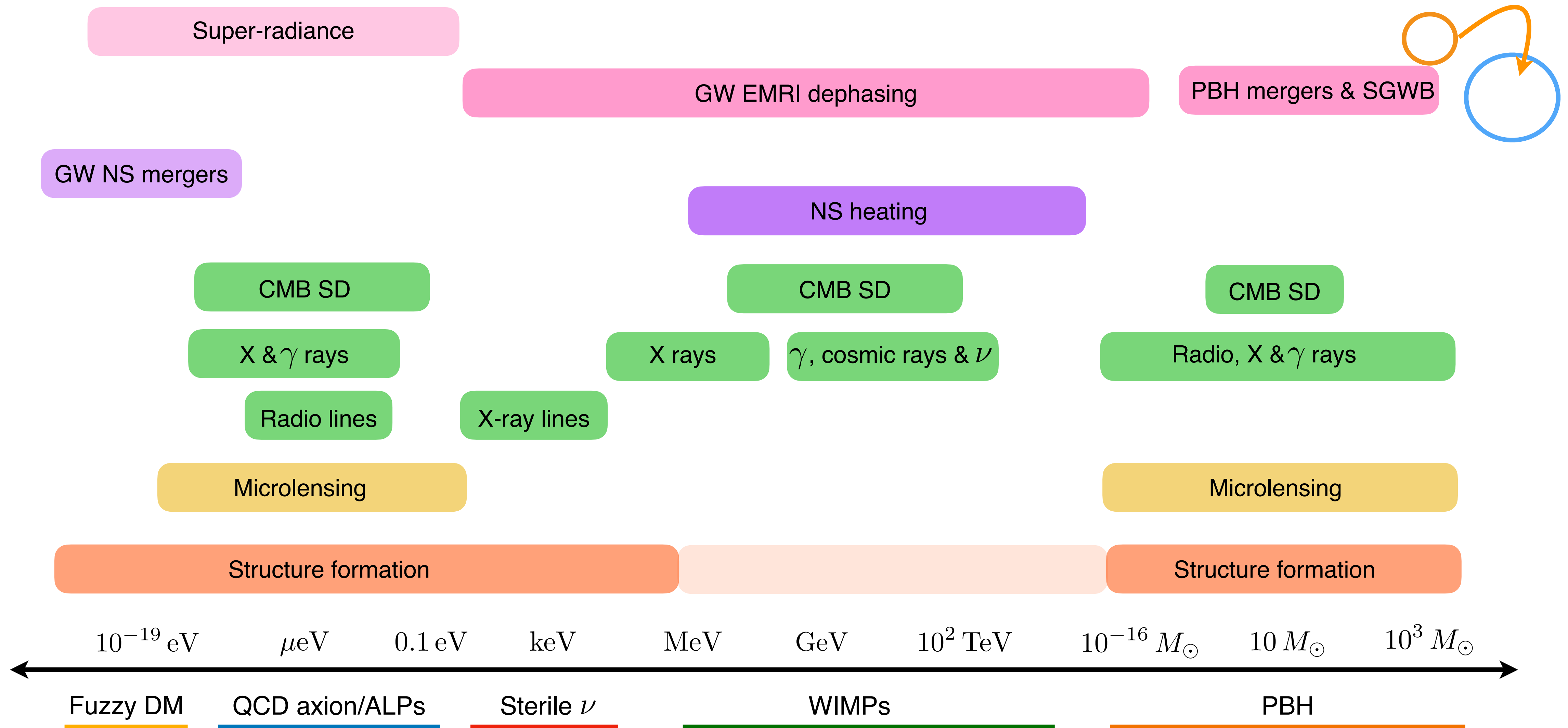
Astroparticle observables for dark matter

EuCAPT White Paper, arXiv:2110.10074



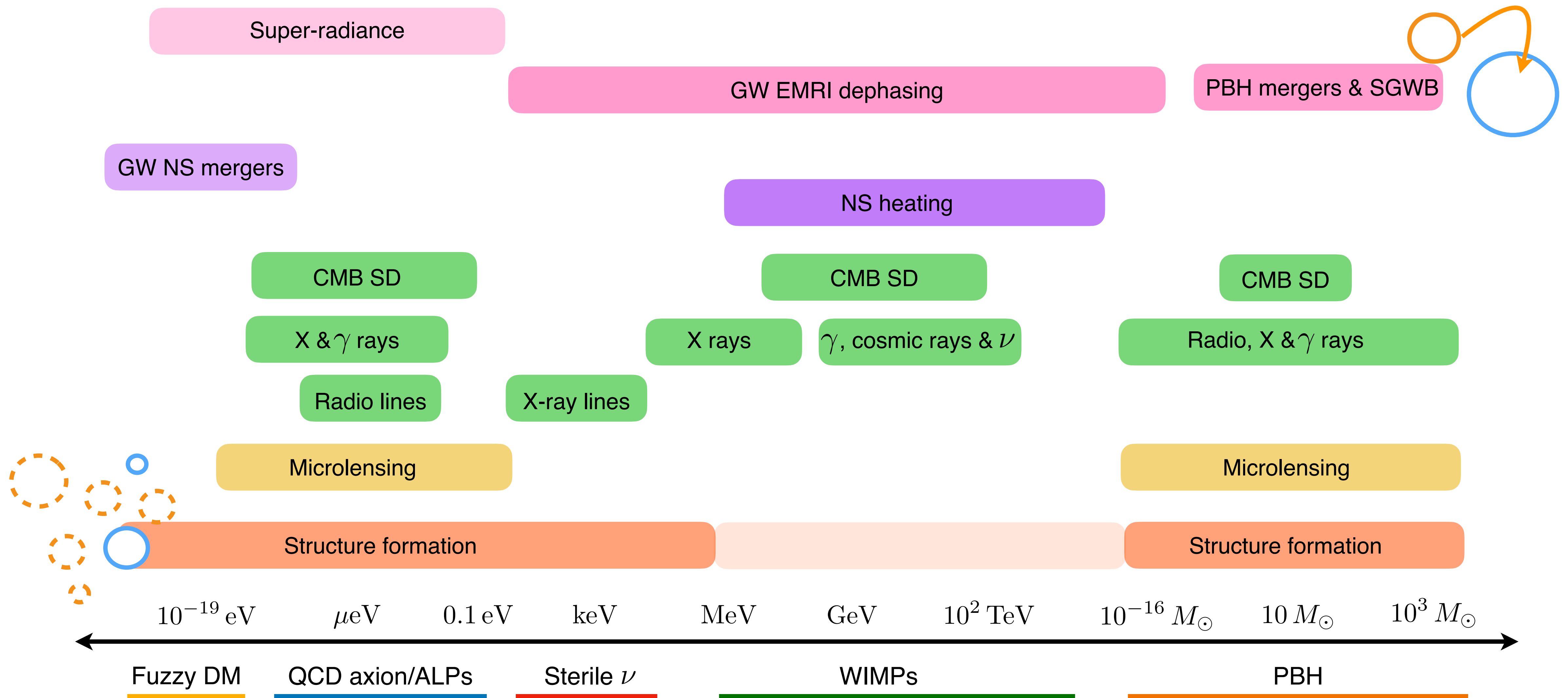
Astroparticle observables for dark matter

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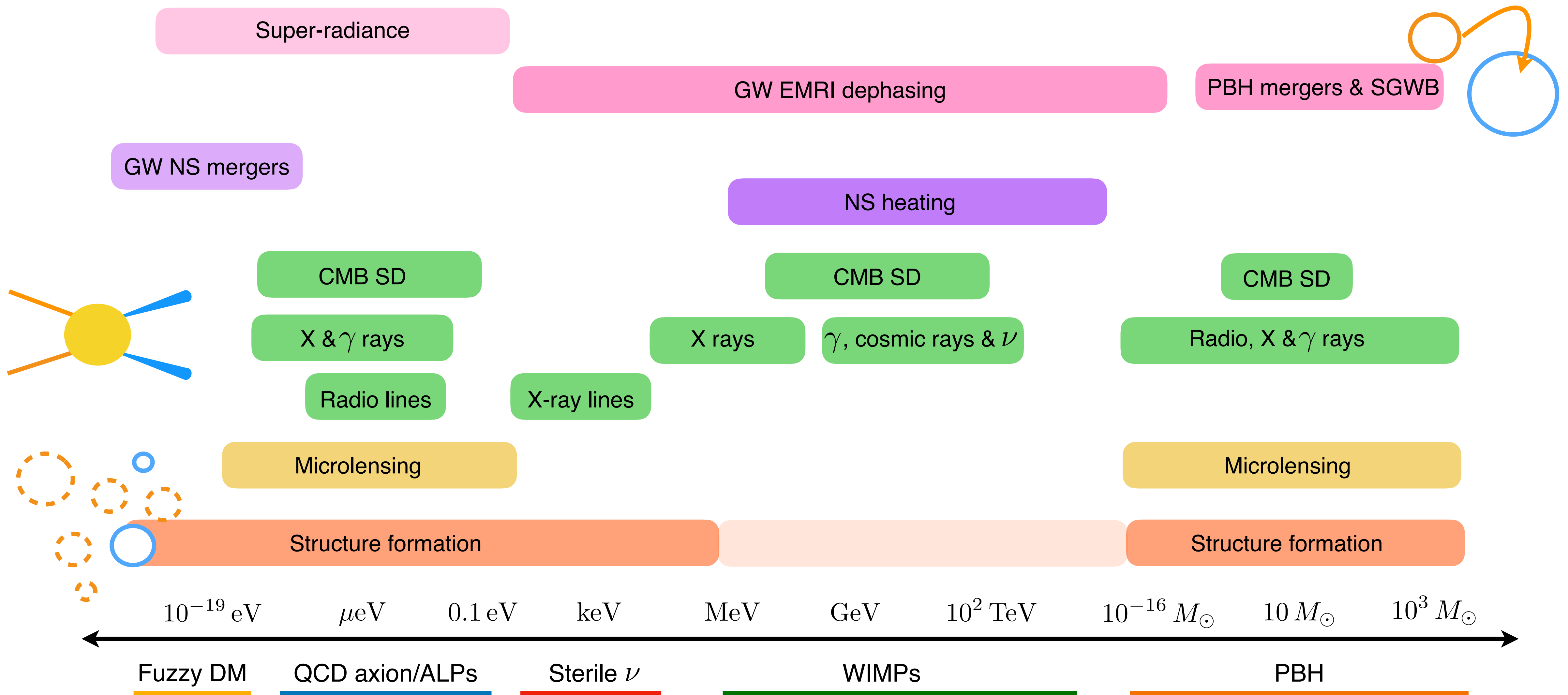
Astroparticle observables for dark matter

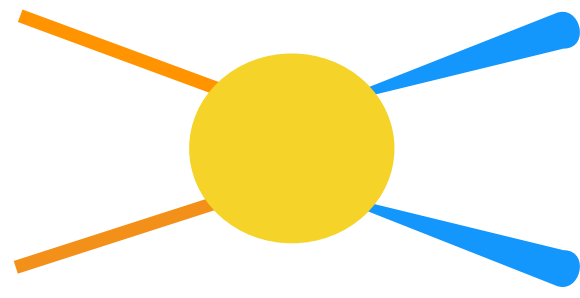
EuCAPT White Paper, arXiv:2110.10074



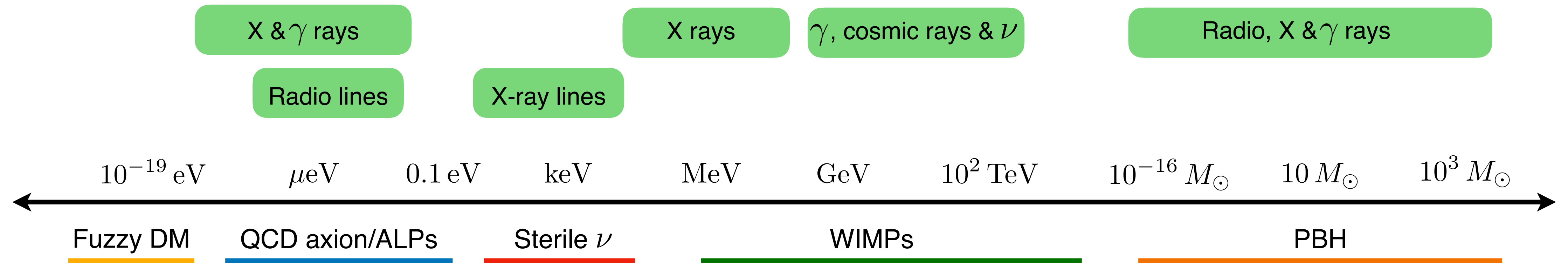
Astroparticle observables for dark matter

EuCAPT White Paper, arXiv:2110.10074





Travelling messengers



Common astrophysical ingredients

$$\Phi_{\text{DM}} \propto n_{\text{DM}}$$

Flux proportional to DM number density

$$\rho_{\text{DM}} \equiv n_{\text{DM}} \times m_{\text{DM}}$$

DM energy density for *non-relativistic* particles (CDM)

$$v_c = (218 - 246) \text{ km/s}$$

Local circular velocity

$$\sigma_v = \sqrt{3/2} v_c$$

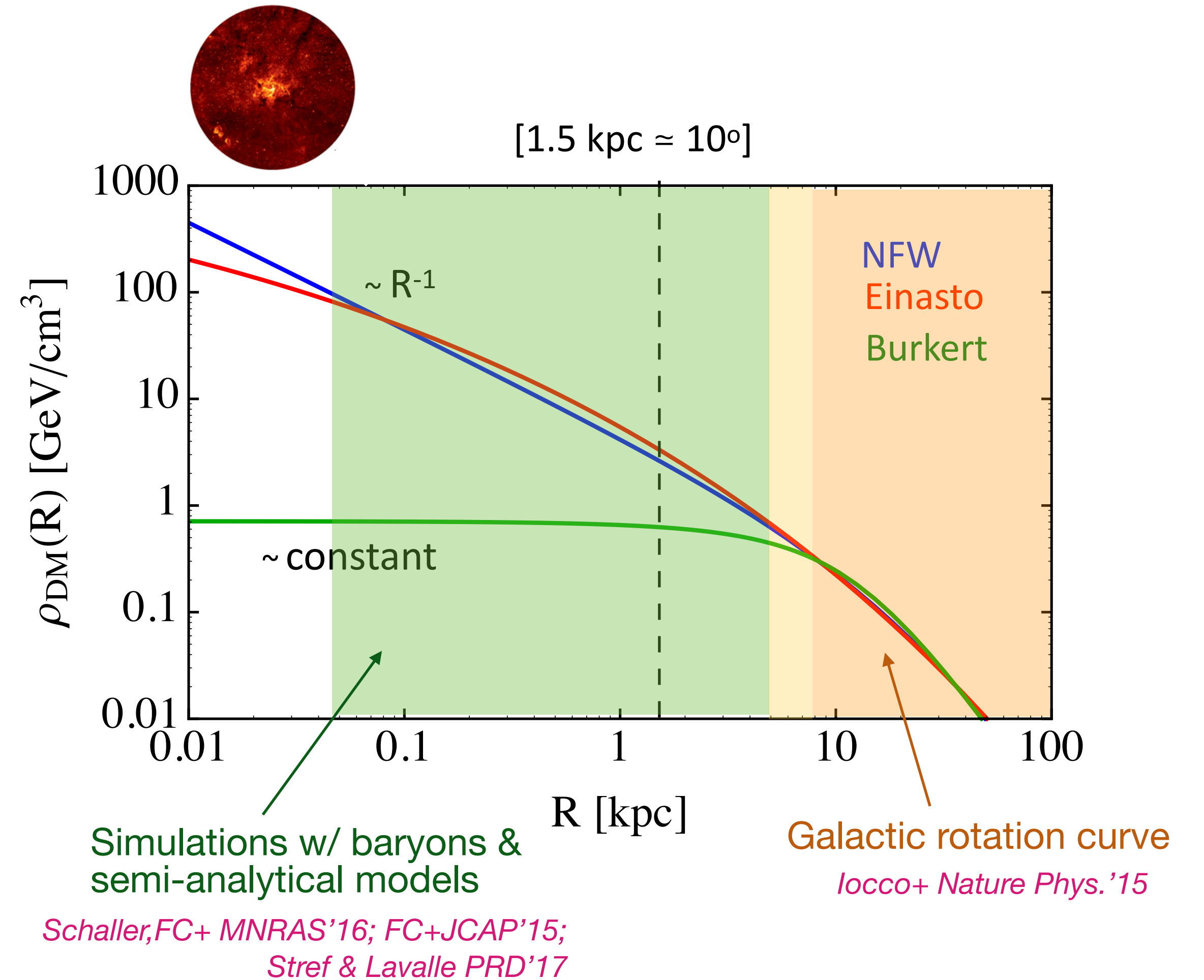
DM velocity dispersion
(Maxwell-Boltzmann
distribution)

=> Decay or self-interactions of DM in the halo of galaxies at $z=0$ occur at rest

Common astrophysical ingredients

Halo dark matter density

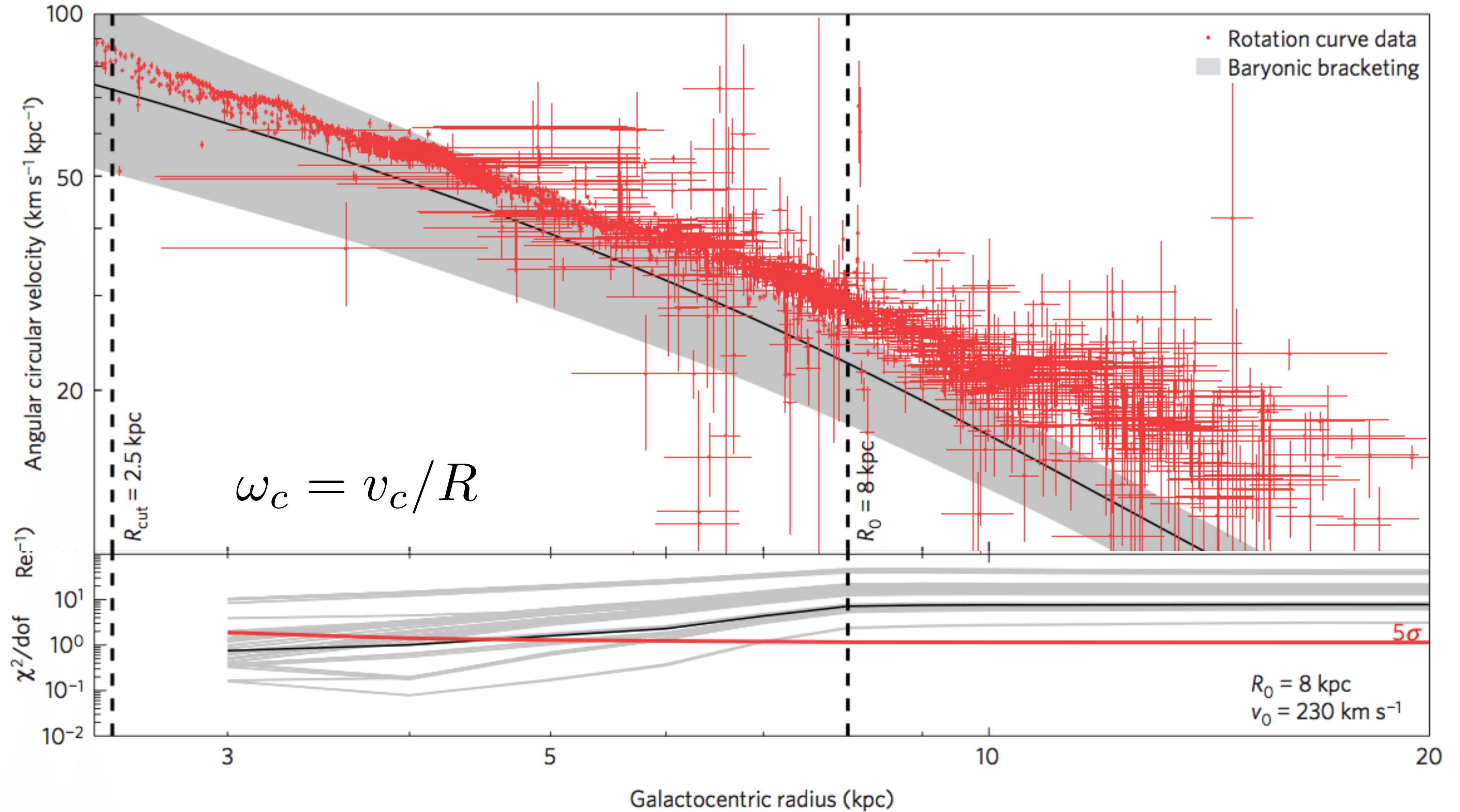
- The distribution of DM in galaxies is affected by large uncertainties
- **Unavoidable modelling uncertainty** to account for when deriving constraints on DM models
- Rely on simulations of galaxy formation to get small-scale density profile and sub-halo distribution
- Less of an issue in dwarf galaxies



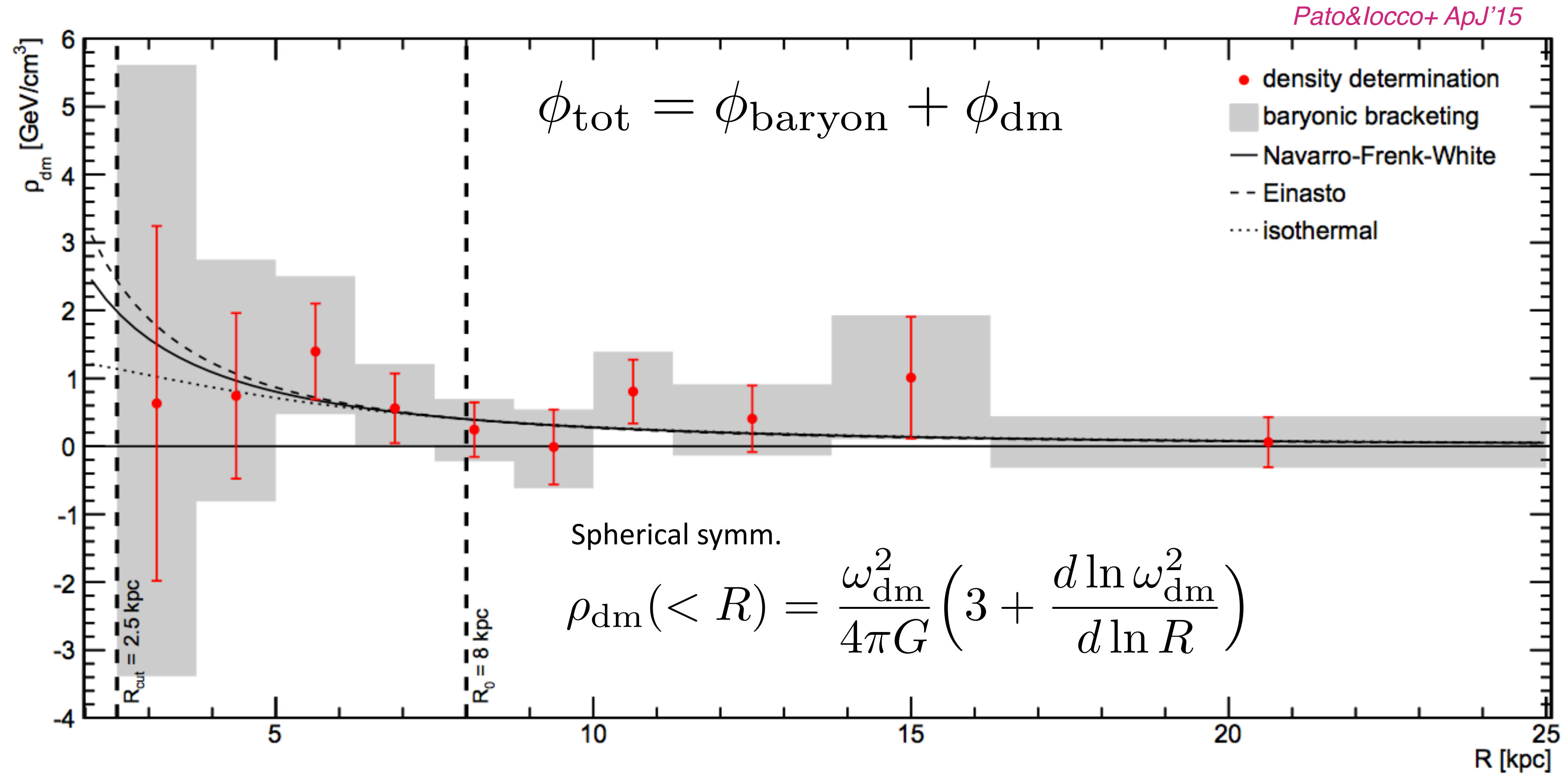
[All profiles normalised by measure of **local DM density** at R_{sun}]

Milky Way rotation curve

locco+ Nature Physics'15

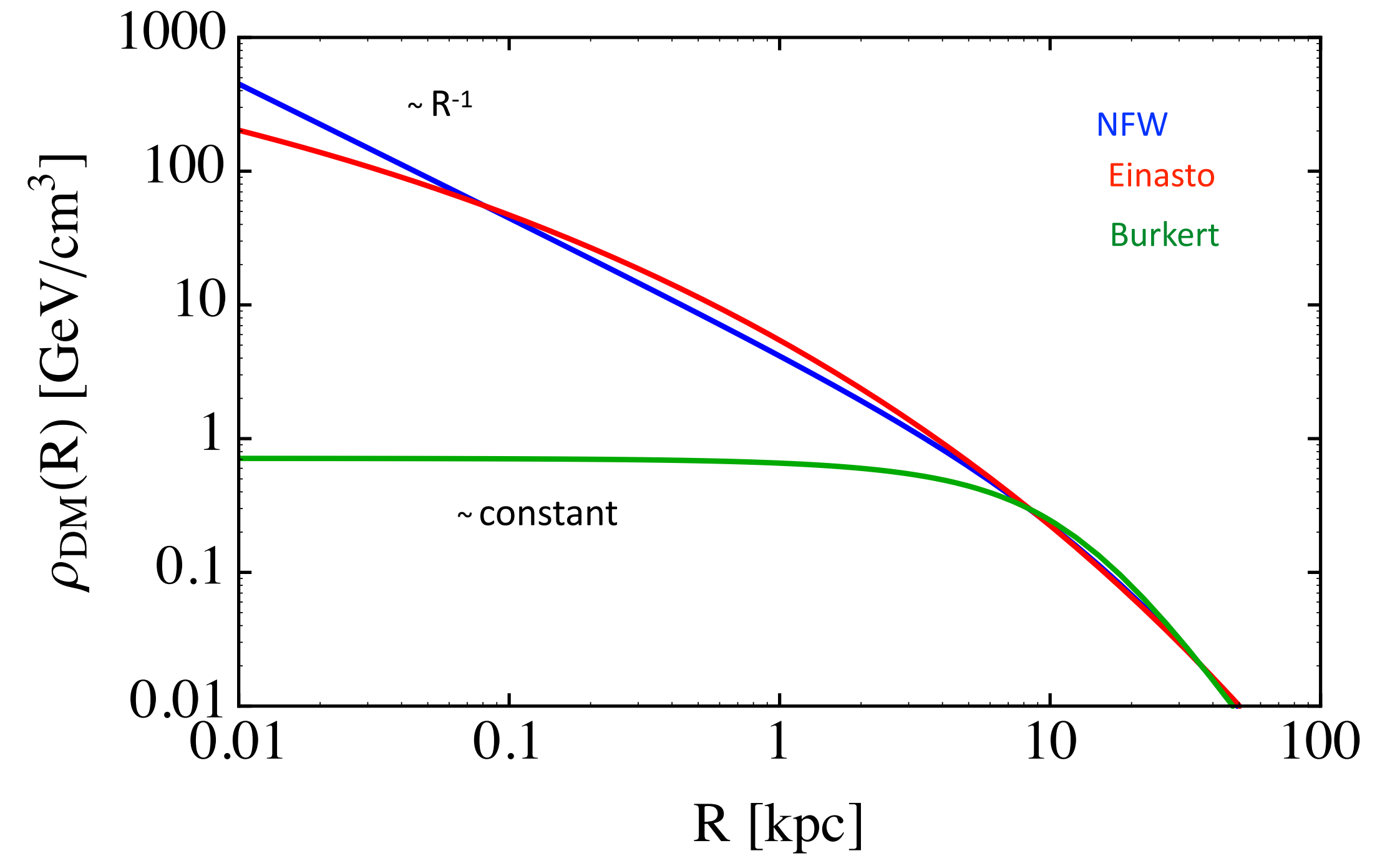
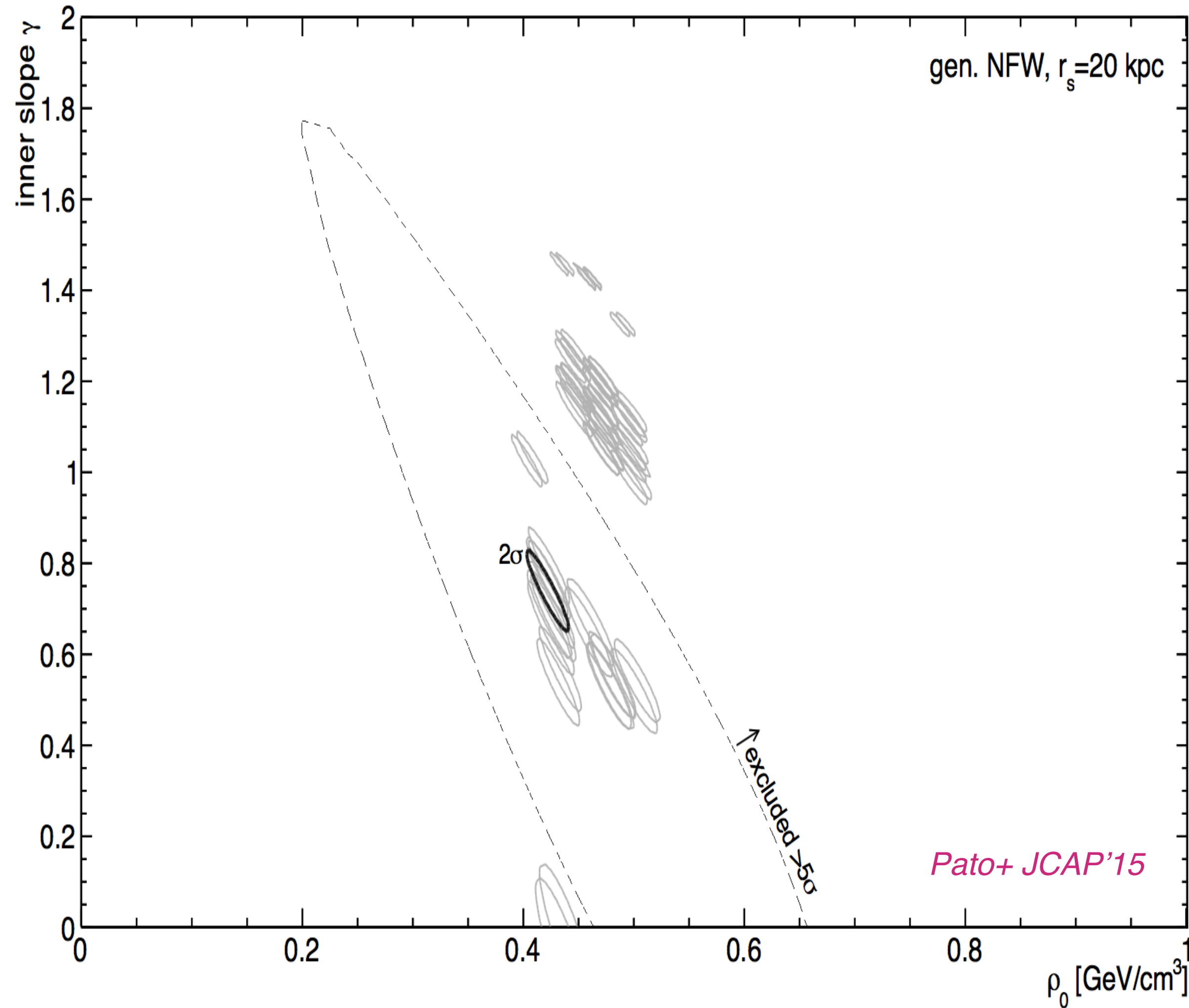


Milky Way DM density



Non-parametric reconstruction: approach free of profile assumptions, but uncertainties are large and hinder discrimination power between different radial behaviours.

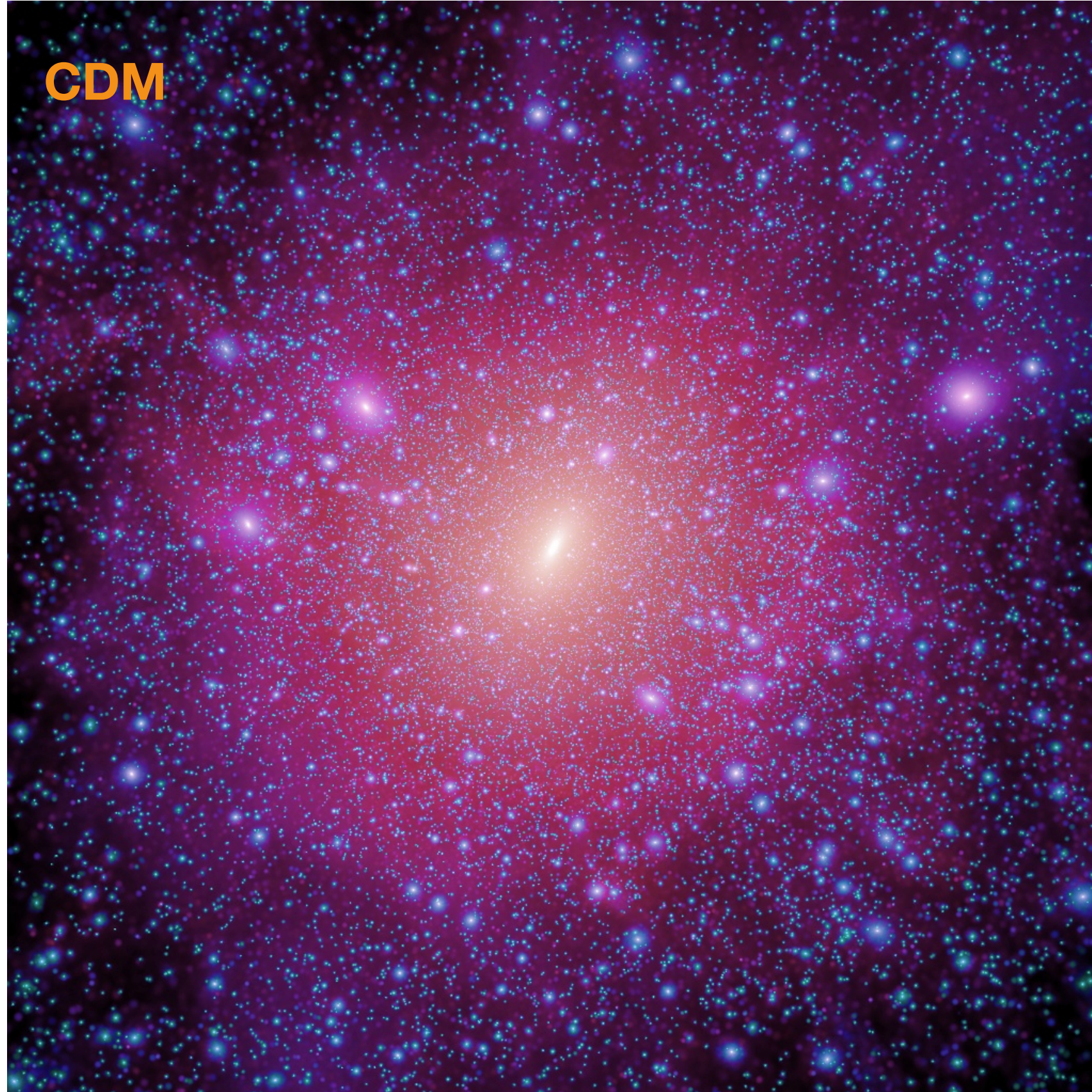
Milky Way DM density



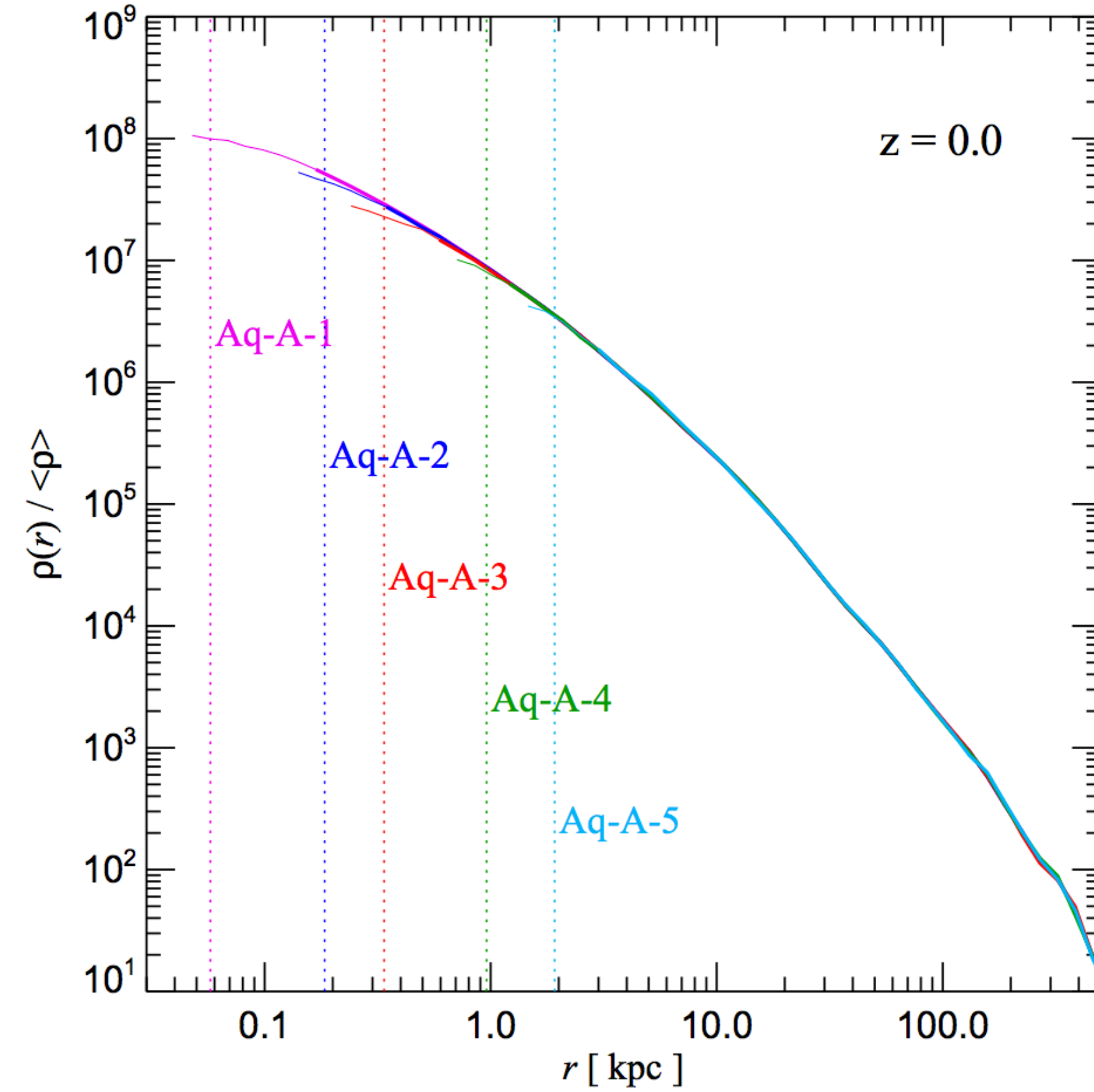
Parametric reconstruction: strong profile assumptions, “global” method to derive local DM density.

e.g: Pato+ JCAP'15; McMillan+ MNRAS'16; Iocco&Benito PDU'17

DM density in haloes



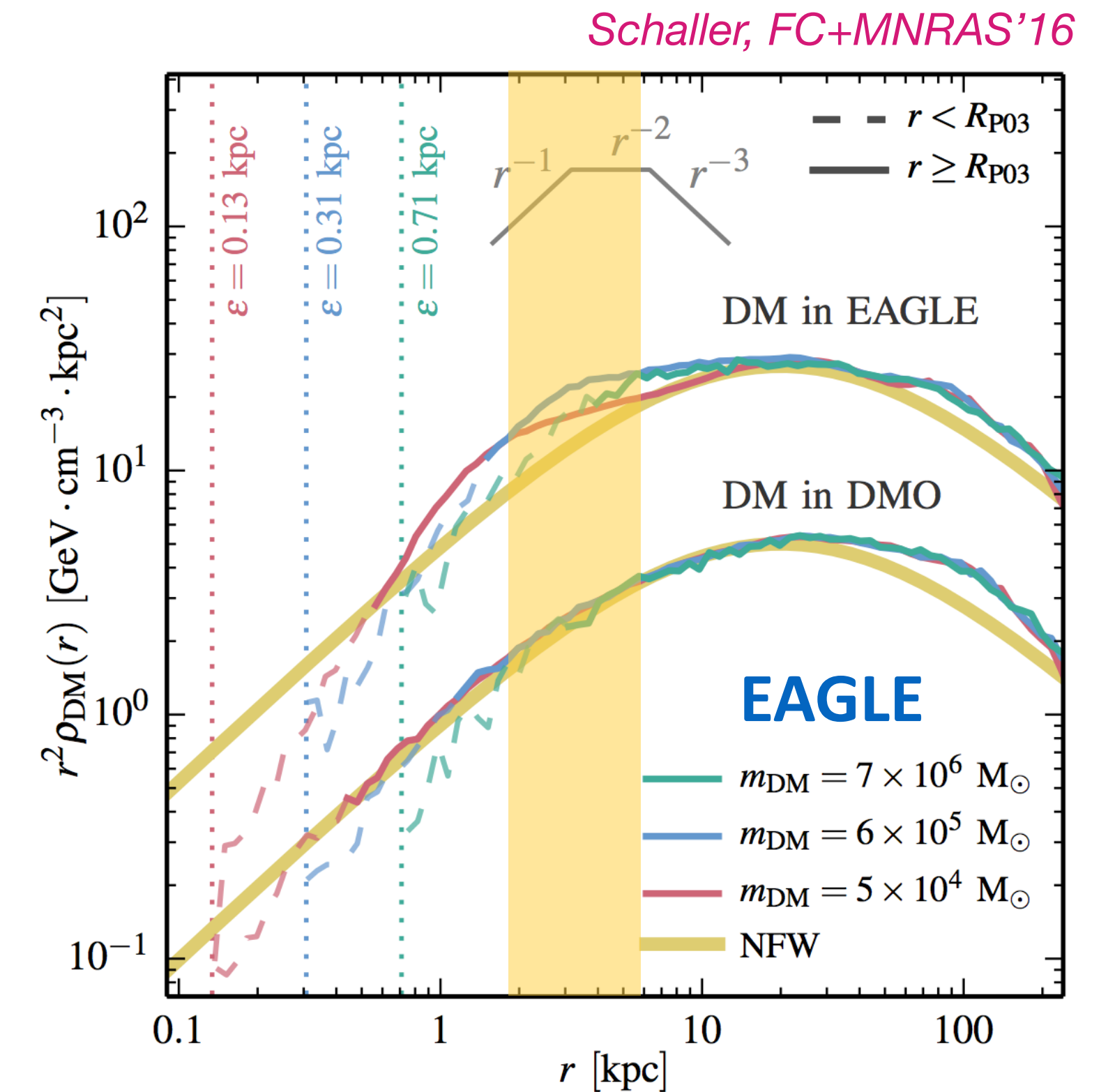
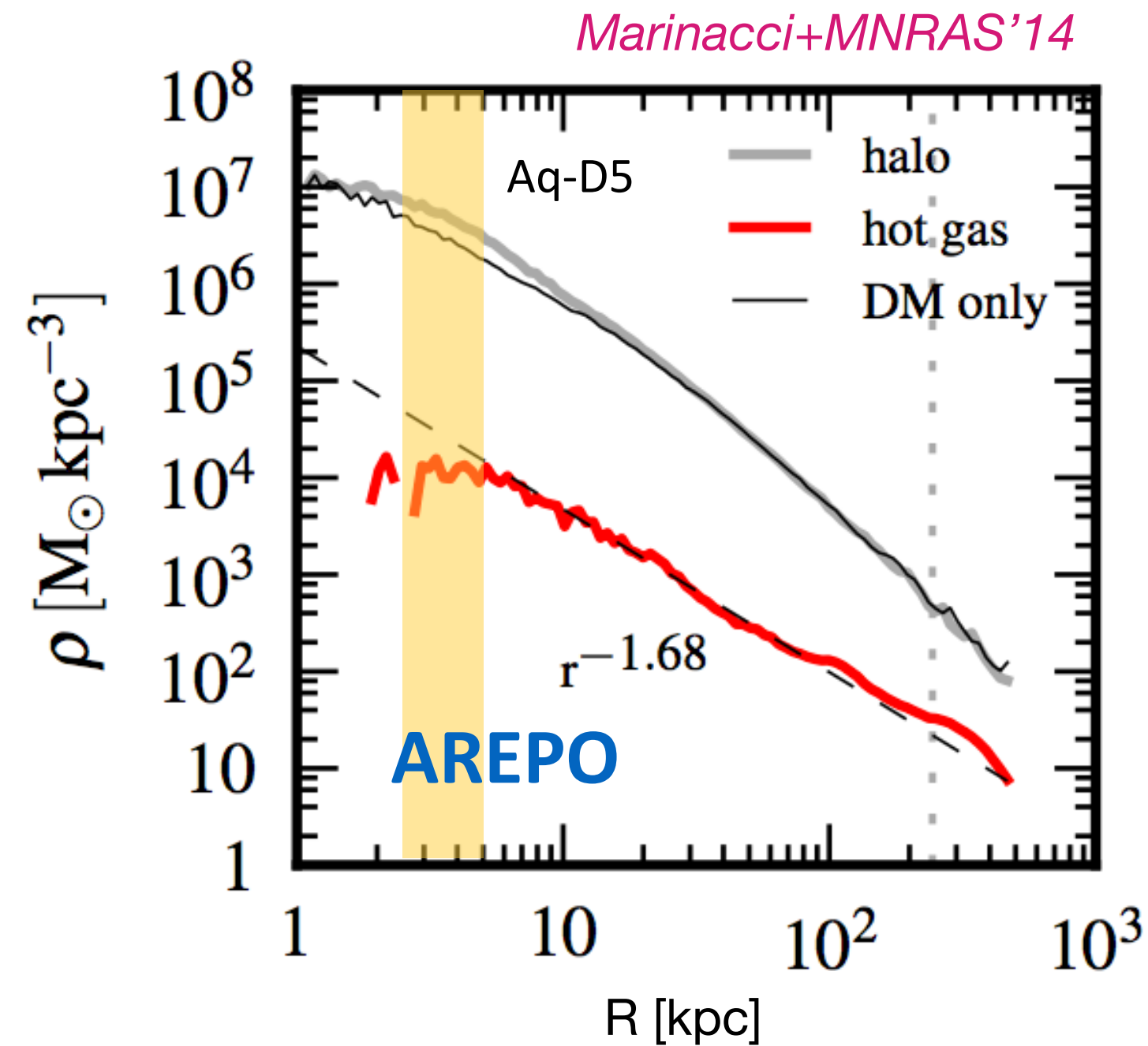
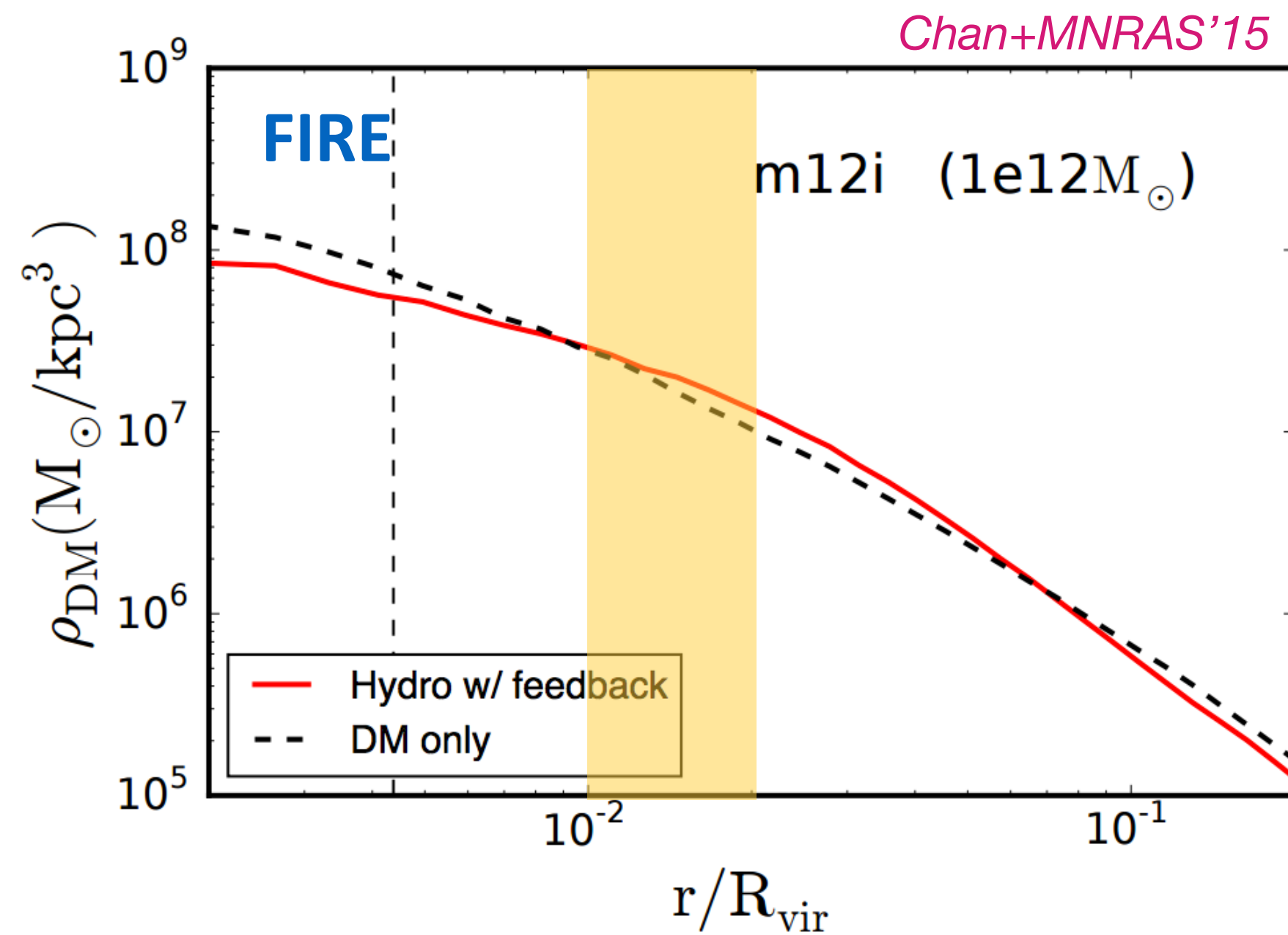
Springel+ MNRAS'08



$$\rho_{\text{tot}}(R) = \rho_s \exp \left\{ -\frac{2}{\alpha} \left[\left(\frac{R}{r_s} \right)^\alpha - 1 \right] \right\}$$

Common astrophysical ingredients

Halo dark matter density: Inner galaxy



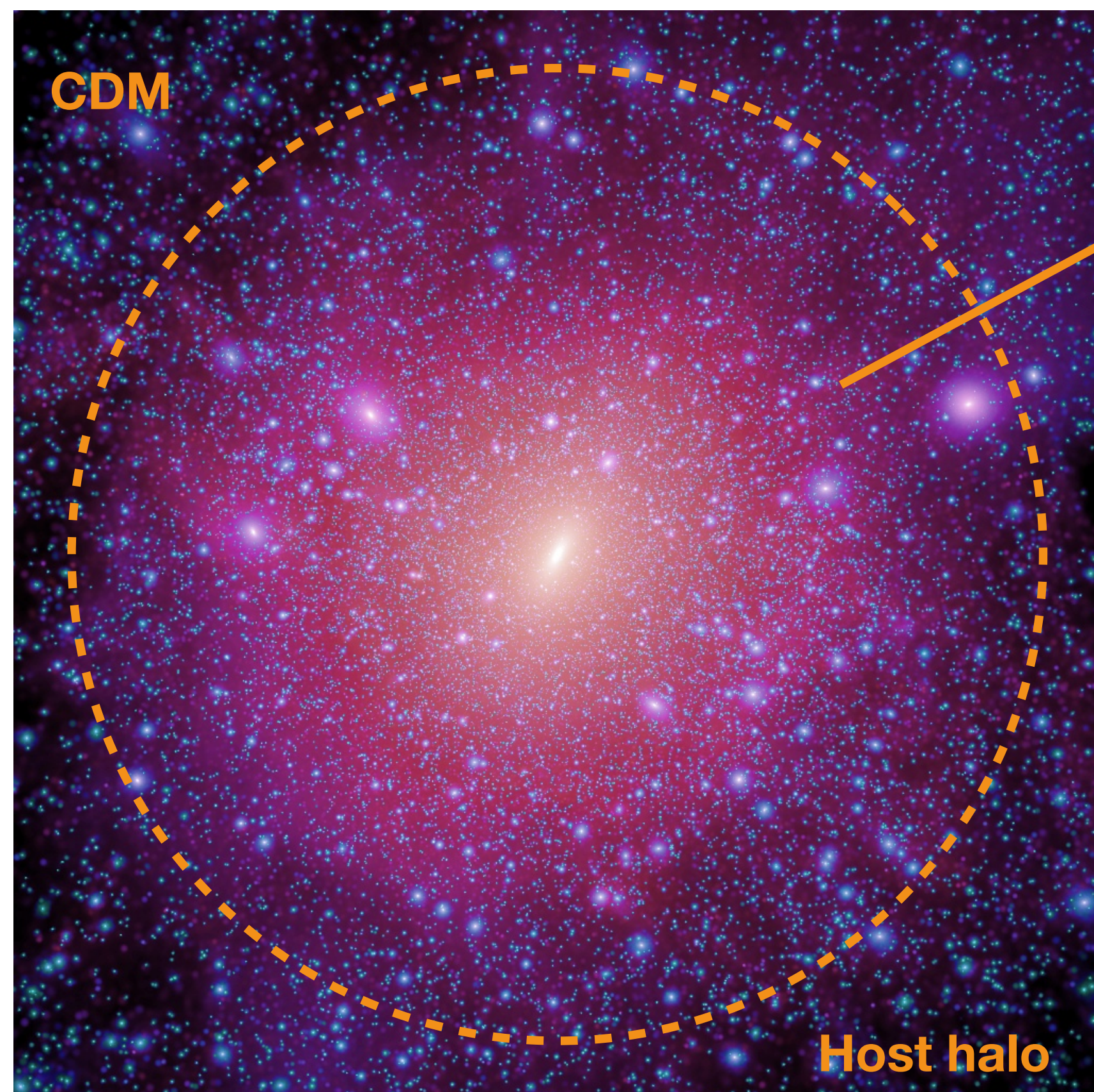
Challenges

- How to select a good MW halo candidate?
- How to extrapolate the profile down to smaller scales?

Wang+MNRAS'12; Gottloeber+2010; FC+JCAP'15

Common astrophysical ingredients

Dark matter substructures



Springel+ MNRAS'08

Dark matter sub-haloes or sub-structures

- Low-mass DM haloes, do not trigger star formation => do not contain stars (**dark haloes**)
- Their **mass distribution** depends on fundamental properties of DM (warm vs cold)
- CDM predicts abundance down to Earth-sized objects ($10^{-6} M_{\text{Sun}}$)*
- Their distribution leads to specific **angular signatures** of the DM signal (sub-halo searches, anisotropies, x-correlation)

* In WIMP and non-thermal axion models

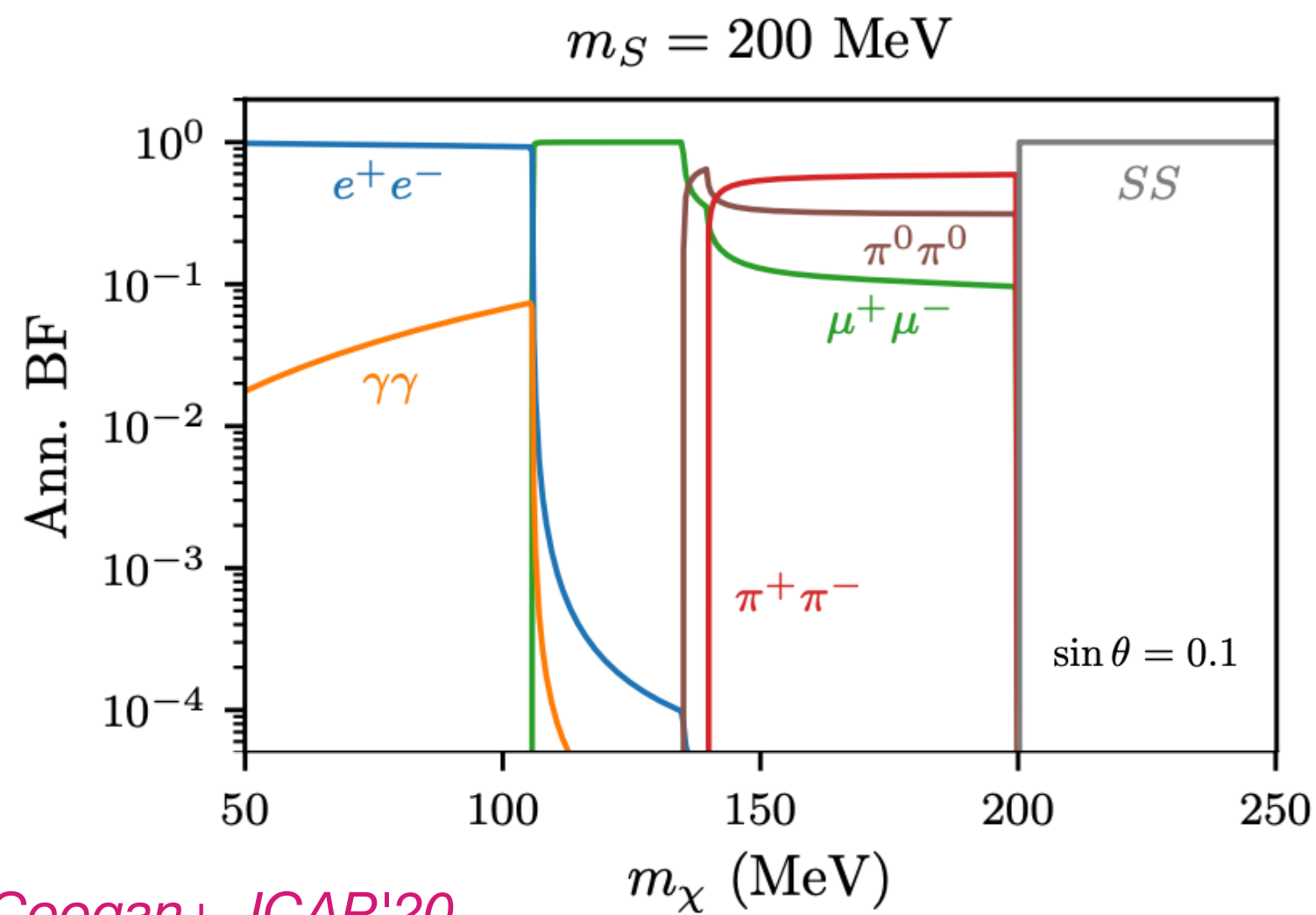
Grand & White MNRAS'21

Particle dark matter

Particle dark matter emission

$$(DM) DM \rightarrow SM SM$$

$$E_{\text{CM}} = N m_{\text{DM}}, \quad N = 1 \text{ (decay)}, 2 \text{ (annih)} \quad \text{Centre of mass energy} \simeq \text{Signal energy}$$



Coogan+ JCAP'20

$$m_{\text{DM}} \lesssim \text{MeV}$$

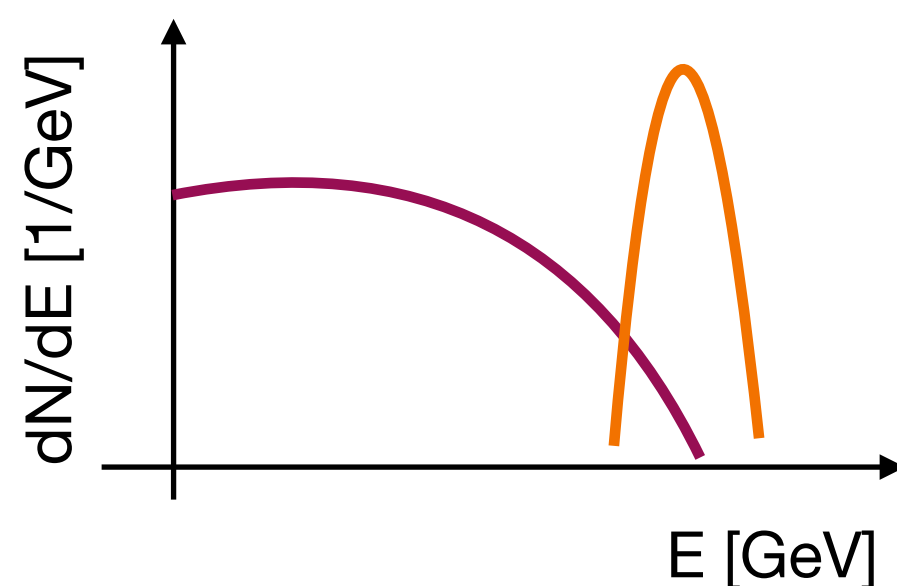
Narrow line signal

$$m_{\text{DM}} \gtrsim \text{MeV}$$

$$E_\gamma = \frac{N m_{\text{DM}}}{2}$$

$$\frac{dN_\gamma}{dE} = 2\delta\left(E - \frac{N m_{\text{DM}}}{2}\right)$$

Broader energy distribution



$$\frac{dN_\gamma}{dE} = \left(\frac{dN_\gamma}{dE}\right)_{\gamma\gamma} + \left(\frac{dN_\gamma}{dE}\right)_{\text{sec}} + \left(\frac{dN_\gamma}{dE}\right)_{\text{FSR}}$$

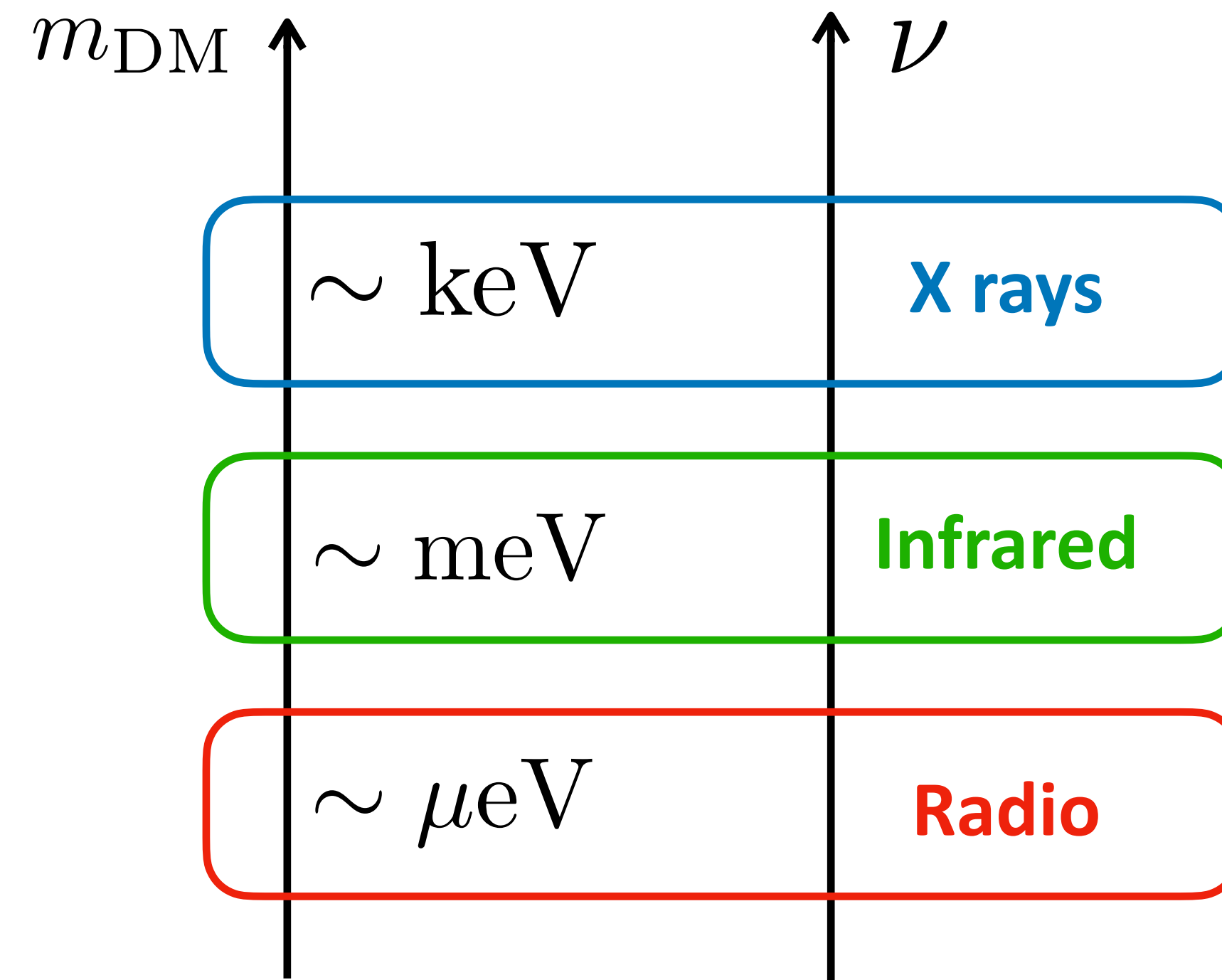
Light dark matter decay: lines

$$m_{\text{DM}} \lesssim \text{MeV}$$

Only allowed final state is into photons emitted back-to-back

$$E_{\gamma} = \frac{m_{\text{DM}}}{2}$$

Narrow line signal @ energy scale of the DM mass



Light dark matter decay: Emissivity and flux

$$Q_{\gamma}^{\text{dec}}(r, E) = \Gamma_{\gamma} \times n(r) \times \frac{dN_{\gamma}}{dE}(E)$$

$$n(r) = \frac{\rho(r)}{m}$$

Local emissivity per unit energy
and volume

$$[\text{GeV}^{-1}\text{s}^{-1}\text{cm}^{-3}]$$

$$\frac{Q_{\gamma}(r, E)}{4\pi}$$

Emissivity per sr (isotropic emission)

$$[\text{GeV}^{-1}\text{s}^{-1}\text{cm}^{-3}\text{sr}^{-1}]$$

$$\Phi_{\gamma}(E) \equiv \frac{d\Phi_{\gamma}}{dE} = \int_{\Delta\Omega} d\Omega(\ell, b) \int_{\text{l.o.s}} ds \frac{Q_{\gamma}(r(s, \ell, b), E)}{4\pi}$$

Gamma-ray (or neutrino)
differential energy flux

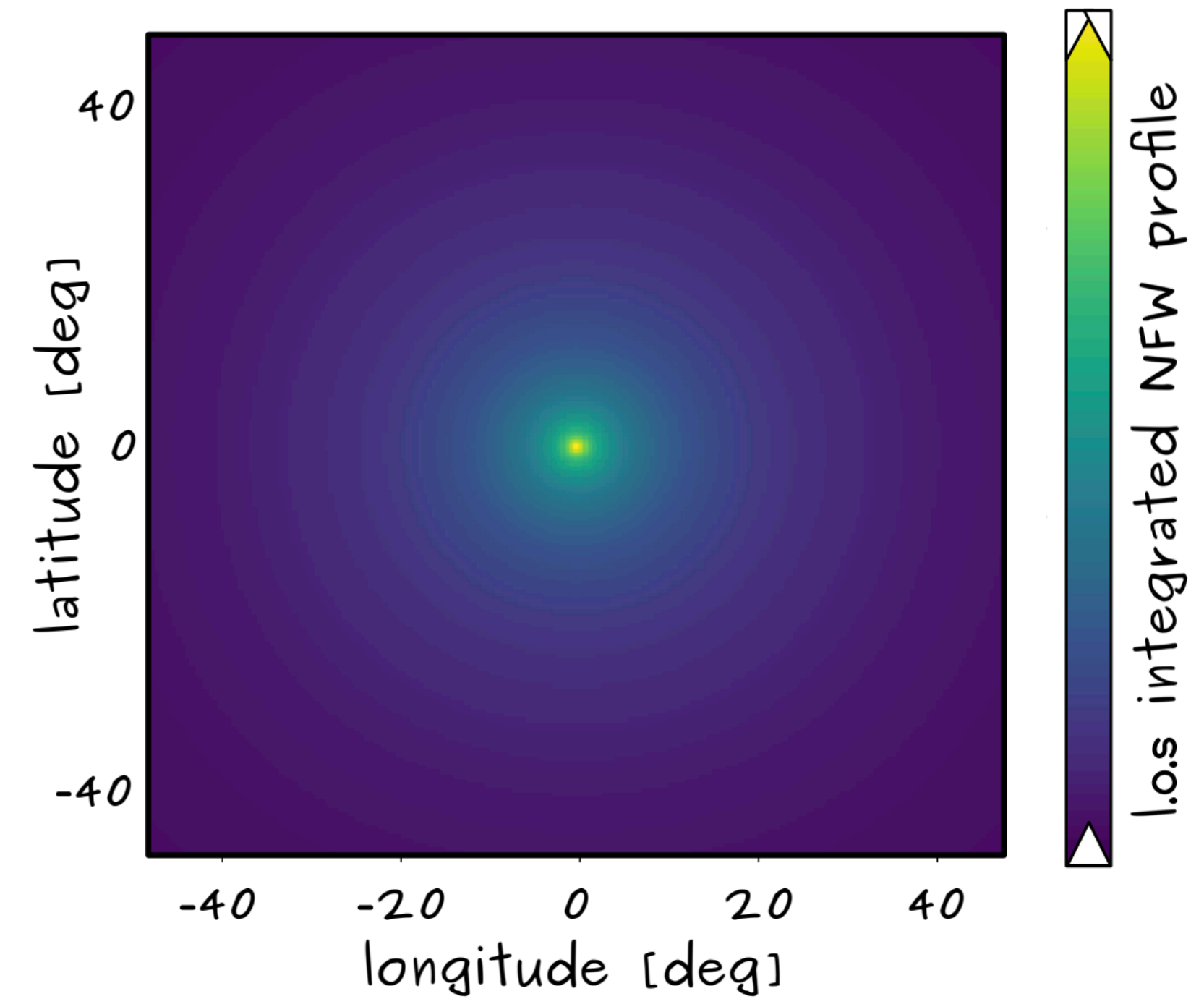
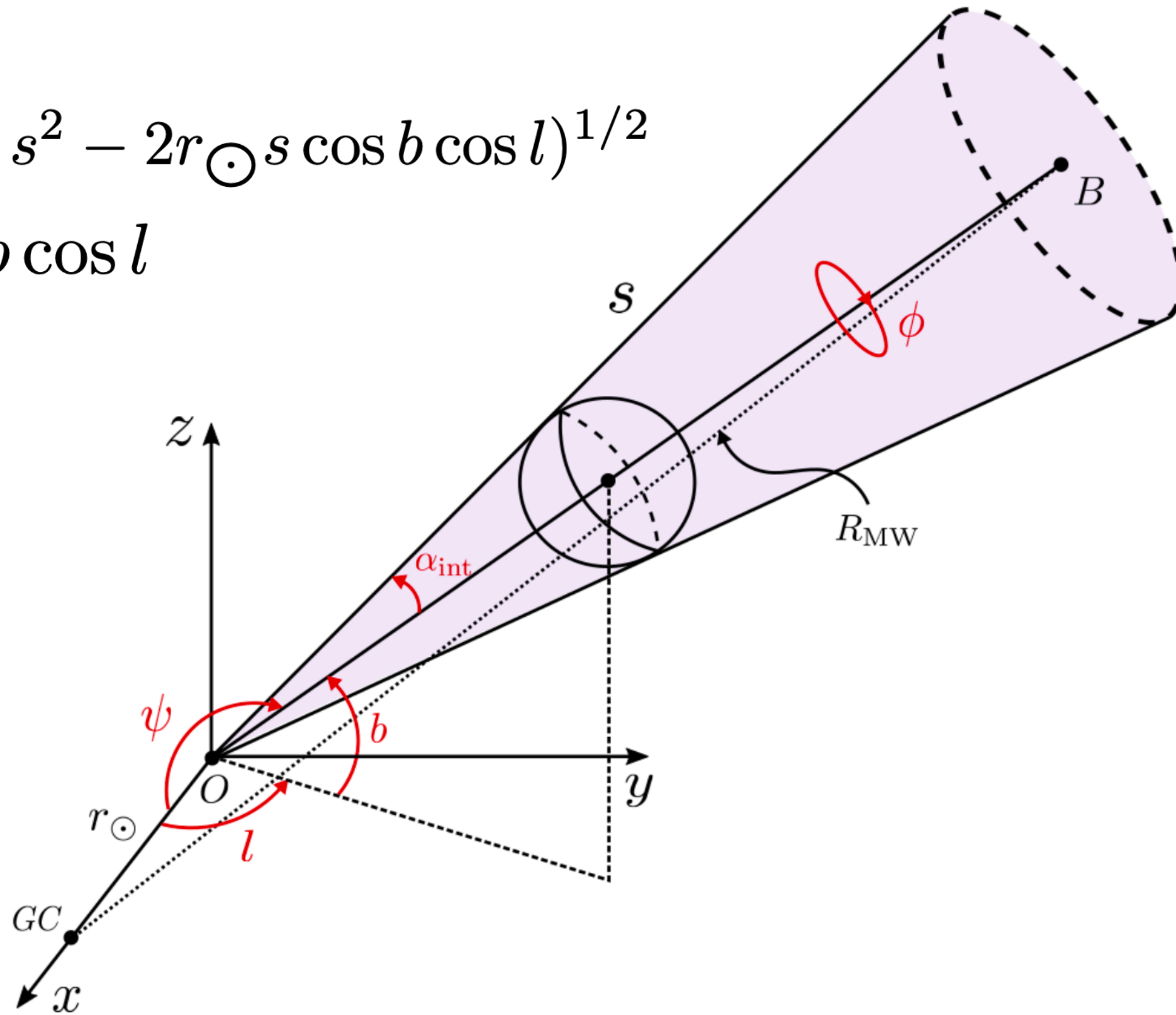
$$[\text{GeV}^{-1}\text{s}^{-1}\text{cm}^{-2}]$$

Light dark matter decay: Flux

$$\frac{d\Phi_\gamma}{dE} = \frac{\Gamma(\text{DM} \rightarrow \gamma\gamma)}{4\pi m_{\text{DM}}} \frac{dN_\gamma}{dE} \times \int_{\Delta\Omega} d\Omega \int_{\text{l.o.s.}} \rho_{\text{DM}}(l) dl$$

$$r(s) = (r_\odot^2 + s^2 - 2r_\odot s \cos b \cos l)^{1/2}$$

$$\cos \psi = \cos b \cos l$$

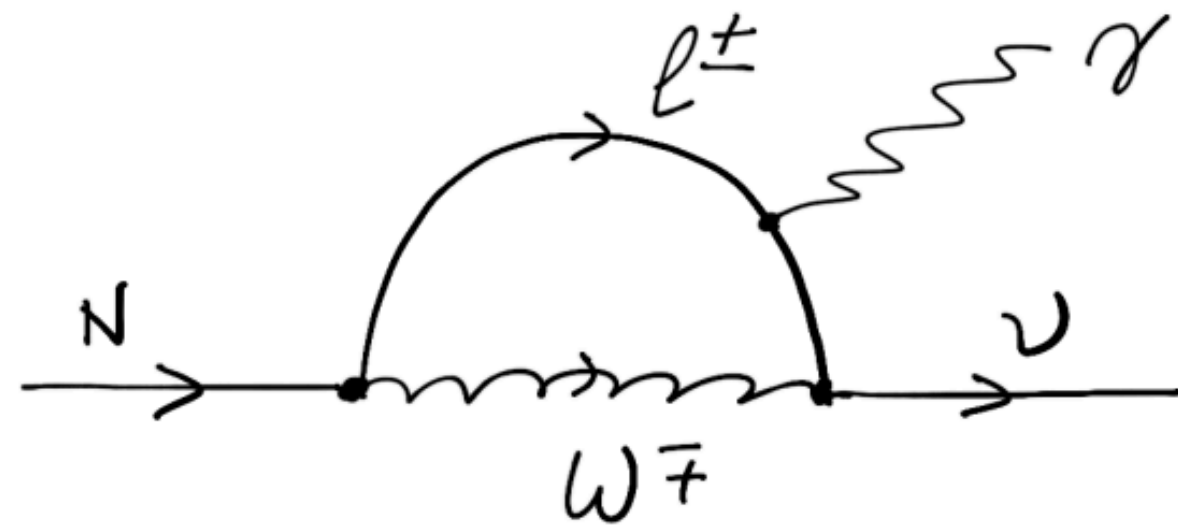


Flux from all galaxies: you need to integrate over redshift and galaxies' redshift distribution

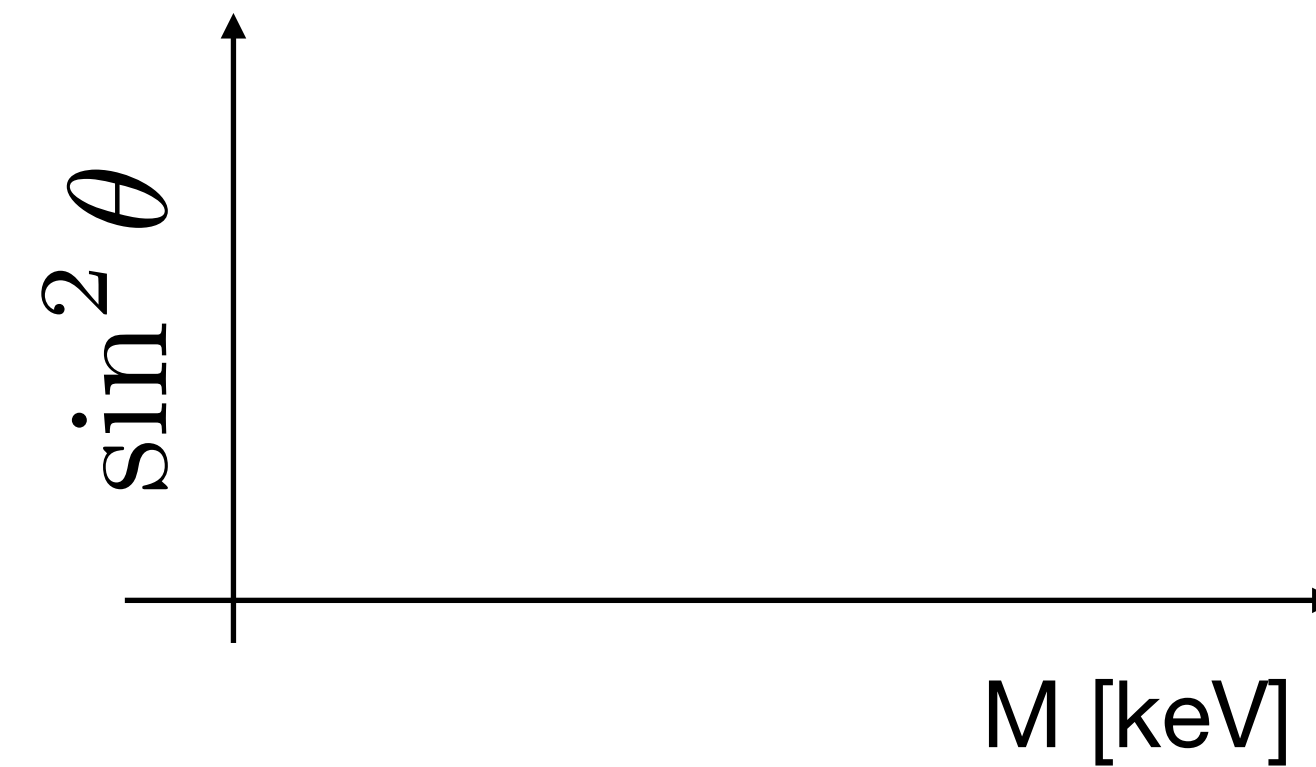
Sterile neutrinos X-ray lines

Model prediction

- X-ray lines for **sterile neutrinos** in the keV to MeV mass range
- Loop mediated radiative decay



$$\Gamma_N \approx 10^{-29} \text{ s}^{-1} \left[\frac{\sin^2(2\theta)}{10^{-7}} \right] \left(\frac{M}{1\text{keV}} \right)^5$$

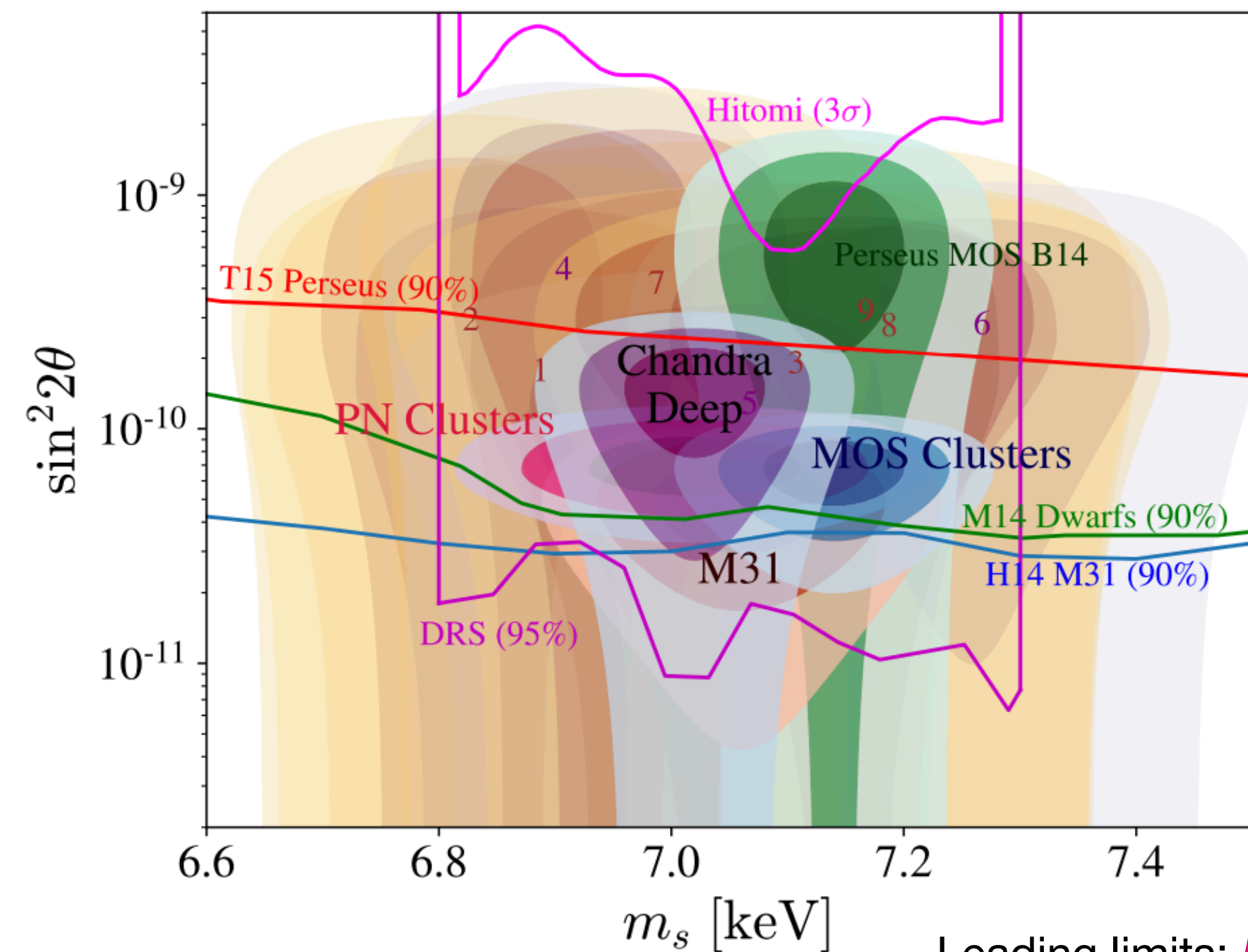
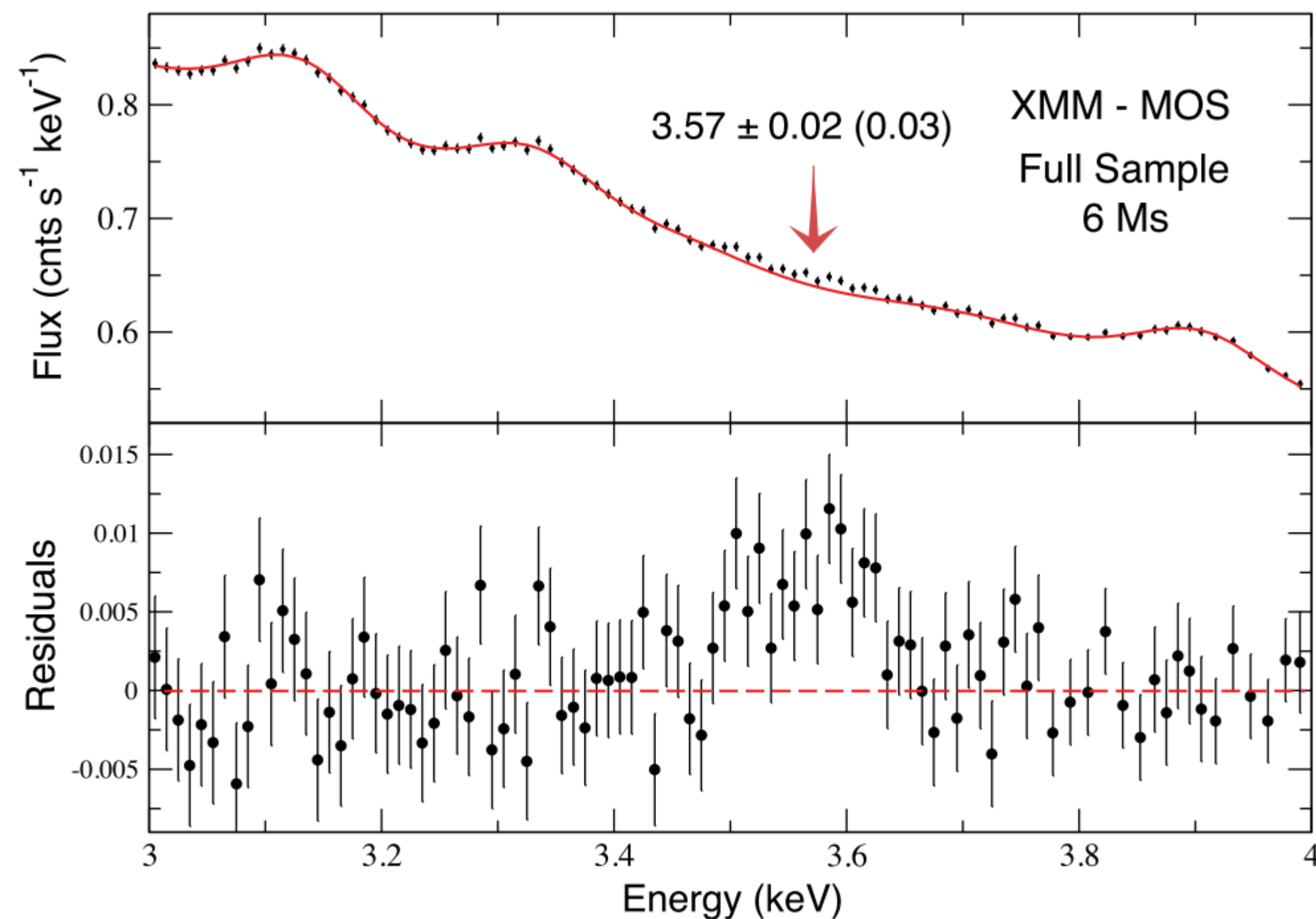


Sterile neutrinos X-ray lines

X-ray telescopes and spectral analysis

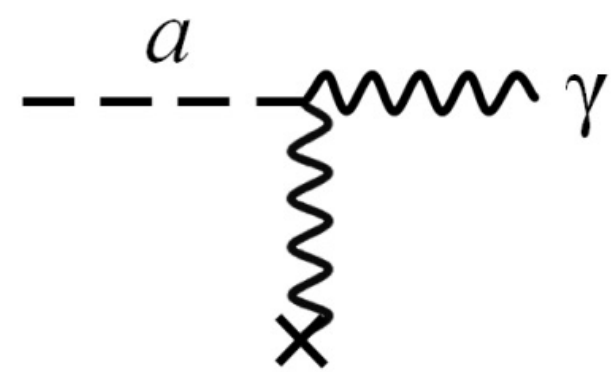
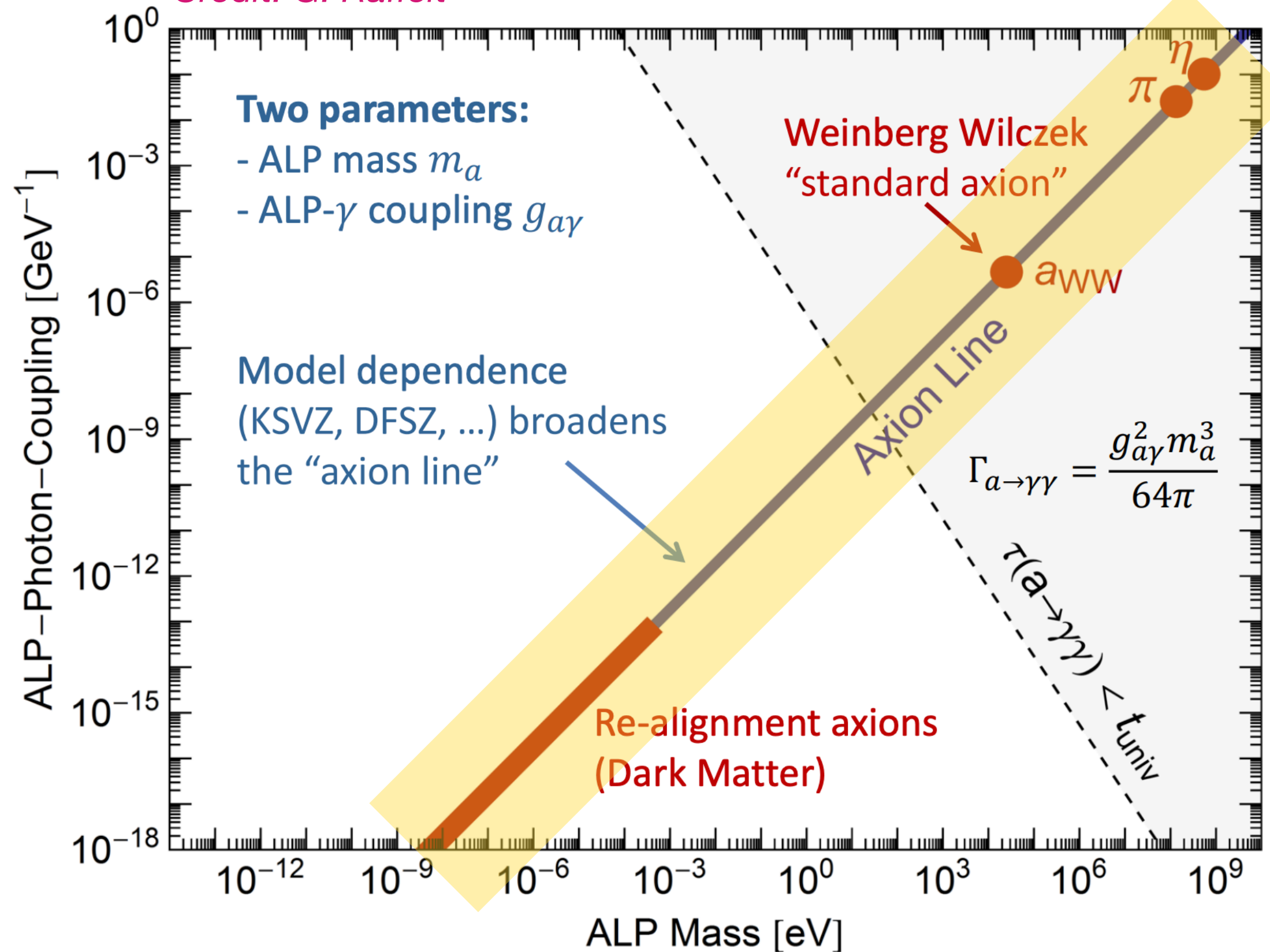
Starting from early 2014:

- ▶ **Detection** of an unidentified line at **3.5 keV**: XMM-Newton (6 Ms) & *Chandra*, Perseus cluster; XMM-Newton, M31; Suzaku, Perseus; etc
- ▶ **Constraints** from *Chandra* M31; XMM-Newton/*Chandra* 80 galaxies; blank field pointings *Chandra* and XMM-Newton, etc

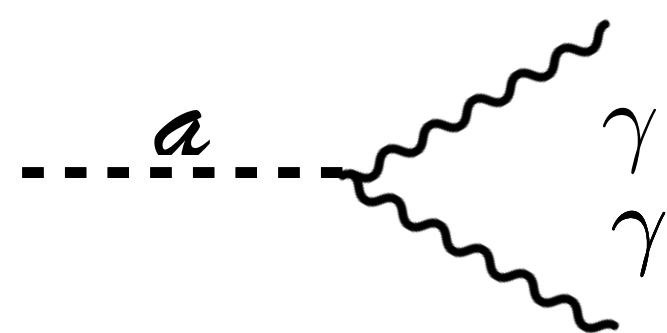


The QCD axion

Credit: G. Raffelt



Primakoff conversion



Decay



Quick ID:

- **Light** pseudo-scalar particle
- Minimal coupling with **gluons** to solve the strong CP problem
- Production through Peccei-Quinn symmetry breaking at the energy scale

$$f_a \approx 10^{10} \text{ GeV} \left(\frac{0.6 \text{ meV}}{m_a} \right)$$

- Viable **cold dark matter candidate** over a large mass range
- Induced coupling with **photons**

$$\mathcal{L}_{a\gamma} = -\frac{1}{4} g_{a\gamma} F_{\mu\nu} \tilde{F}^{\mu\nu} a = g_{a\gamma} \mathbf{E} \cdot \mathbf{B} a$$

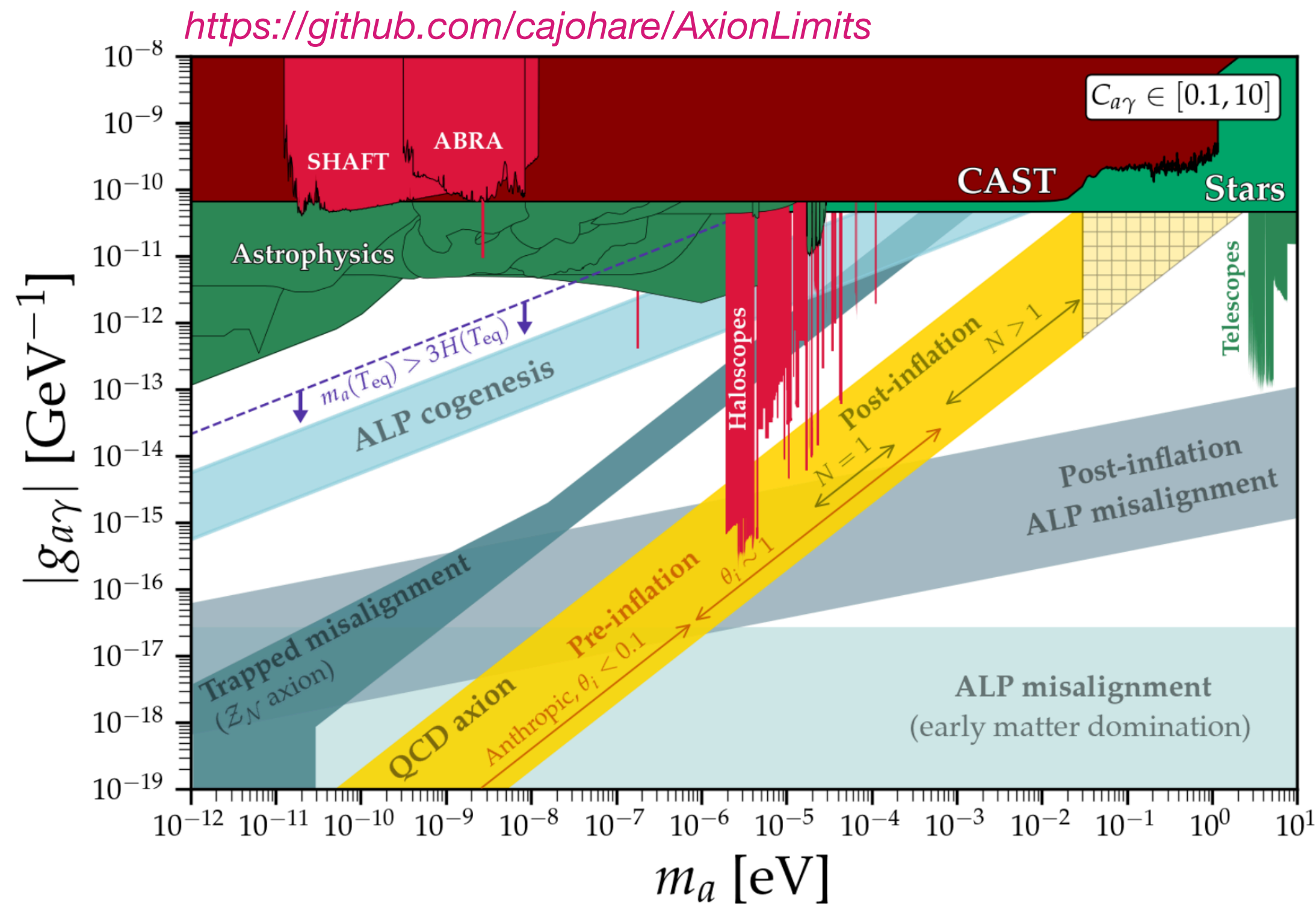
$$g_{a\gamma} = \frac{\alpha_{\text{em}}}{2\pi f_a} \left(\frac{E}{N} - 1.92 \right)$$

Broadening the landscape

Axion-like particles

Axion-like particles: (pseudo-)scalar particles, masses as low as ZeV, very weak couplings with SM, coupled with photons as QCD axions

Chang+ PRD 2000; Turok PRL 1996; Arvanitaki+ PRD'10

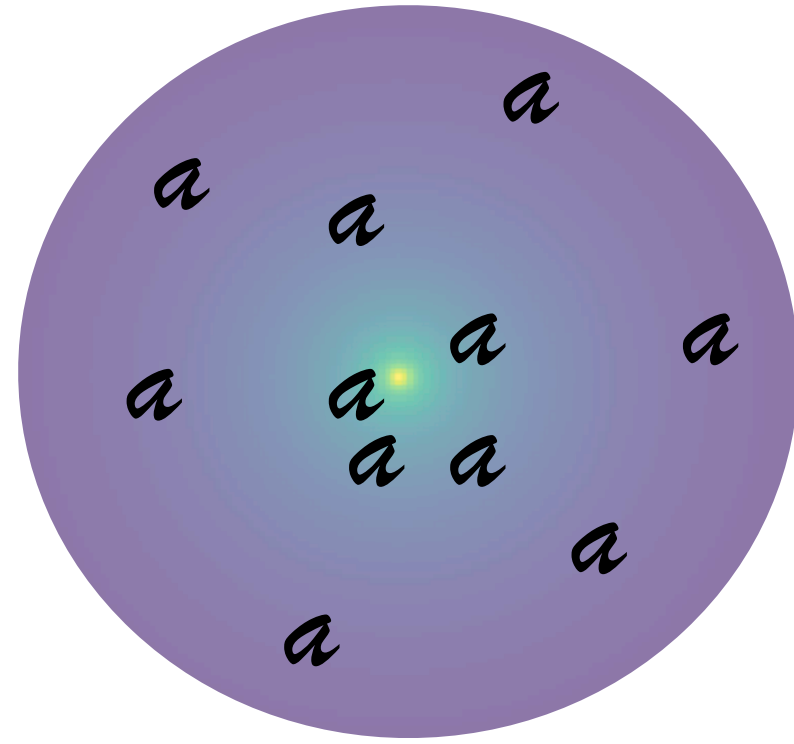


Quick ID:

- **Common** in many **extensions of the SM** from the spontaneous breaking of approximate global symmetries
- **Mass is not determined** by QCD effects
- They do **not solve the strong CP problem**
- They can be **viable dark matter candidates**

[Wave-like DM is even broader: more generally, light scalars or vectors]

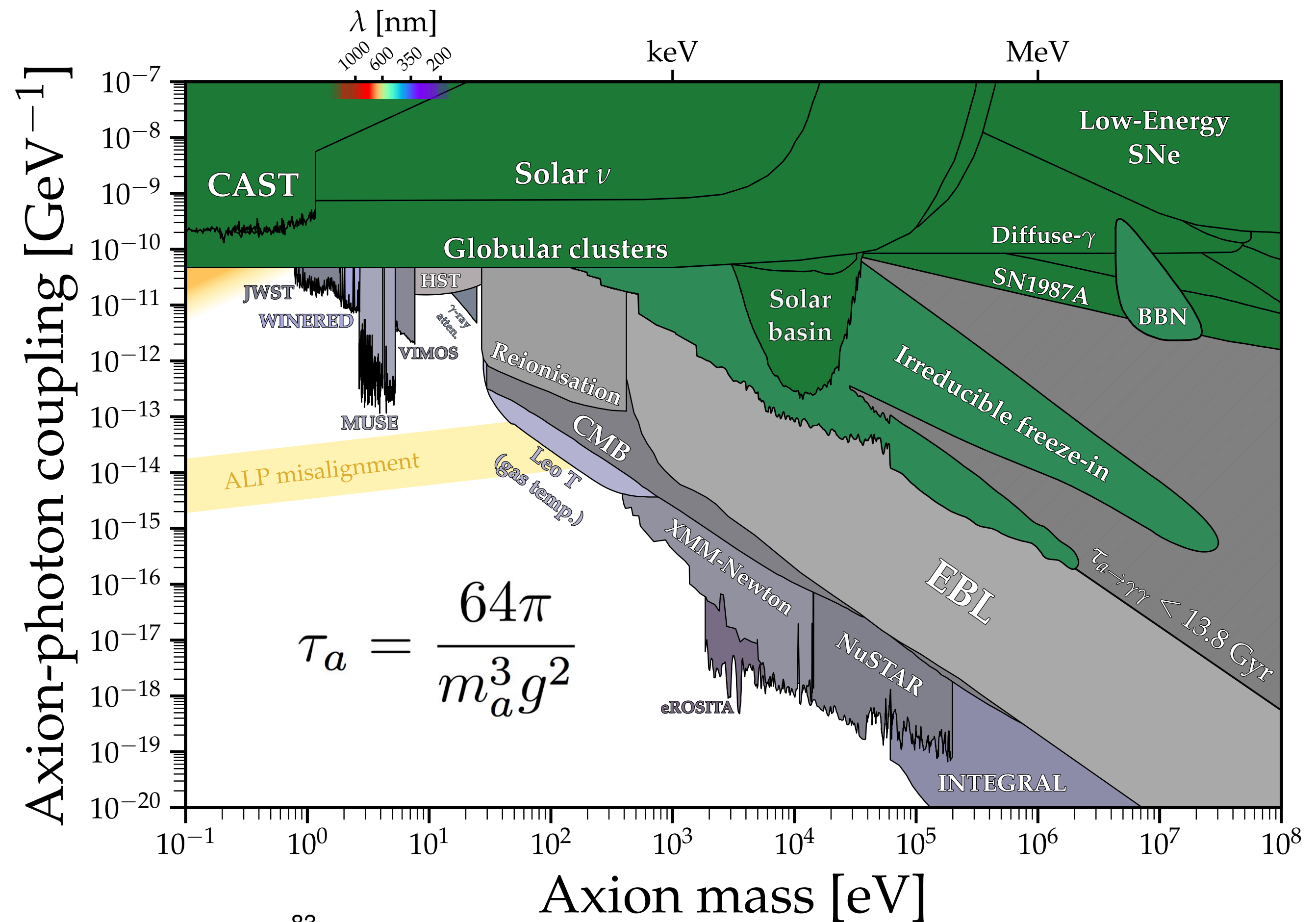
ALPs dark matter



- If DM, ALPs distributed in galaxies according standard DM density distributions (e.g. NFW)
- Search for narrow lines in DM-rich environments

- ALPs can be good DM candidates in some portions of the parameter space

Preskill+ PLB 1983; Sikivie International Journal of Modern Physics '10

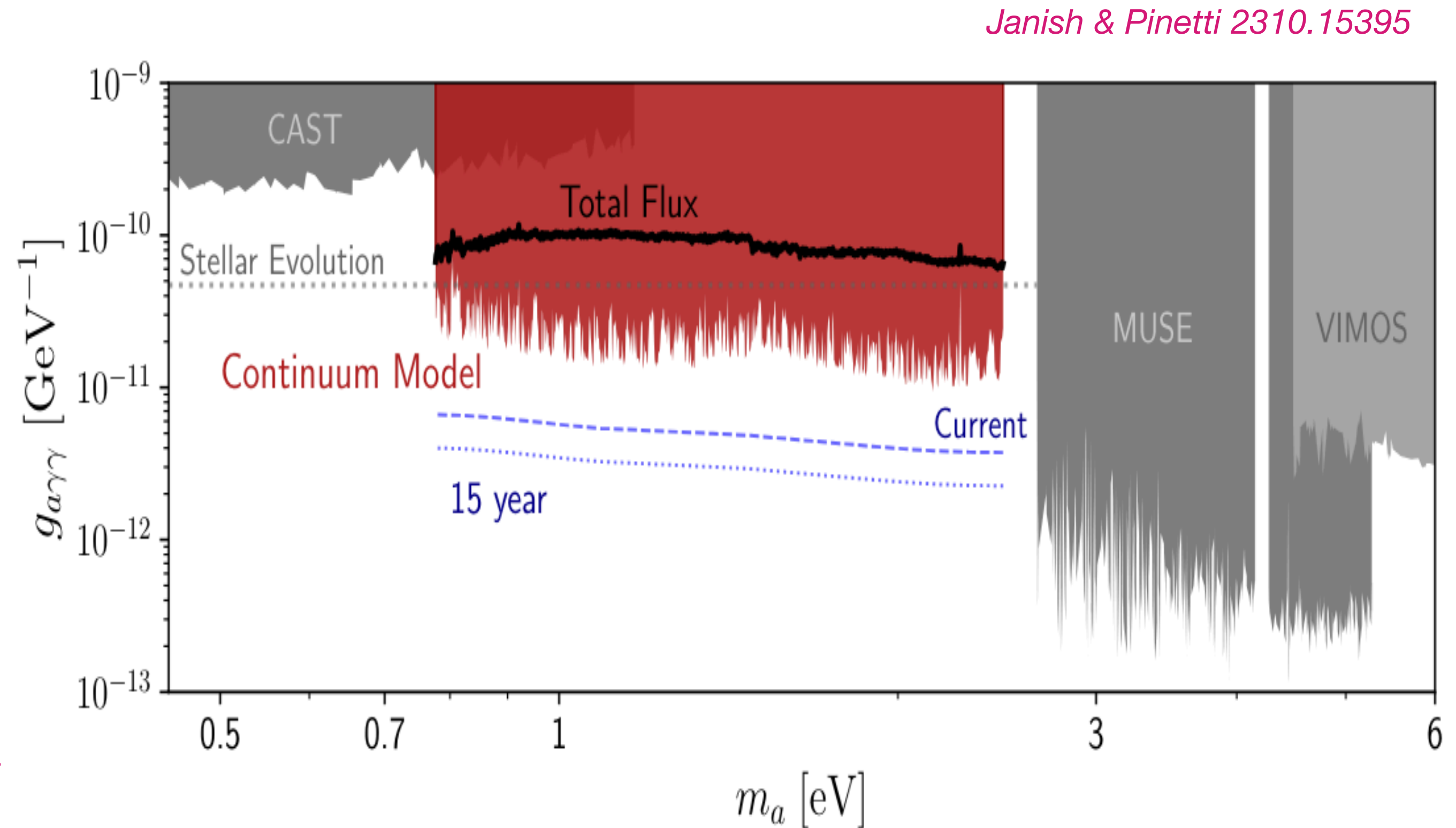


Constraints on eV ALPs

IR - optical wavelengths

Search for **narrow lines** in IR and optical data

- **MUSE**: search in the direction of 5 known dwarf galaxies
Todarello+ JCAP'24
- **VIMOS (Visible Multi-Object Spectrograph)**: galaxy clusters Abell 2667 and 2390
Grin+ PRD'06
- **JWST**: public blank sky observations from the NIRSpec IFU
Janish & Pinetti 2310.15395



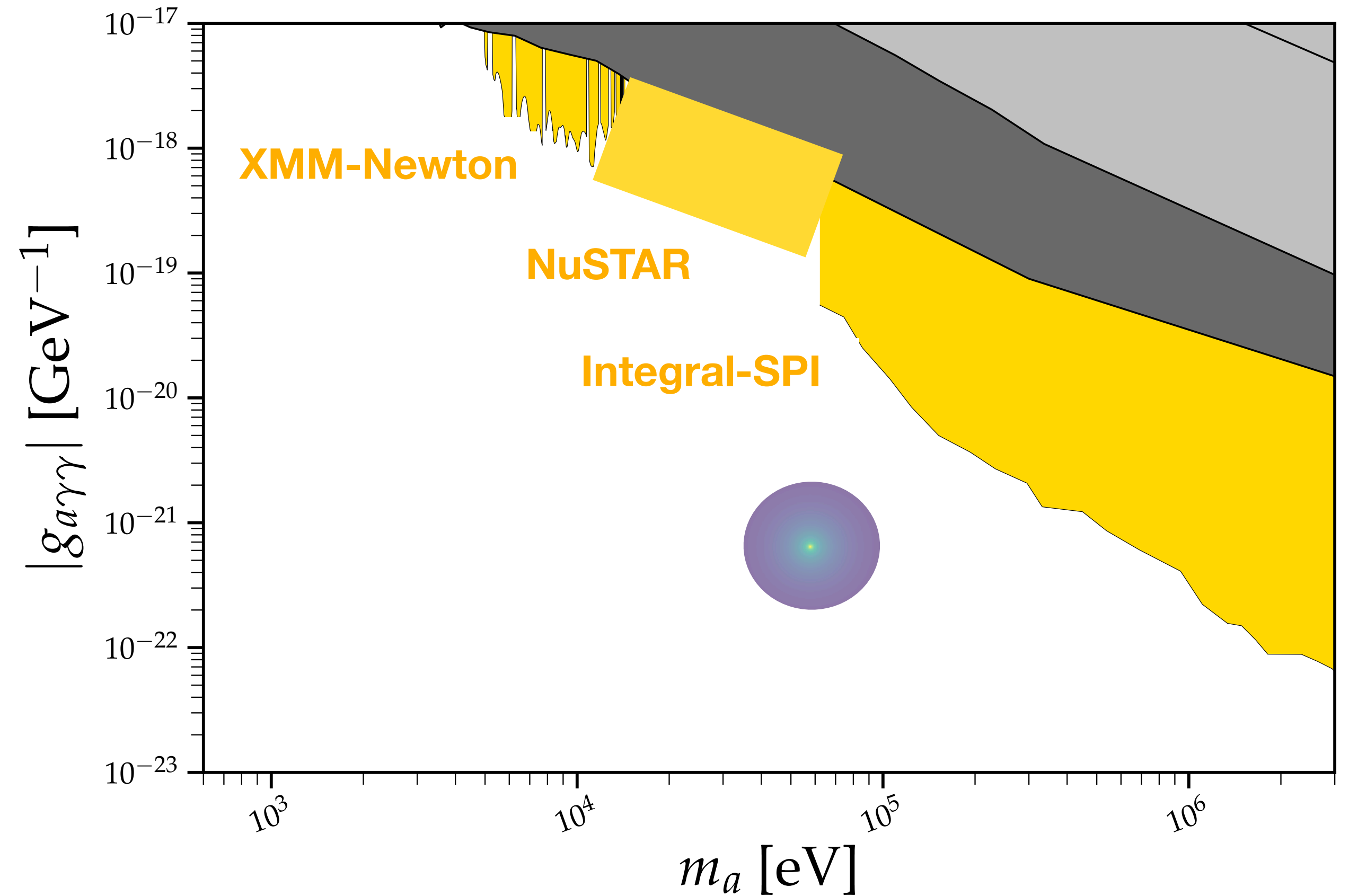
Constraints on keV - MeV ALPs

X-ray and soft gamma rays energies

Heavy ALPs DM decay

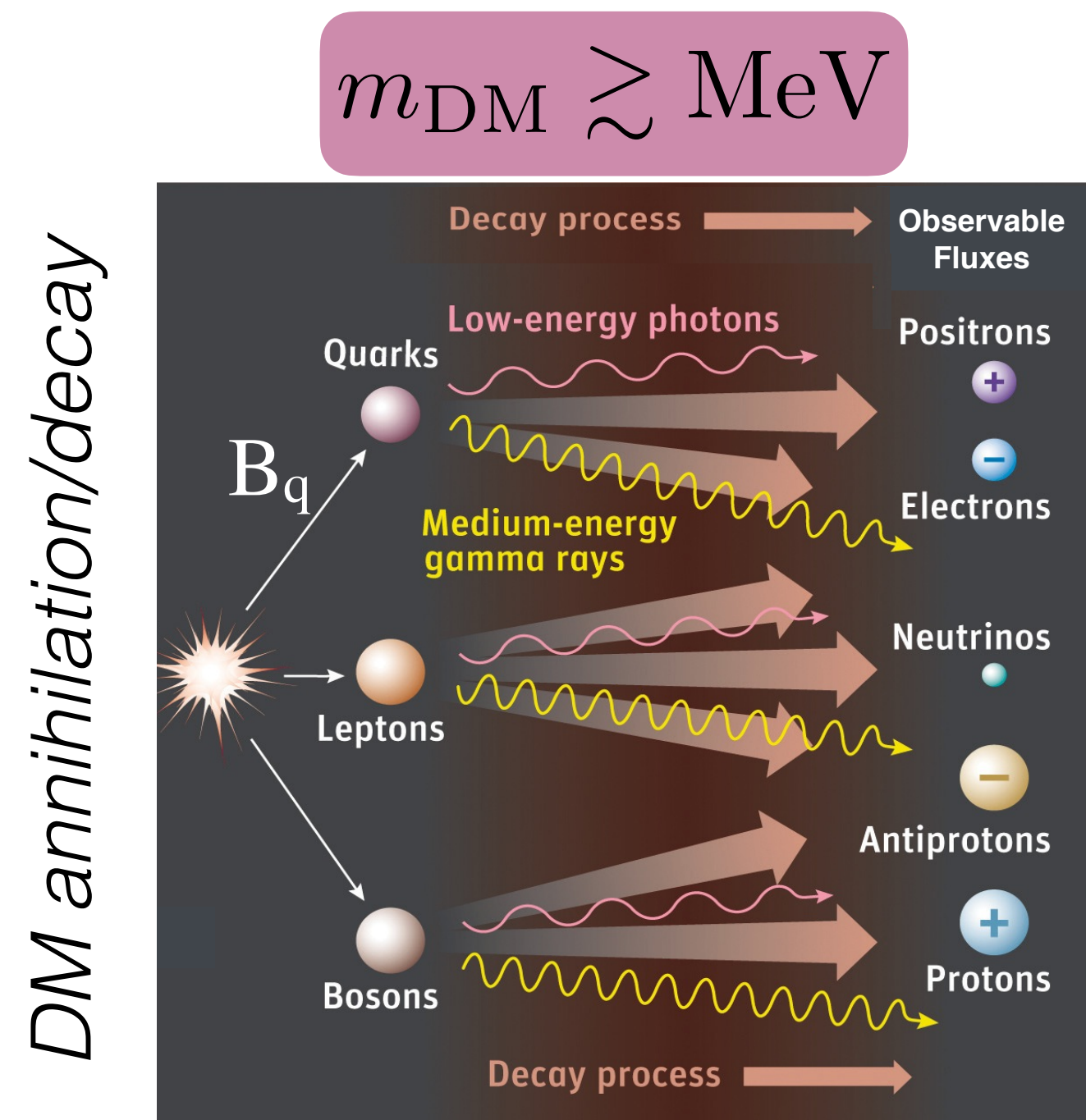
Search for **narrow lines** in X and gamma-ray data

- **XMM-Newton**: 5-16 keV, archival data
=> No evidence found for unassociated X-ray lines
Foster+ PRL'21
- **NuSTAR**: 7-Ms/detector deep blank-sky exposures
Roach+ PRD'23
- **Integral-SPI**: new analysis of 16yr data with dedicated search for DM component in continuum Galactic emission
Berteaud, FC+PRD'22; FC+ MNRAS'23



Particle dark matter emission*

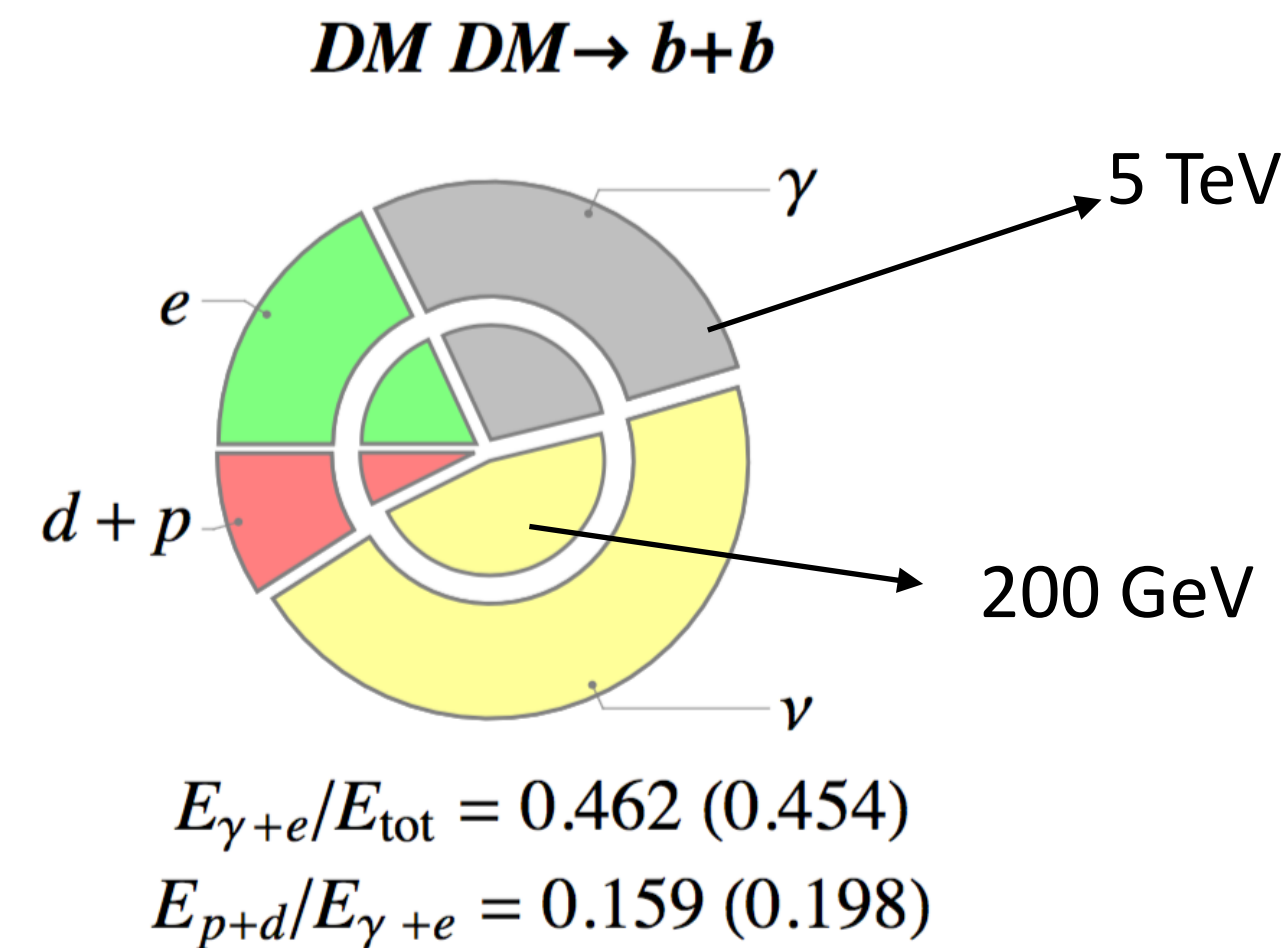
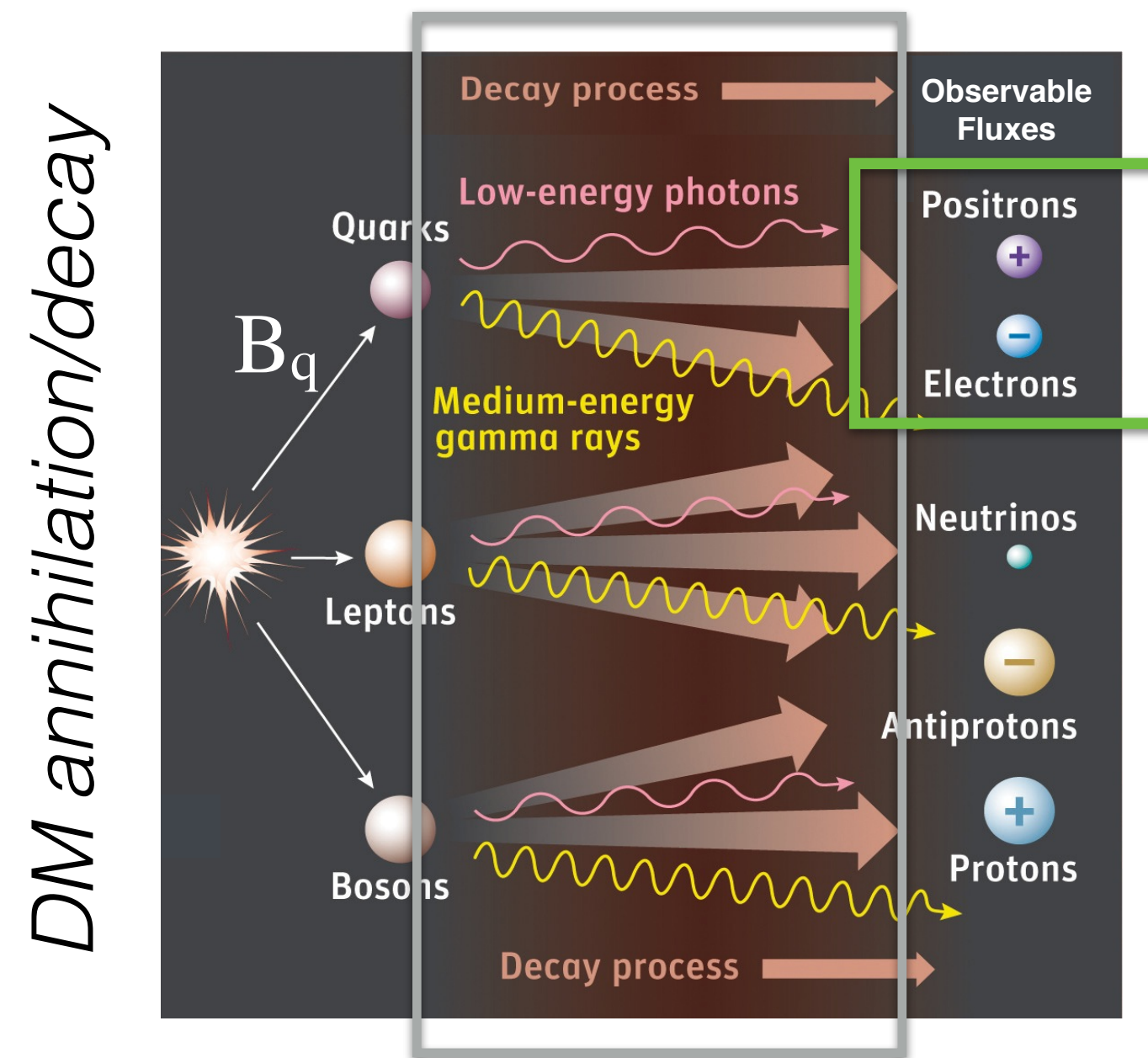
$$E_{\text{CM}} = N m_{\text{DM}}, \quad N = 1 \text{ (decay)}, 2 \text{ (annih)} \quad \text{Centre of mass energy} \simeq \text{Signal energy}$$



[*Dark matter = Weakly interacting massive particles, **WIMPs**]

Particle dark matter emission

$$E_{\text{CM}} = N m_{\text{DM}}, \quad N = 1 \text{ (decay)}, 2 \text{ (annih)} \quad \text{Centre of mass energy} \simeq \text{Signal energy}$$



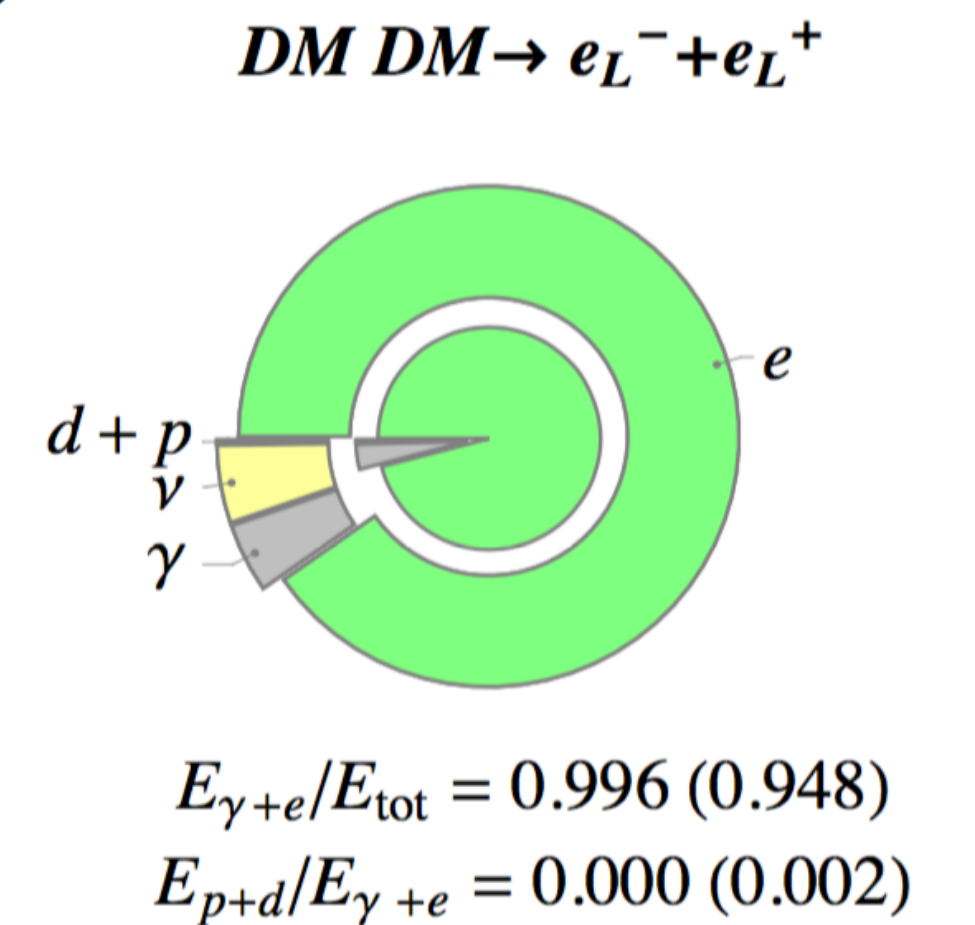
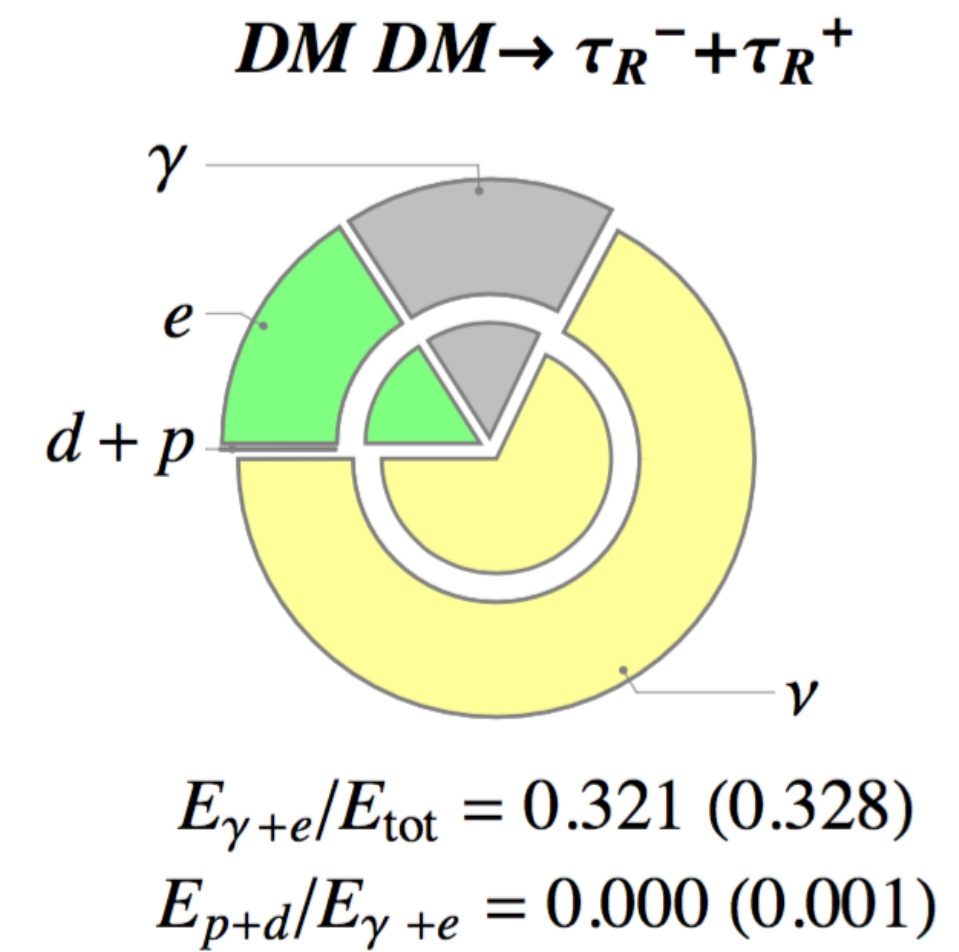
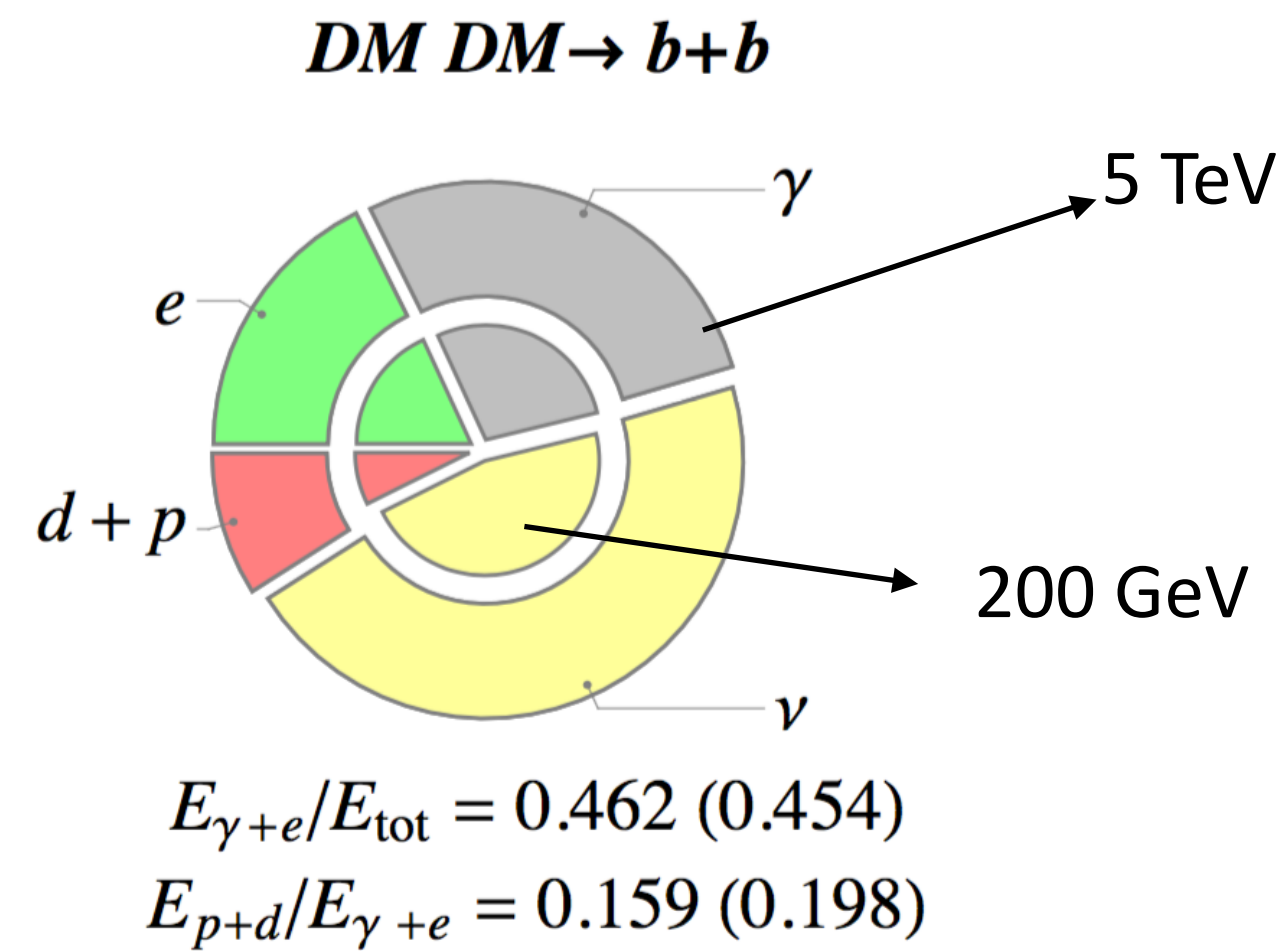
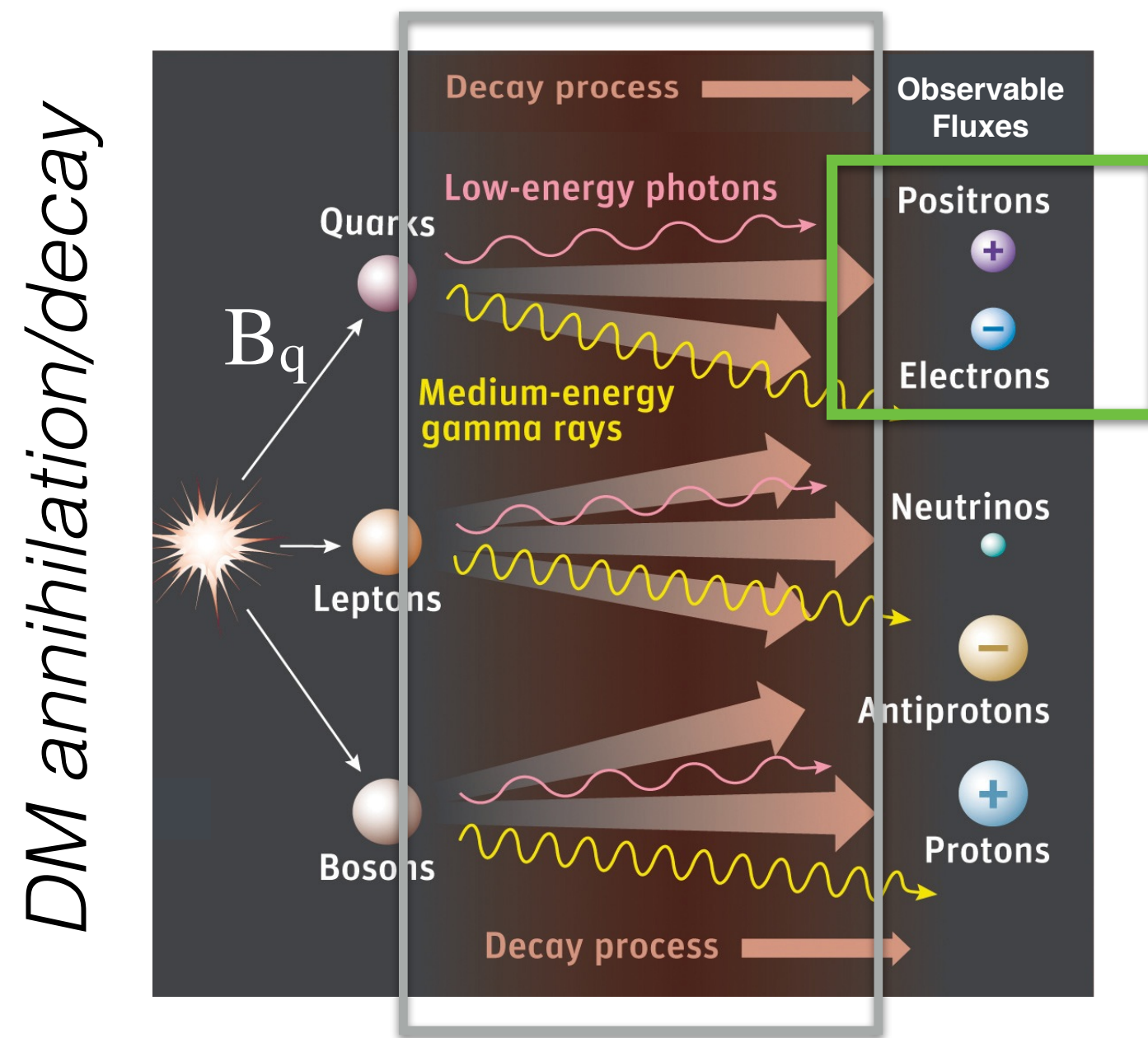
Prompt emission of observable particles *i*

$$\sum_f B_f \frac{dN_i^f}{dE}(E)$$

Particle dark matter emission

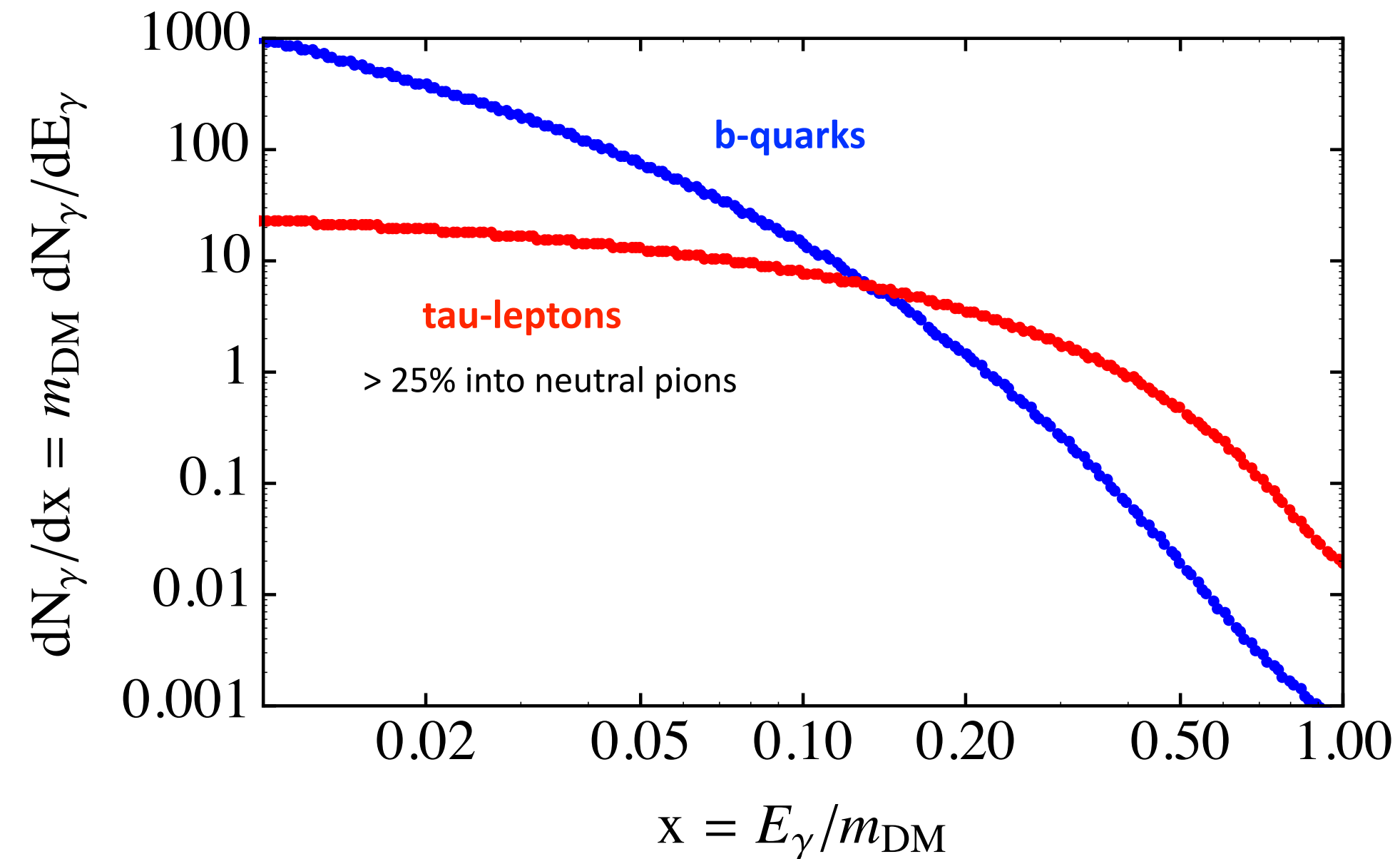
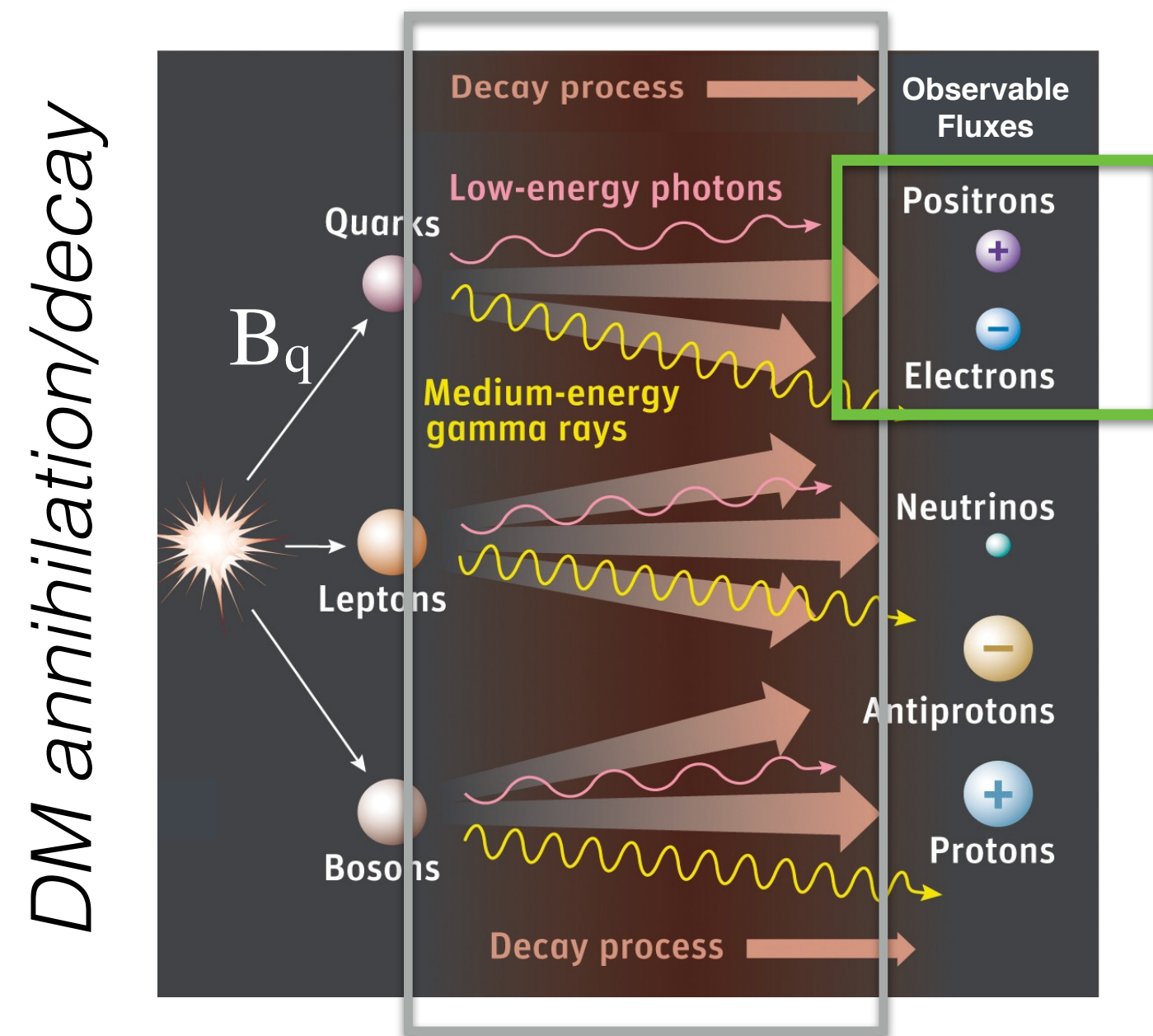
$$E_{\text{CM}} = N m_{\text{DM}}, \quad N = 1 \text{ (decay)}, 2 \text{ (annih)}$$

Centre of mass energy \simeq Signal energy



Particle dark matter emission

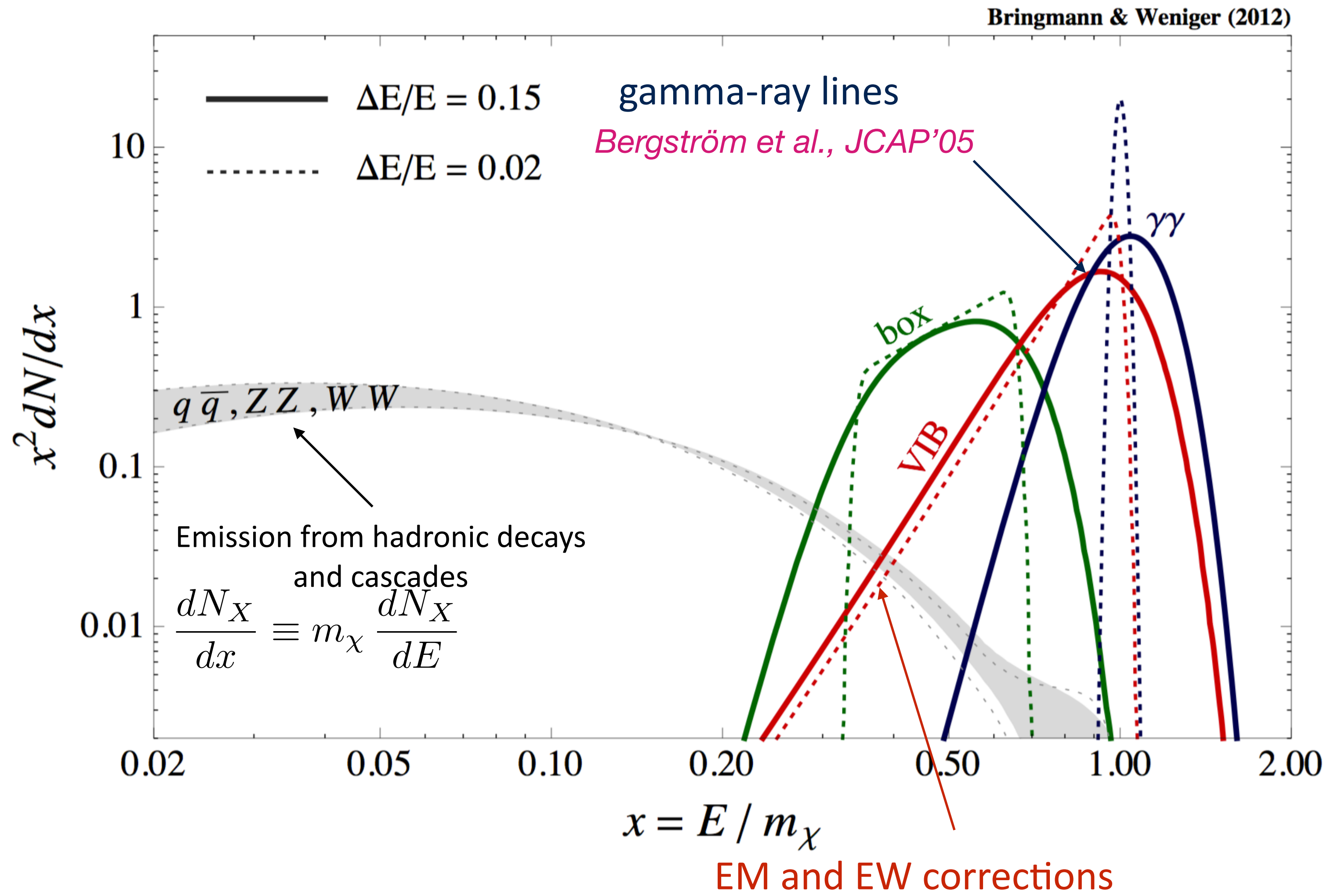
$$E_{\text{CM}} = N m_{\text{DM}}, \quad N = 1 \text{ (decay)}, 2 \text{ (annih)} \quad \text{Centre of mass energy} \simeq \text{Signal energy}$$



$$x \equiv \frac{E_X}{m_\chi}$$

$$\frac{dN_X}{dx} \equiv m_\chi \frac{dN_X}{dE}$$

Dark matter spectrum: prompt photons



Bringmann, FC+ JHEP'08; Ciafaloni+ JCAP'11; Bringmann&FC PRL'14

Numerical codes for computation of DM spectra:

DarkSUSY <http://www.fysik.su.se/~edsjo/darksusy/>

Gondolo+ JCAP'04

MicrOMEGAs <https://lapth.cnrs.fr/micromegas/>

Belanger+ JCAP'05

PPC 4 DM ID <http://www.marcocirelli.net/PPPC4DMID.html>

Cirelli+ JCAP 2012

Analytical fitting functions:

Fornengo, Pieri, Scopel, PRD 2004

Cembranos et al., PRD 2011

For dependence on event Monte Carlo generators see, e.g.,

Cembranos+ JHEP'13

Dark matter annihilation: Rate and flux

$$Q_{\gamma}^{\text{ann}}(r, E) = \langle \sigma v \rangle \times n_{\text{pairs}}(r) \times \sum_f B_f \frac{dN_{\gamma}^f}{dE}(E)$$

$$n_{\text{pairs}}(r) = s \times \frac{\rho^2(r)}{m^2}$$

$$s = \left\{ \begin{array}{l} \frac{1}{2} \\ \frac{1}{4} \end{array} \right\}$$

Majorana       Dirac

Emissivity per unit energy and volume

$$[\text{GeV}^{-1} \text{s}^{-1} \text{cm}^{-3}]$$

$$\frac{d\Phi_{\gamma}}{dE} = s \frac{\langle \sigma v \rangle}{4\pi m_{\text{DM}}^2} \sum_f B_f \frac{dN_{\gamma}^f}{dE} \times \int_{\Delta\Omega} d\Omega \int_{\text{l.o.s.}} \rho_{\text{DM}}^2(\ell) d\ell$$

Gamma-ray (or neutrino) differential energy flux

$$[\text{GeV}^{-1} \text{s}^{-1} \text{cm}^{-2}]$$

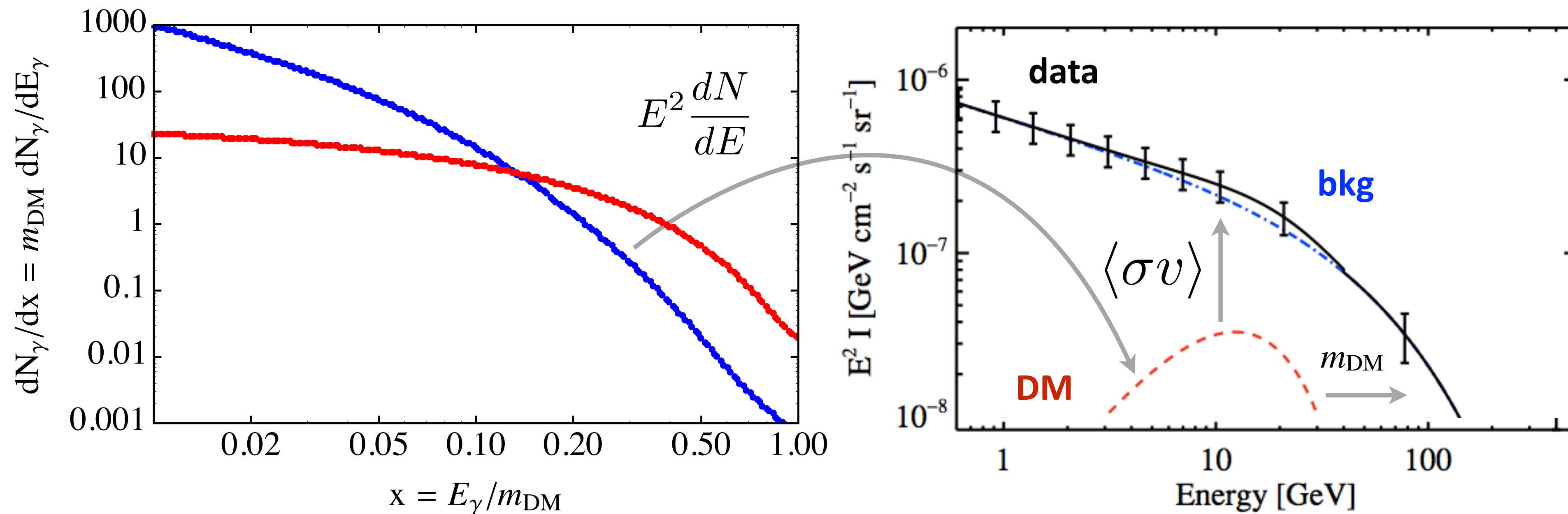
Be careful! This equation holds in the s-wave approximation, i.e. of velocity dependent cross-section

$$\langle \sigma v \rangle = a + bv^2 + \mathcal{O}(v^4), \quad v/c \sim 10^{-3}$$

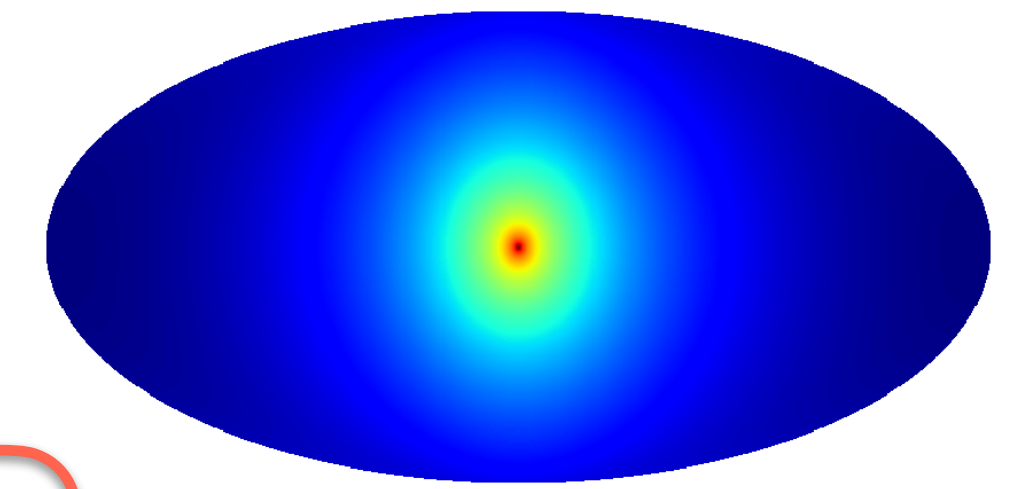
Non-relativistic regime: if present, s-wave is dominant

Dark matter annihilation flux

$$\frac{d\Phi_\gamma}{dE} = s \frac{\langle\sigma v\rangle}{4\pi m_{\text{DM}}^2} \sum_f B_f \frac{dN_\gamma^f}{dE} \times \int_{\Delta\Omega} d\Omega \int_{\text{l.o.s.}} \rho_{\text{DM}}^2(\ell) d\ell$$

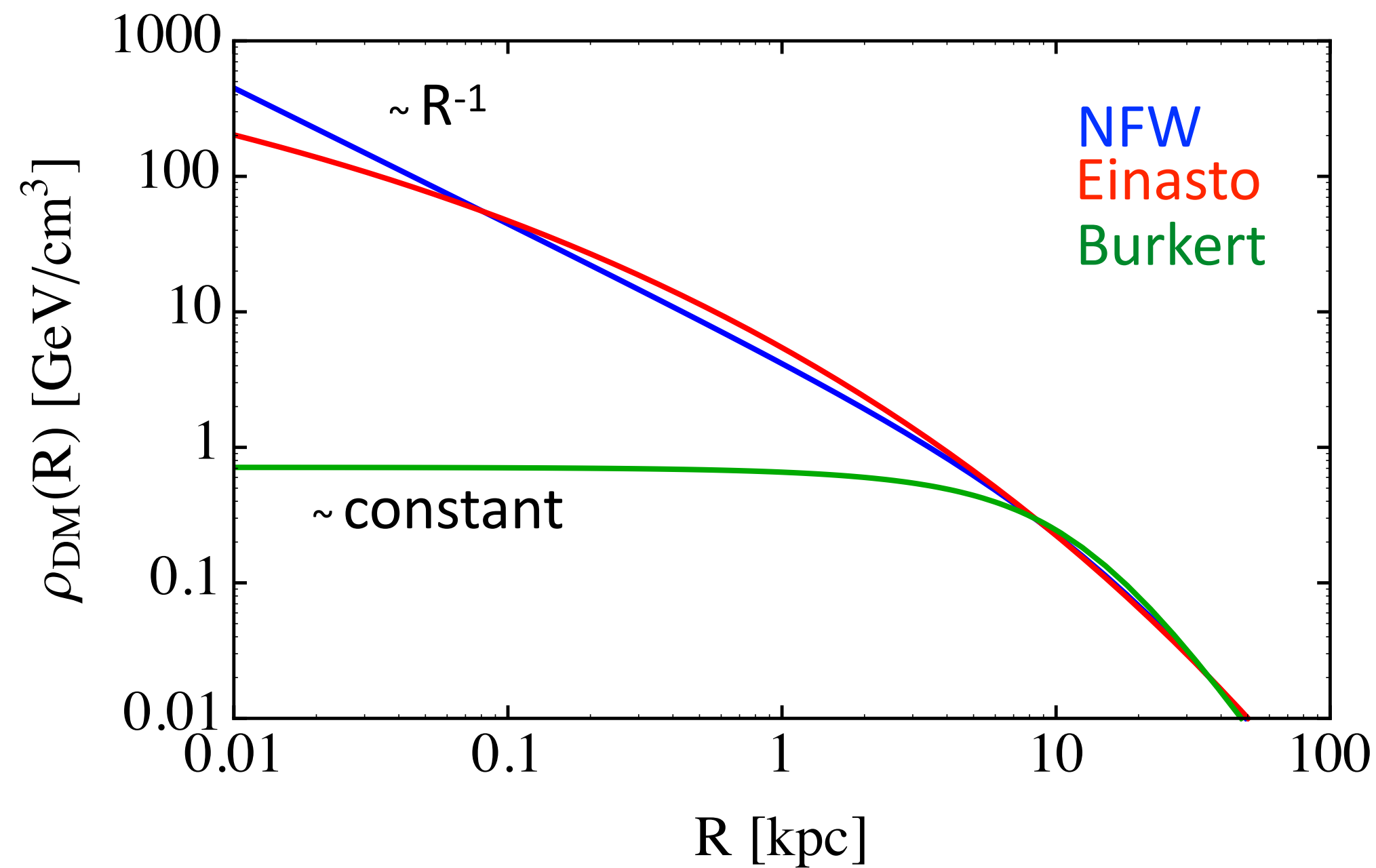


Dark matter annihilation flux

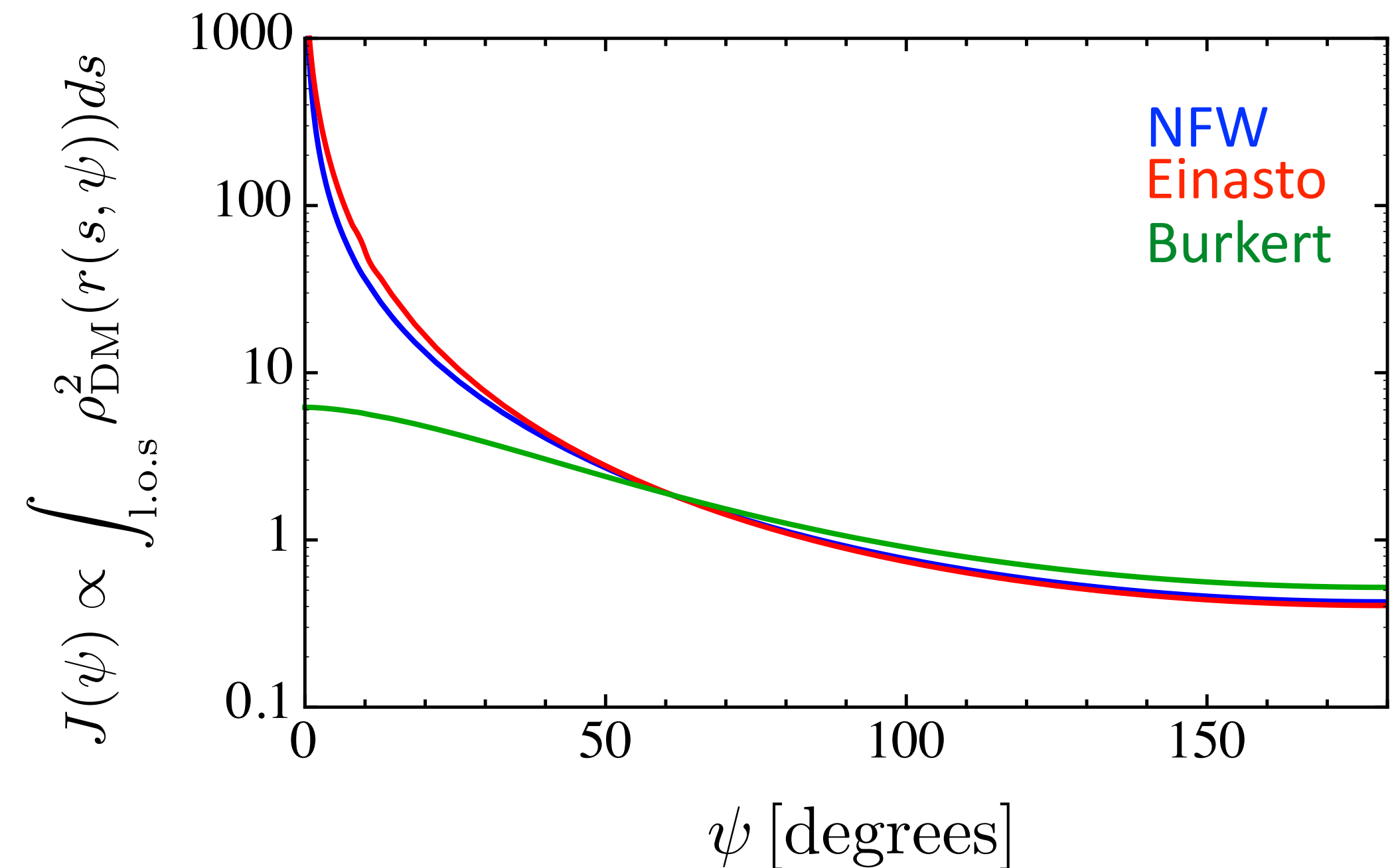


$$\frac{d\Phi_\gamma}{dE} = s \frac{\langle \sigma v \rangle}{4\pi m_{\text{DM}}^2} \sum_f B_f \frac{dN_\gamma^f}{dE} \times \int_{\Delta\Omega} d\Omega \int_{\text{l.o.s.}} \rho_{\text{DM}}^2(\ell) d\ell$$

Dark matter density profiles

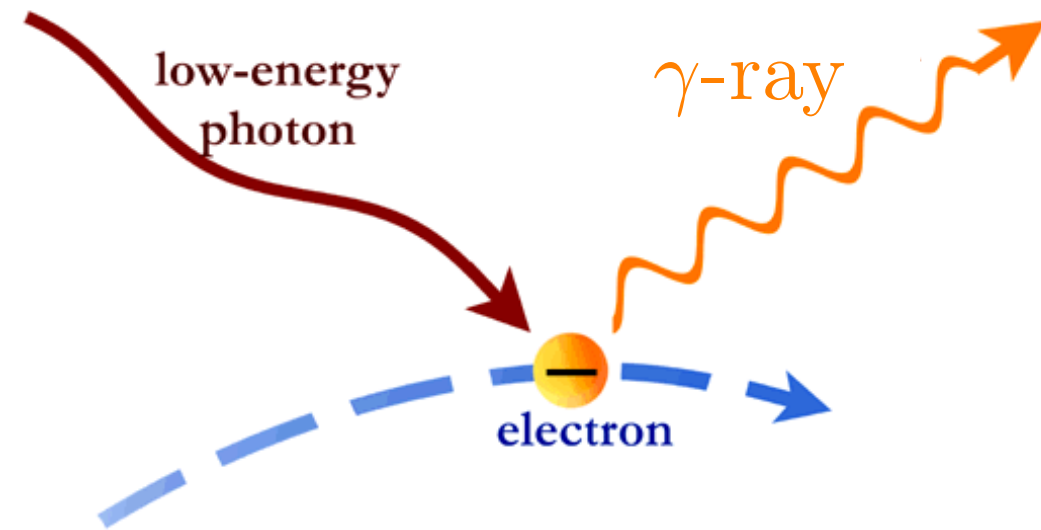


Spatial distribution of the signal

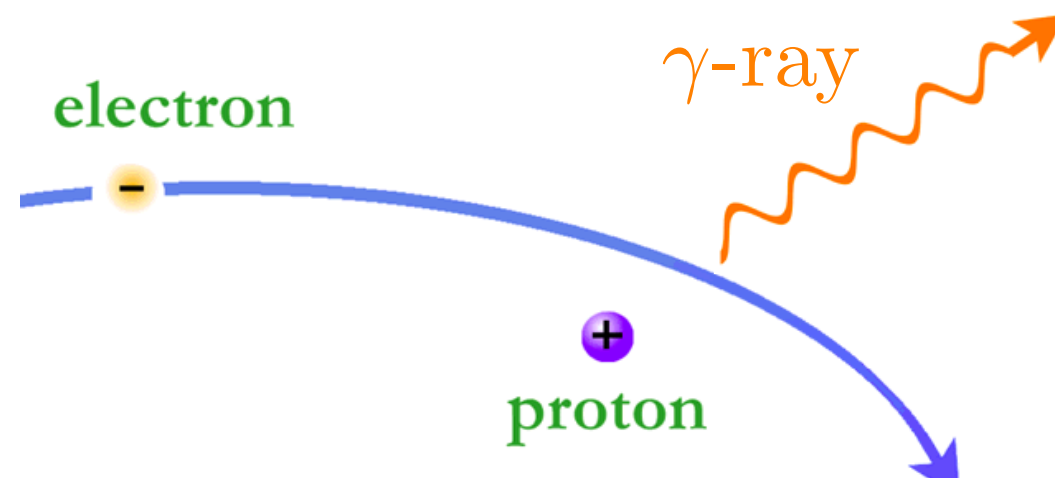


Radiative emission from leptons

$$\chi\chi \rightarrow \left\{ \begin{array}{l} ZZ, W^+W^-, \gamma\gamma \\ q\bar{q}, l^+l^-, \nu\bar{\nu} \end{array} \right\} \xrightarrow[\text{decays}]{\text{hadronization}} \gamma, e^\pm, \mu^\pm, p/\bar{p}, \pi^\pm, \nu/\bar{\nu}, \dots$$



Inverse Compton scattering
on CMB, star-light, infrared-light



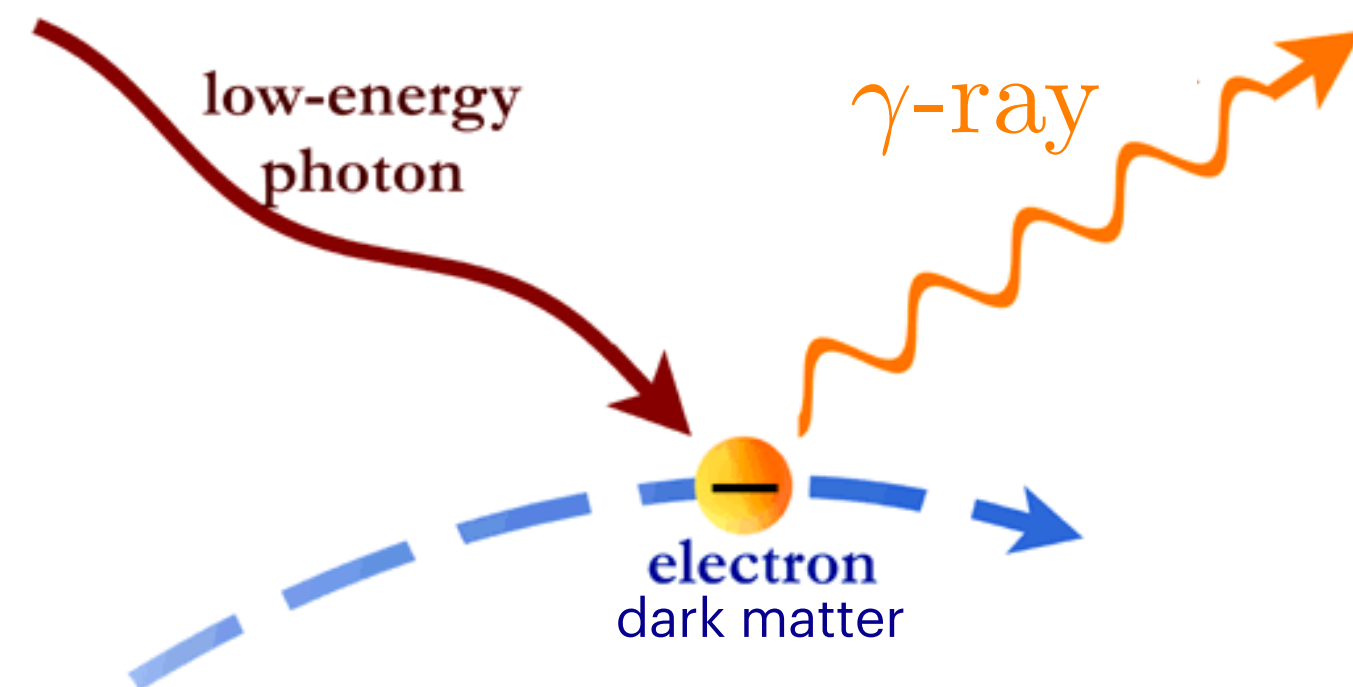
Bremsstrahlung
onto gas of interstellar medium



Synchrotron radiation
magnetic field μG for GeV-TeV electrons
 \Rightarrow MHz-GHz radio signal

Radiative emission from leptons

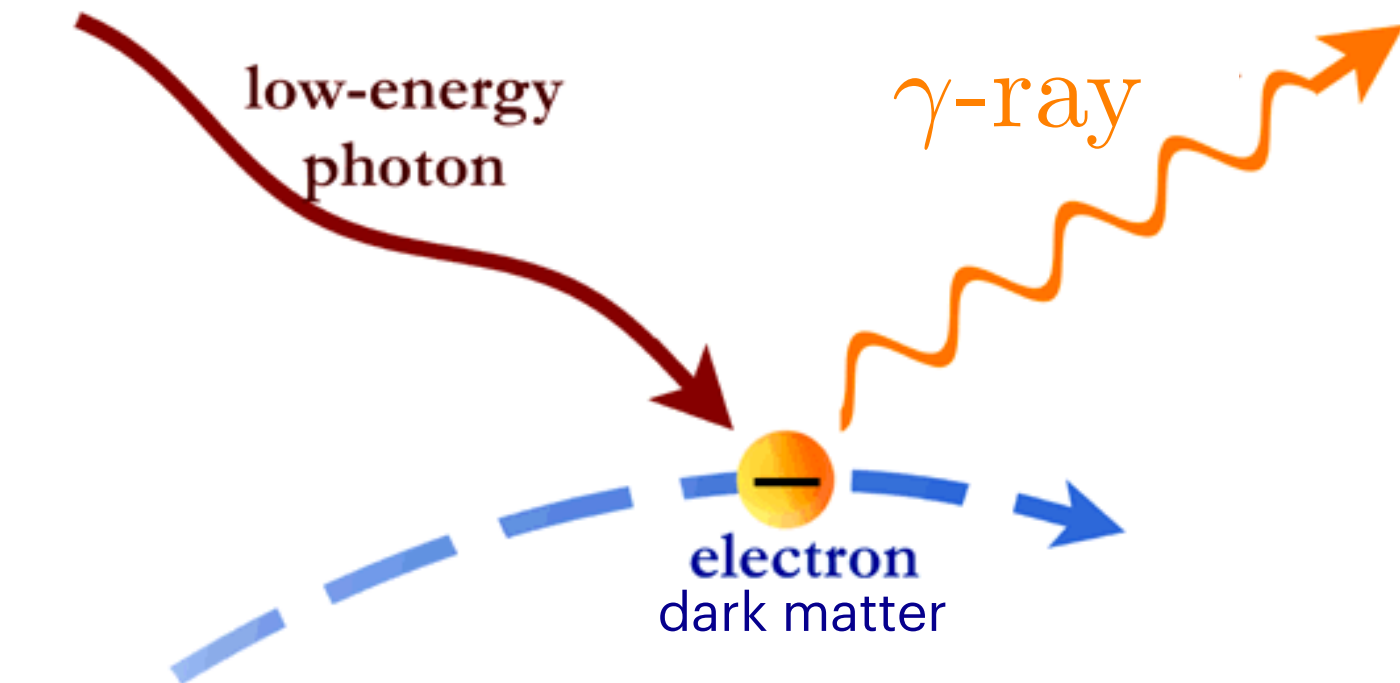
Inverse Compton scattering



$$E_\gamma \approx \gamma^2 \varepsilon \quad ; \quad \varepsilon E_e \ll m_e^2$$
$$E_\gamma \approx 30 \left(\frac{\varepsilon}{\text{eV}} \right) \left(\frac{E_e}{\text{GeV}} \right)^2 \text{ MeV}$$

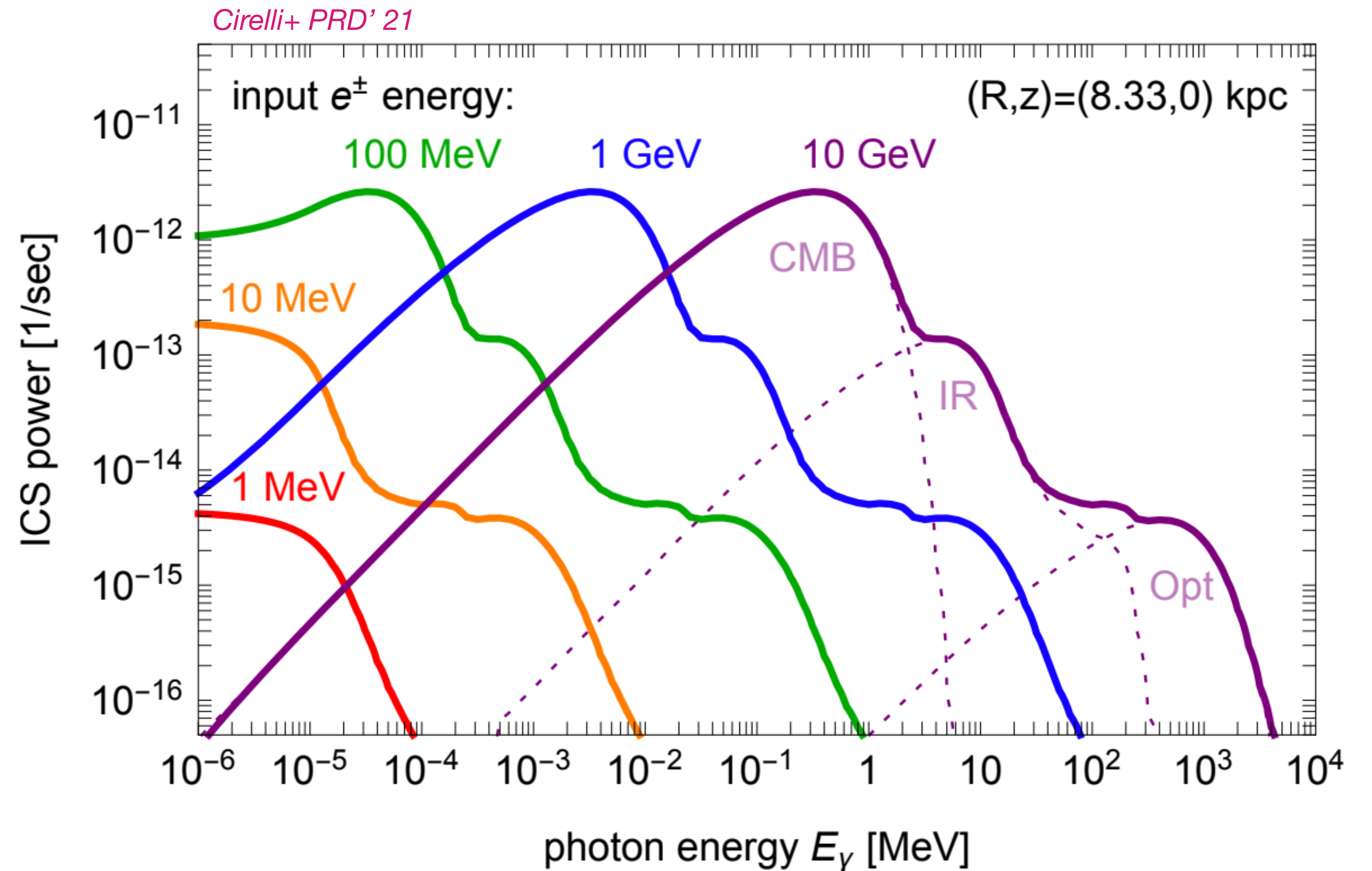
Radiative emission from leptons

Inverse Compton scattering



$$E_\gamma \approx \gamma^2 \varepsilon \quad ; \quad \varepsilon E_e \ll m_e^2$$

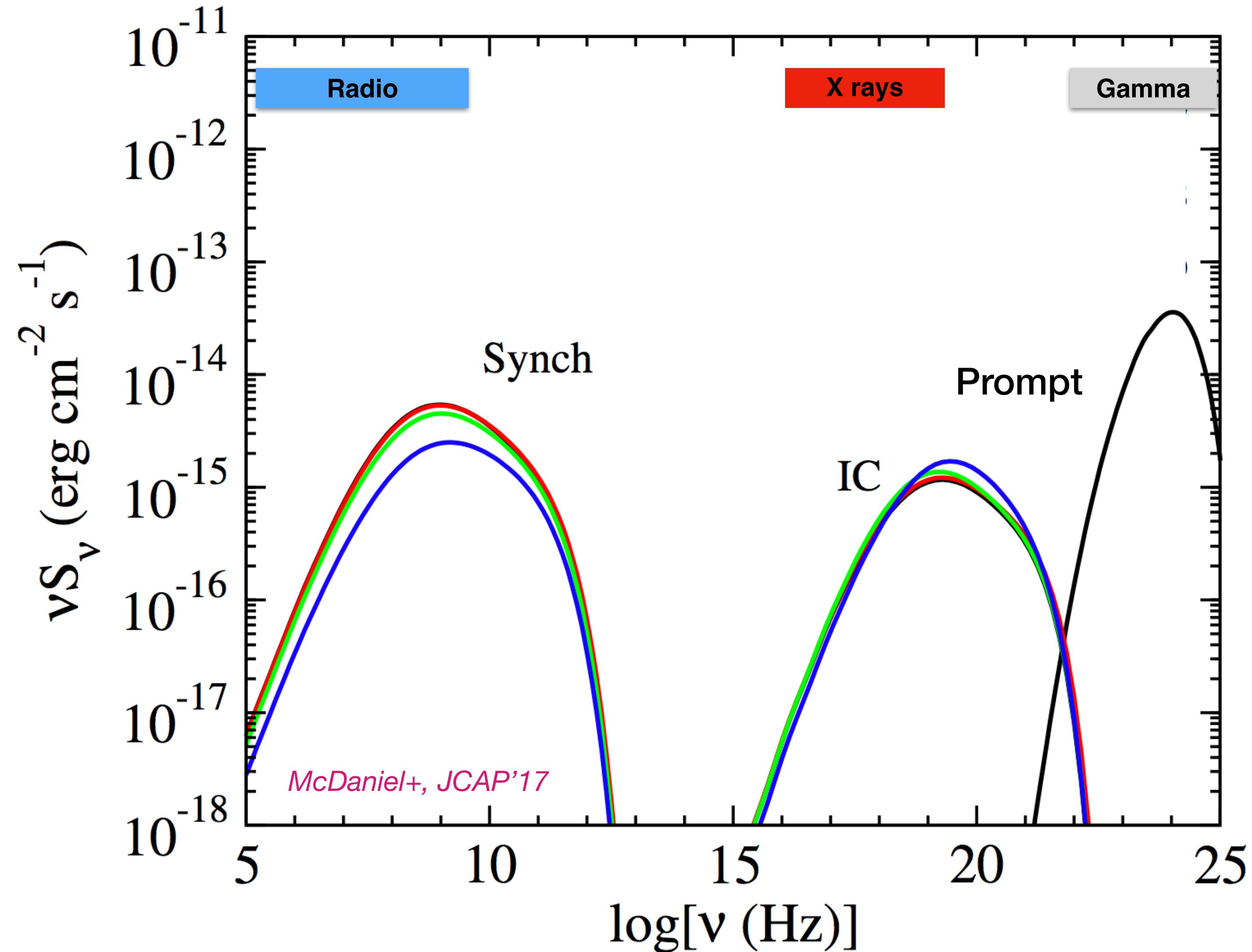
$$E_\gamma \approx 30 \left(\frac{\varepsilon}{\text{eV}} \right) \left(\frac{E_e}{\text{GeV}} \right)^2 \text{ MeV}$$



Secondary emission processes allow us to probe DM at much higher masses than prompt energy scales

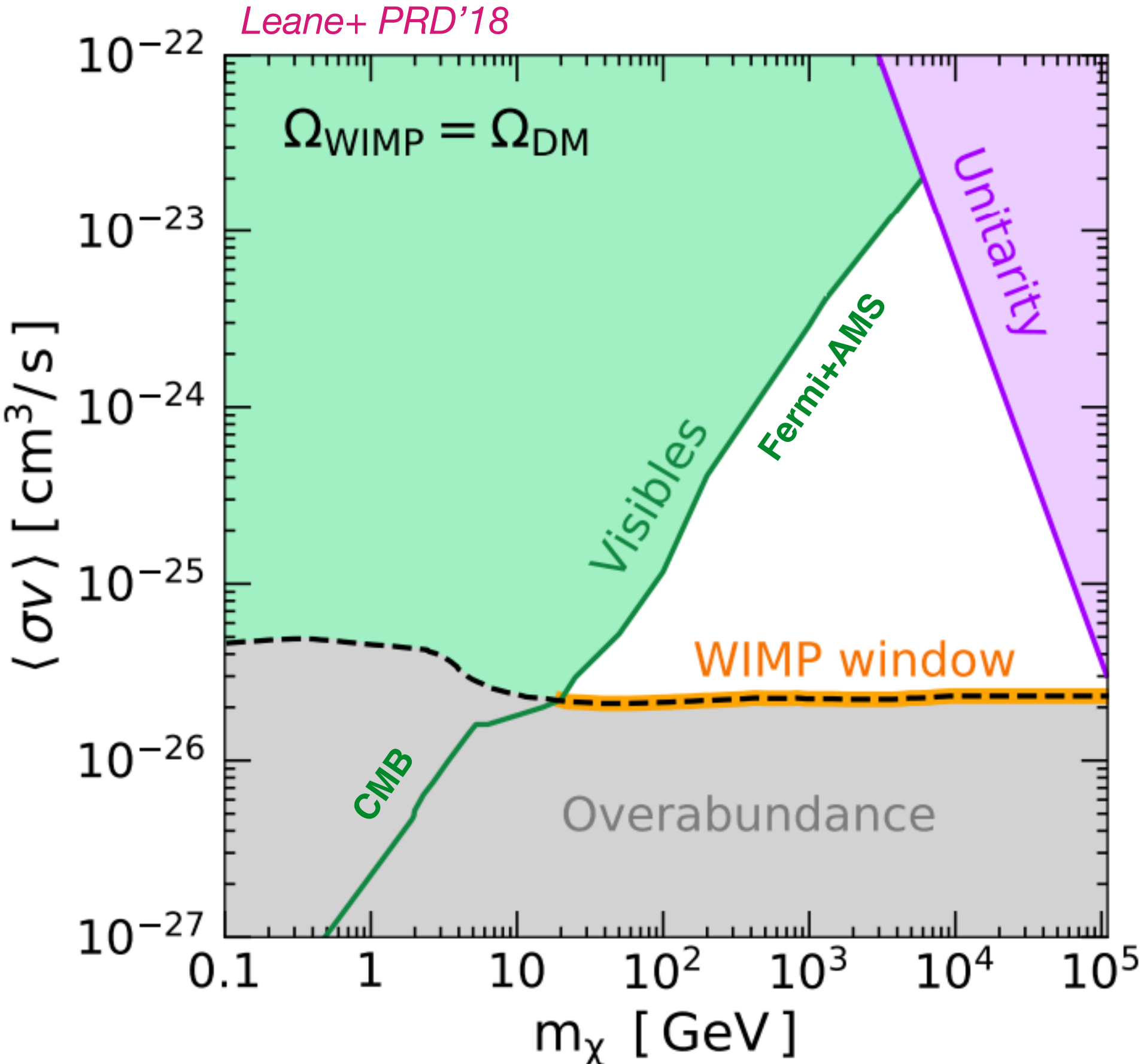
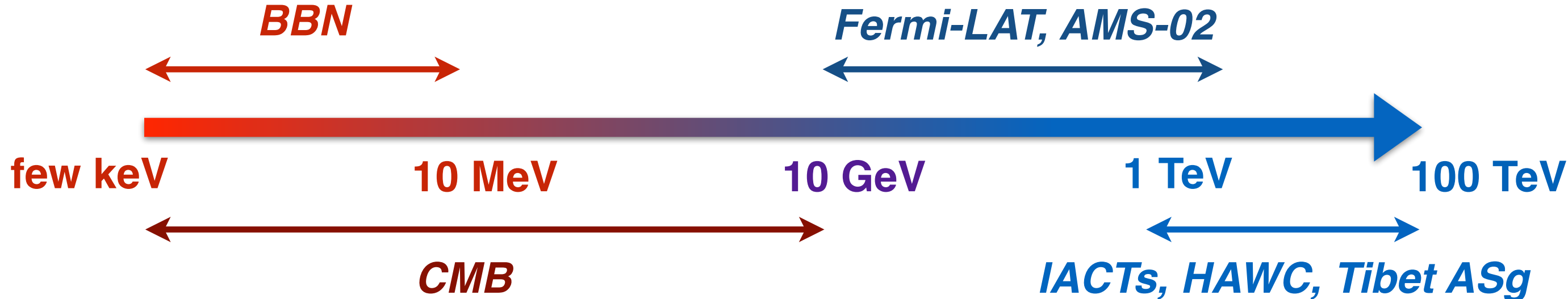
Multi-wavelength particle DM emission

“Secondary” radiative emission induced by leptonic particles interacting with the environment



*Example: Annihilation into b-quarks;
 $m_{\text{DM}} = 100$ GeV*

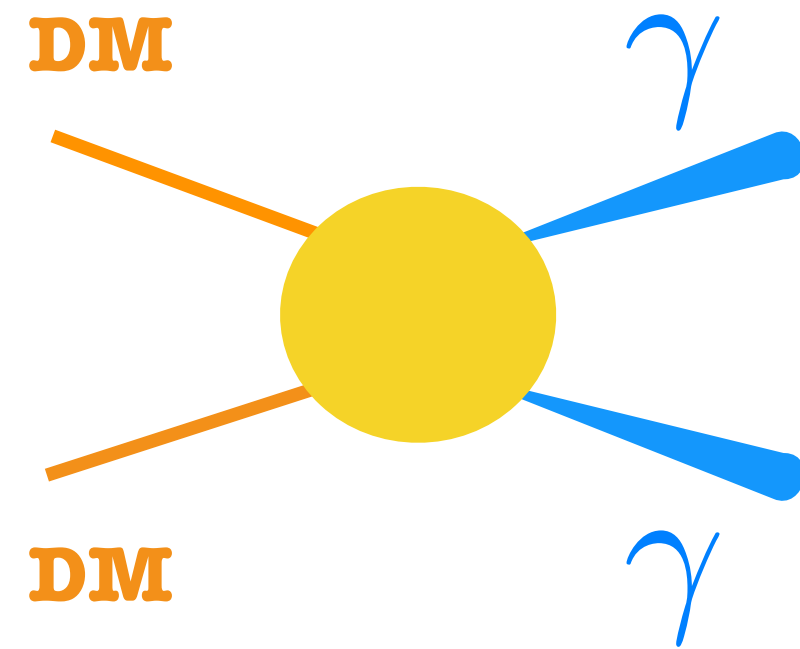
WIMP annihilation window



- **Total cross-section sets relic abundance**
 - **Indirect detection** provides model-independent UL on annihilation **cross-section for a given final state**
- ➔ Consistent and conservative interpretation of the data in the context of the generic thermal WIMP

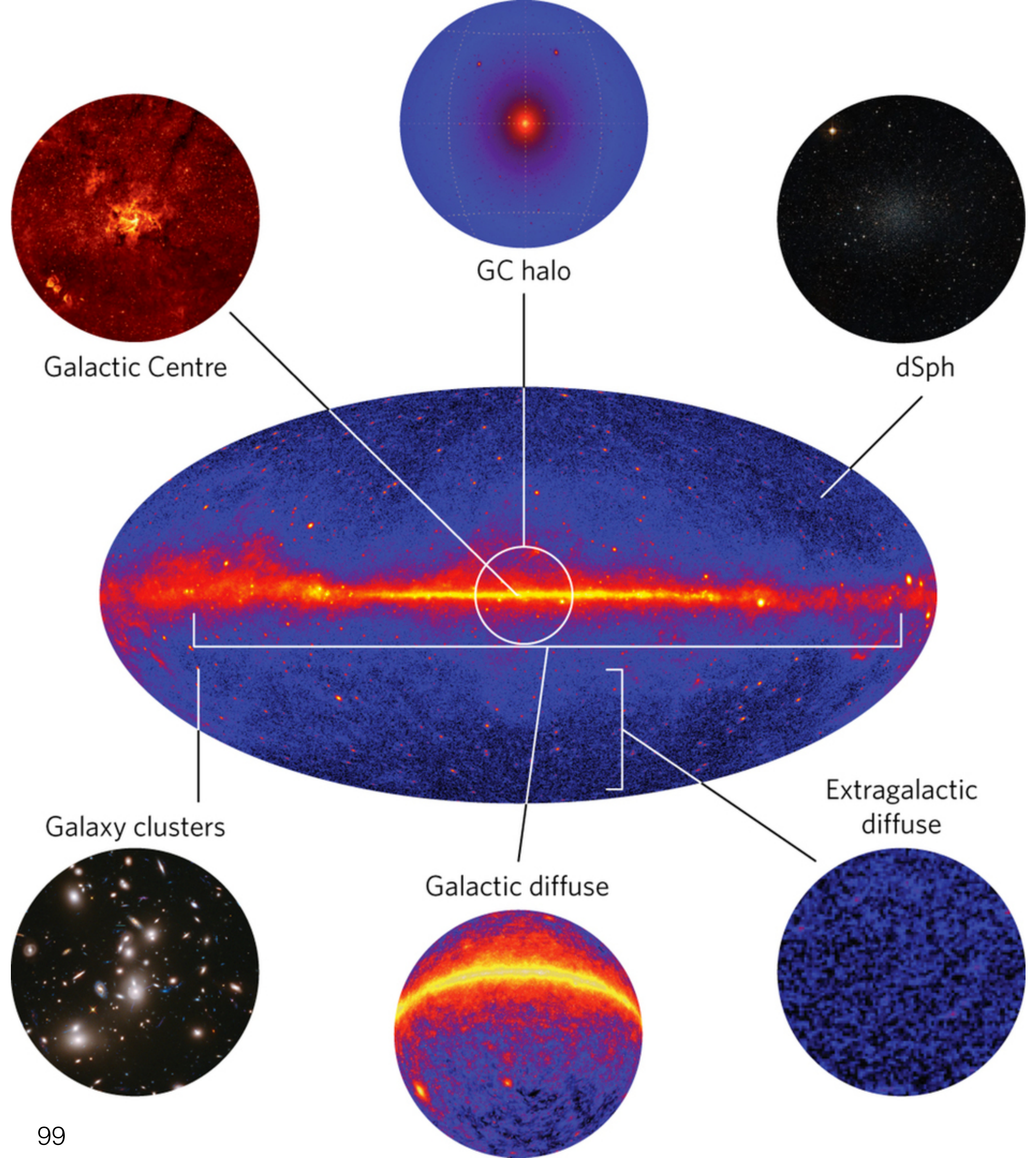
[Low DM masses constrained by energy injection at early times and CMB observations *Slatyer & Wu, PRD'17*]

Targets for WIMP gamma-ray searches



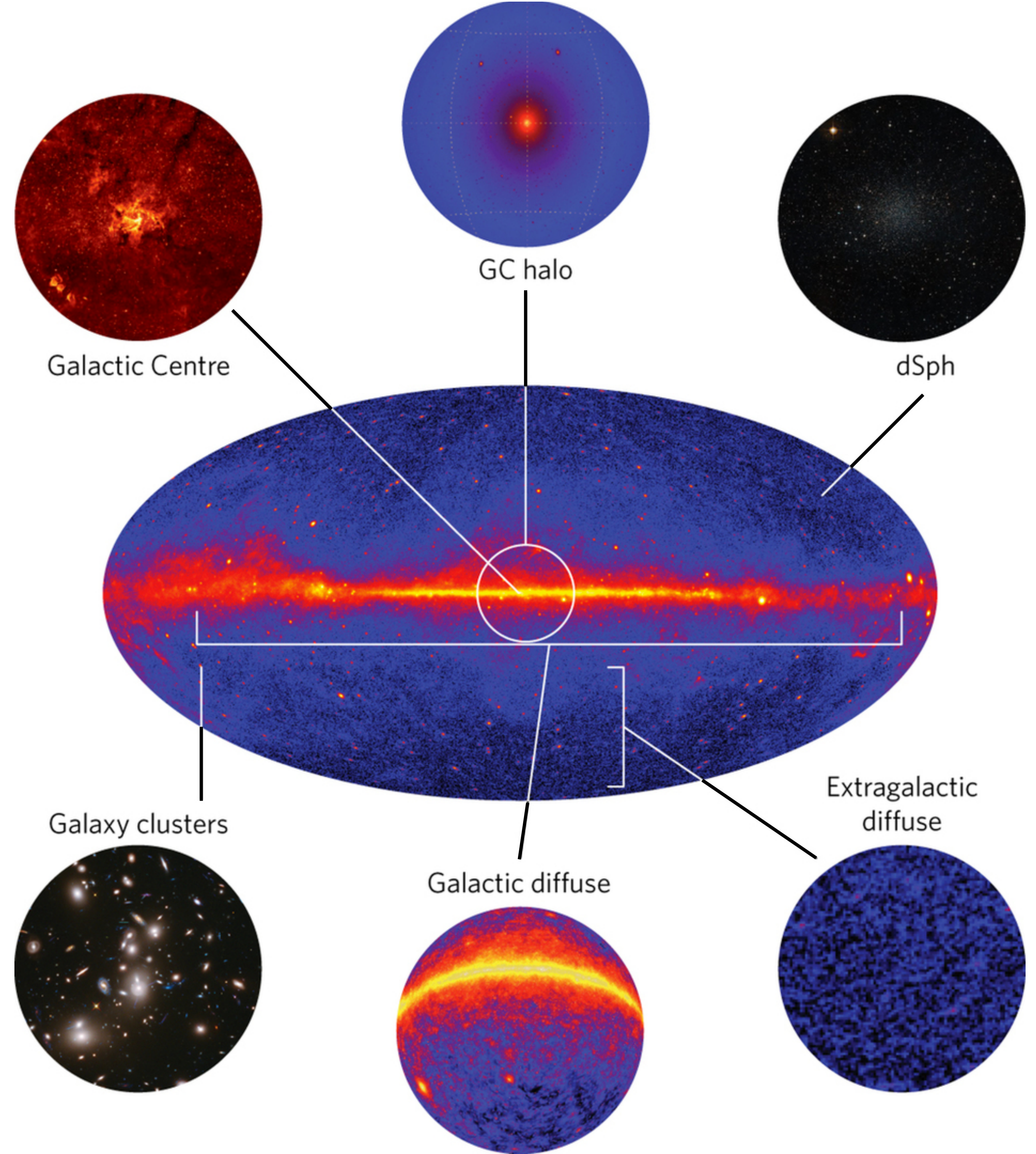
$$J \propto \int dl \rho [r(l, \psi)]^2$$

- + dedicated searches for gamma-ray lines
- + similar targets for radio searches (synchrotron)



Astrophysical Backgrounds

- **Guaranteed gamma-ray emission** from astrophysical sources and cosmic-ray interactions with gas and ambient photons
- **Modelling uncertainty** from cosmic-ray propagation conditions and target distributions (gas, radiation fields)
- Careful assessment is S/B and systematic modelling uncertainties



Feature-based searches

Different WIMP searches leverage on different WIMP (generic) features

- Generally, the signal looks like a **smooth bump** (from decay/hadronization products) from the main, smooth, **Galactic halo** => Importance of astrophysical background modelling
- **Spectral features**: look for sharp (or less sharp) features at the high-end of the energy spectrum
- **Spatial (angular) features**: look for specific DM-dominated targets and/or for angular correlations in the sky (anisotropies/cross-correlation)

=> Example: **Searches for DM towards dwarf spheroidal galaxies**

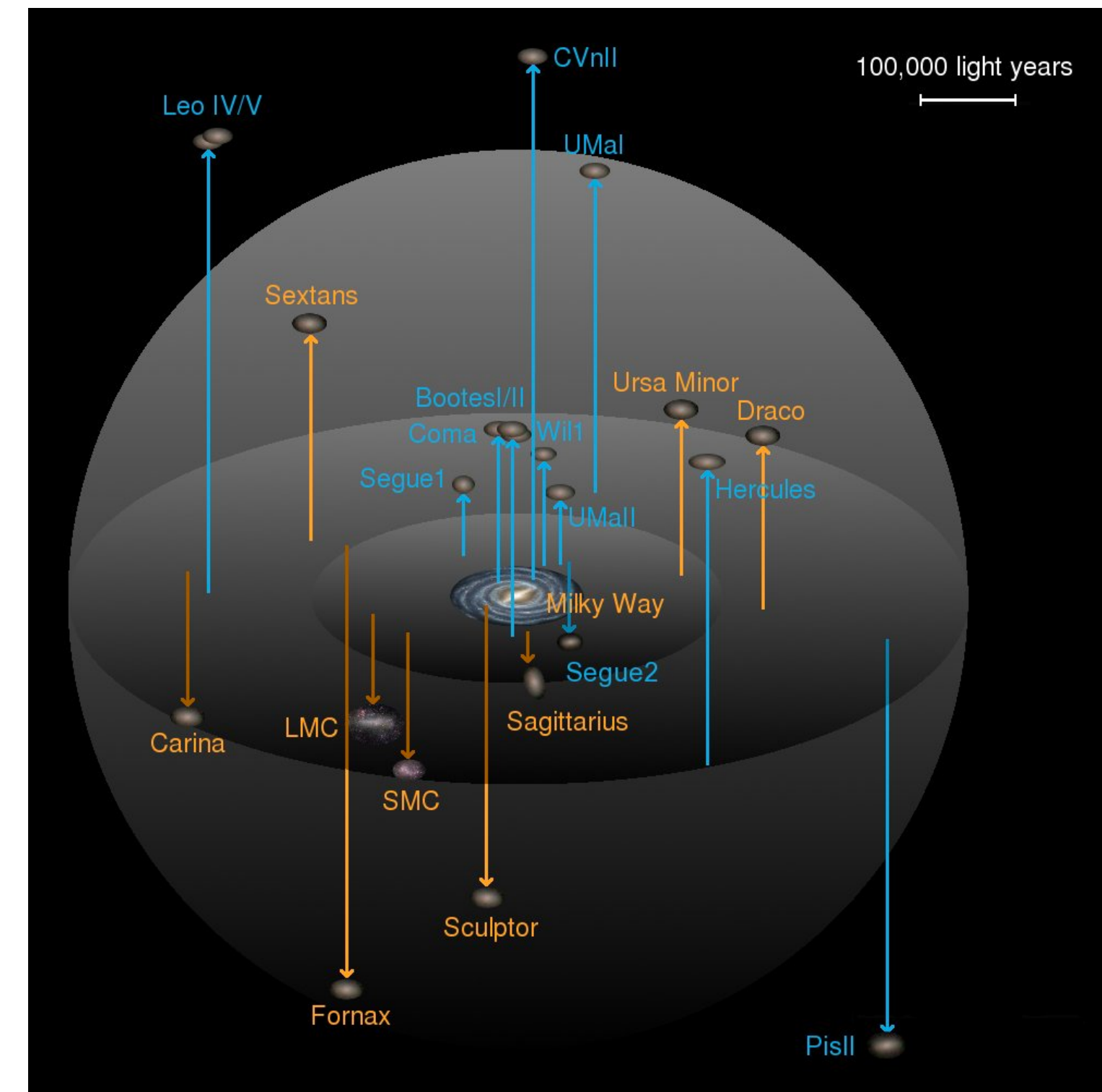
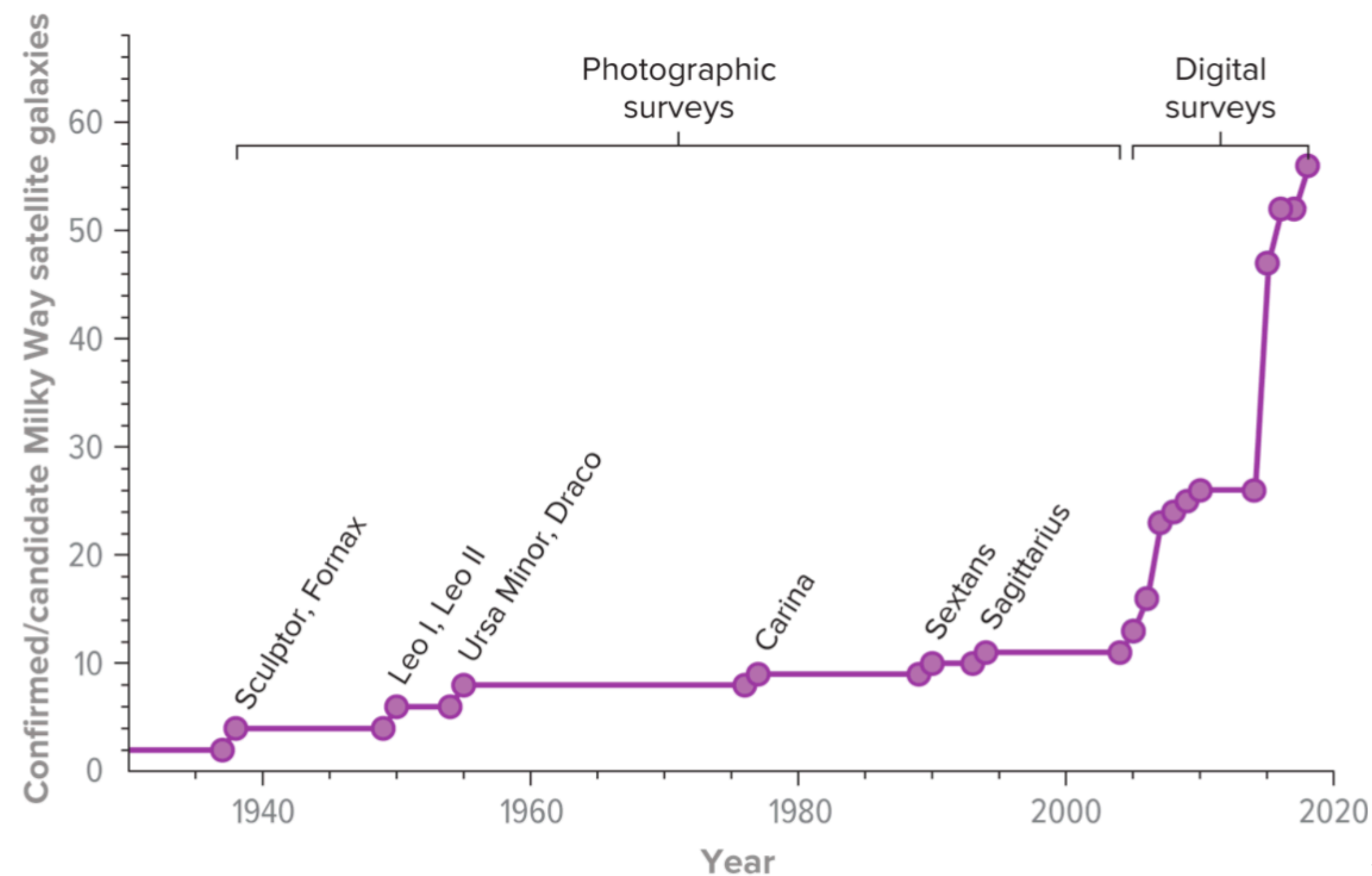
[NB: Models are predictive => can link different observables (direct/indirect/collider) and break degeneracies]

Dwarf spheroidal galaxies

Known satellites of the Milky Way at ~100 kpc from Earth

“Clean” target for DM searches, high mass-to-light ratio and little astrophysical emission

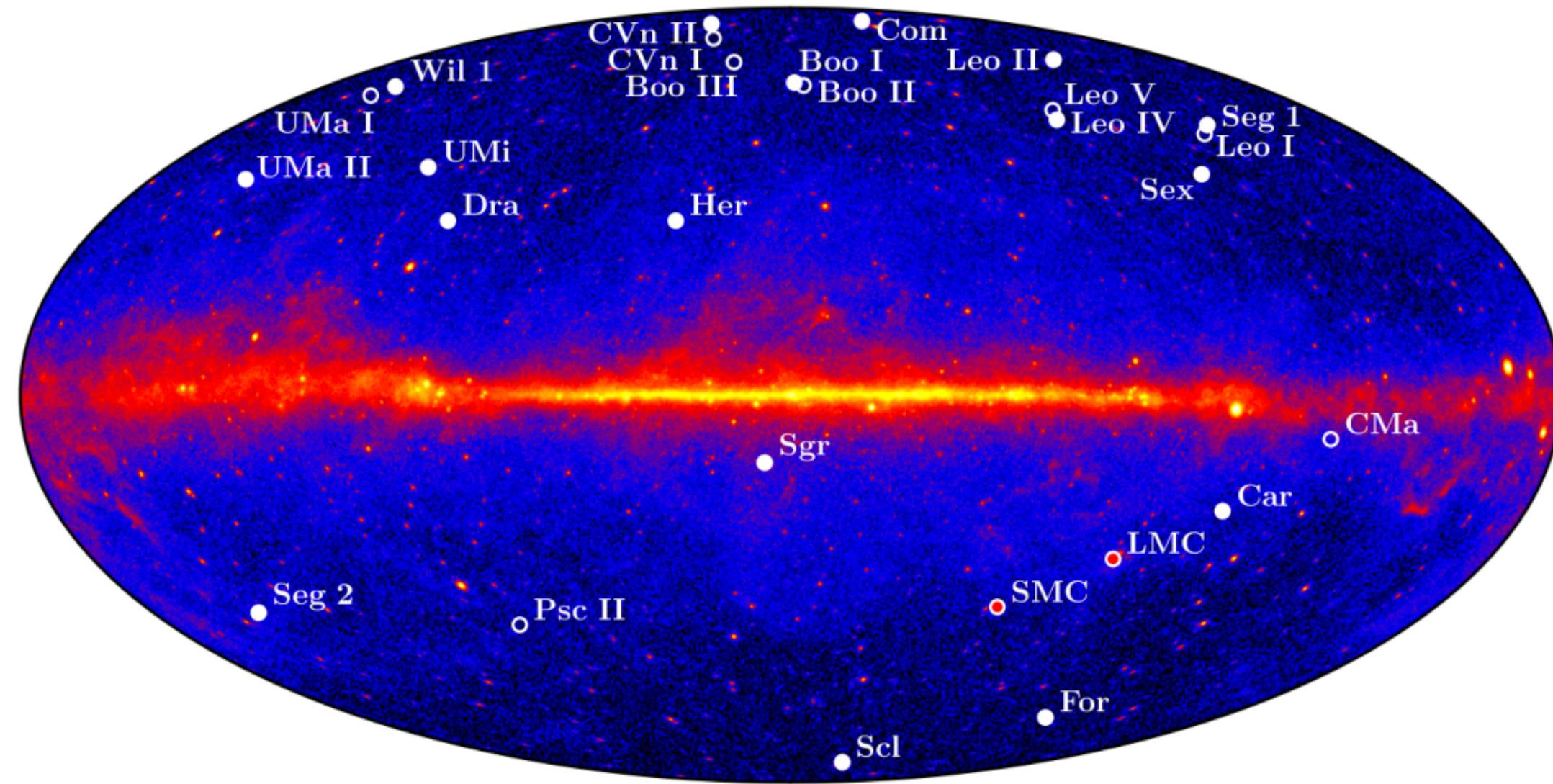
Winter+ ApJ'16



A growing Galactic crowd
> 50 satellites
(SDSS, PanSTARRS, DES)

Credit: J.D. Simon / AR Astronomy and Astrophysics

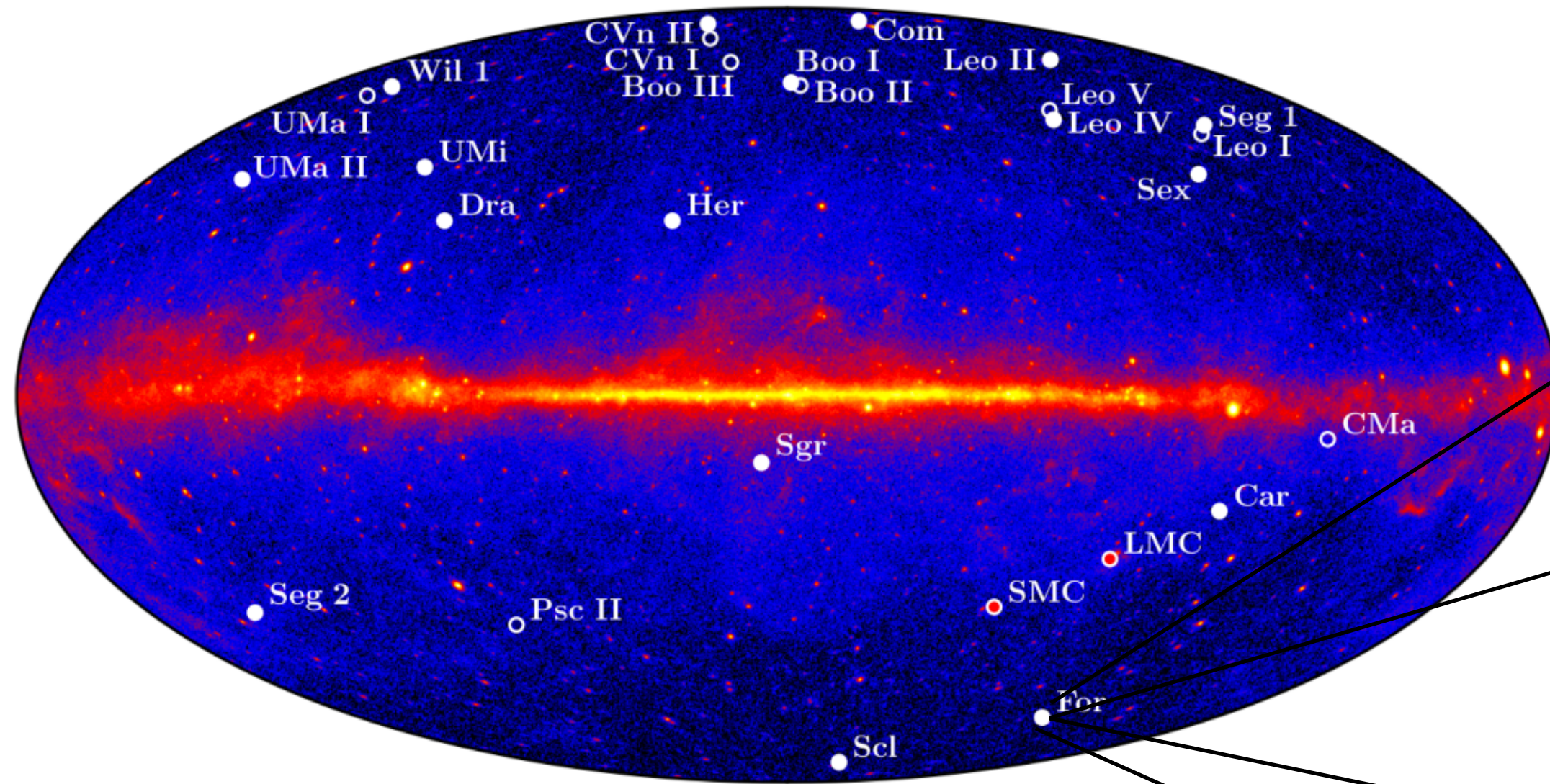
Limits from dwarf spheroidal galaxies



$$J \propto \int dl \rho [r(l, \psi)]^2$$

Fermi-LAT Collaboration, PRL'11

Limits from dwarf spheroidal galaxies

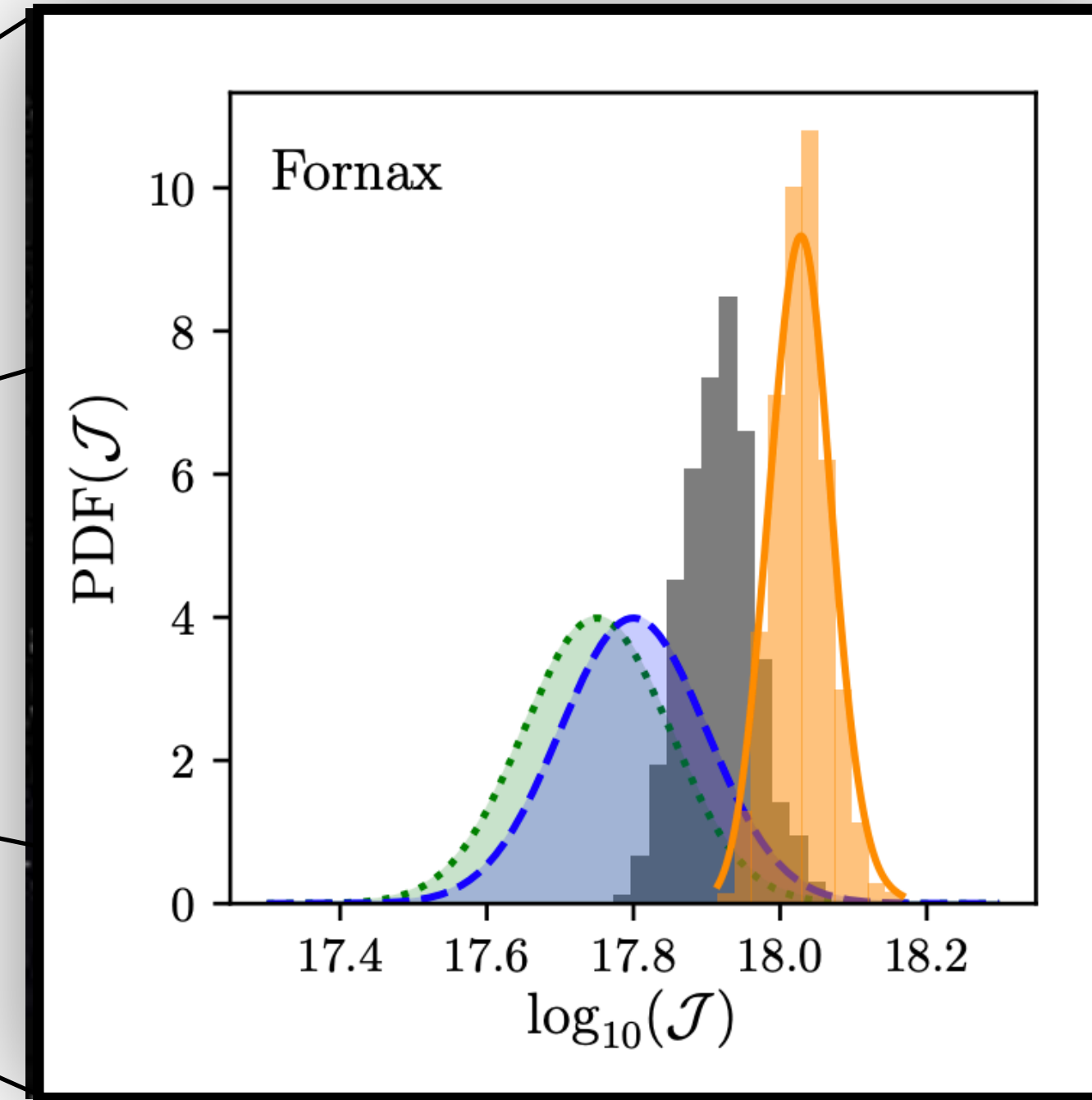
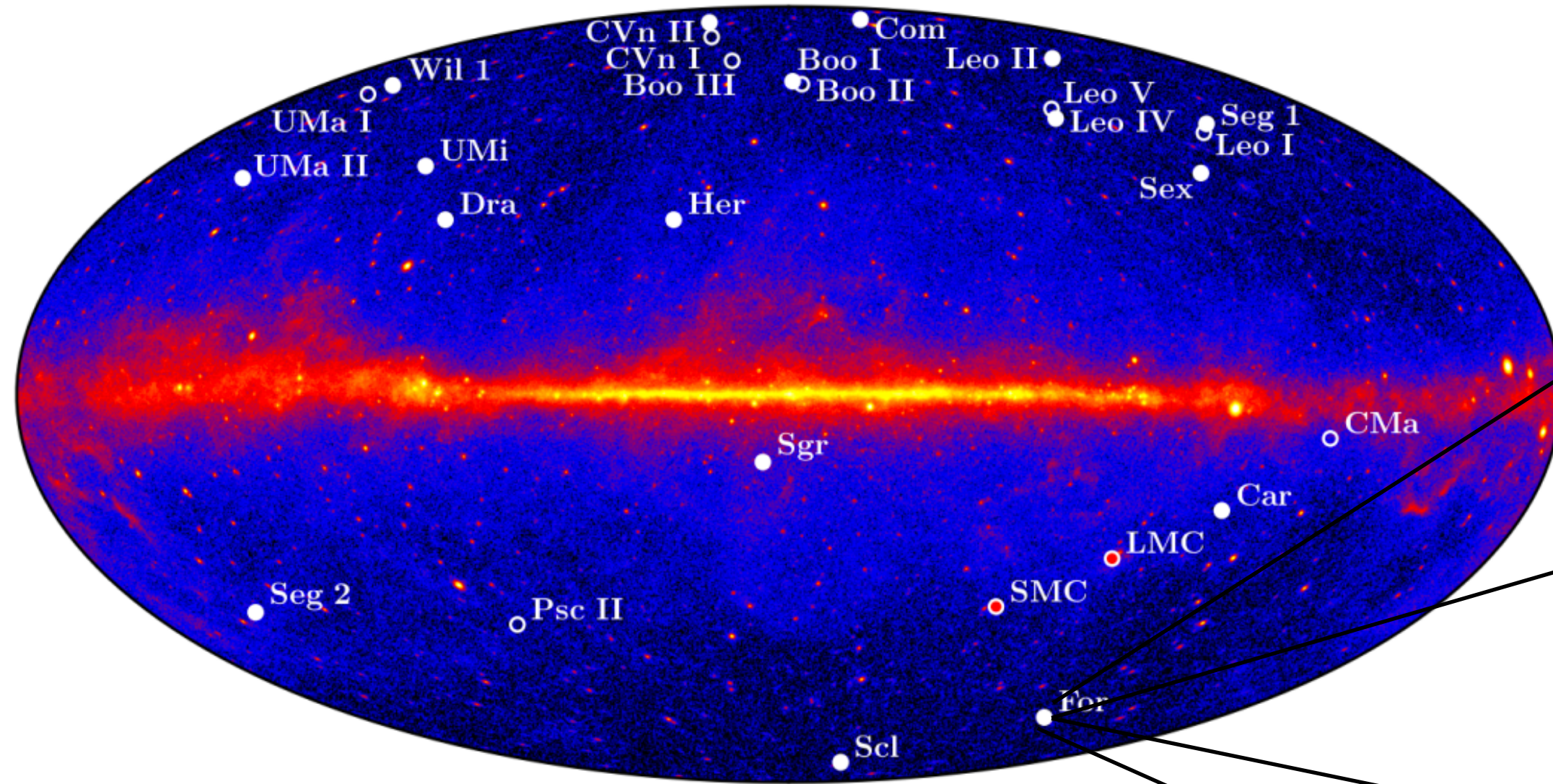


$$J \propto \int dl \rho [r(l, \psi)]^2$$

Fermi-LAT Collaboration, PRL'11

Credit: ESO/Fornax galaxy

Limits from dwarf spheroidal galaxies

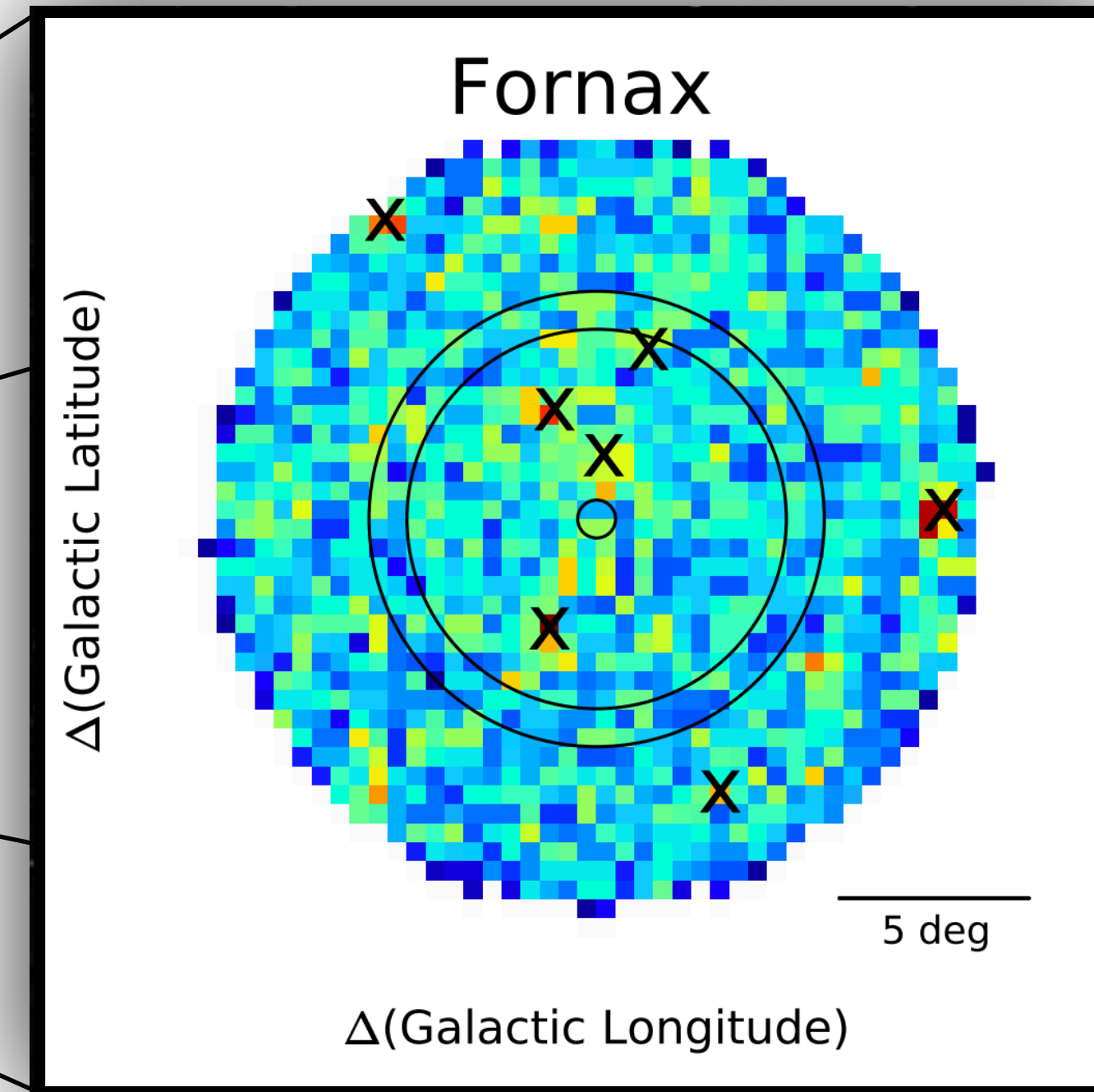
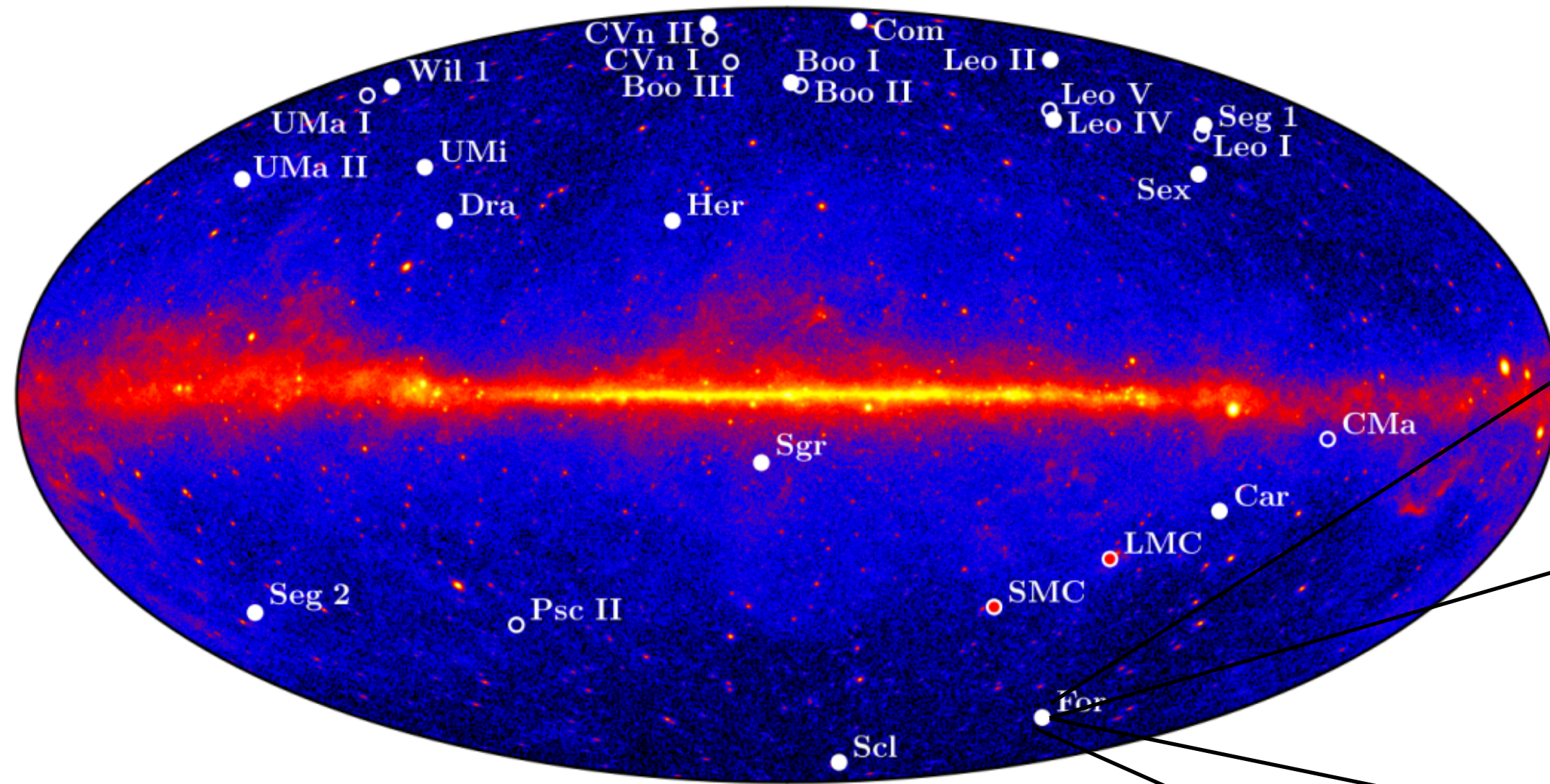


$$J \propto \int dl \rho [r(l, \psi)]^2$$

Fermi-LAT Collaboration, PRL'11

*GRAVSPHERE
Alvarez, FC+ JCAP'20*

Limits from dwarf spheroidal galaxies

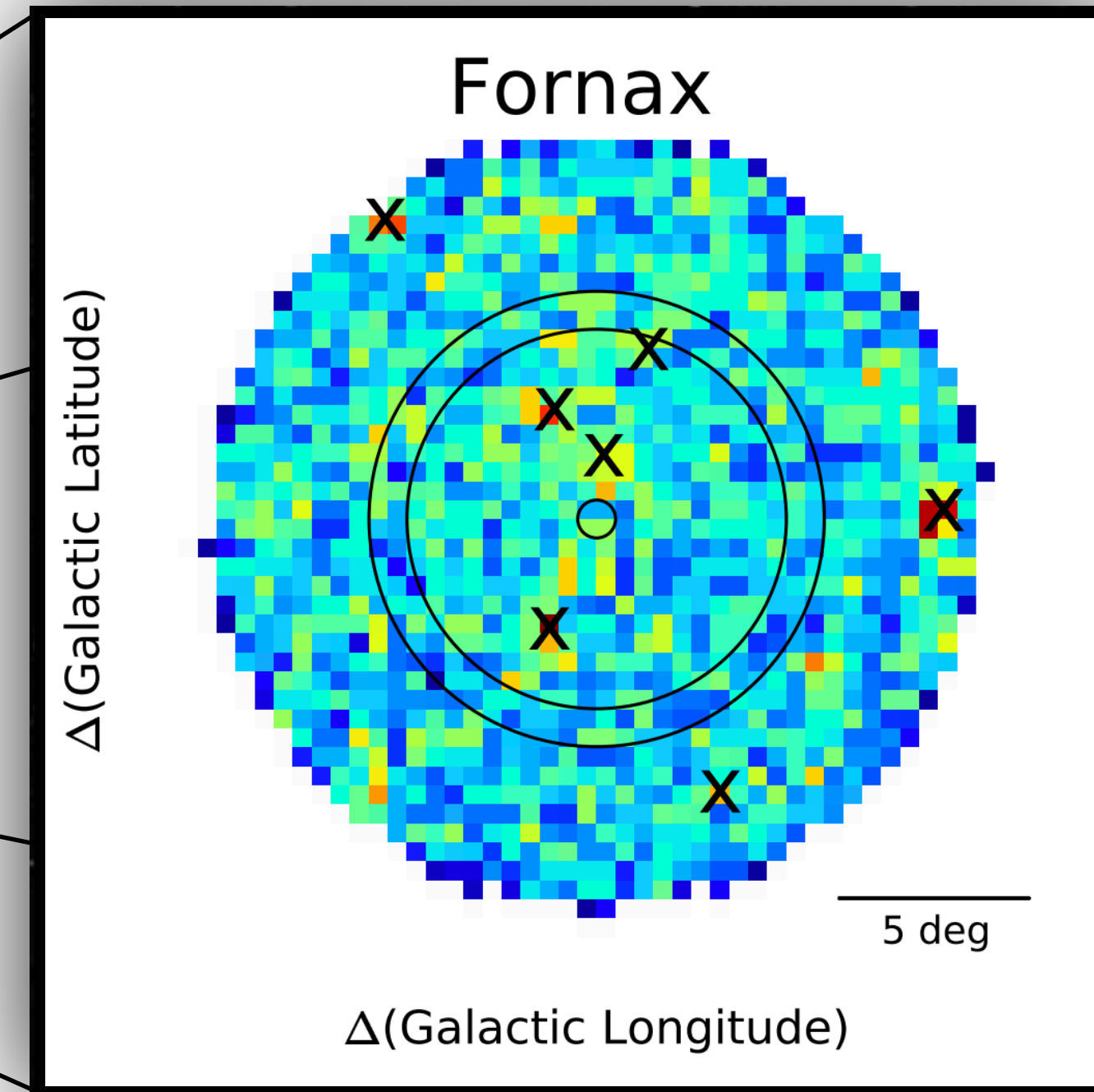
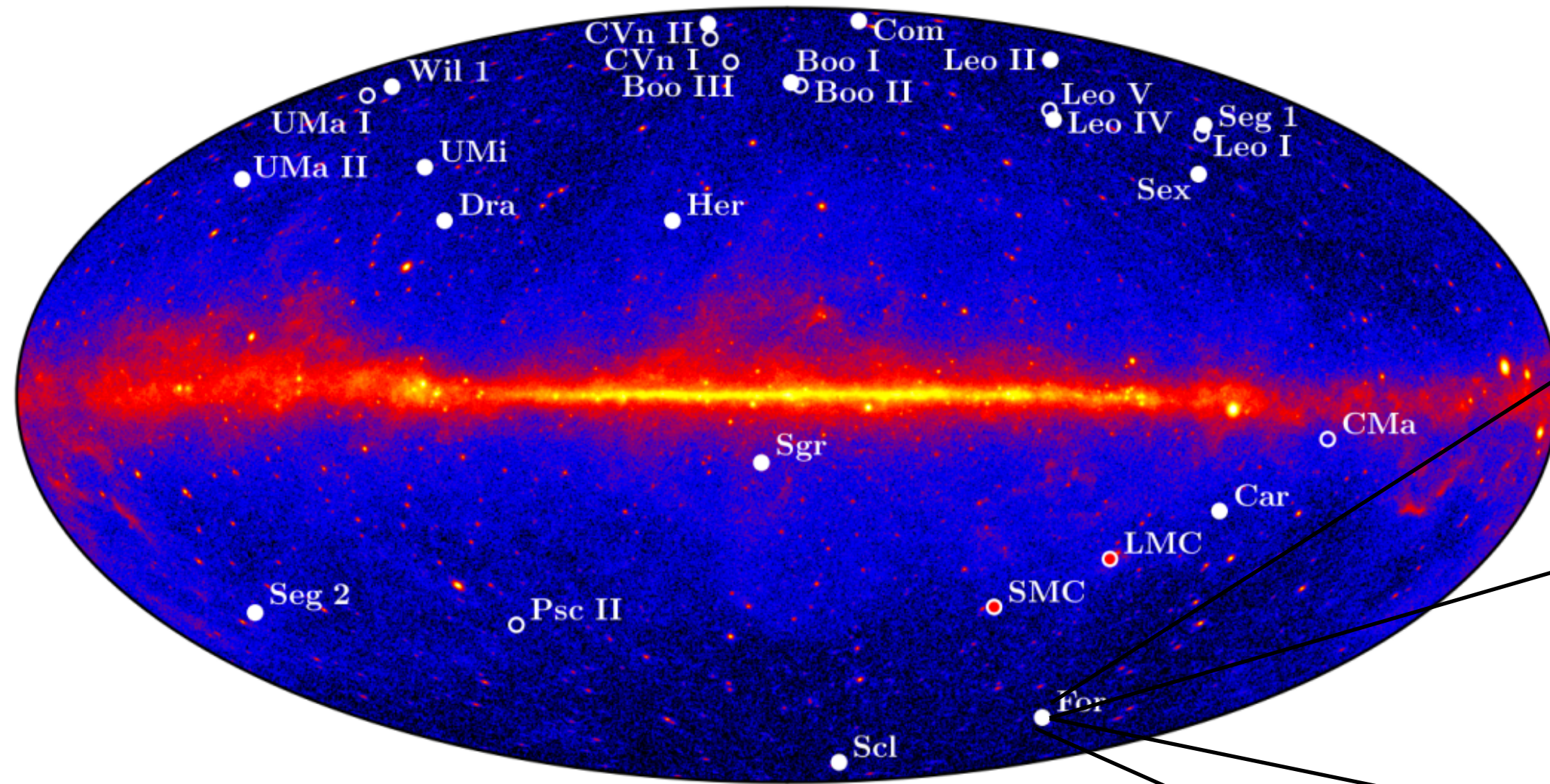


$$J \propto \int dl \rho [r(l, \psi)]^2$$

Fermi-LAT Collaboration, PRL'11

Mazziotta+Astrop. Phys.'12

Limits from dwarf spheroidal galaxies



$$J \propto \int dl \rho [r(l, \psi)]^2$$

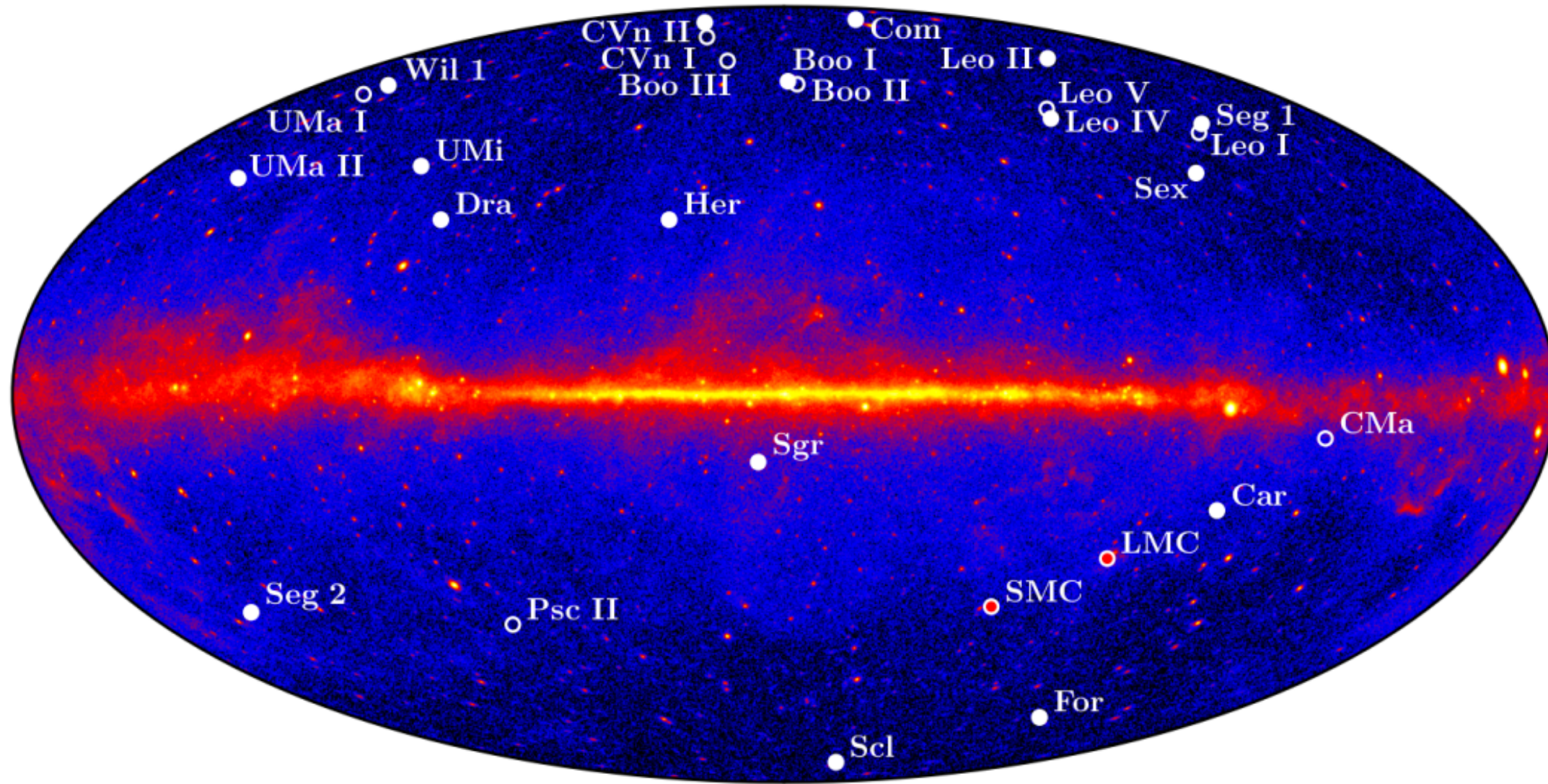
Fermi-LAT Collaboration, PRL'11

Mazziotta+Astrop. Phys.'12

Evidence for additional **DM signal?**

$$\frac{d\Phi_\gamma}{dE}(\ell, b) = \left(\frac{d\Phi_\gamma}{dE} \right)_{\text{diffuse}} + \left(\frac{d\Phi_\gamma}{dE} \right)_{\text{PS}} + \left(\frac{d\Phi_\gamma}{dE} \right)_{\text{DM signal}}$$

Limits from dwarf spheroidal galaxies

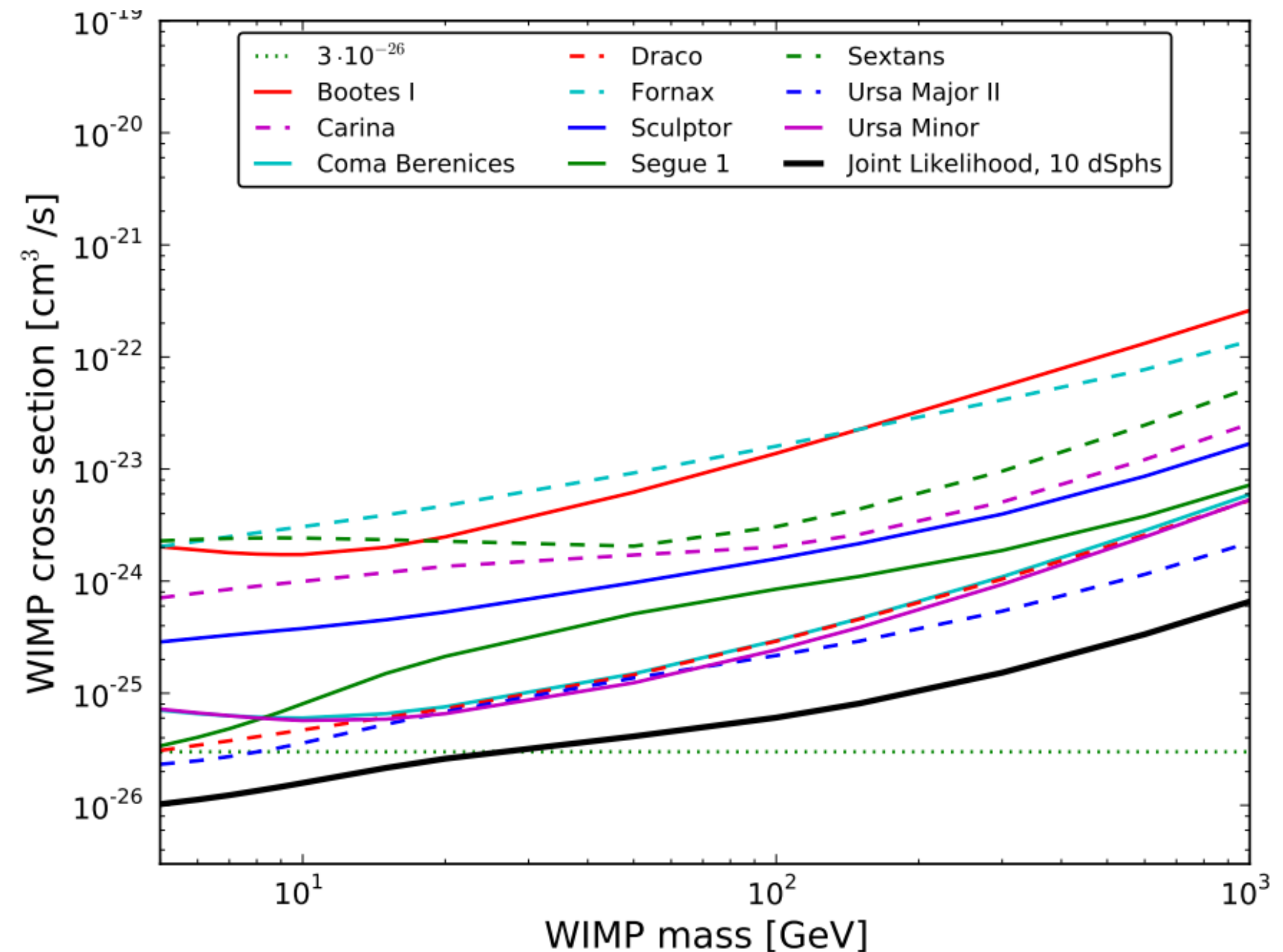


$$L(D|\mathbf{p}_W, \{\mathbf{p}\}_i) = \prod_i L_i^{\text{LAT}}(D|\mathbf{p}_W, \mathbf{p}_i) \times \frac{1}{\ln(10) J_i \sqrt{2\pi\sigma_i}} e^{-[\log_{10}(J_i) - \overline{\log_{10}(J_i)}]^2 / 2\sigma_i^2}$$

Analysing dSphs as a group results in sensitivity competitive with other targets => **Stacking technique**

Fermi-LAT Collaboration, PRL'11

$$J \propto \int dl \rho [r(l, \psi)]^2$$

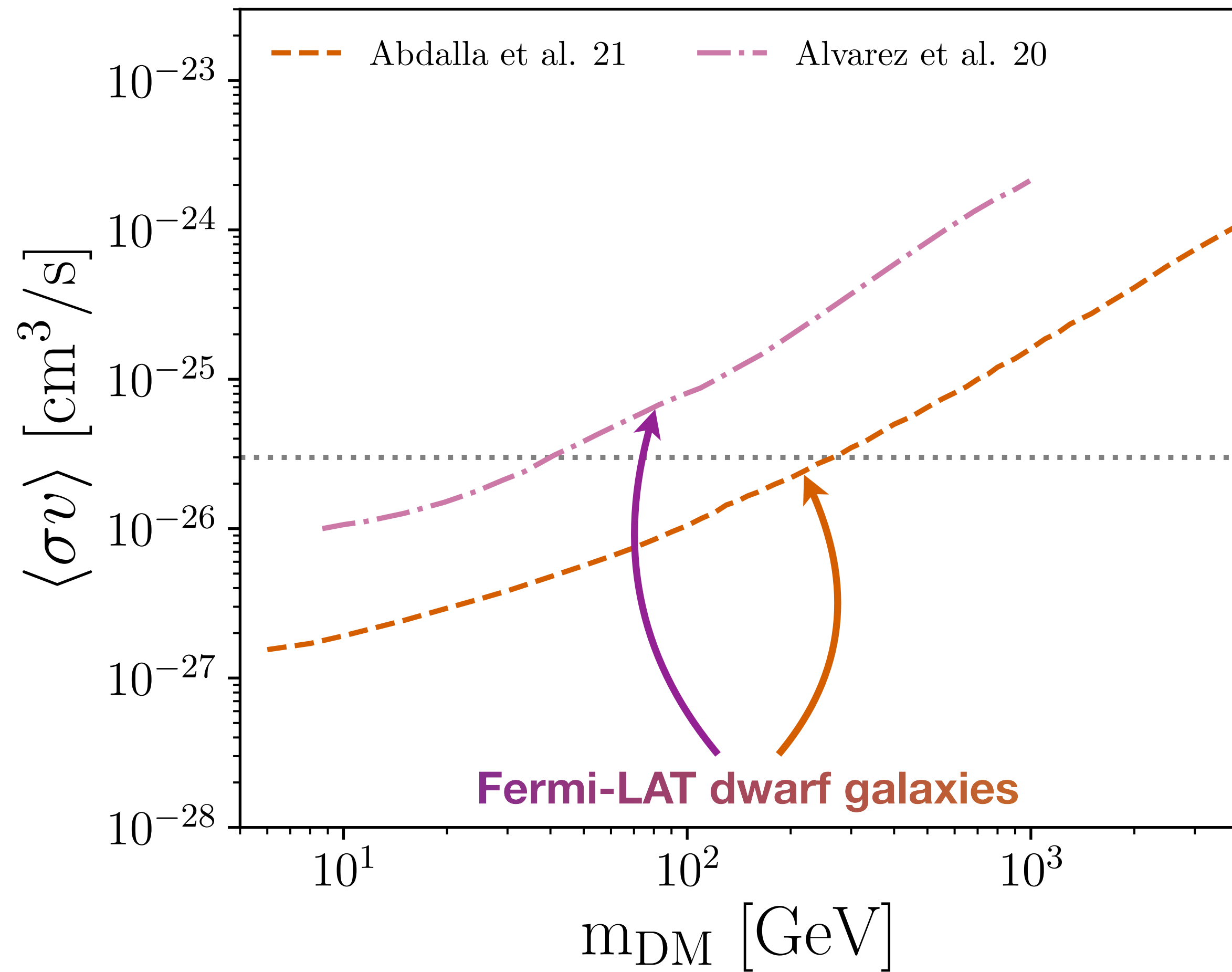


Limits on annihilating WIMPs



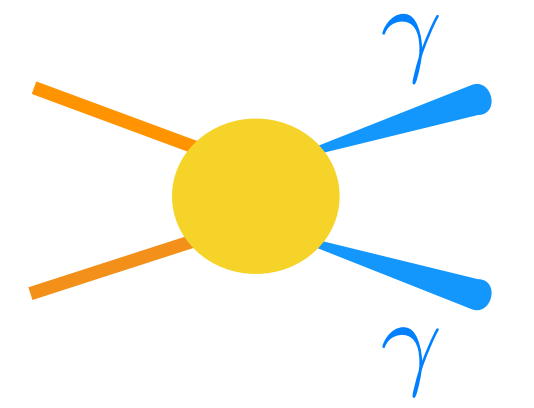
Gamma rays: Dwarf spheroidal galaxies

~ a few GeV — few TeV

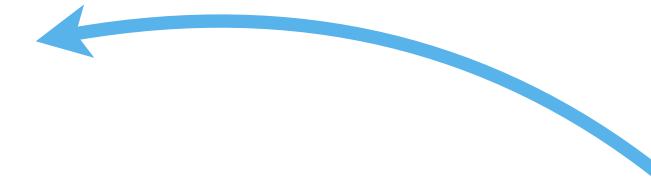
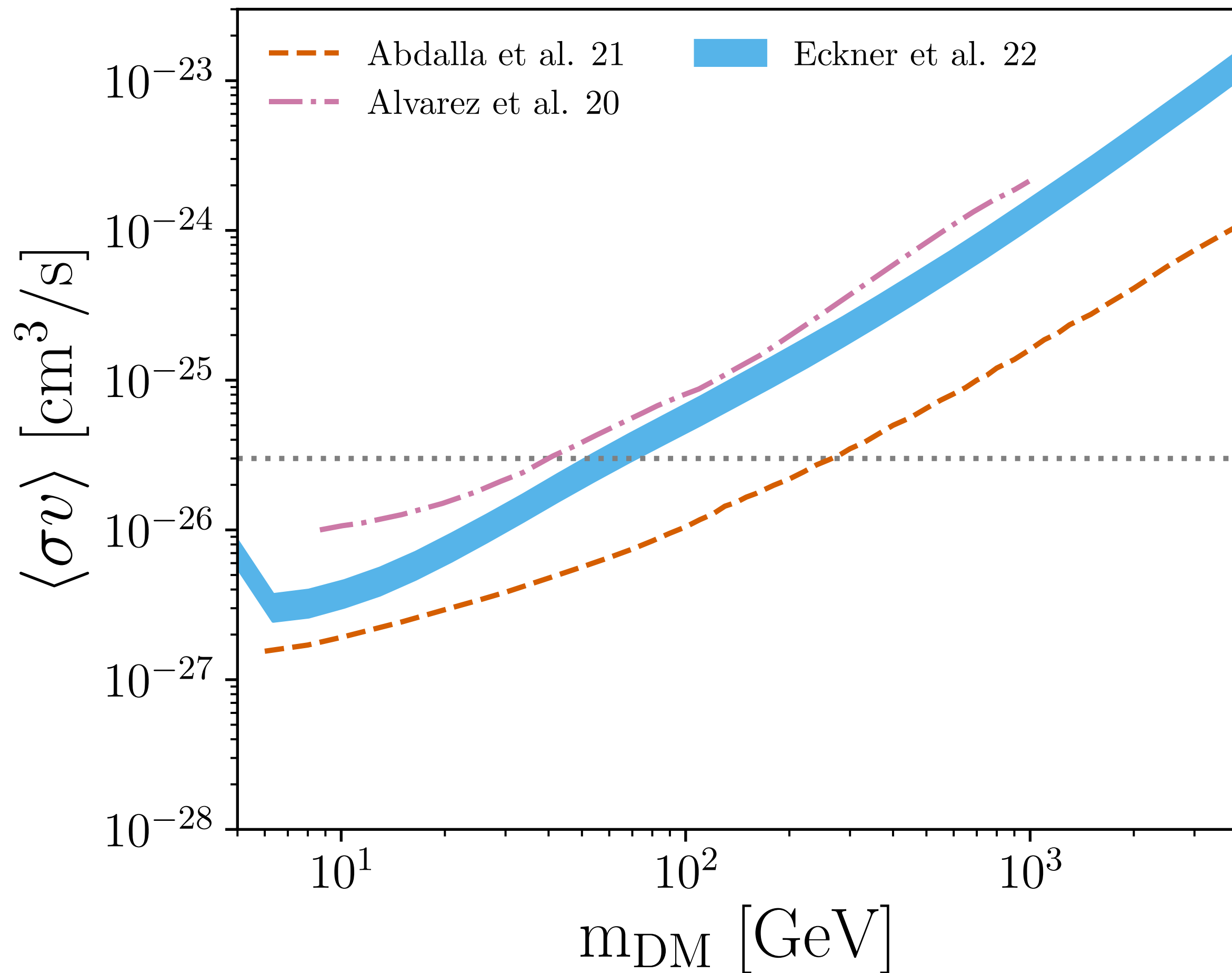


Limits on annihilating WIMPs

Gamma rays: High-latitude MW halo

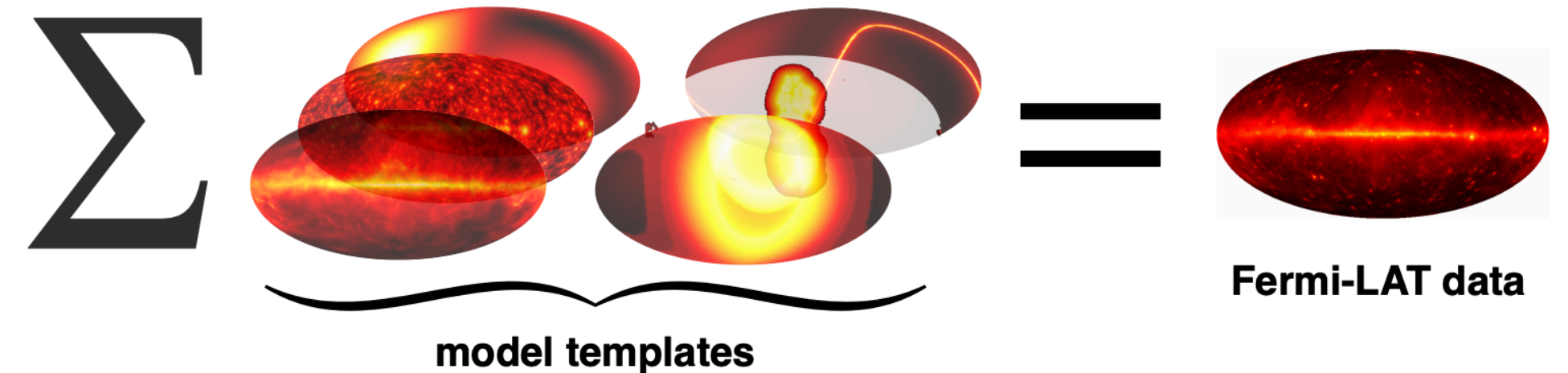
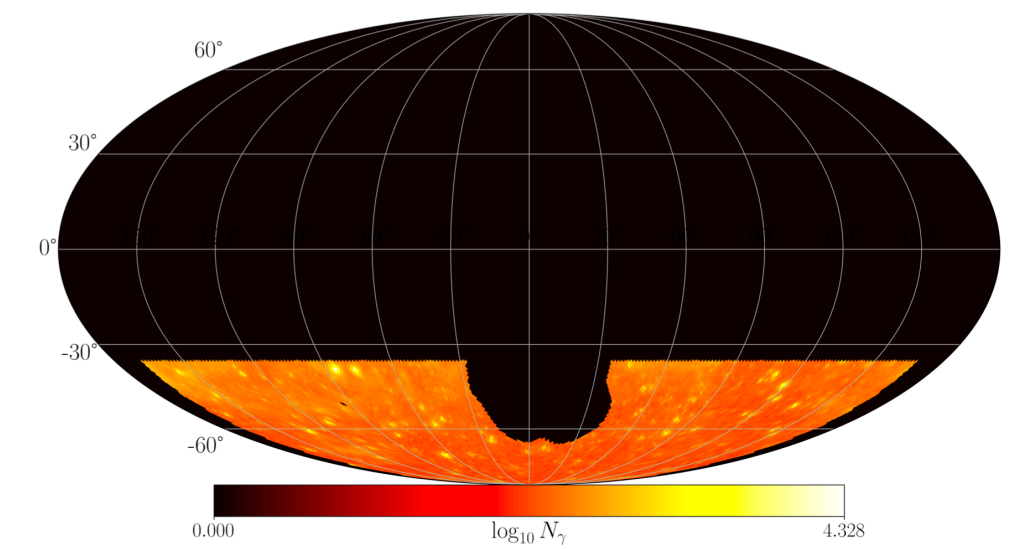


~ a few GeV — few TeV



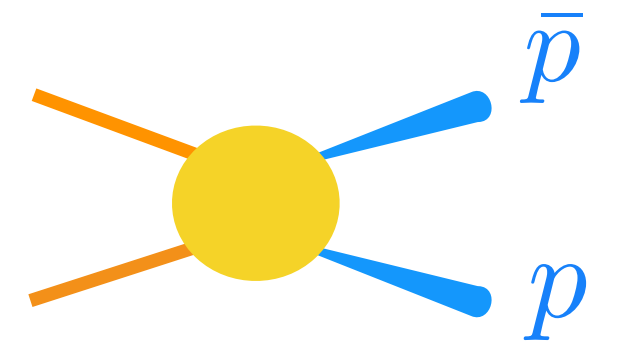
High-latitude
Fermi-LAT sky

Eckner, FC+ MNRAS'22
Zechlin+ PRD'18
Chang+ PRD'18

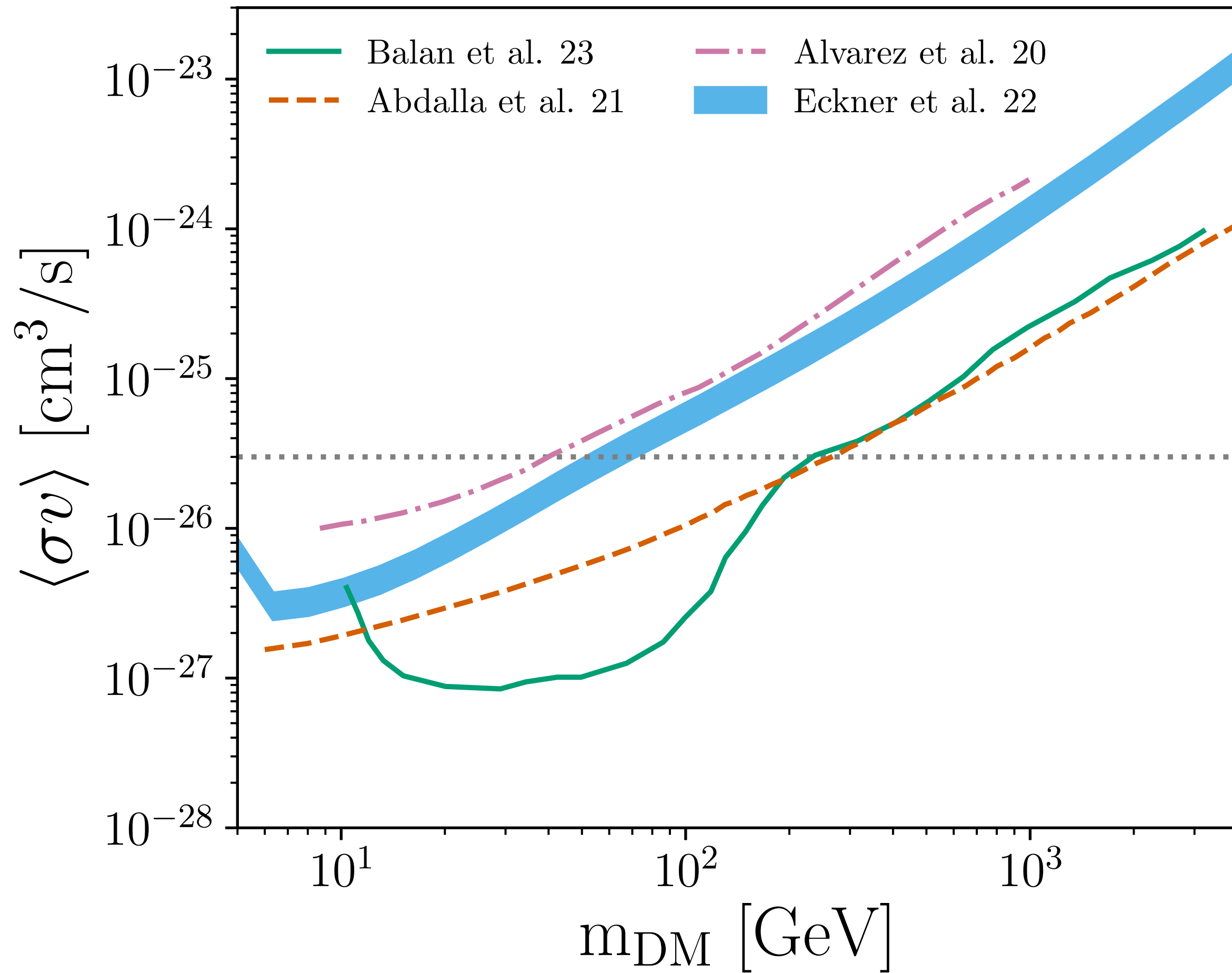


Limits on annihilating WIMPs

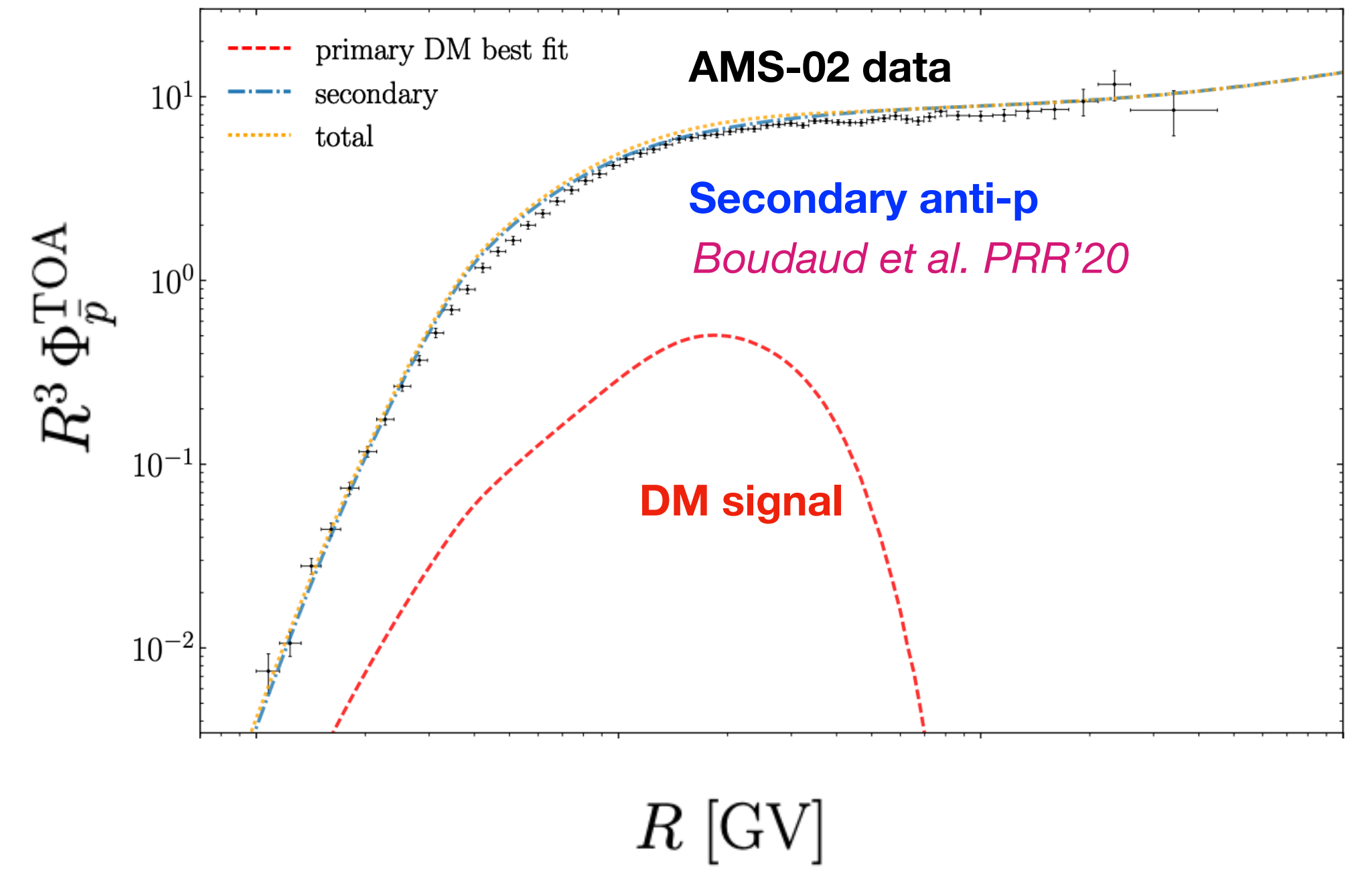
Anti-protons flux from the MW



~ a few GeV — few TeV

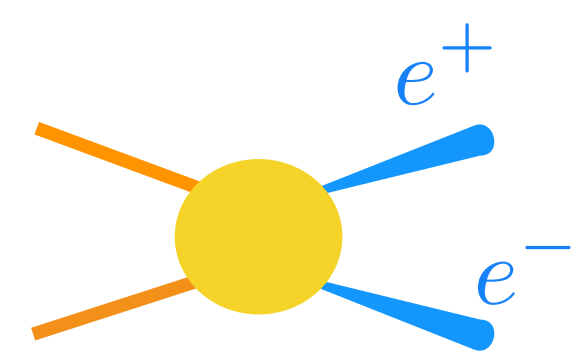


Anti-protons
Balan+ arXiv:2303.07362
Di Mauro+ PRD'21,
FC+ SciPost'22

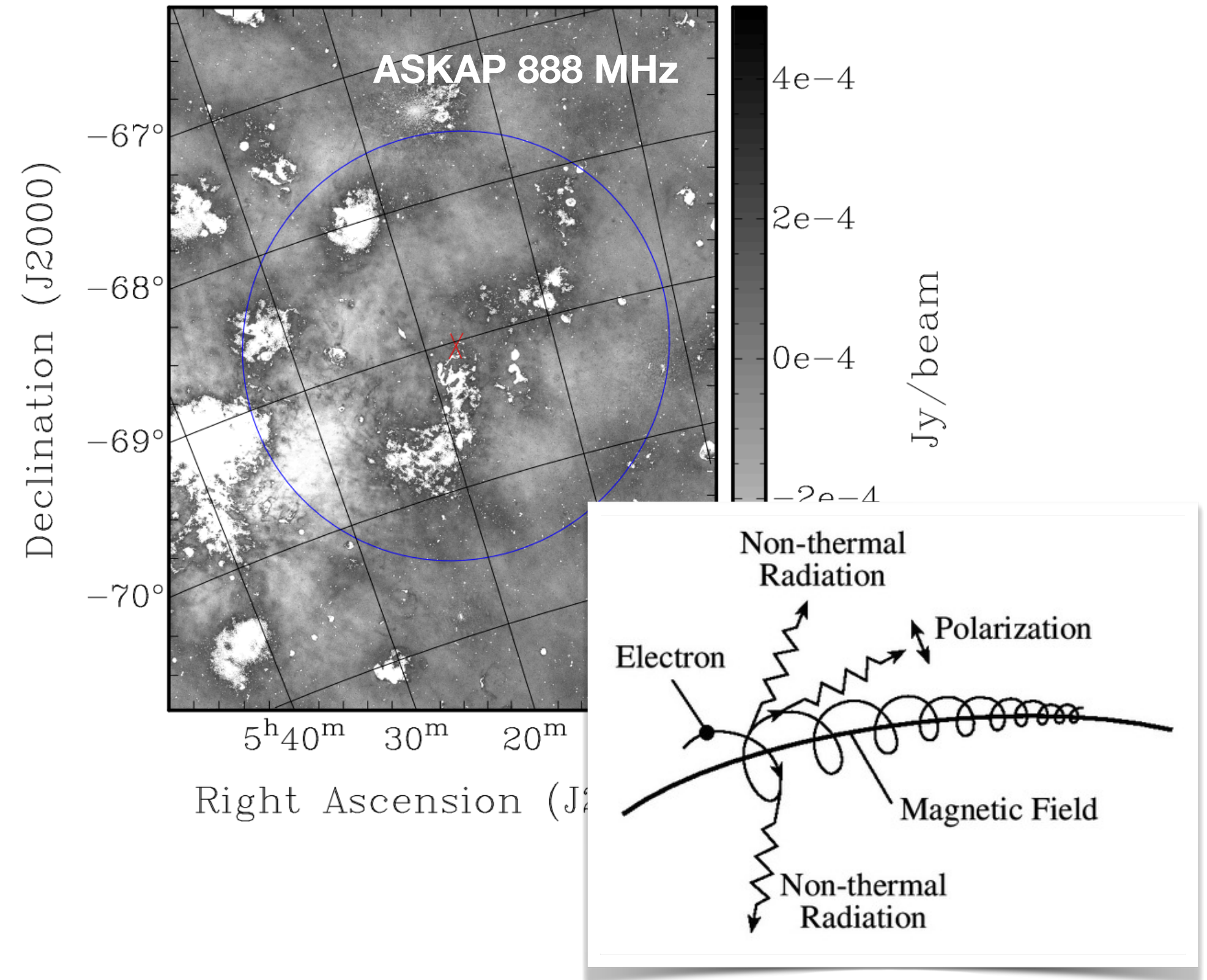
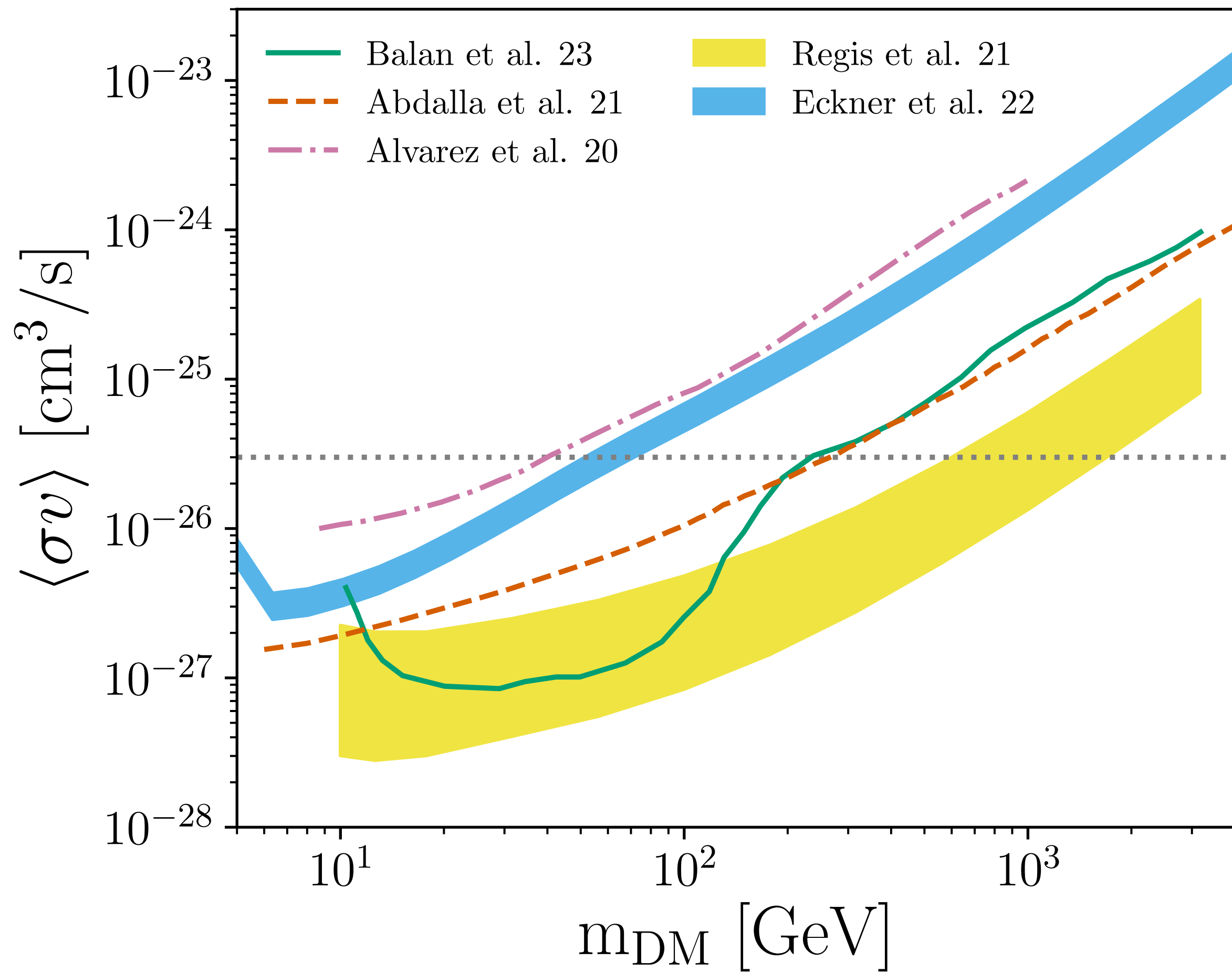


Limits on annihilating WIMPs

Synchrotron emission in the Large Magellanic Cloud



~ a few GeV — few TeV

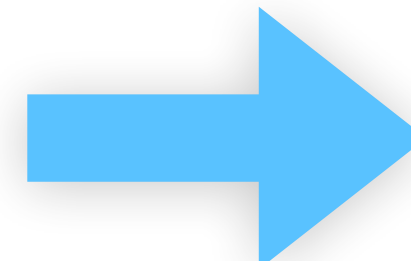


Radio LMC
Regis+ JCAP'21

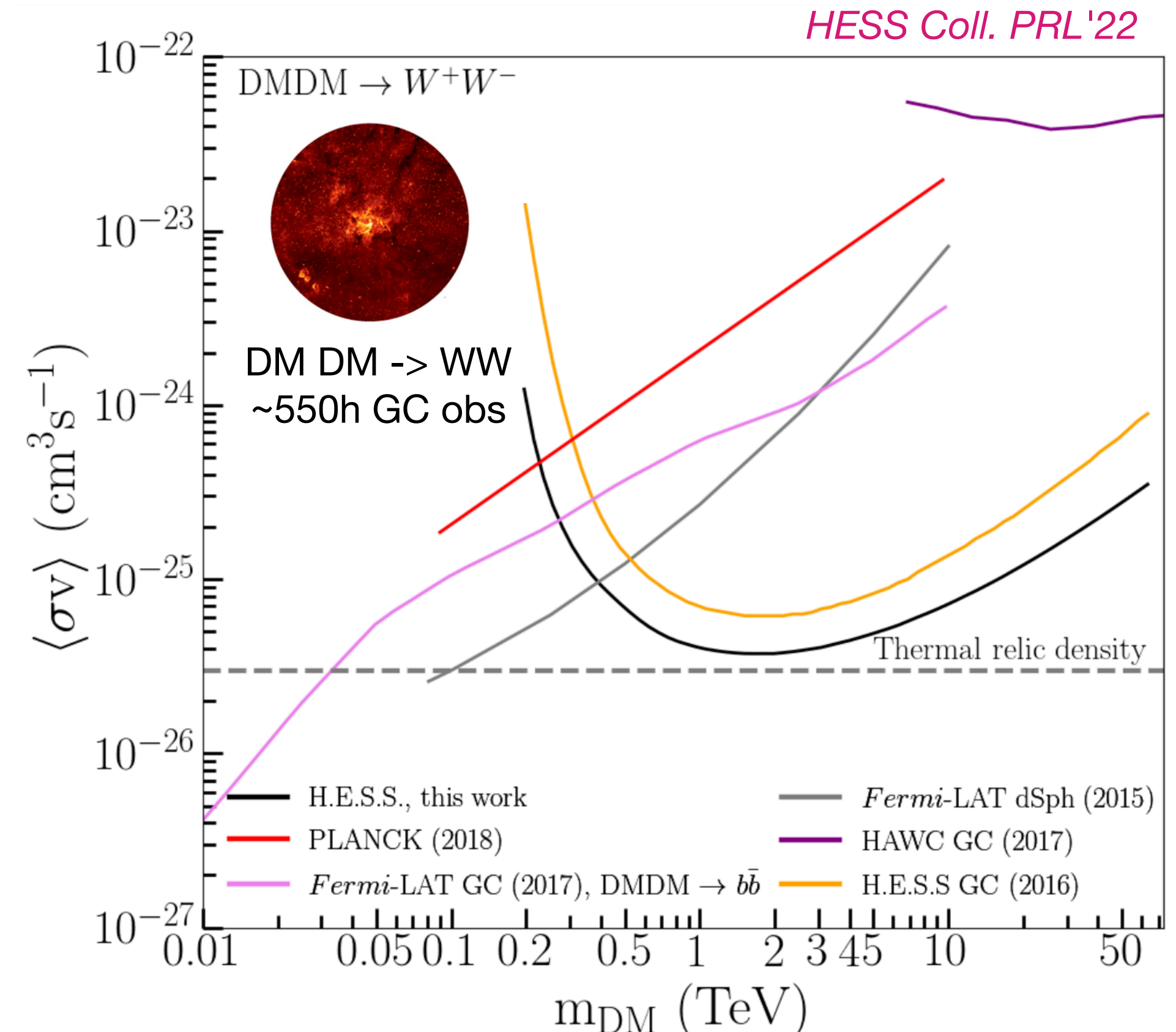
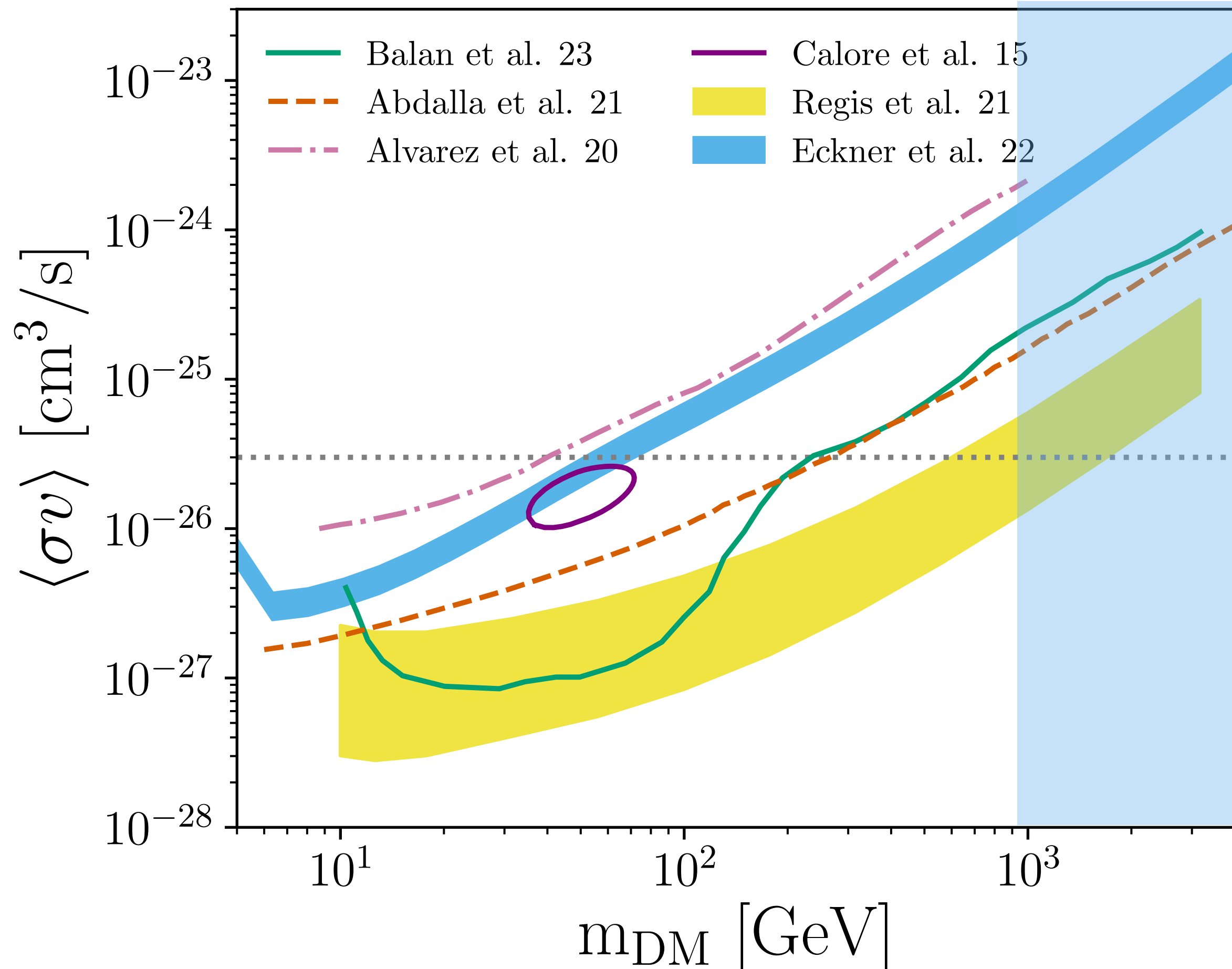
Limits on annihilating WIMPs

Summary of multi-targets and MW constraints

~ a few GeV — few TeV



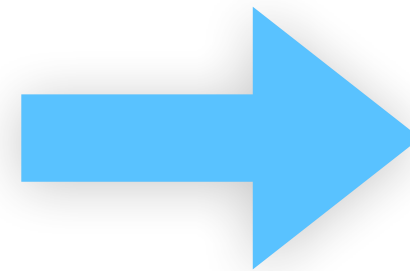
0.2 TeV — 50 TeV



Prospects and opportunities

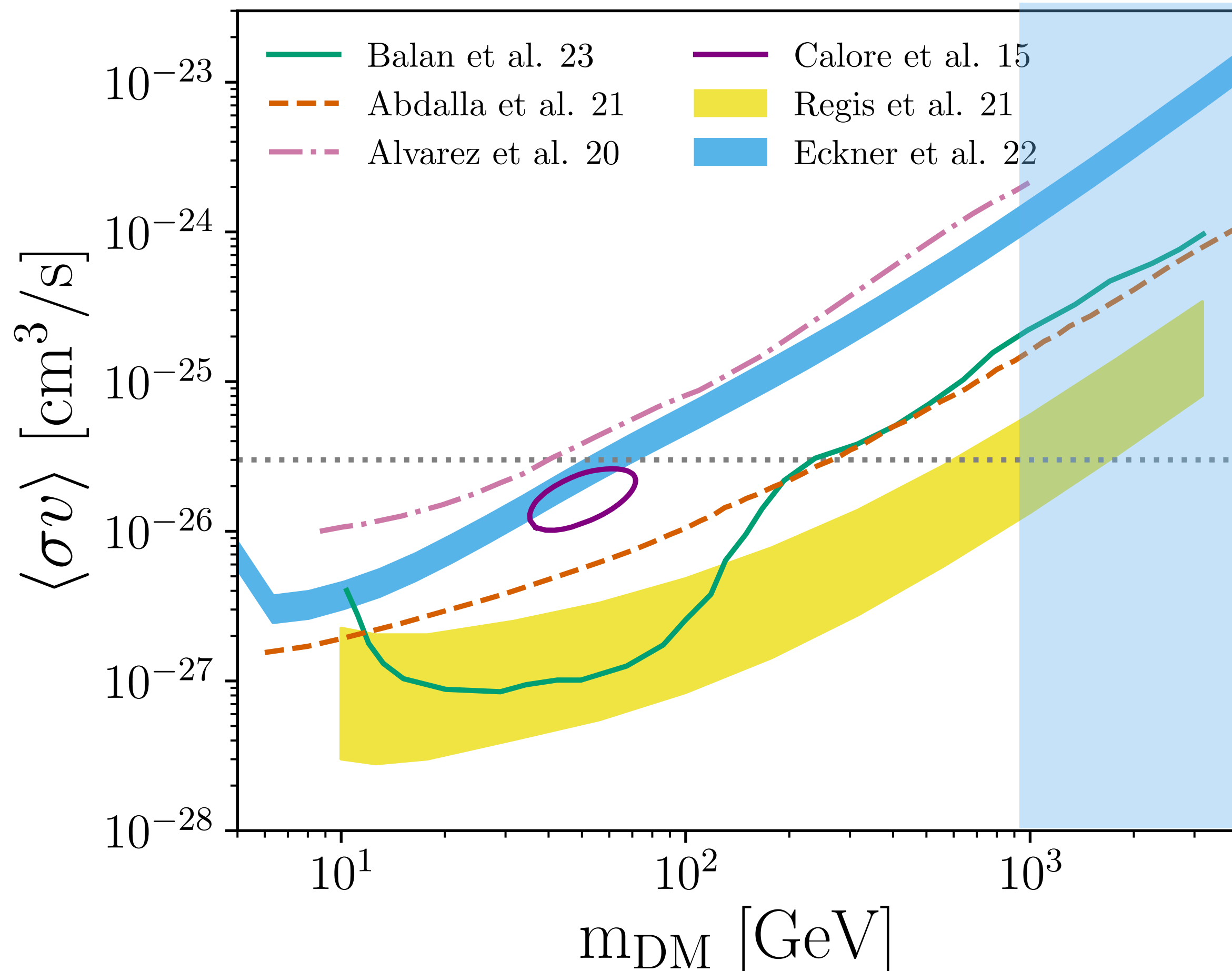
Extending the energy/mass scale

~ a few GeV — few TeV



Sub-PeV frontier

LHAASO, Tibet ASg



- Cannot be thermally produced (WIMpy) DM, since you hit the unitarity bound

Griest & Kamionkowski, PRD' 90

- Viable production mechanisms for PeV DM exist, e.g. inflation decay in low-scale reheating scenarios

Harigaya+ 1402.2846

- The signal should come through decay and should appear in neutrino fluxes even before gamma rays

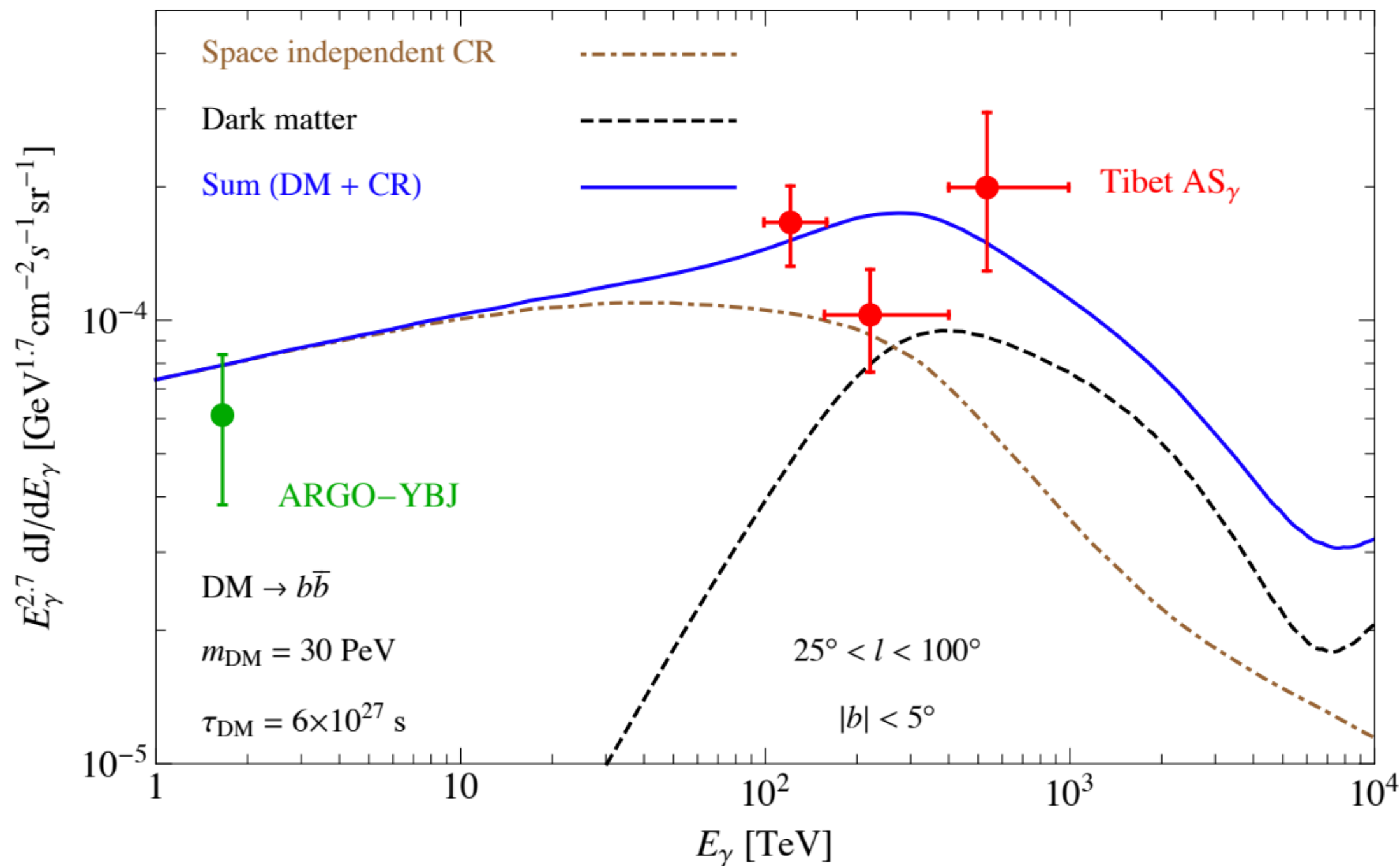
Feldstein+ PRD'13; Esmaili & Serpico, JCAP'13; Chianese+ arXiv:2108.01678

- These data often provide *best bounds* to heavy DM lifetime

Esmaili & Serpico, PRD'21; Chianese+ arXiv:2108.01678

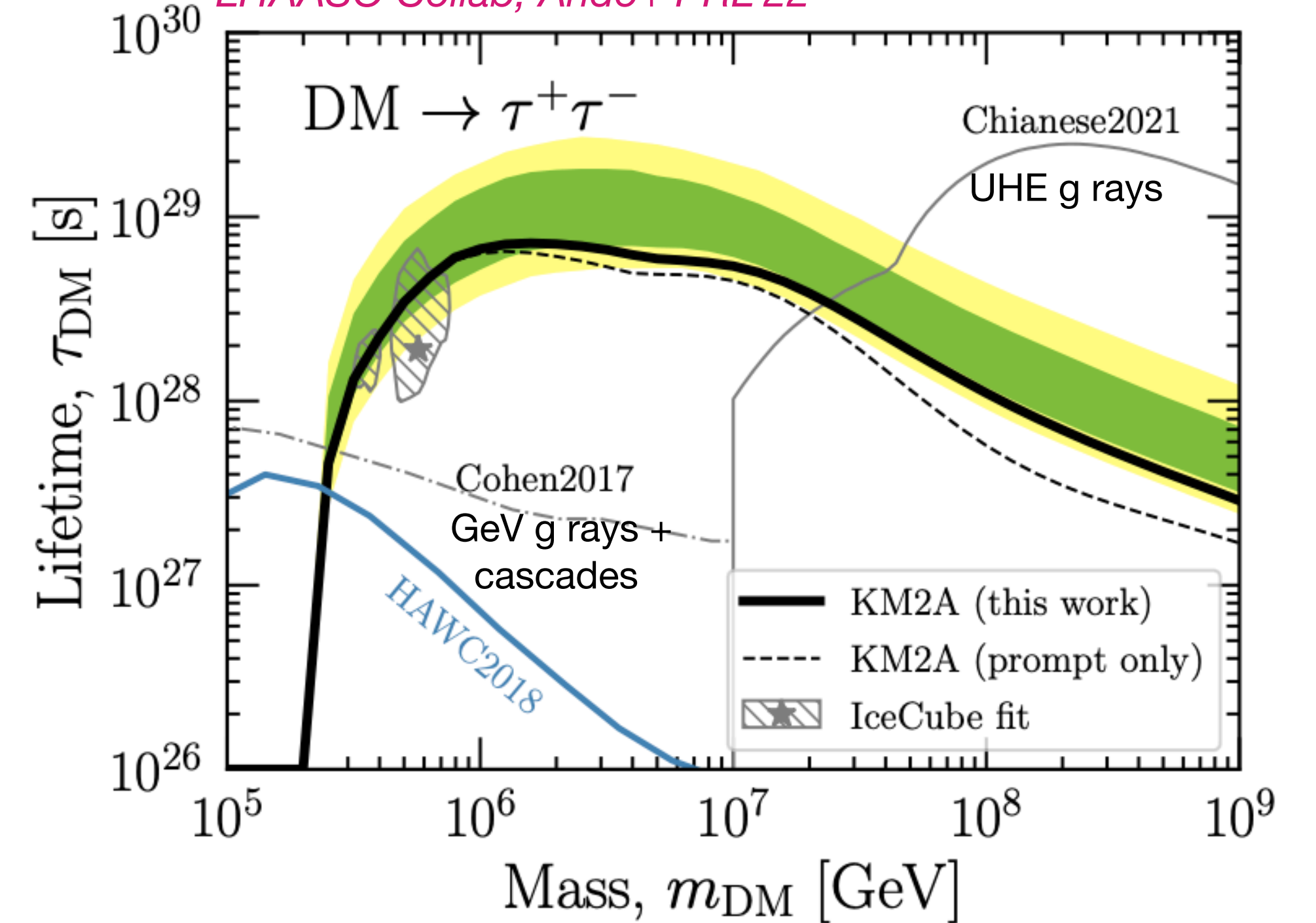
VHE gamma rays: PeV dark matter

Esmaili & Serpico, PRD Letters'21



Important degeneracies broken by **angular distribution of arrival photon directions**

LHAASO Collab, Ando+ PRL'22



Dedicated LHAASO DM search at slightly higher latitudes → Strongest constraints on PeV DM

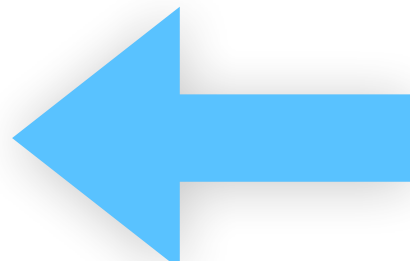
Challenges: 1. Production and propagation of UHE gamma rays; 2. DM spectra at production for $E_{CM} \gg E_{LHC}$

Bauer+ JHEP'21

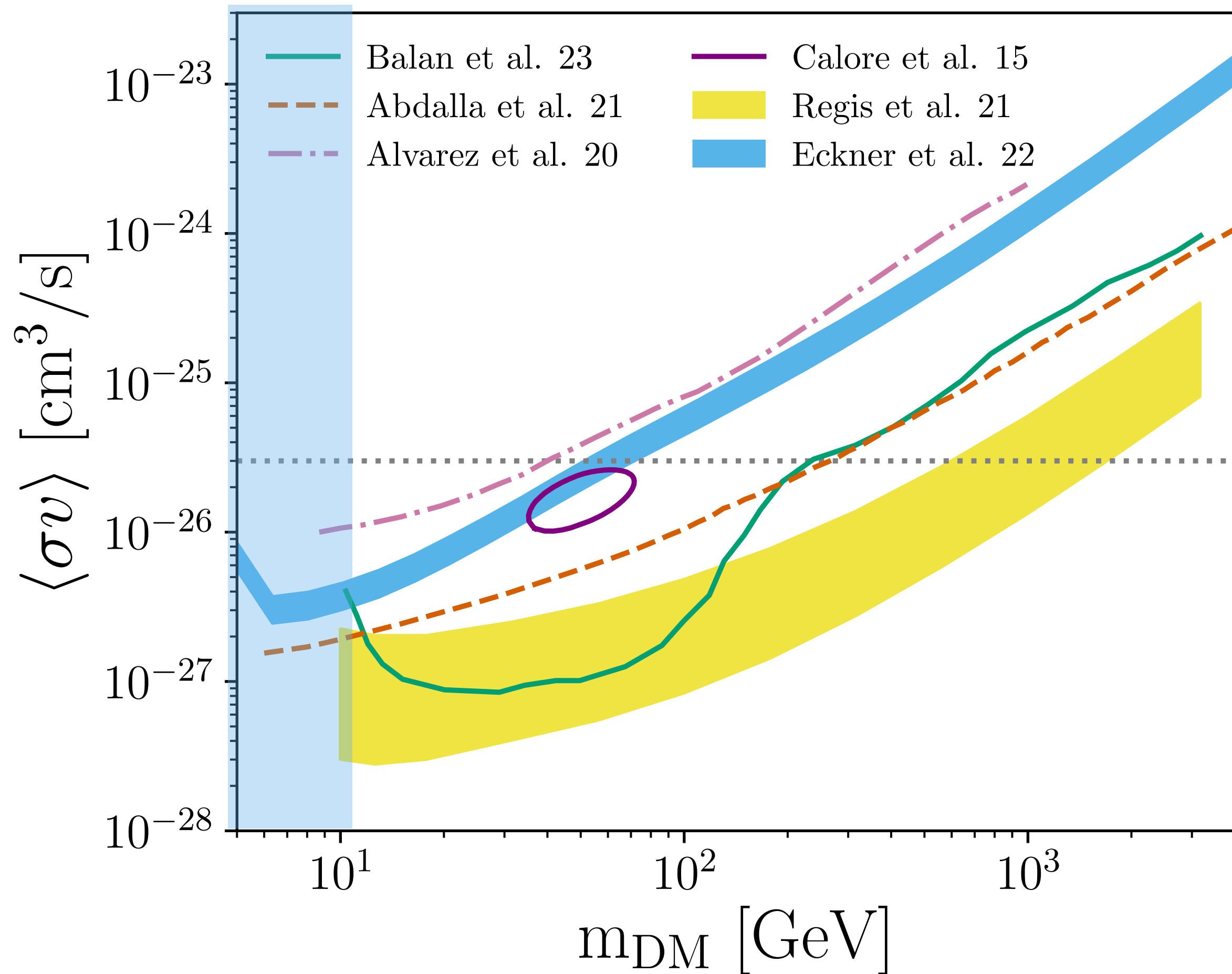
Application to **ALPs** in HE neutrino sources *Eckner, FC+ PRD'22; Mastrotoaro+ Eur.Phys.J.C'22*

Prospects and opportunities

Extending the energy/mass scale

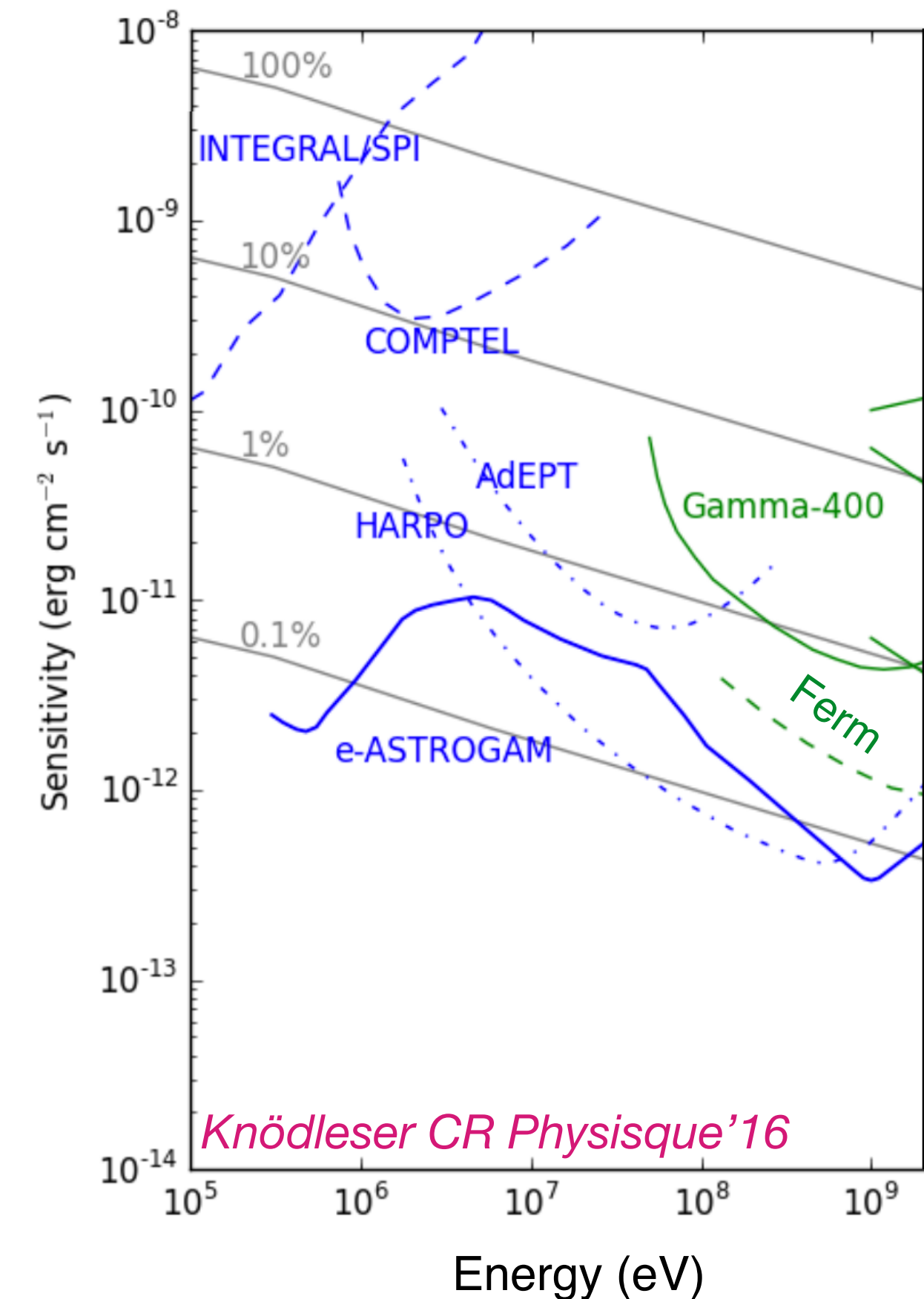


~ a few GeV — few TeV



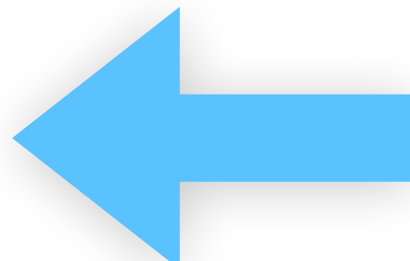
MeV sensitivity gap

Amego, e-ASTROGAM, GECCO, GRAMS, COSI, MeVCube, etc

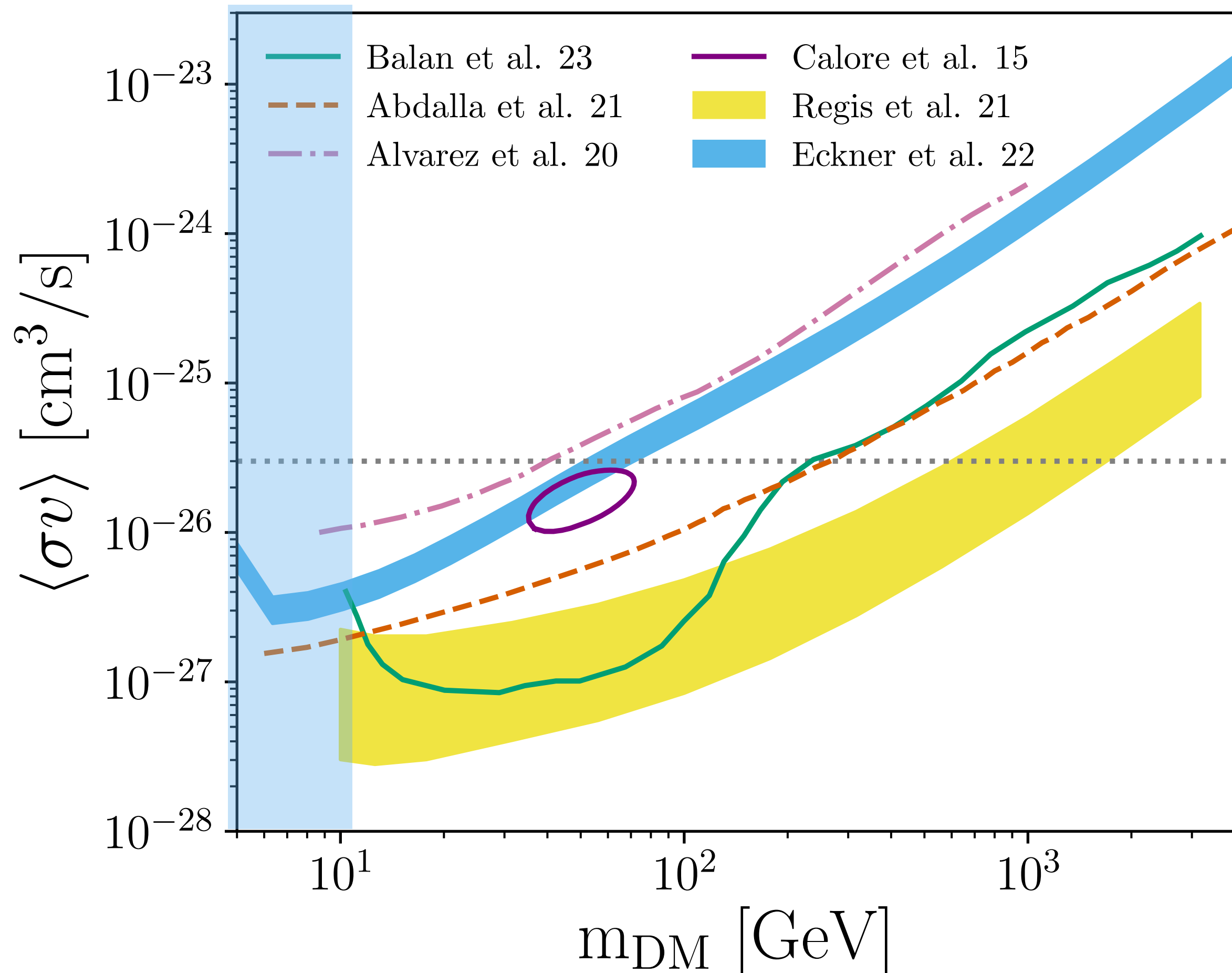


Prospects and opportunities

Extending the energy/mass scale

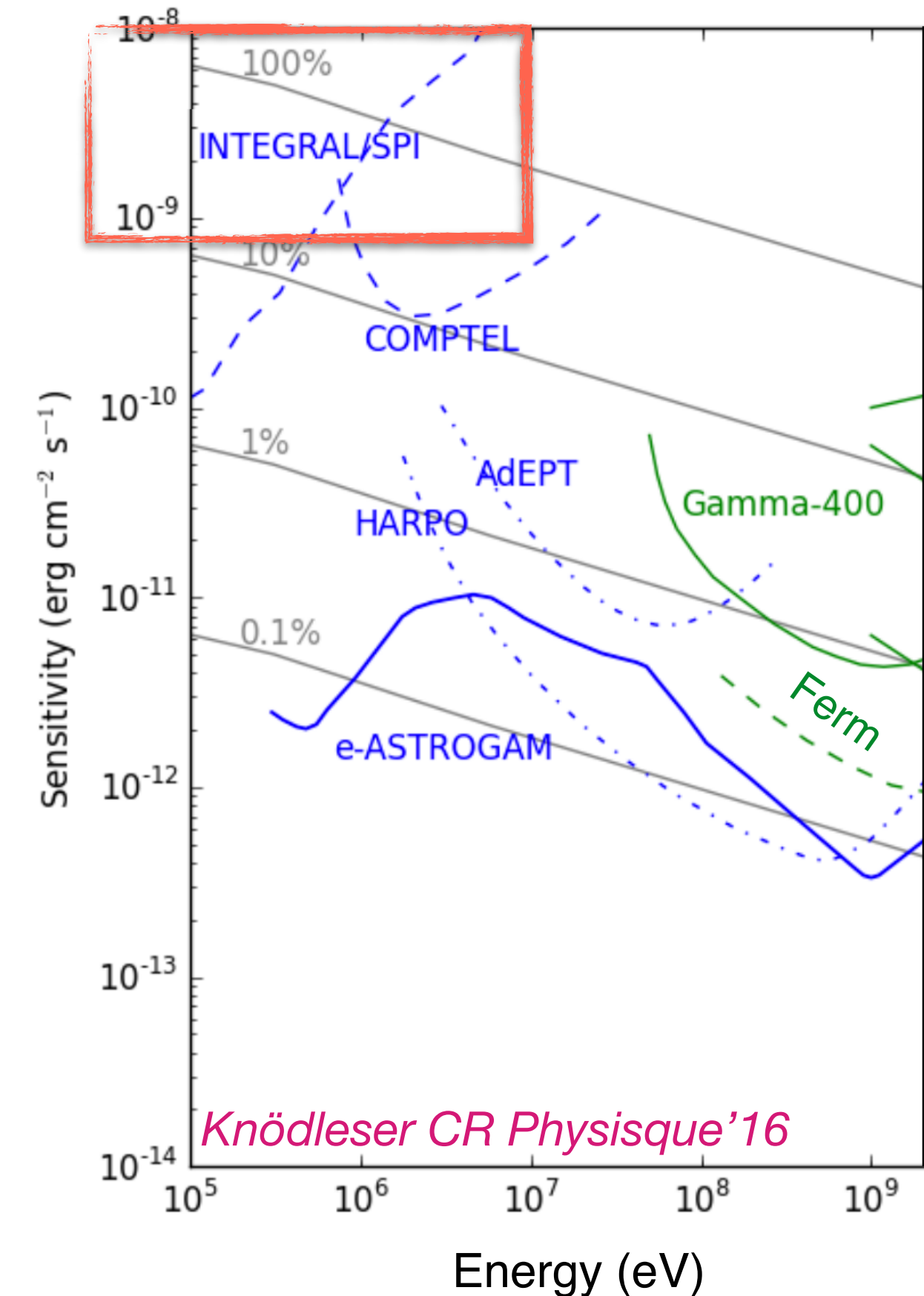


~ a few GeV — few TeV



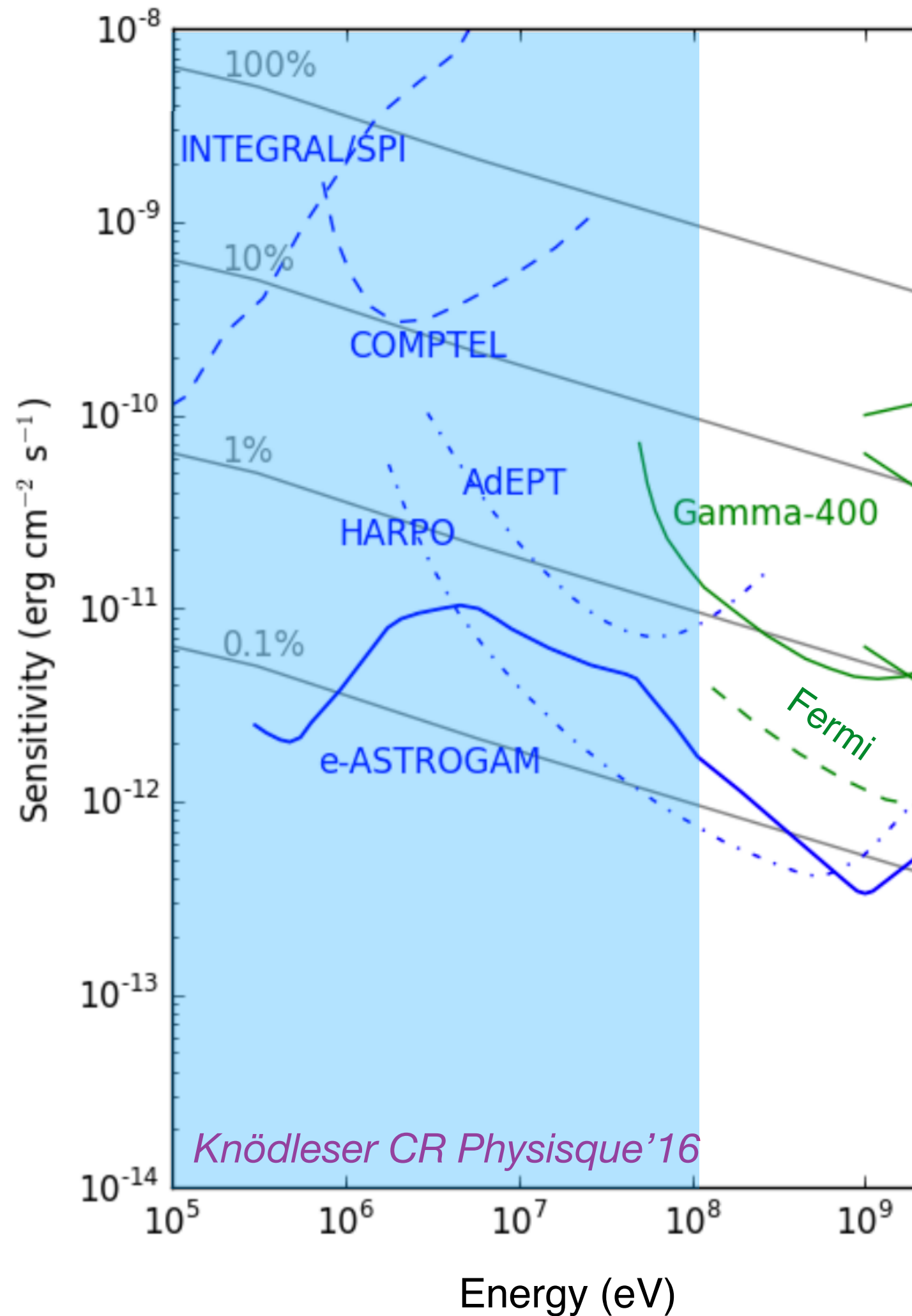
MeV sensitivity gap

Amego, e-ASTROGAM, GECCO, GRAMS, COSI, MeVCube, etc



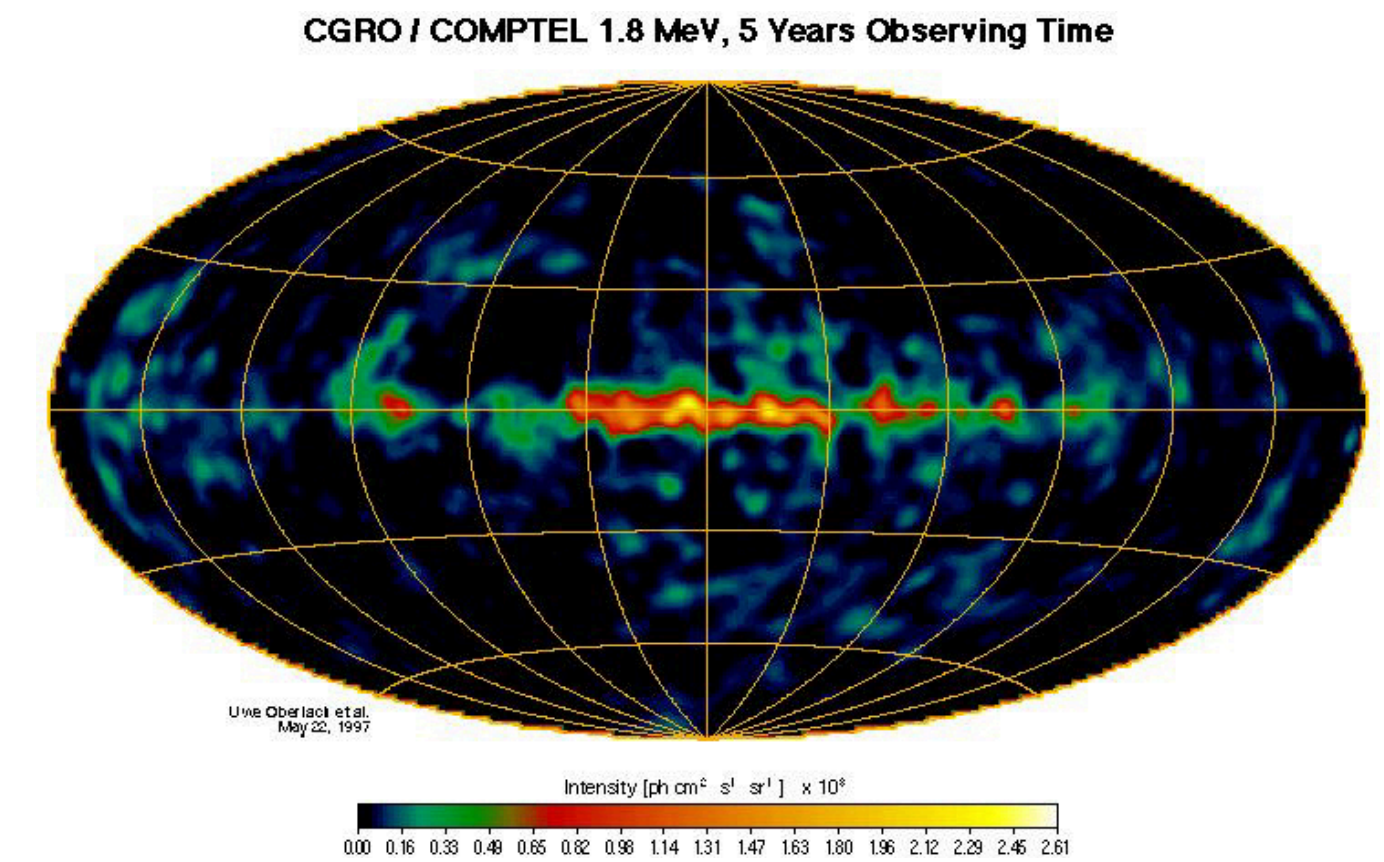
The data landscape

Hard X rays and gamma-ray sky



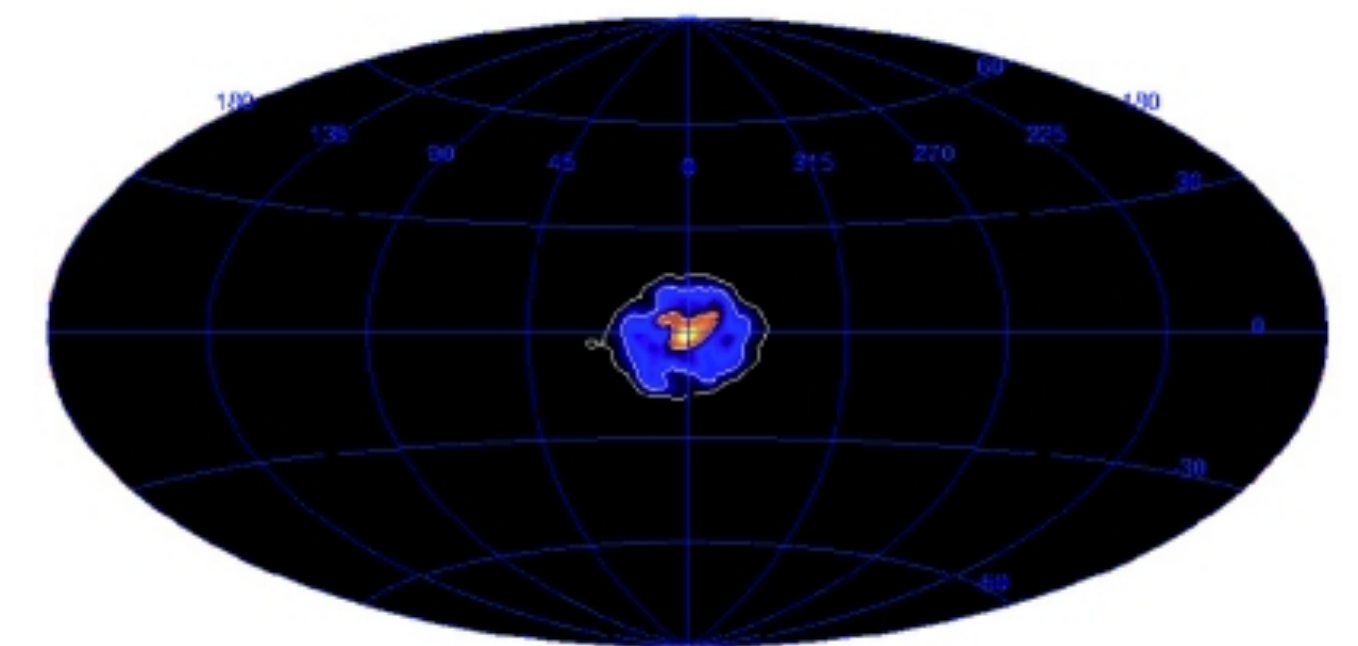
CGRO Compton Telescope (COMPTEL)

Lifetime: 1991 — 2000
Energy: 0.8 MeV - 30 MeV
Large FoV: 1 sr
Angular res: 1 deg



INTEGRAL Spectrometer (SPI)

Lifetime: 2002 — present
Energy: 20 keV - 8 MeV
Good energy res
Angular res: 2.5 deg

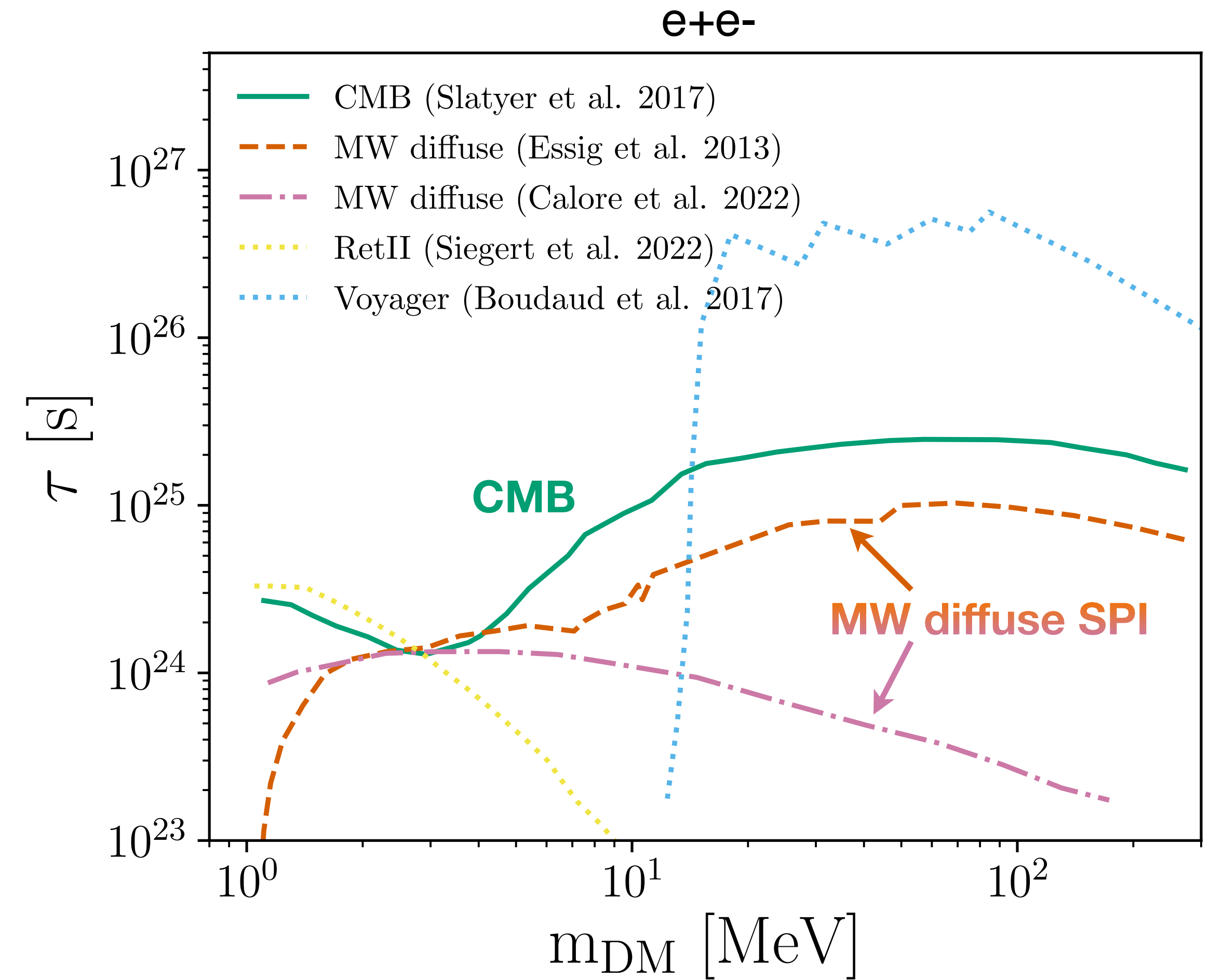
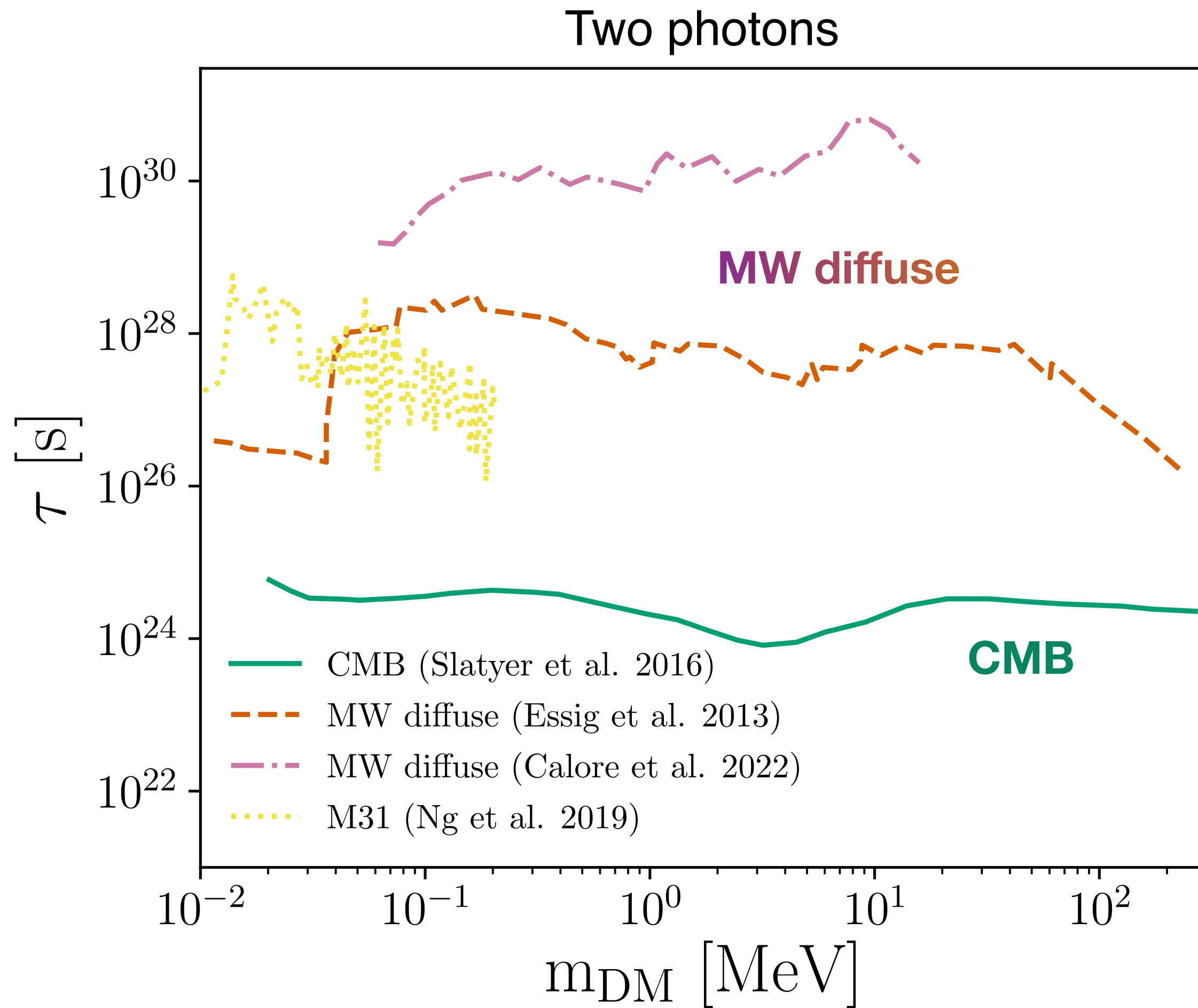


511 keV electron-positron line

Constraints on sub-GeV dark matter

Summary: decay

FC FIPs2022 Proceedings

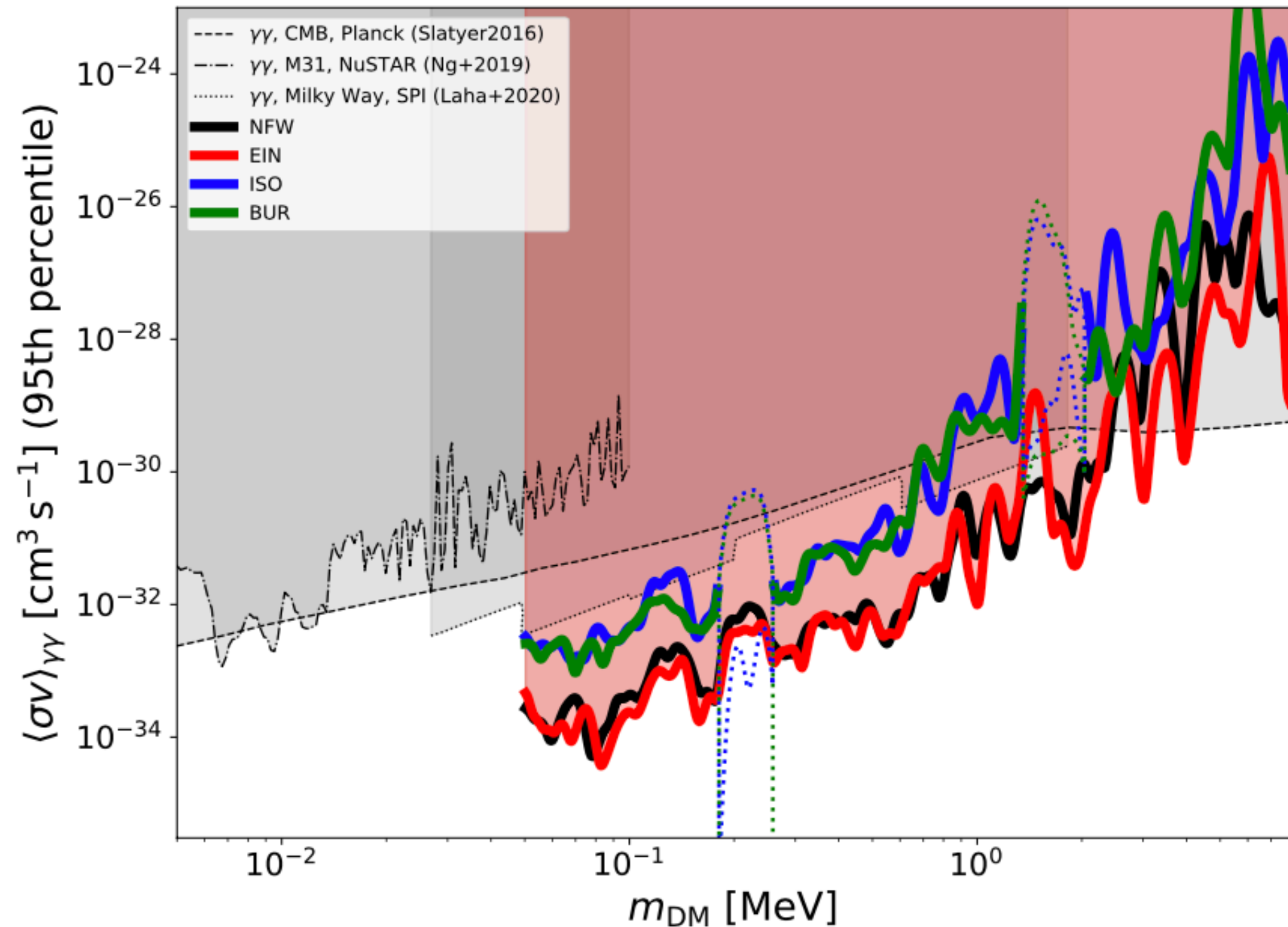


For e+e- channel, very relevant limits from Voyager 1 data *Boudaud+ PRL'17*

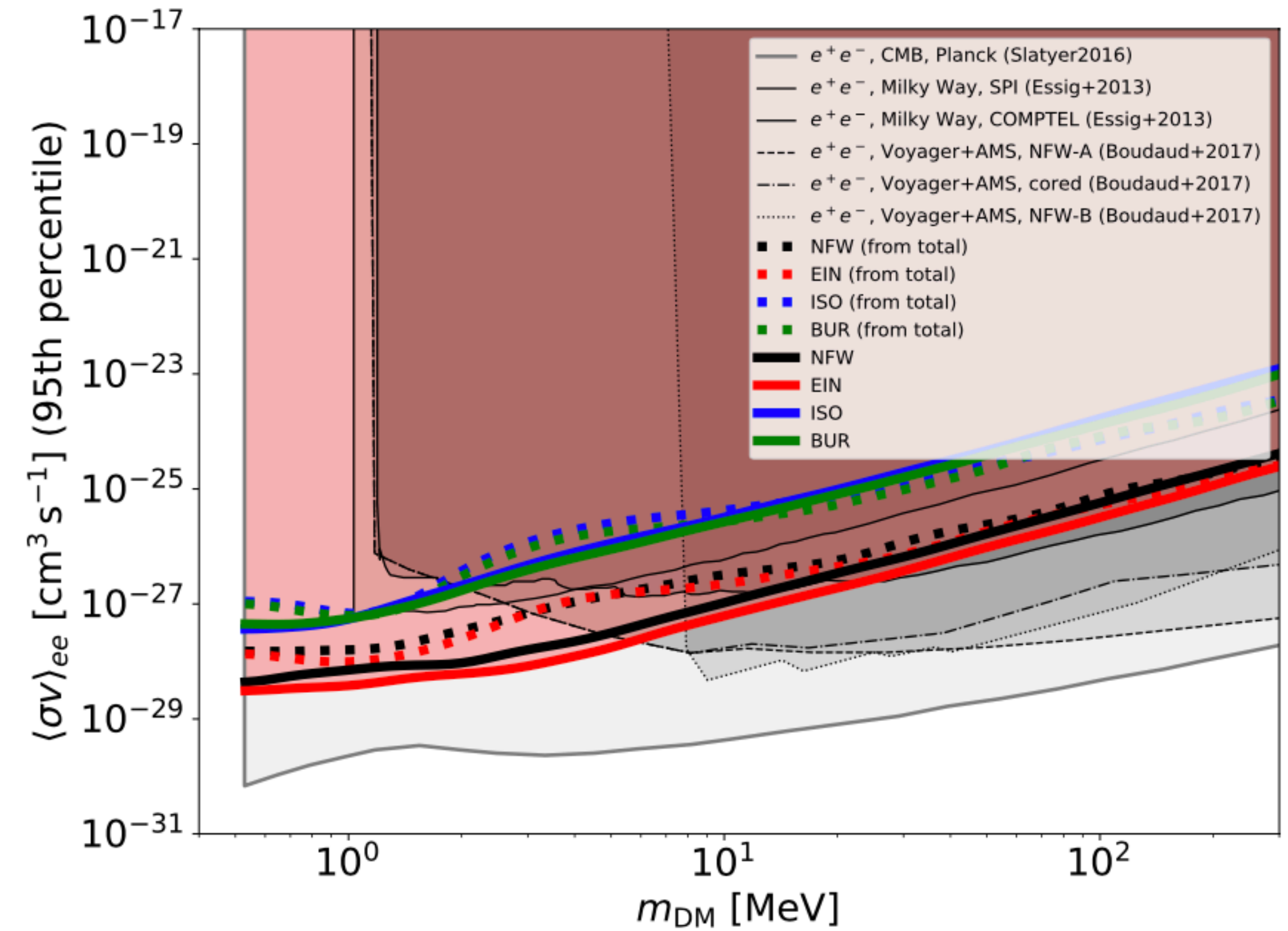
Constraints on annihilating dark matter

Summary from prompt emission

Two photons



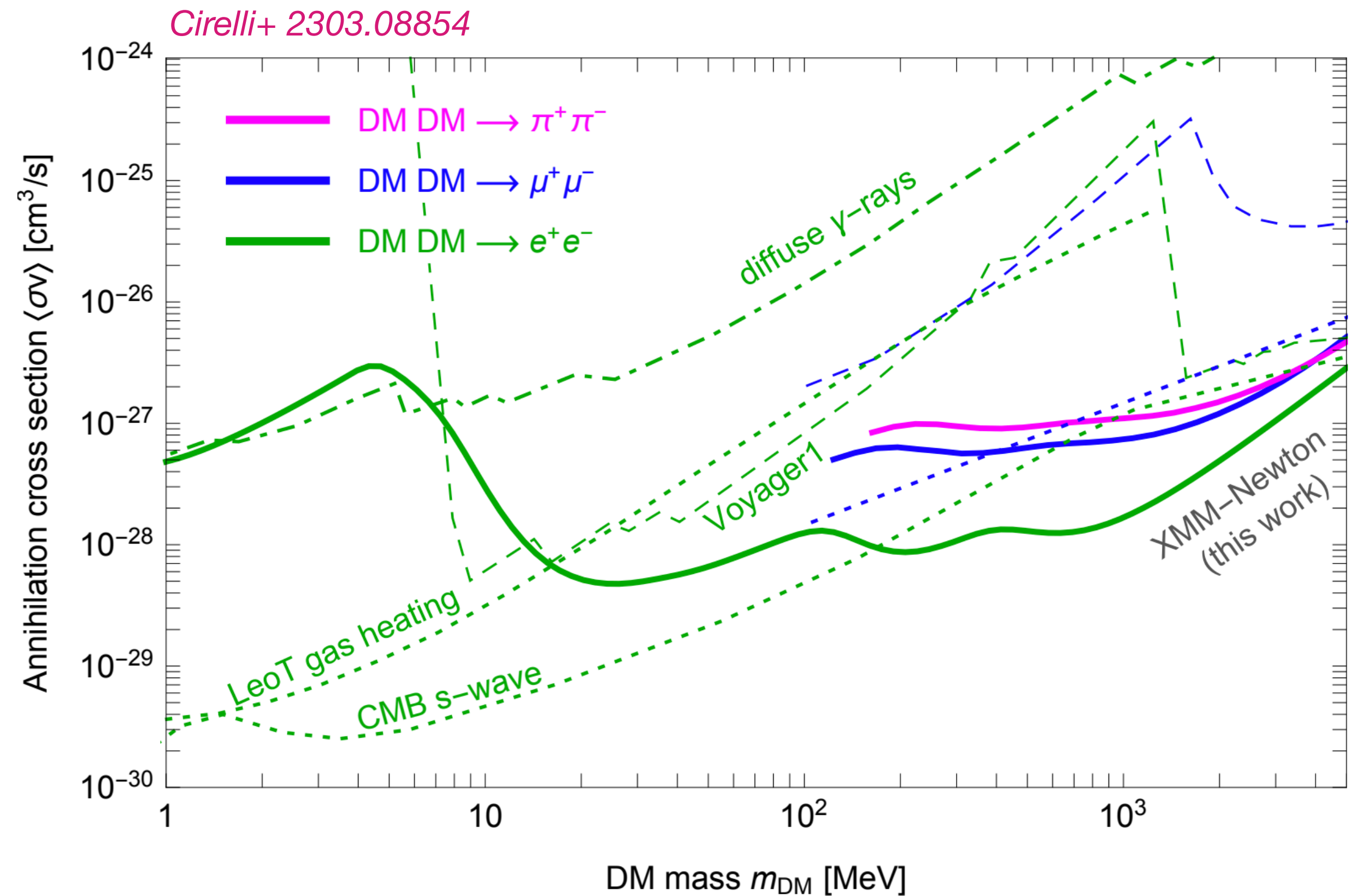
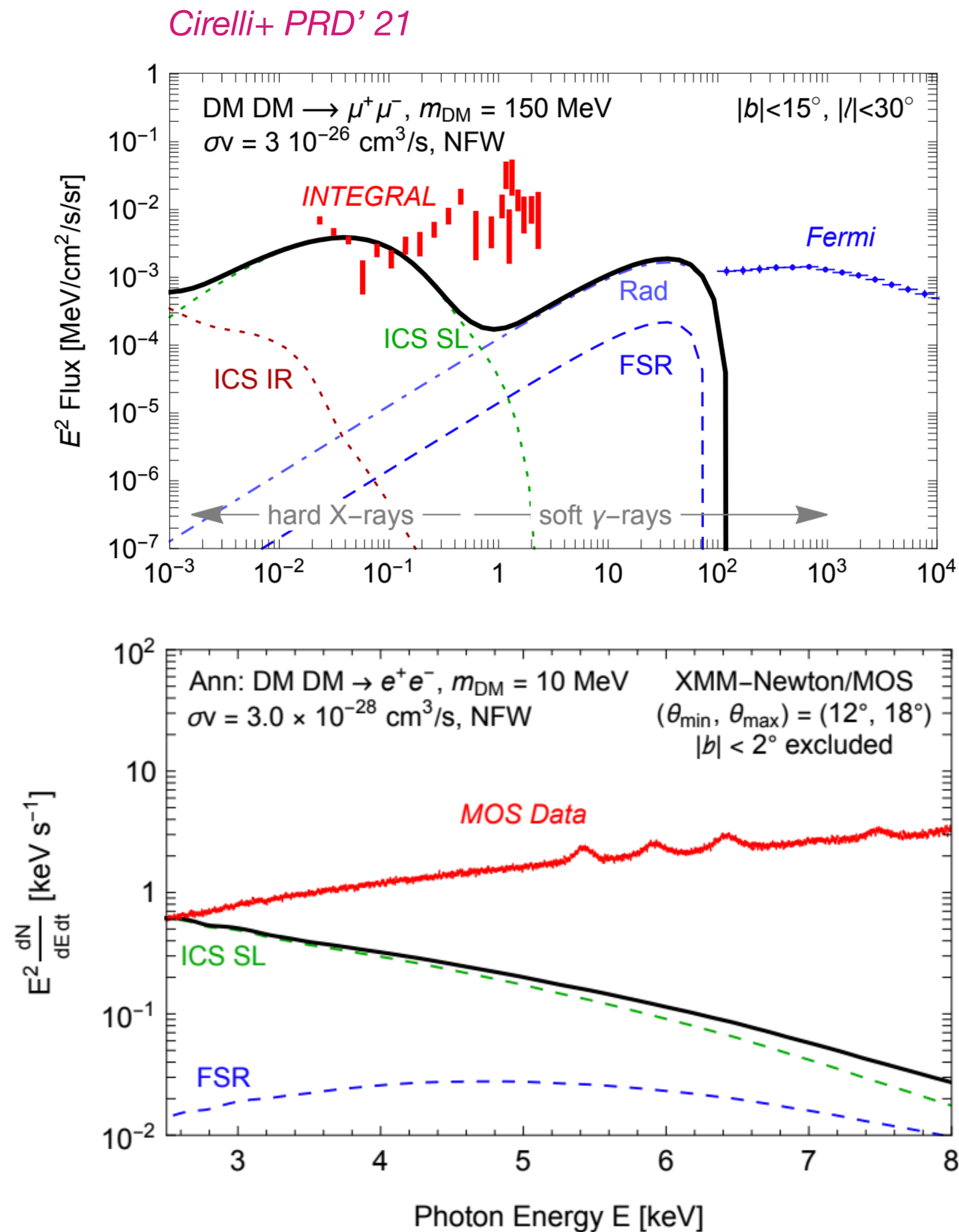
e+e-



For e^+e^- channel, very relevant limits from Voyager 1 data *Boudaud+ PRL'17*

Constraints on annihilating dark matter

Constraining higher dark matter masses with inverse Compton



Limits on IC induced gamma-ray emission from XMM-Newton blank-sky observations

Supplemental material

(Primordial black holes)

PBH: non-particle DM candidates

Black holes formed via the **collapse of large overdensities from inflation** in the early universe, before matter-radiation equality

Review & Refs in Carr+PRD'16, Green & Kavanagh J. Phys. G'21



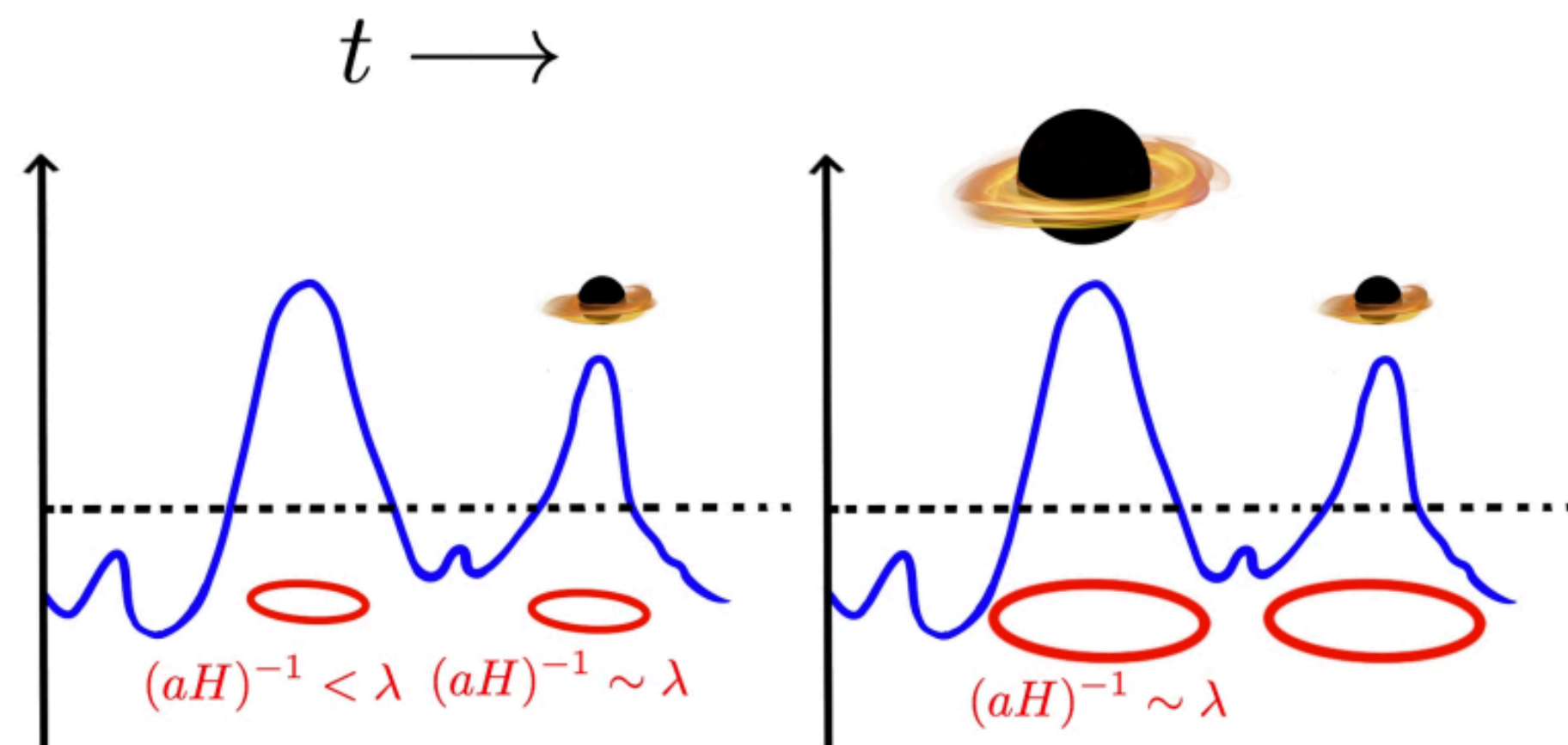
Quick ID:

- PBHs evaporate through Hawking radiation

Page & Hawking ApJ'76

$$\frac{dM_{\text{BH}}}{dt} \propto \frac{1}{(GM_{\text{BH}})^2} \longrightarrow \tau \propto (GM_{\text{BH}})^3$$

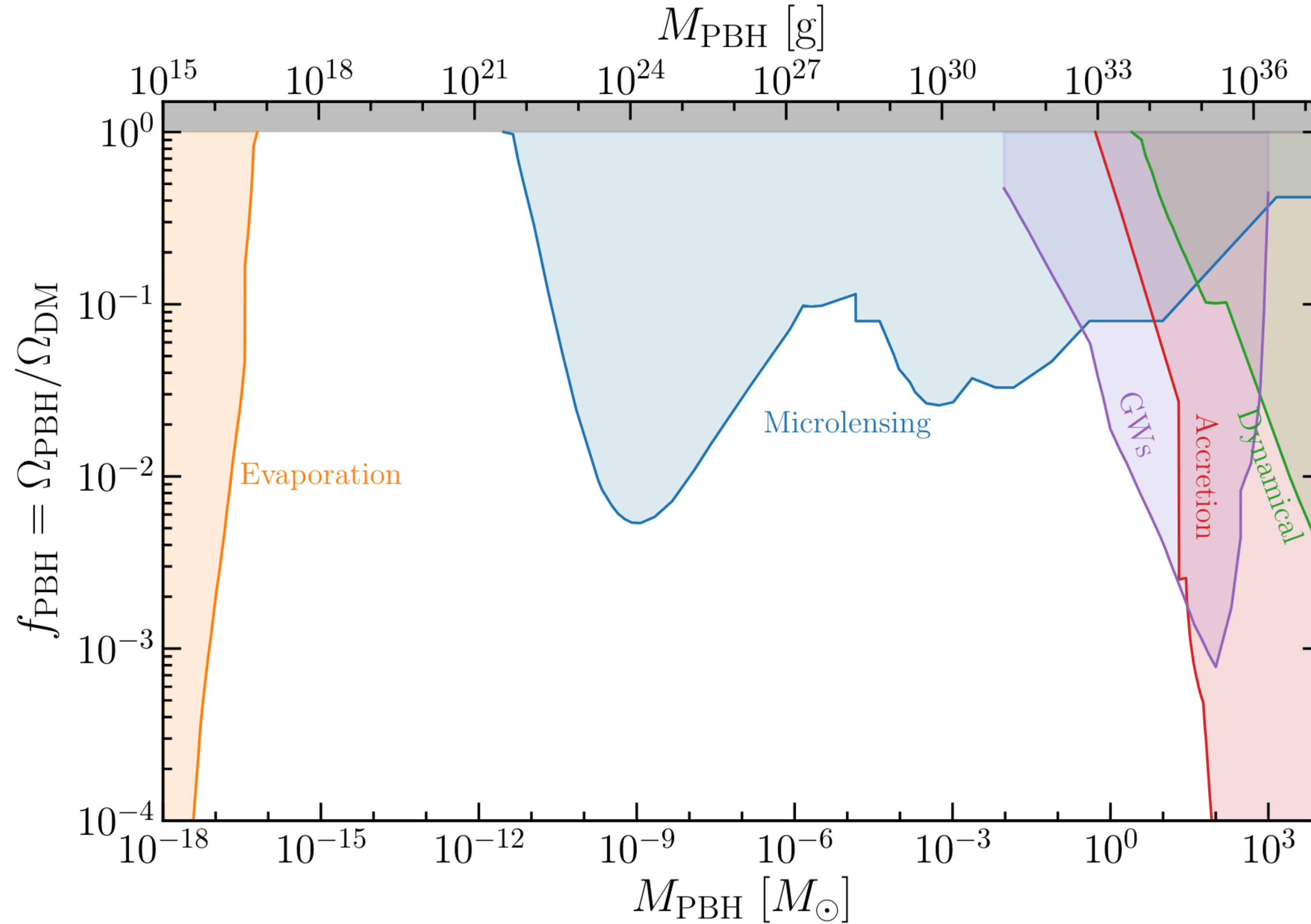
- If $M > 10^{14}$ g, PBHs have a lifetime longer than the age of the universe
- On cosmological scales PBH DM behaves like **cold** particle DM
- Granularity observable at galactic scales



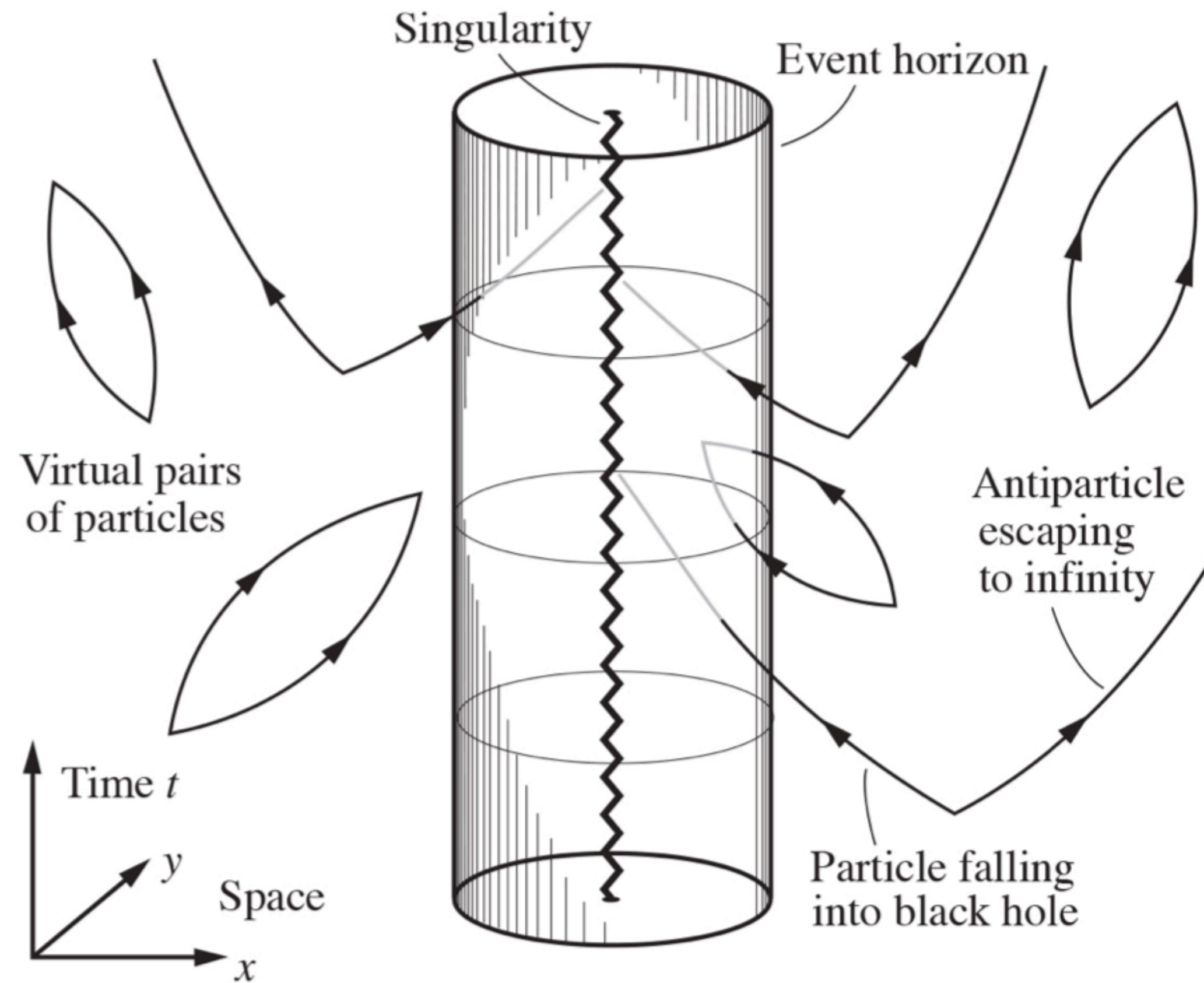
Villanueva-Domingo+ FrASS'21

[Intermediate between unitarity bound and Planck scale: ultra-heavy dark matter]

Limits on primordial black holes



Particle emission from primordial BHs



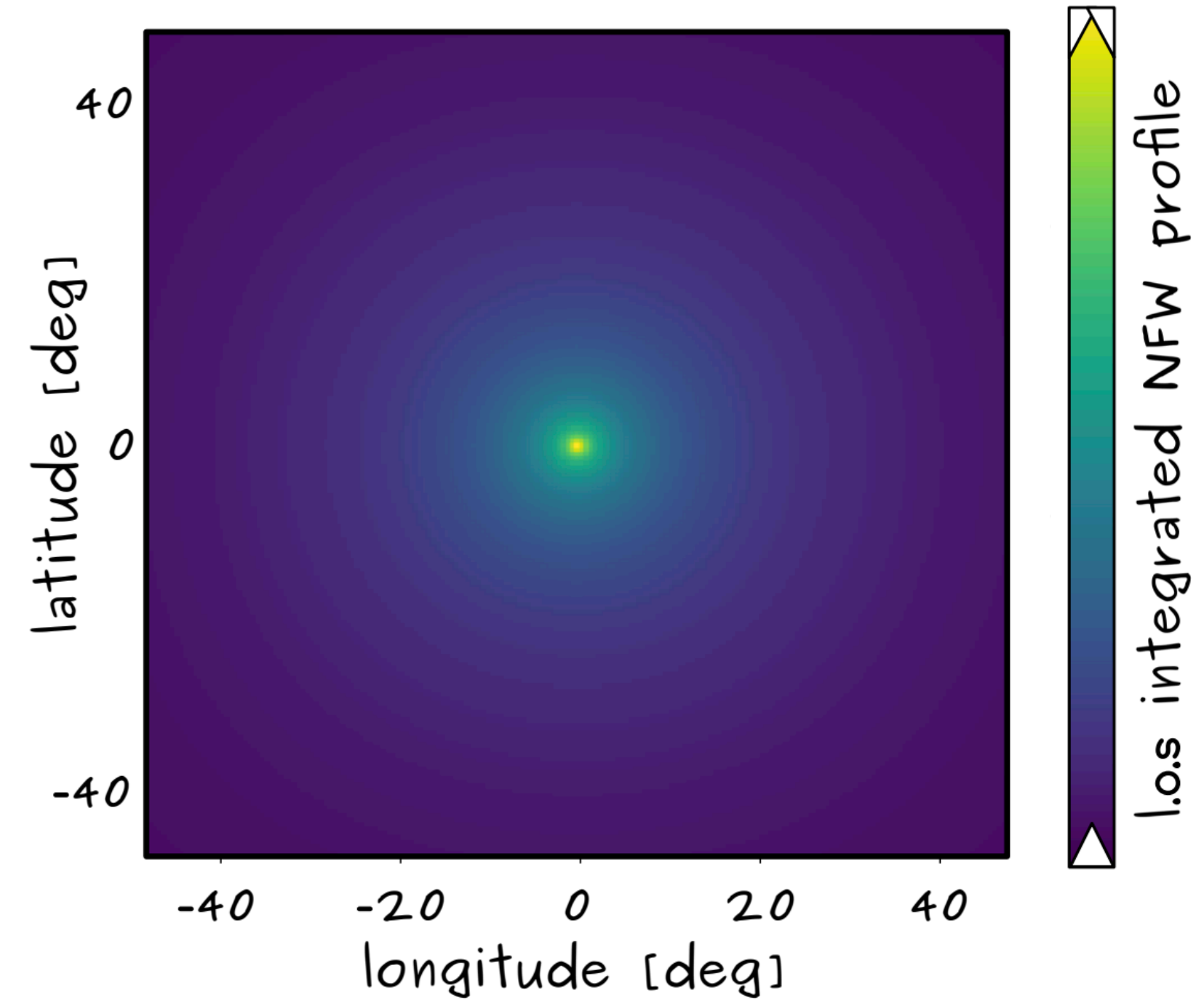
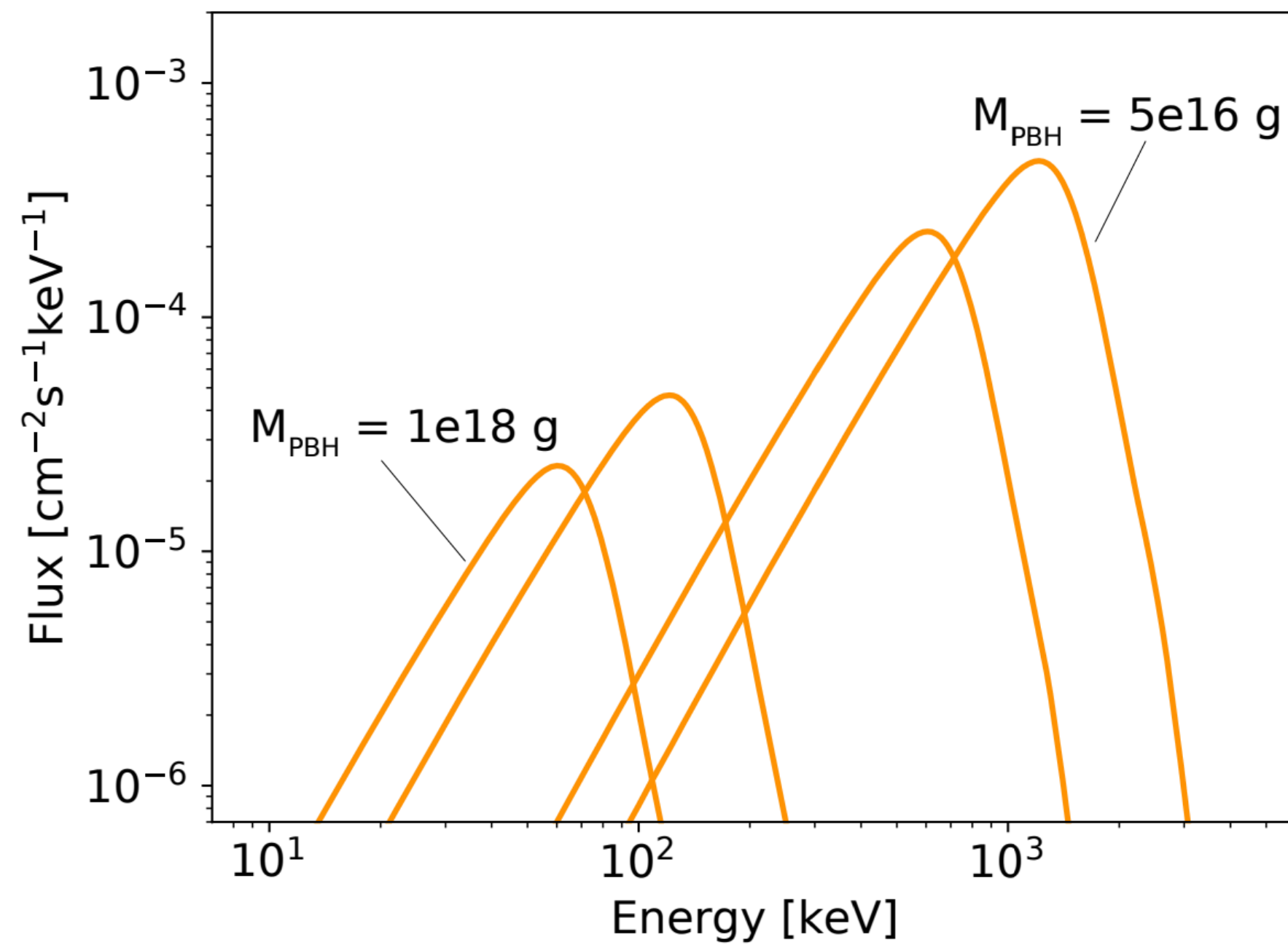
Hawking radiation

$$\frac{d^2 N}{dE dt} = \frac{1}{2\pi} \frac{\Gamma(E, M_{\text{BH}})}{\exp(E/T_{\text{BH}}) - (-1)^{2S}}$$

$$T_{\text{BH}} = \frac{1}{k_{\text{B}}} \frac{\hbar c^3}{8\pi G M_{\text{BH}}}$$

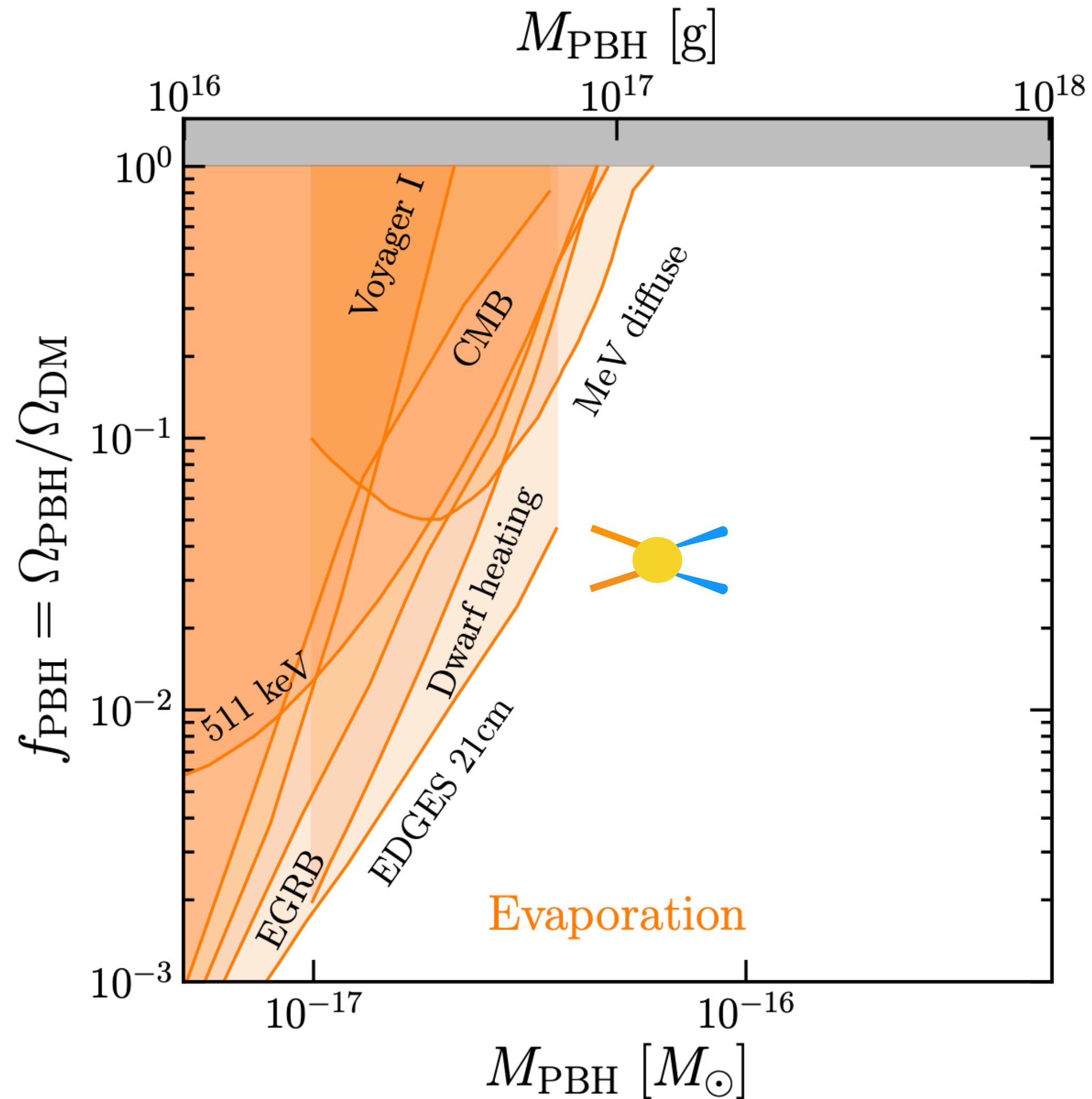
Photons from PBHs

$$\frac{d\Phi_\gamma}{dE}(l, b) = \frac{f_{\text{PBH}}}{4\pi M_{\text{PBH}}} \frac{d^2 N_\gamma}{dE dt} \int_{\text{l.o.s.}} ds \rho(r(s, l, b))$$



Limits on primordial black holes

Evaporation of PBH and cosmic backgrounds



Green & Kavanagh *J. Phys. G*'19

- Dominant limits from Galactic MeV diffuse emission from INTEGRAL/SPI

Berteaud, FC+ PRD'22

- Photon contribution to the extragalactic gamma-ray and X-ray backgrounds

Carr+ PRD'10; Ballesteros+ PLB'20; Iguaz+ PRD'21

- Unconstrained mass range $\sim 10^{17} - 10^{22}$ g, the so-called *asteroid mass gap* where f_{PBH} can be 1

Anomalies and excesses

Low-energy excess in antiproton data?

Signal:

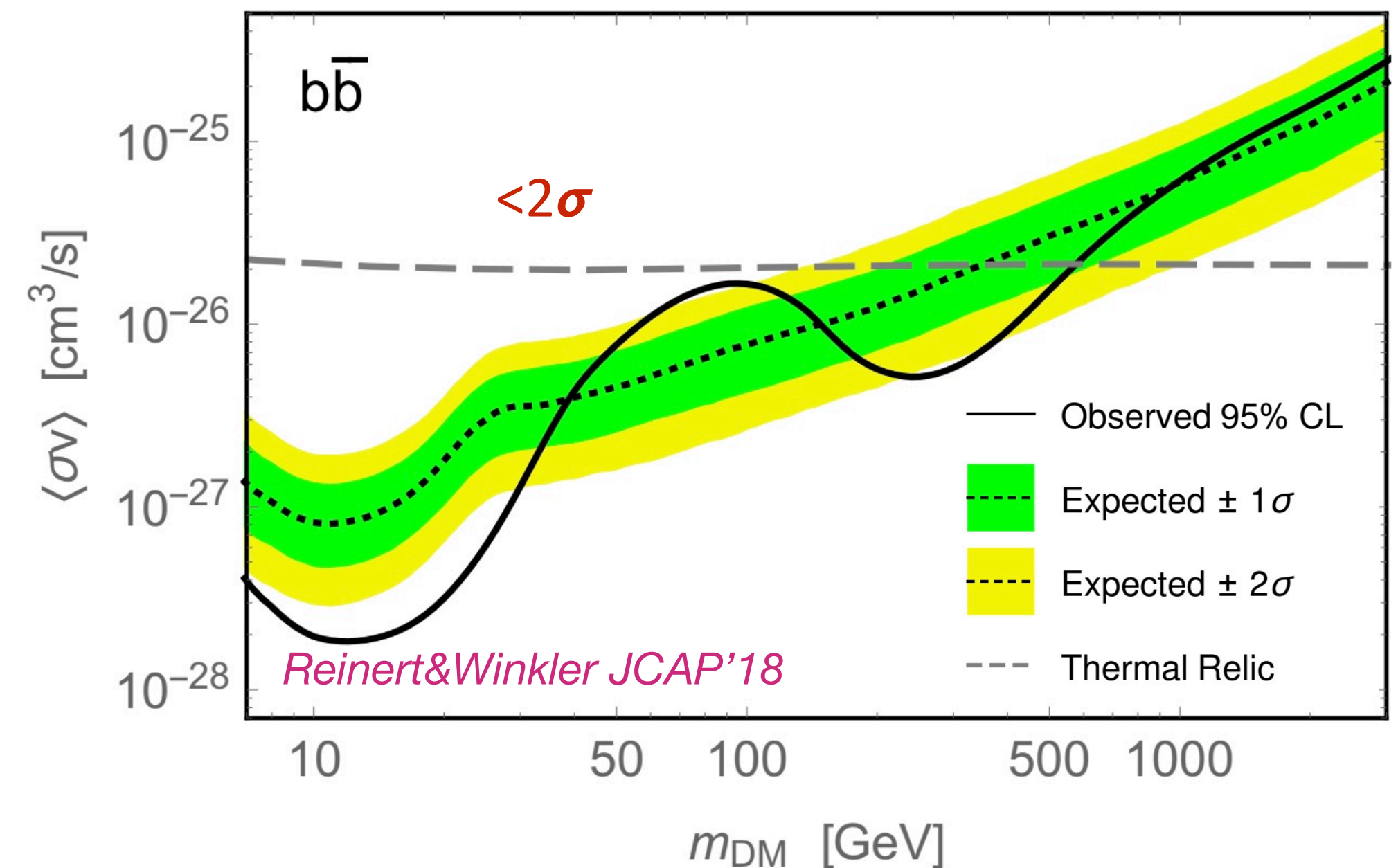
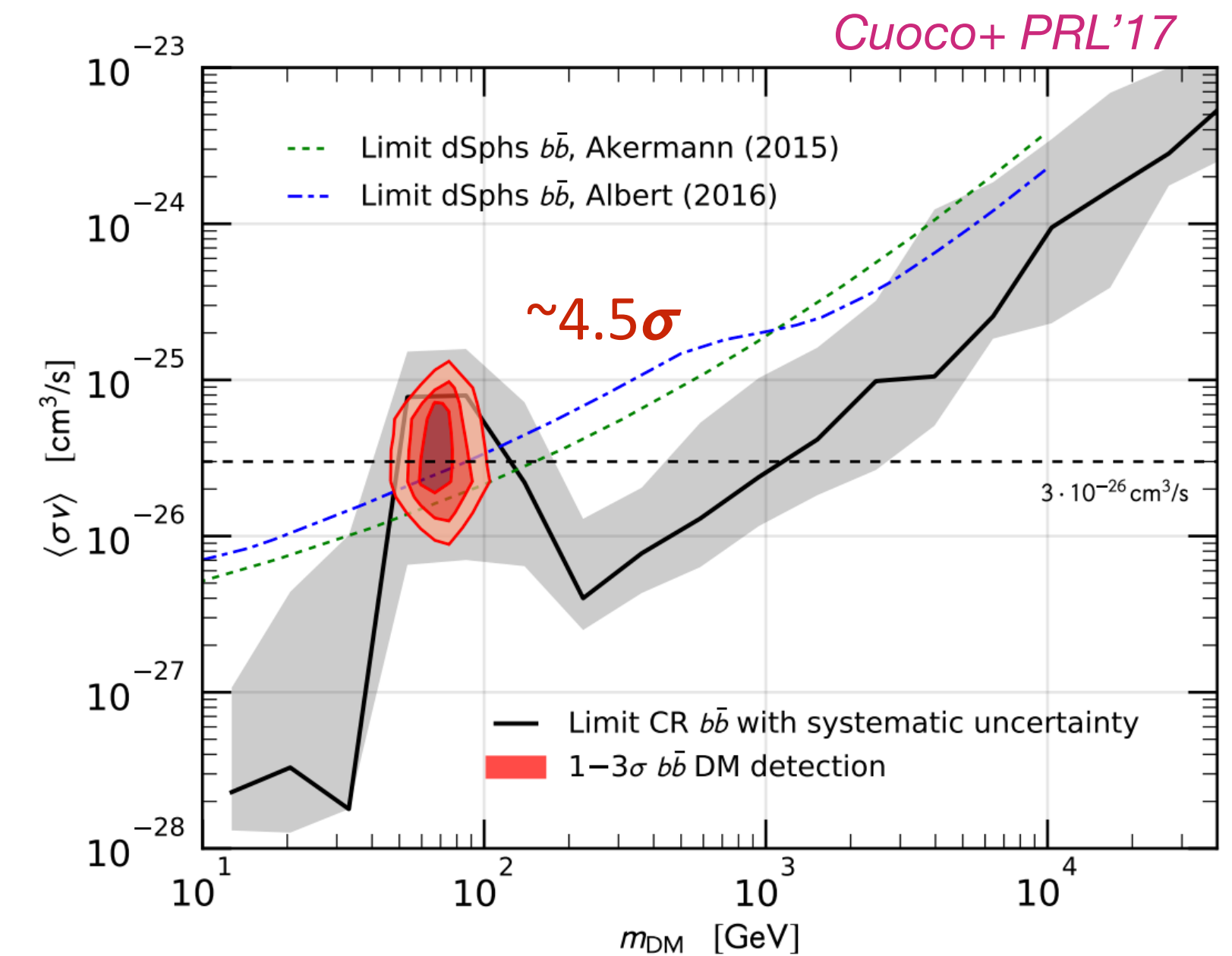
- Excess in AMS-02 cosmic-ray antiprotons @ 10 – 20 GeV
Cuoco+ PRL'17; Cui+ PRL'17; Cholis+ '17
- Accounting for covariance of various systematics the significance drops $< 2\sigma$

Reinert&Winkler JCAP'18; Boschini+ ApJ'17

Interpretations:

- Dark matter annihilation with mass $\sim 40 - 130$ GeV (consistent with GeV excess) *Cuoco+ JCAP'17*
- However, simple propagation scenarios cannot explain *all* CR data
- Syst. uncertainties still large: pbar production cross section? Effects of solar modulation? Cosmic-ray propagation models?

➔ Refined treatment of uncertainties is needed!



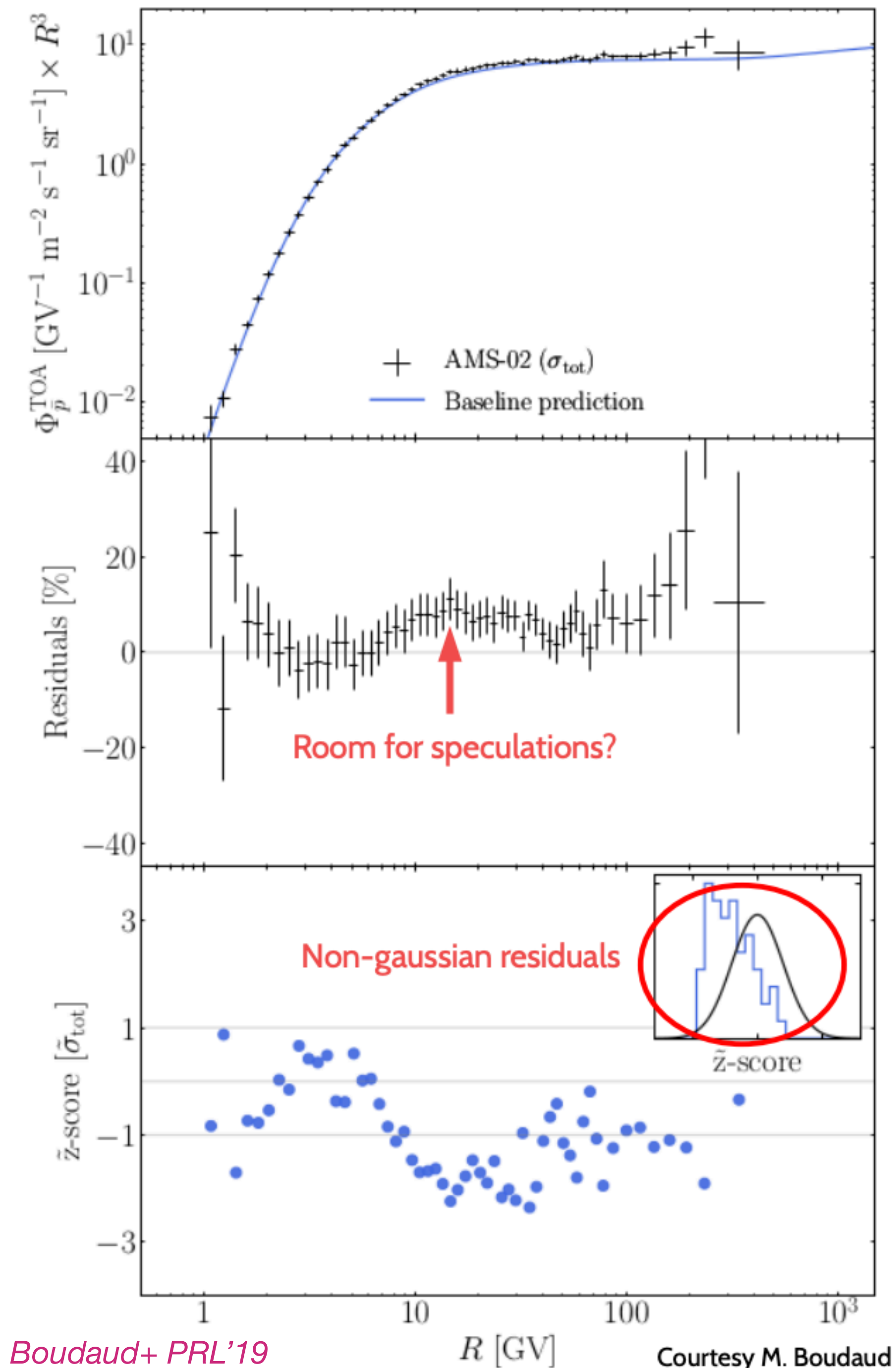
Antiproton uncertainties

Data: AMS02 antiproton from 2016

Model: semi-analytical

Comparison with data => discrepancy ~ few 10GV

**New Physics?
Or sys uncertainties?**



Antiproton uncertainties

Data: AMS02 antiproton from 2016

Model: semi-analytical

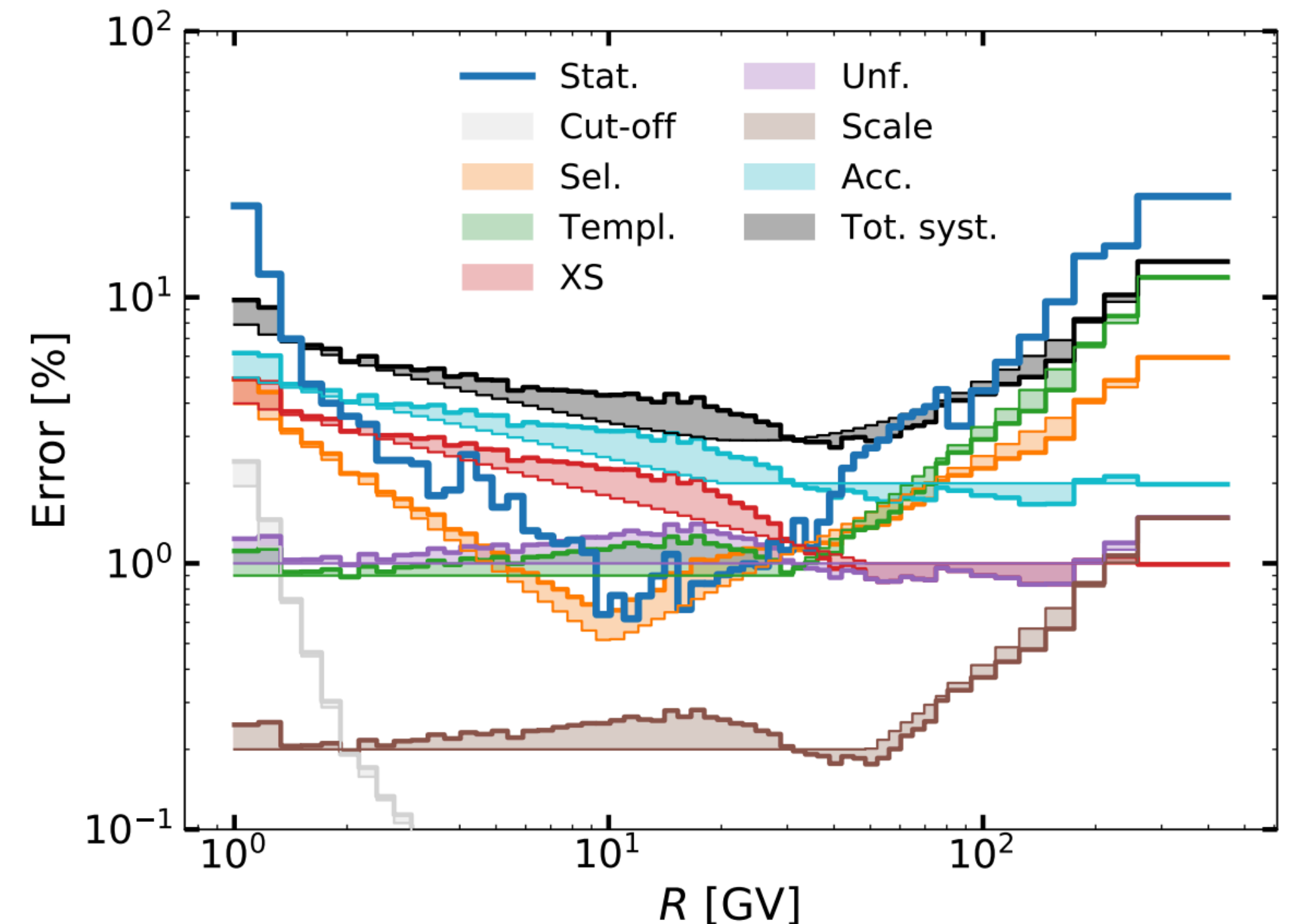
Comparison with data => discrepancy ~ few 10GV

New Physics? Or sys uncertainties?

Errors on the **data**: Covariance matrix estimated from detector info

Errors on the **model**:

1. Pbar production cross-sections → Updated parameterisation and uncertainties
2. Transport → Updated transport models and uncertainties
3. Parents → Updated fit and contribution of high- γ elements



Antiproton uncertainties

Data: AMS02 antiproton from 2016

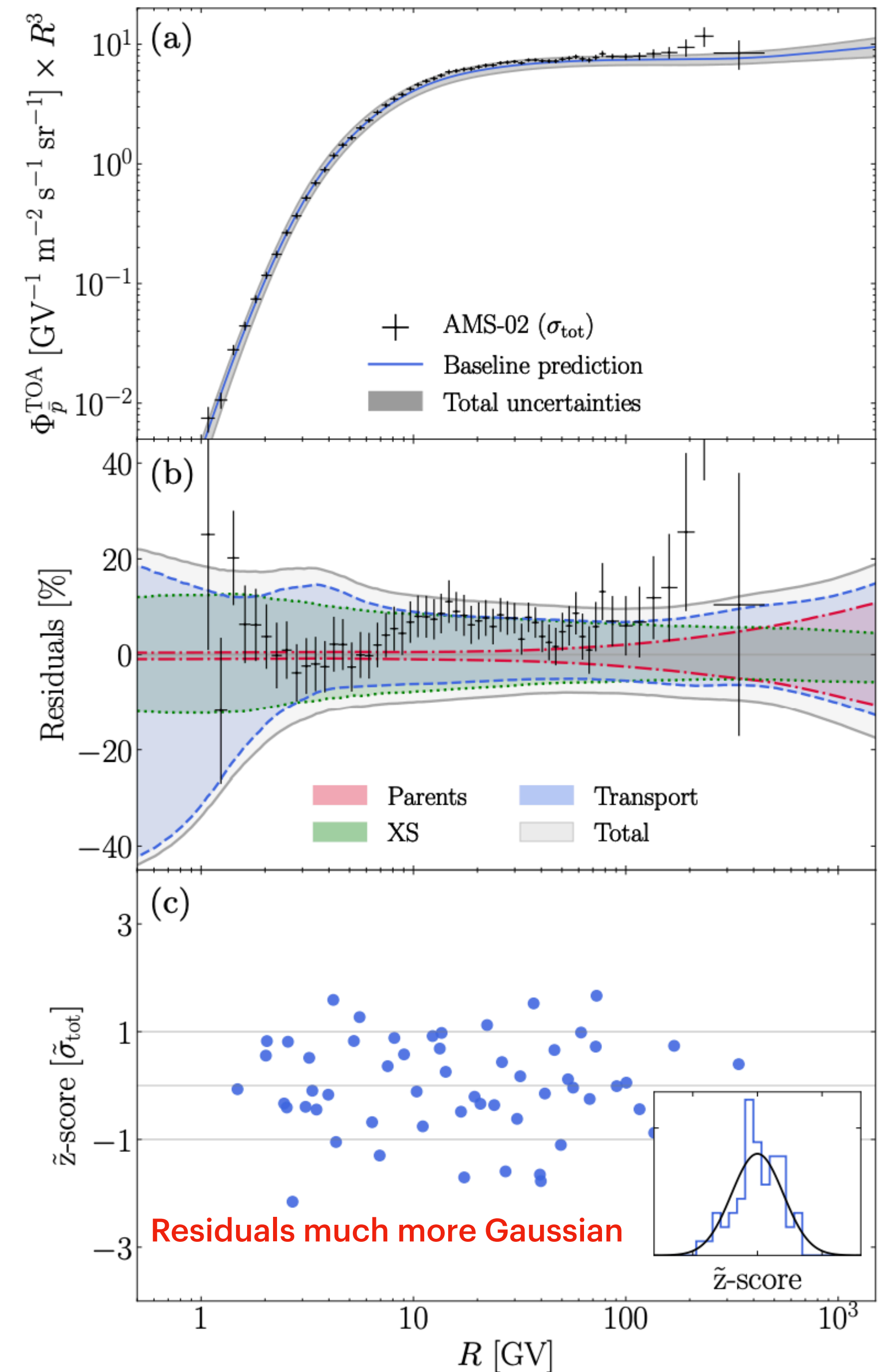
Model: semi-analytical

Comparison with data => discrepancy ~ few 10GV

**New Physics?
Or sys uncertainties?**

AMS-02 antiprotons are consistent with a secondary astrophysical origin

$$\chi^2 = (\text{data} - \text{model})^T (\mathcal{C}^{\text{model}} + \mathcal{C}^{\text{data}})^{-1} (\text{data} - \text{model})$$



No room for dark matter

Likelihood ratio:

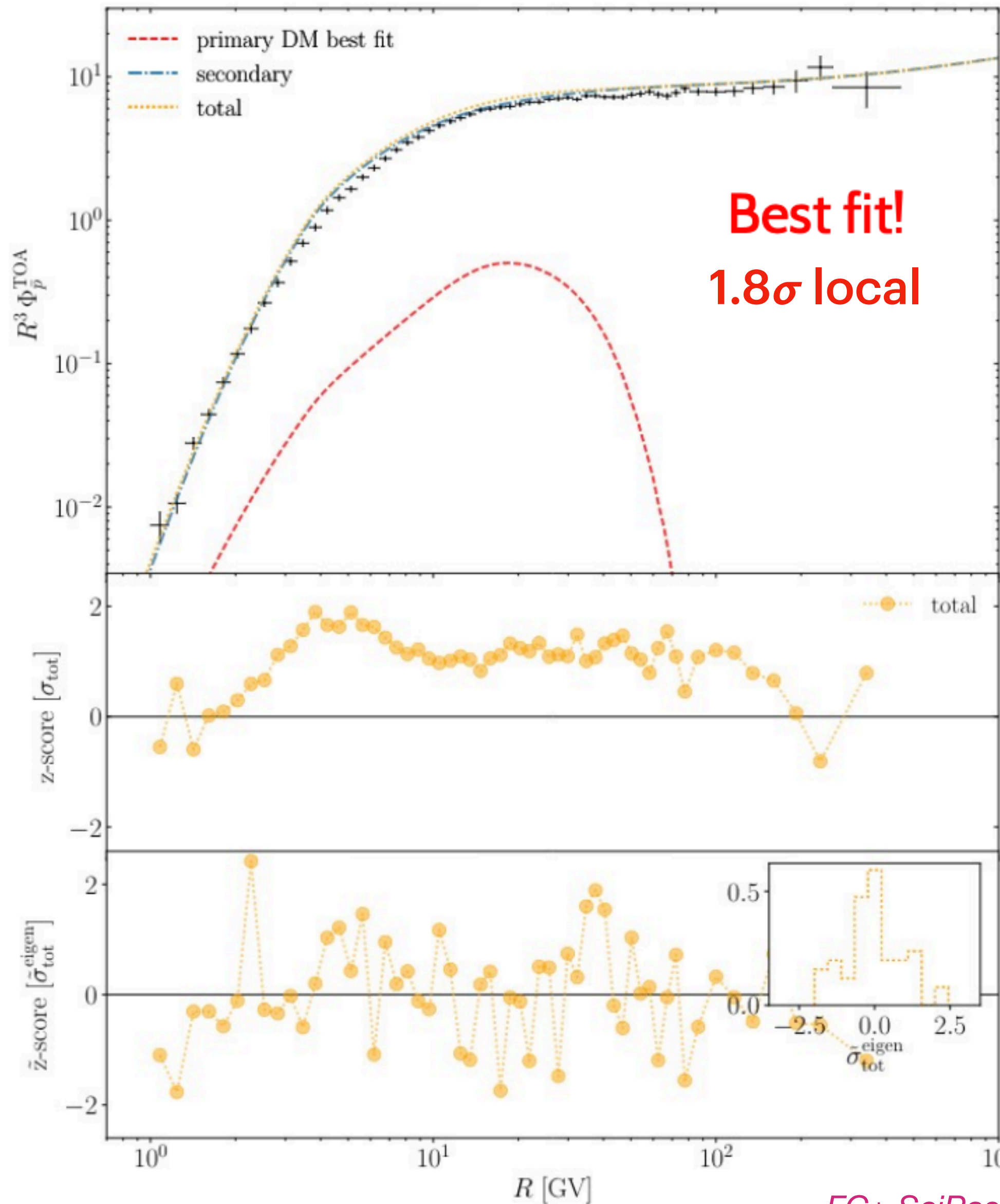
$$LR(\mu_0) = -2 \ln \frac{\sup_{\lambda \in \Lambda} \mathcal{L}(\lambda, \mu_0)}{\sup_{\{\lambda, \mu\} \in \Lambda \cup M} \mathcal{L}(\lambda, \mu)}$$

CR-specific parameters vs DM-specific ones

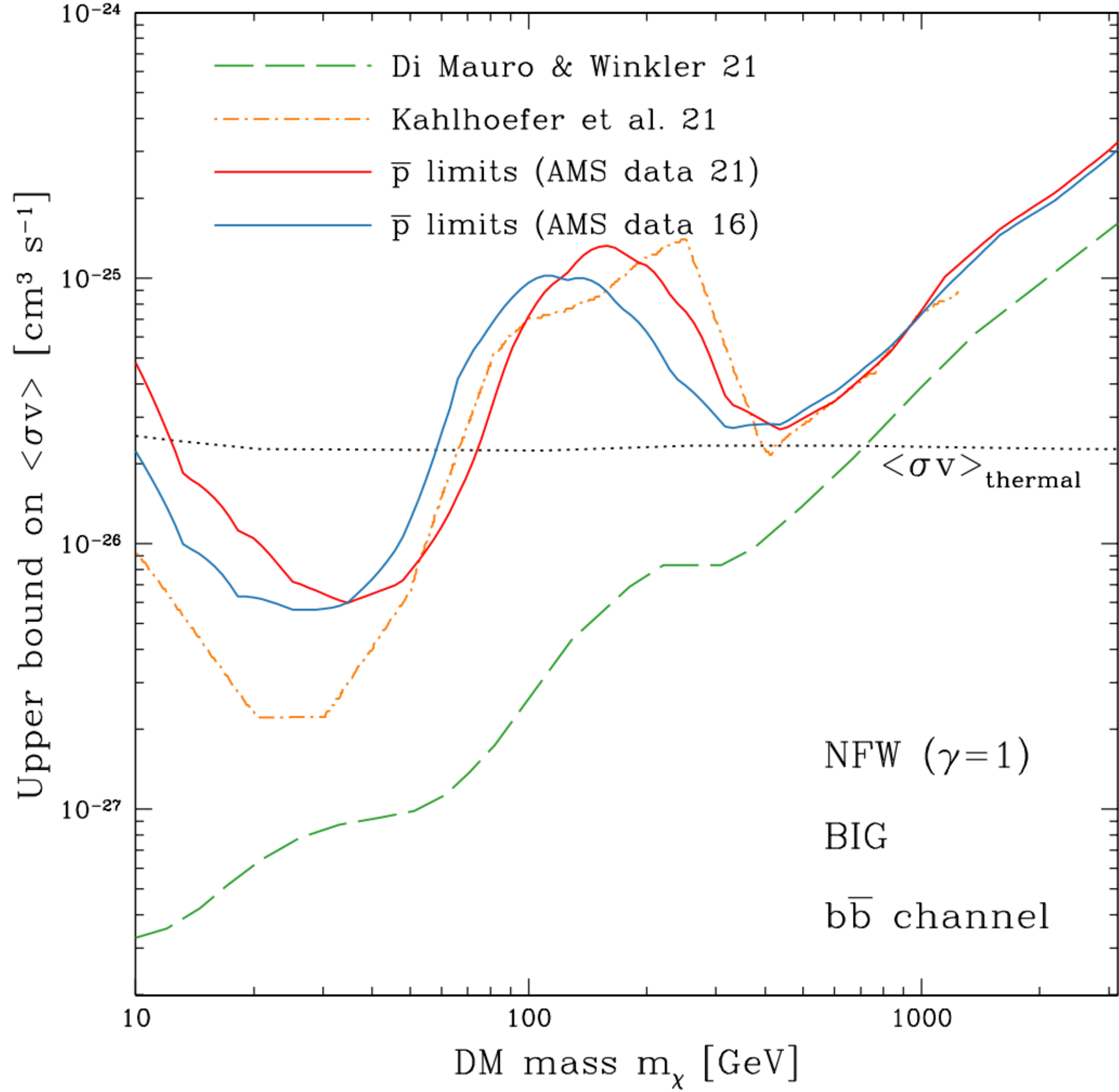
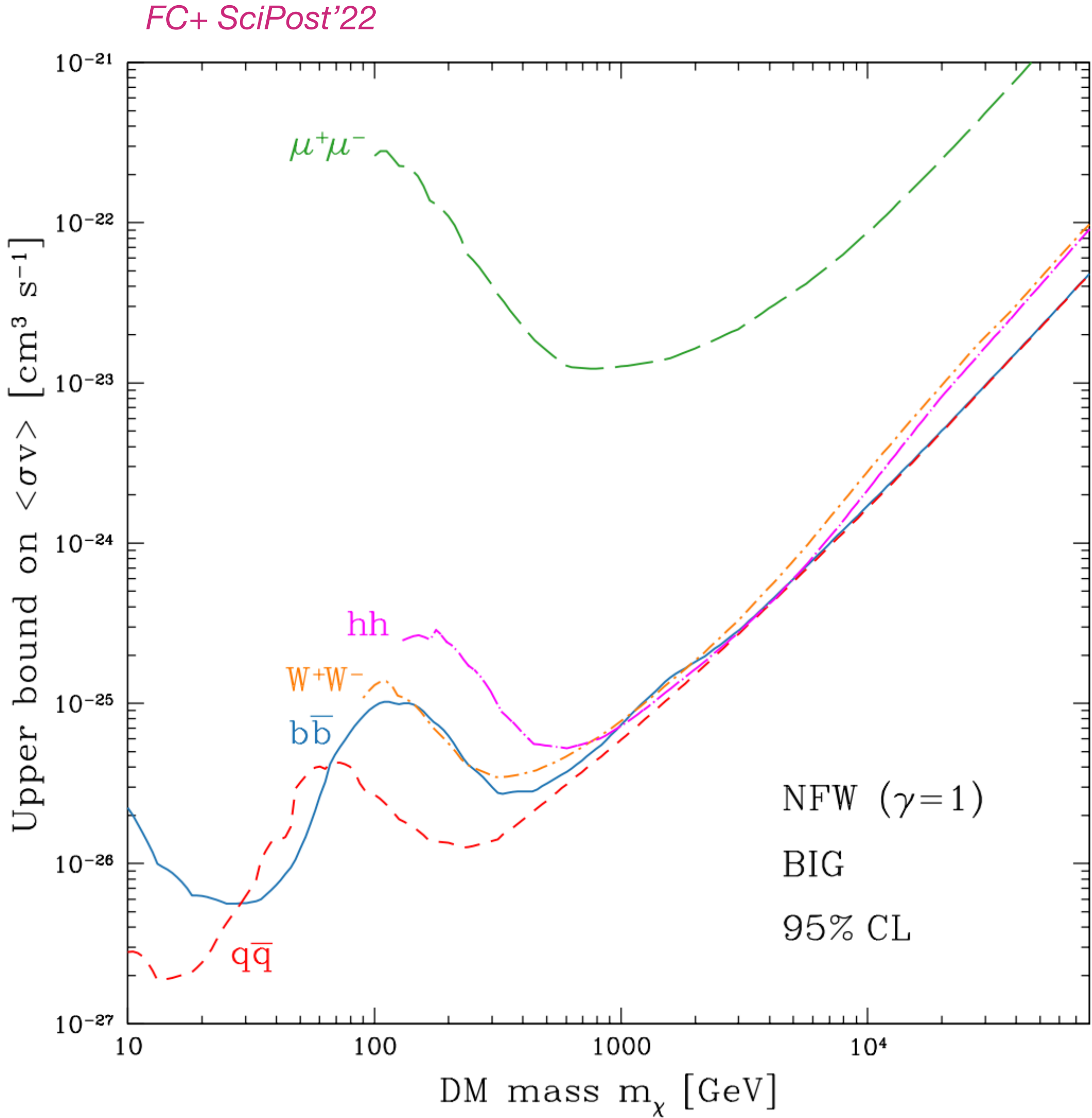
$$-2 \ln \mathcal{L}(\lambda, \mu) \equiv \chi_{\text{LiBeB}}^2(\lambda) + \chi_{\bar{p}}^2(\lambda, \mu)$$

1. CR parameters derived from LiBeB, are good for anti-protons
2. DM does not alter best-fit for propagation parameters since subdominant
3. Uncertainty on primary antiproton flux dominated by the size of the diffusive halo, L

$$-2 \ln \mathcal{L}(\lambda, \mu) \equiv -2 \ln \mathcal{L}(L, \mu) = \left\{ \frac{\log L - \log \hat{L}}{\sigma_{\log L}} \right\}^2 + x_i (\mathcal{C}^{-1})_{ij} x_j$$



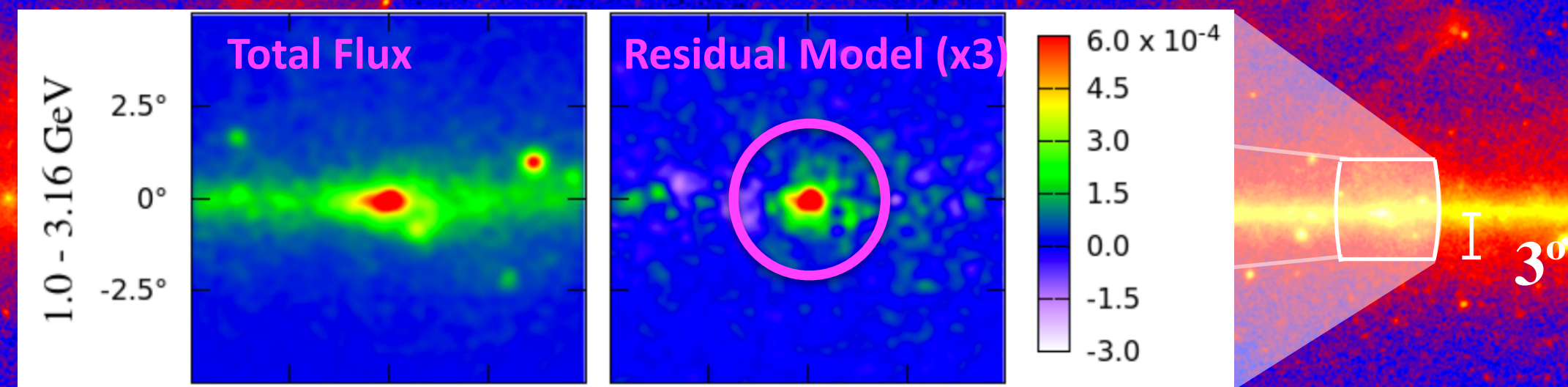
Antiproton constraints on WIMPs



The Fermi-LAT GeV excess

The Galactic centre GeV excess

The Galactic centre ROI



Daylan+ PDU'14

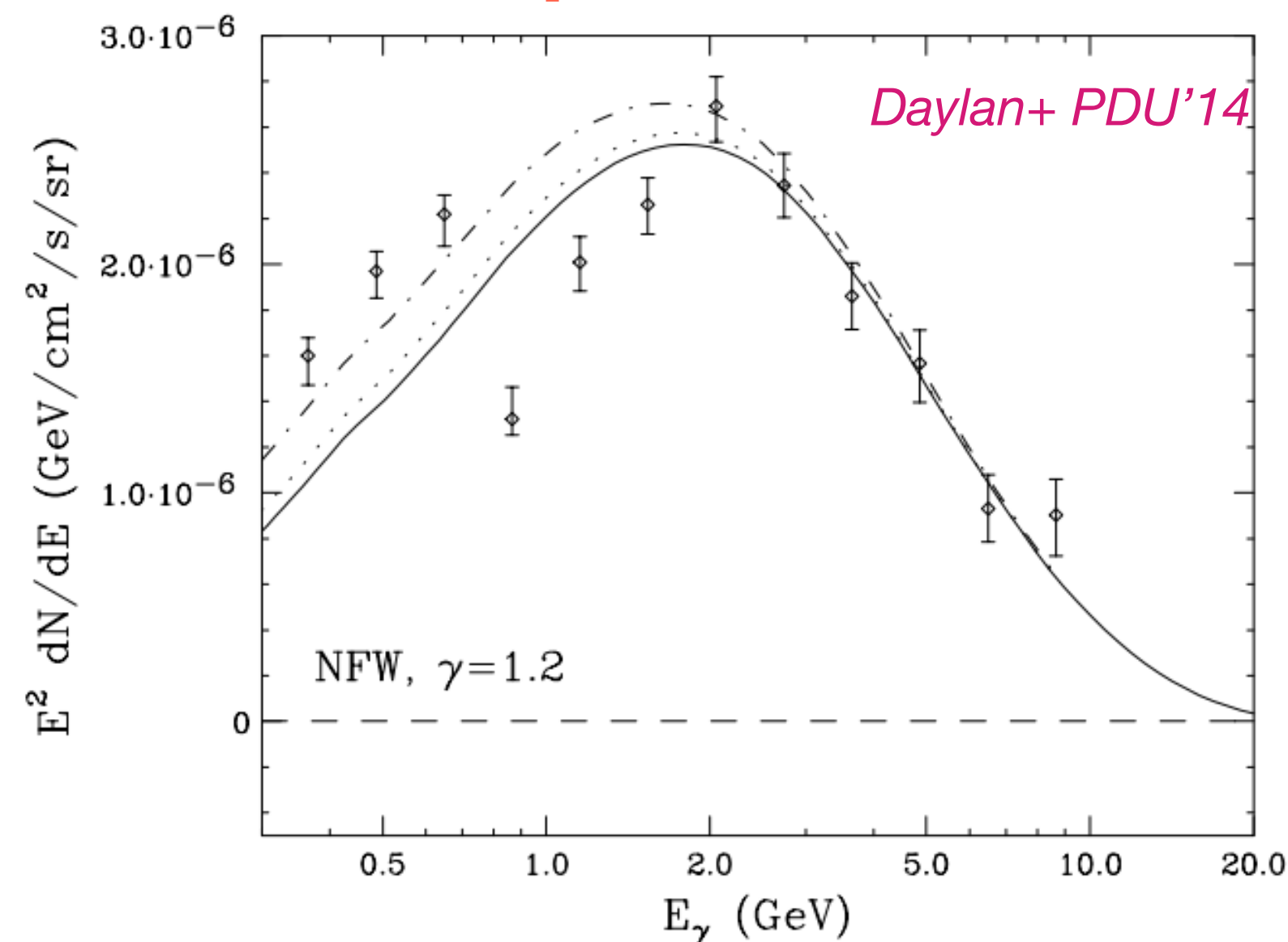
Hooper&Goodenough '09; Vitale&Morselli '09;
Hooper&Linden PRD'11;
Hooper&Goodenough PLB'11;
Boyarsky+ PLB'11;
Abazajian&Kaplinghat PRD'12;
Macias&Gordon PRD'14;
Abazajian+ PRD'14; Daylan+ '14;
Huang+ '15; Carlson+ '15; Ajello+15;
Casandjian Fermi Symp.'14;
de Boer+'16; Macias+'16; etc.

The GeV excess

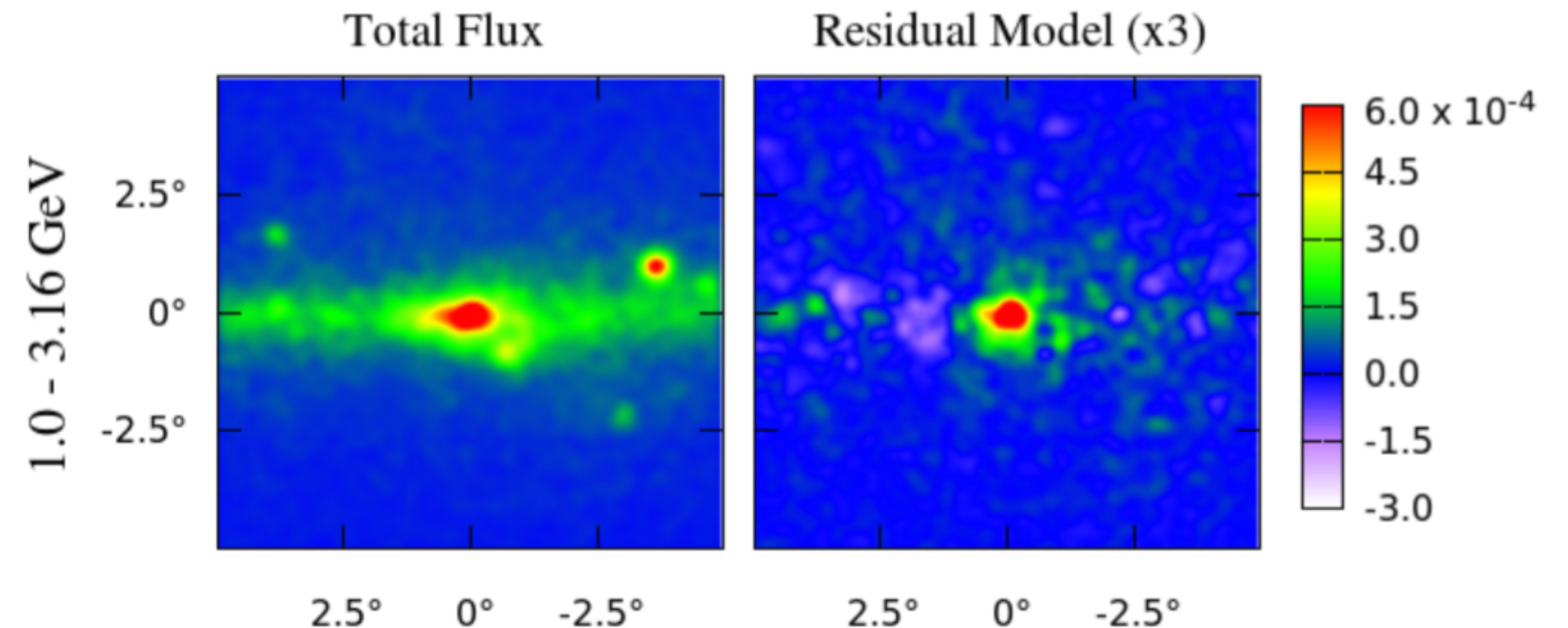
Galactic centre characterisation

$$|\ell|, |b| \lesssim 2^\circ$$

Spectrum



Morphology



✓ **Extended excess emission** above: model for diffuse emission, Sgr A* and other point sources

✓ The **spectrum** might strongly suffer from **background modelling**

✓ Compatible to be **spherically symmetric** about the Galactic centre

✓ Connection with HESS TeV GC ridge

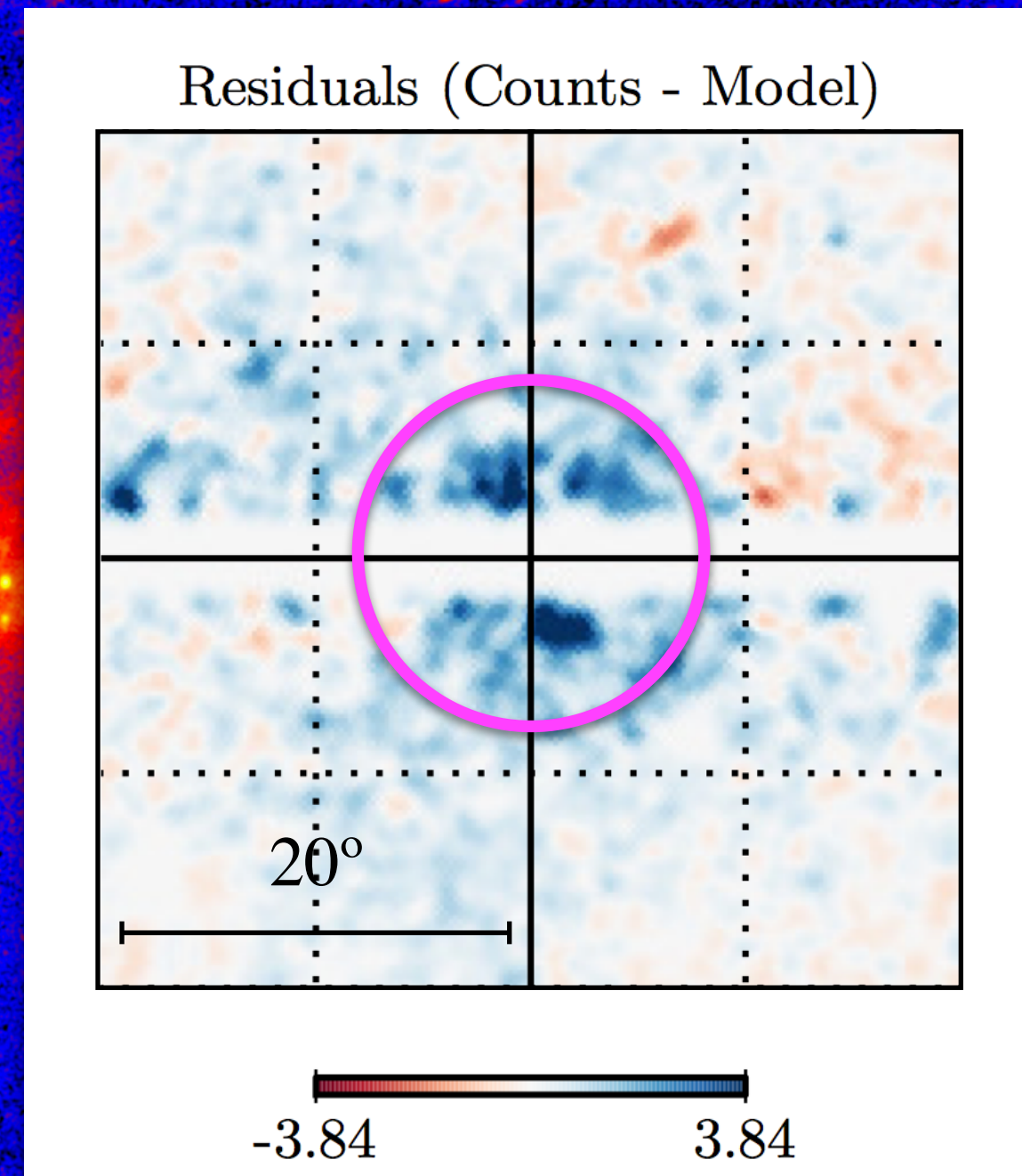
Abazjian+ PRD'14

$$\frac{dn}{dV} \sim r^{-\Gamma} \quad \Gamma \sim 2.6$$

Macias&Gordon PRD'14; Macias+ MNRAS'15

The Galactic centre GeV excess

The inner Galaxy ROI



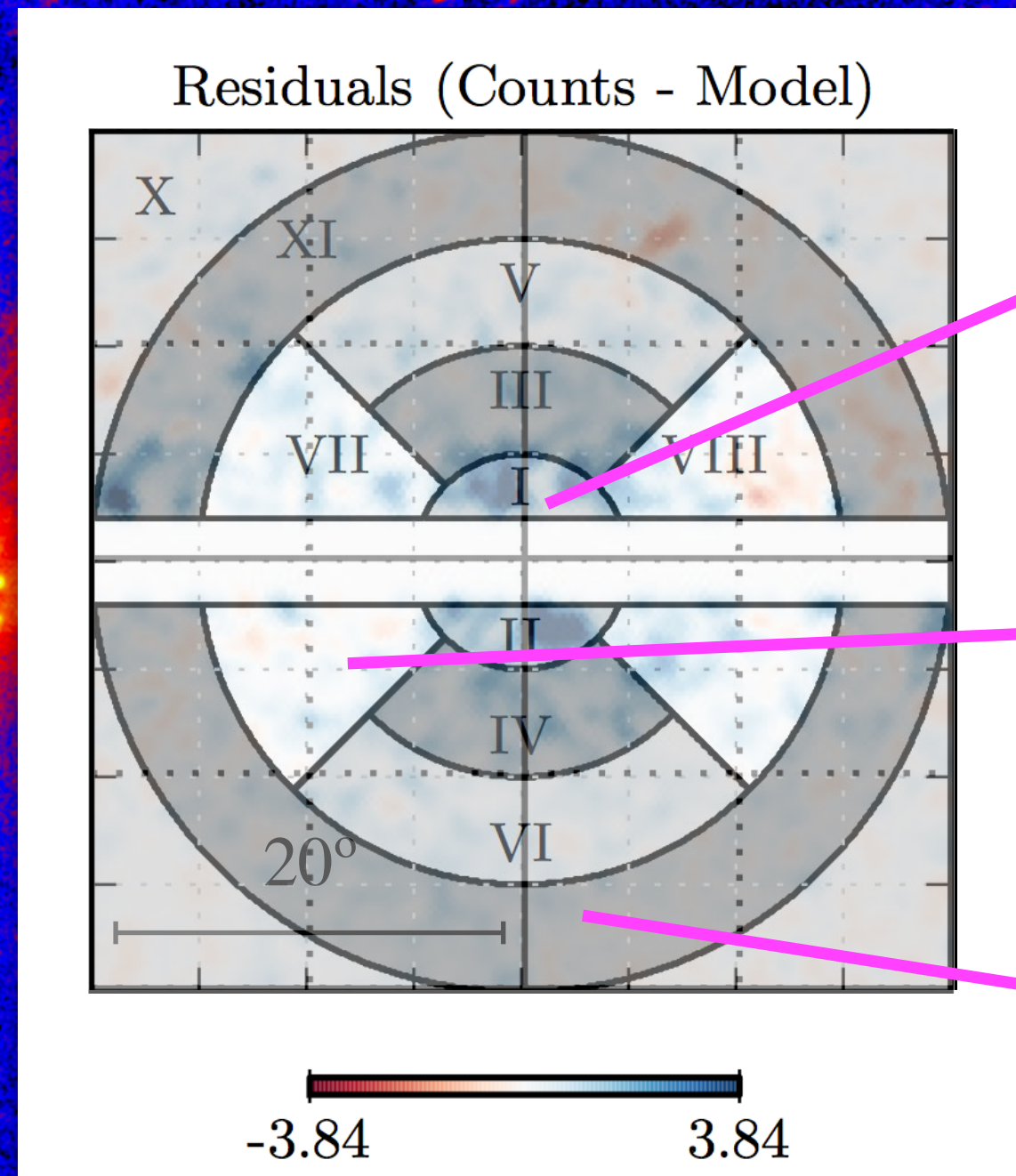
FC+ JCAP'15

20°

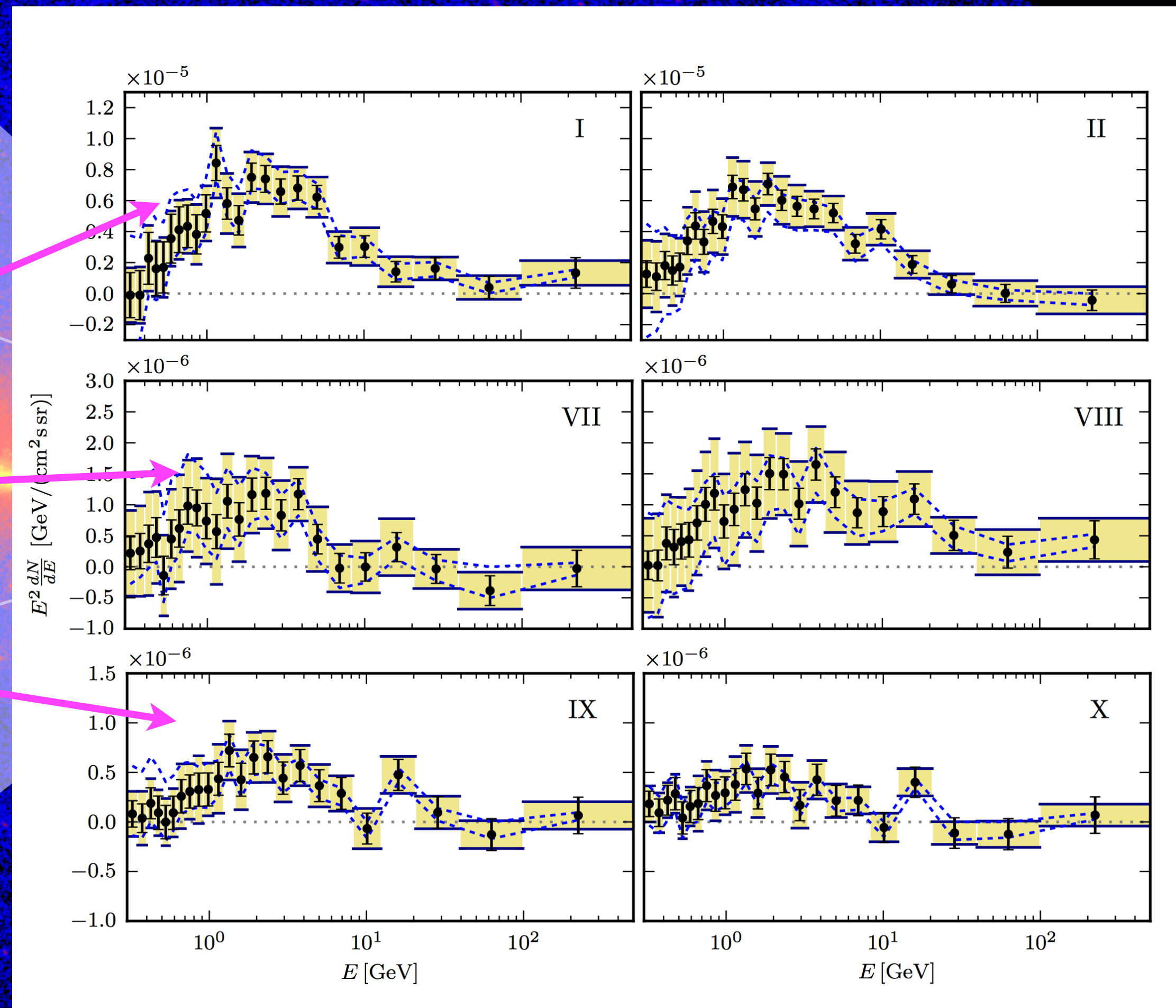
Hooper&Slatyer PDU'13; Huang+ JCAP'13;
Zhou+ PRD'15; Daylan+ PDU'14; FC+ JCAP'15;
Gaggero+ JCAP'15; Ajello+ 2015; Huang+JCAP '15
Linden+PRD'16; Horiuchi+'16; Ackermann+ApJ'17; Ackermann+2017; etc.

The Galactic centre GeV excess

The inner Galaxy ROI



FC+ JCAP'15



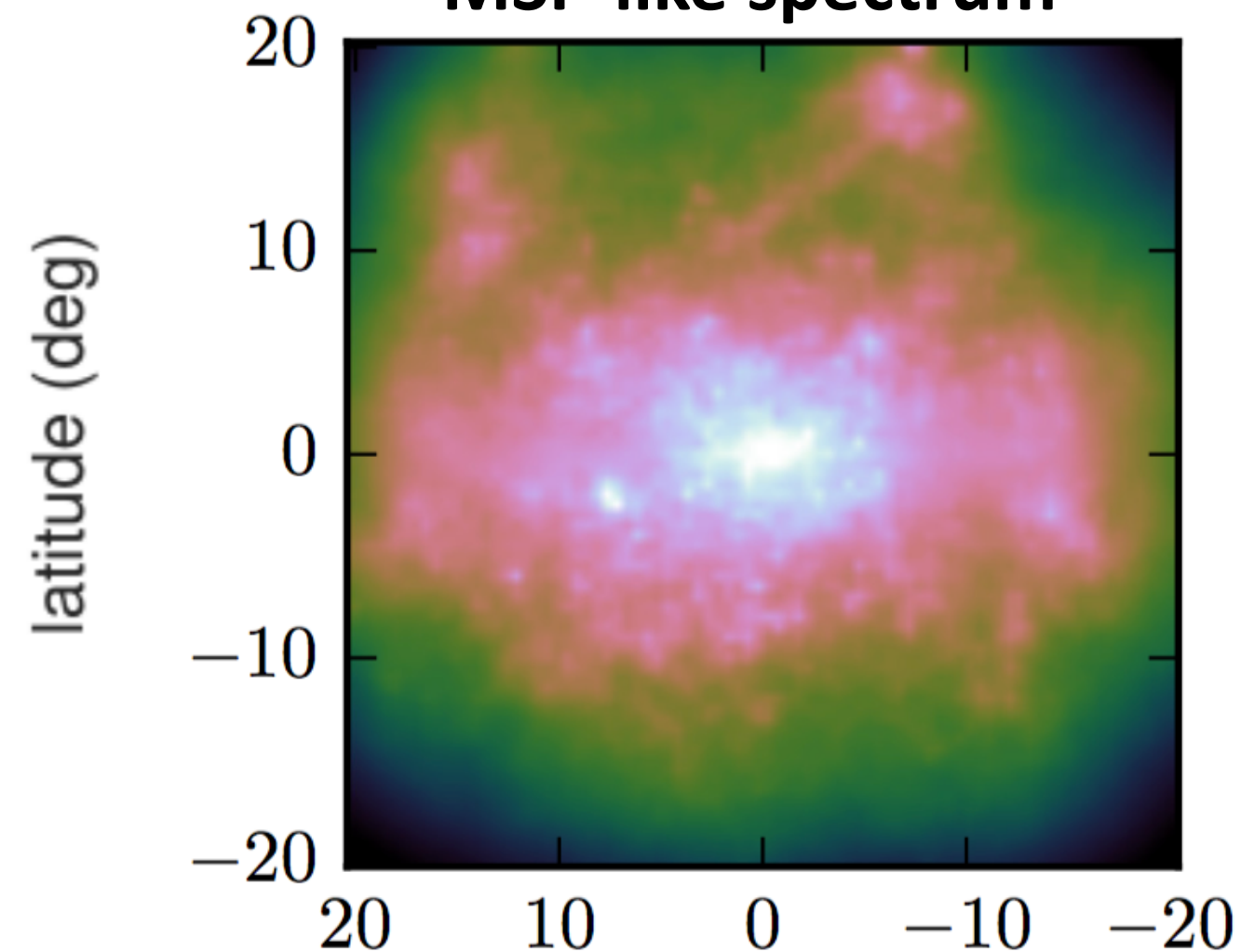
1. Almost uniform spectrum peaked at ~ 2 GeV
2. Extended at least up to 10 degrees

Hooper&Slatyer PDU'13; Huang+ JCAP'13;
Zhou+ PRD'15; Daylan+ PDU'14; FC+ JCAP'15;
Gaggero+ JCAP'15; Ajello+ 2015; Huang+JCAP '15

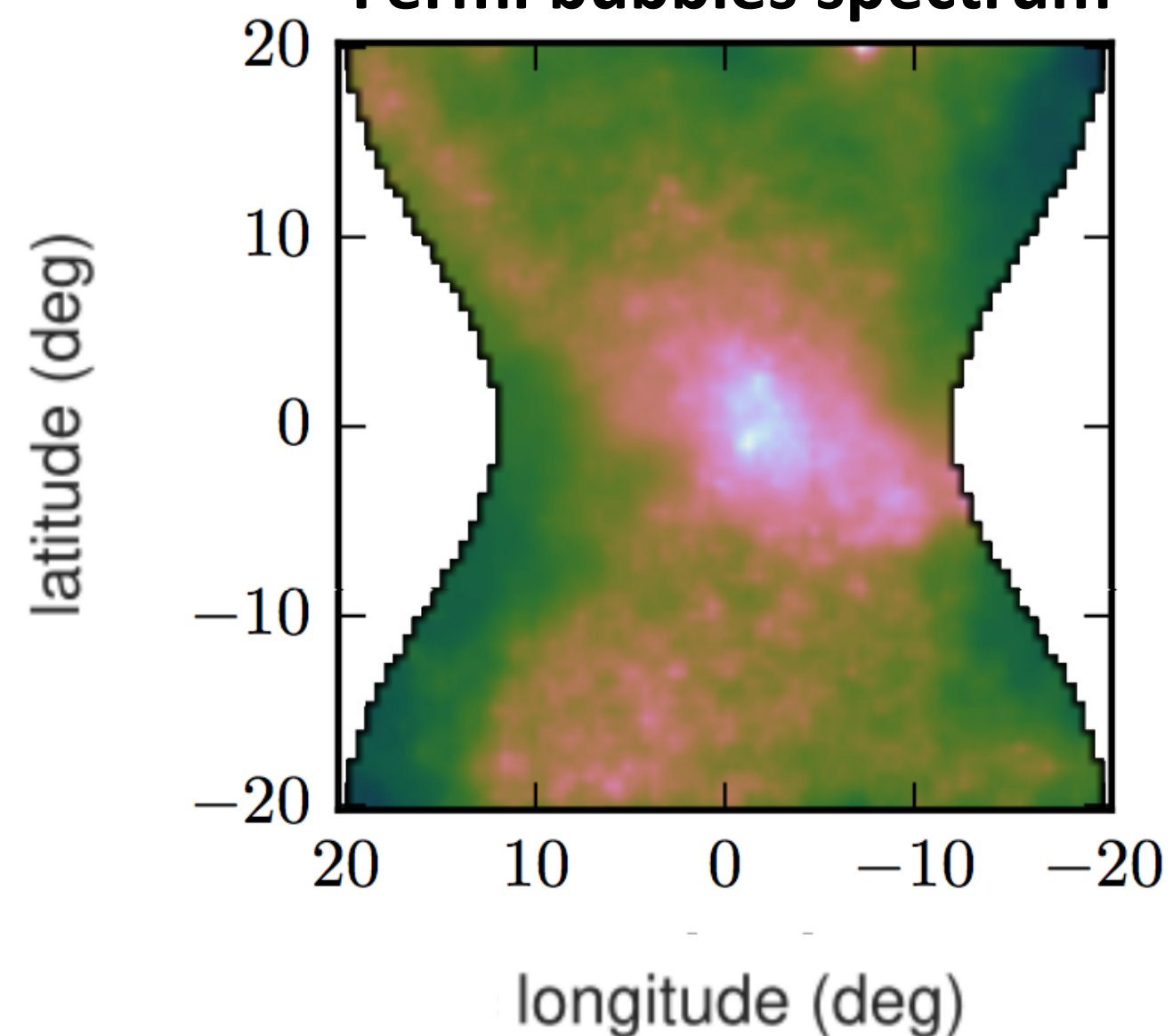
Linden+PRD'16; Horiuchi+'16; Ackermann+ApJ'17; Ackermann+2017; etc.

The GeV excess emission

MSP-like spectrum



Fermi bubbles spectrum



- Established evidence for an excess emission above **known** astrophysical backgrounds (diffuse emission + point-like sources)
- **Several independent techniques** find analogous results (template fitting, spectral decomposition, image reconstruction)

Hooper+ PDU'13; Huang+ JCAP'13; Daylan+ '14; FC+ JCAP'15; Ajello+ ApJ'15; Gaggero+ JCAP'15; etc

Selig+ A&A'14; Huang+ JCAP'16; de Boer+'16

Storm, Weniger & FC JCAP'17

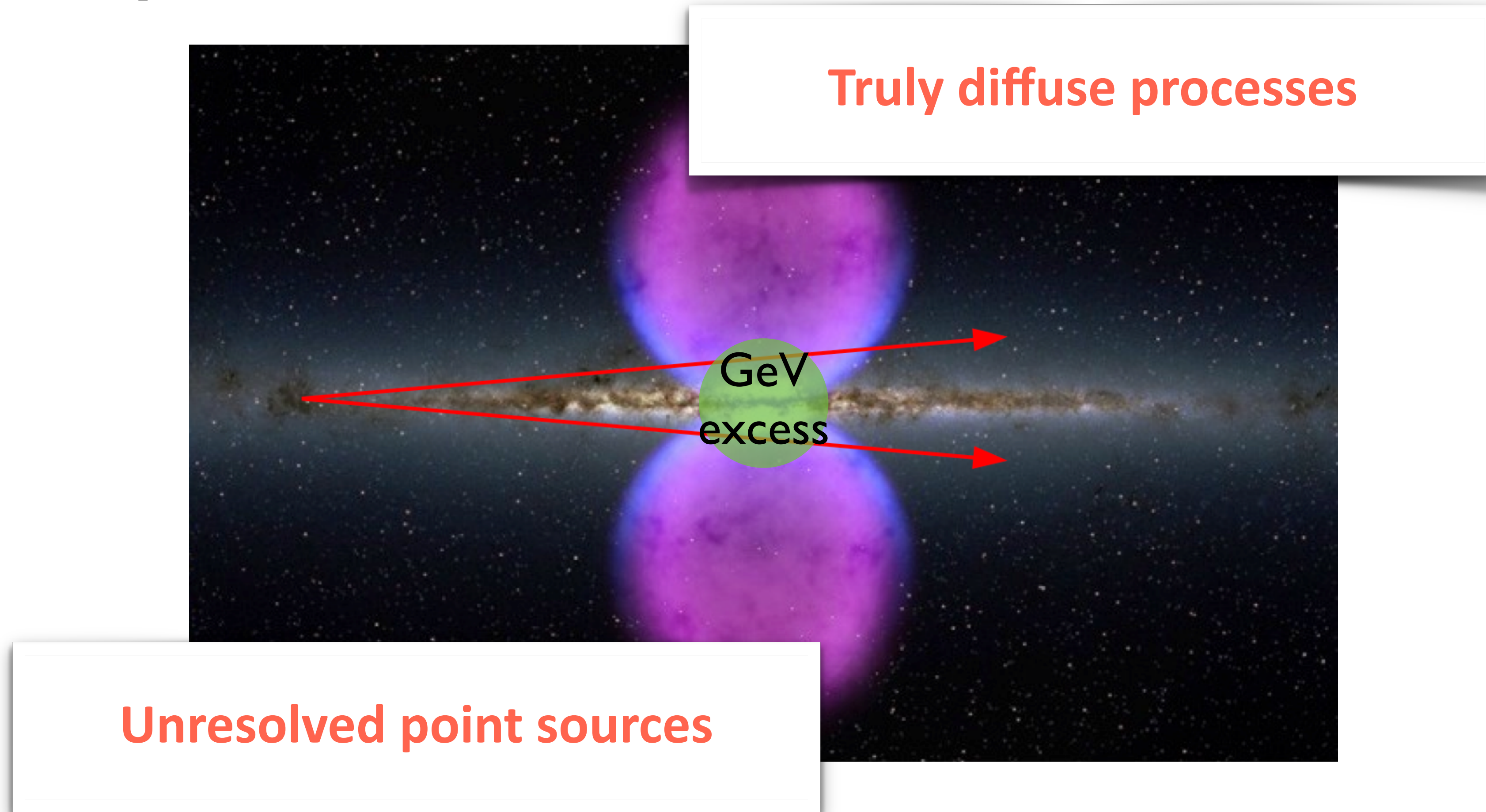
- **Template fitting - image reconstruction hybrid approach** (SKYFACT) has been proved very powerful in disentangling gamma-ray emission components

Storm, Weniger & FC JCAP'17

- **Residuals reduced significantly** when (realistic) nuisance parameters are included in the fit

What is the origin of the GeV excess?

Possible interpretations



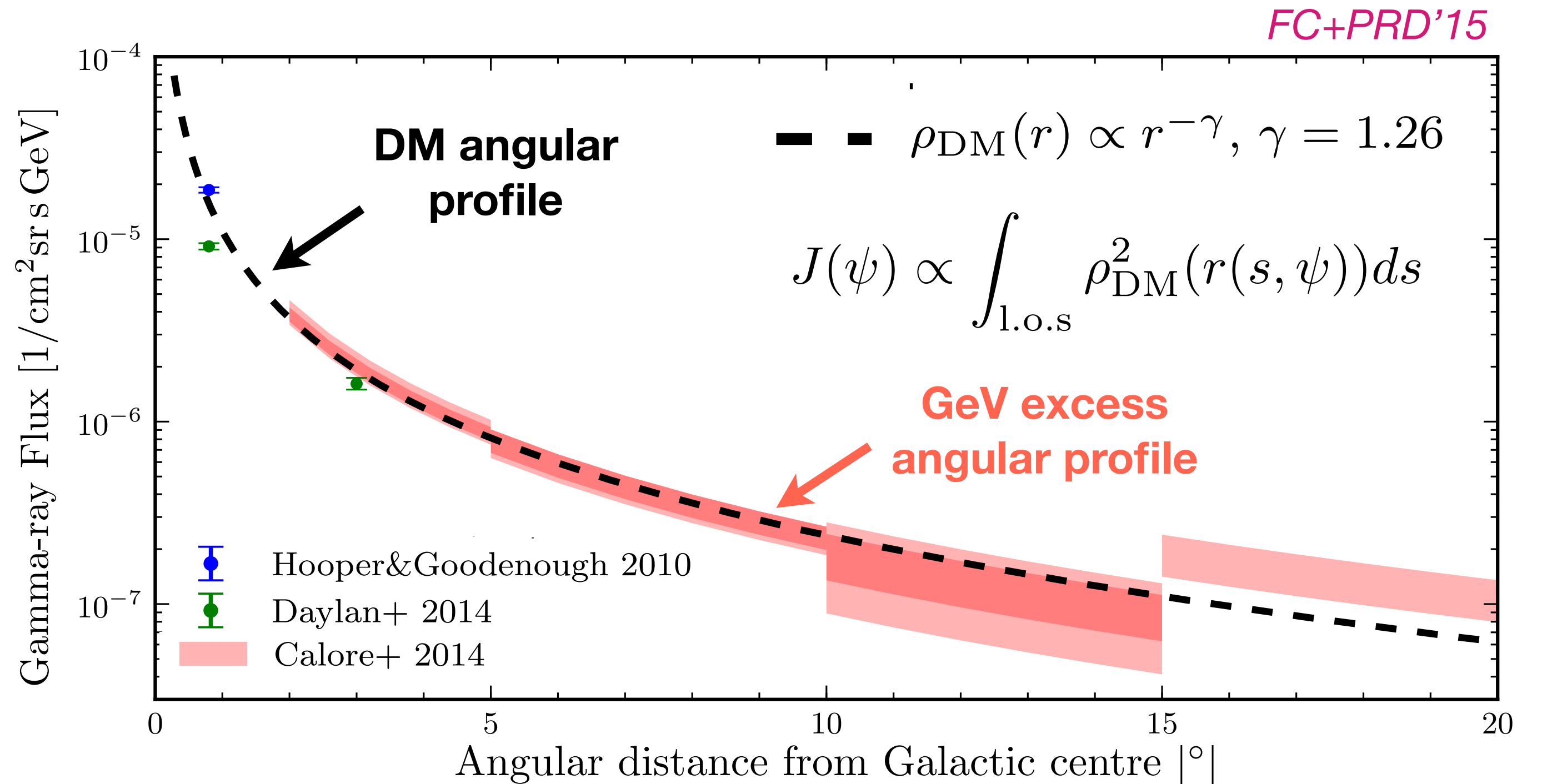
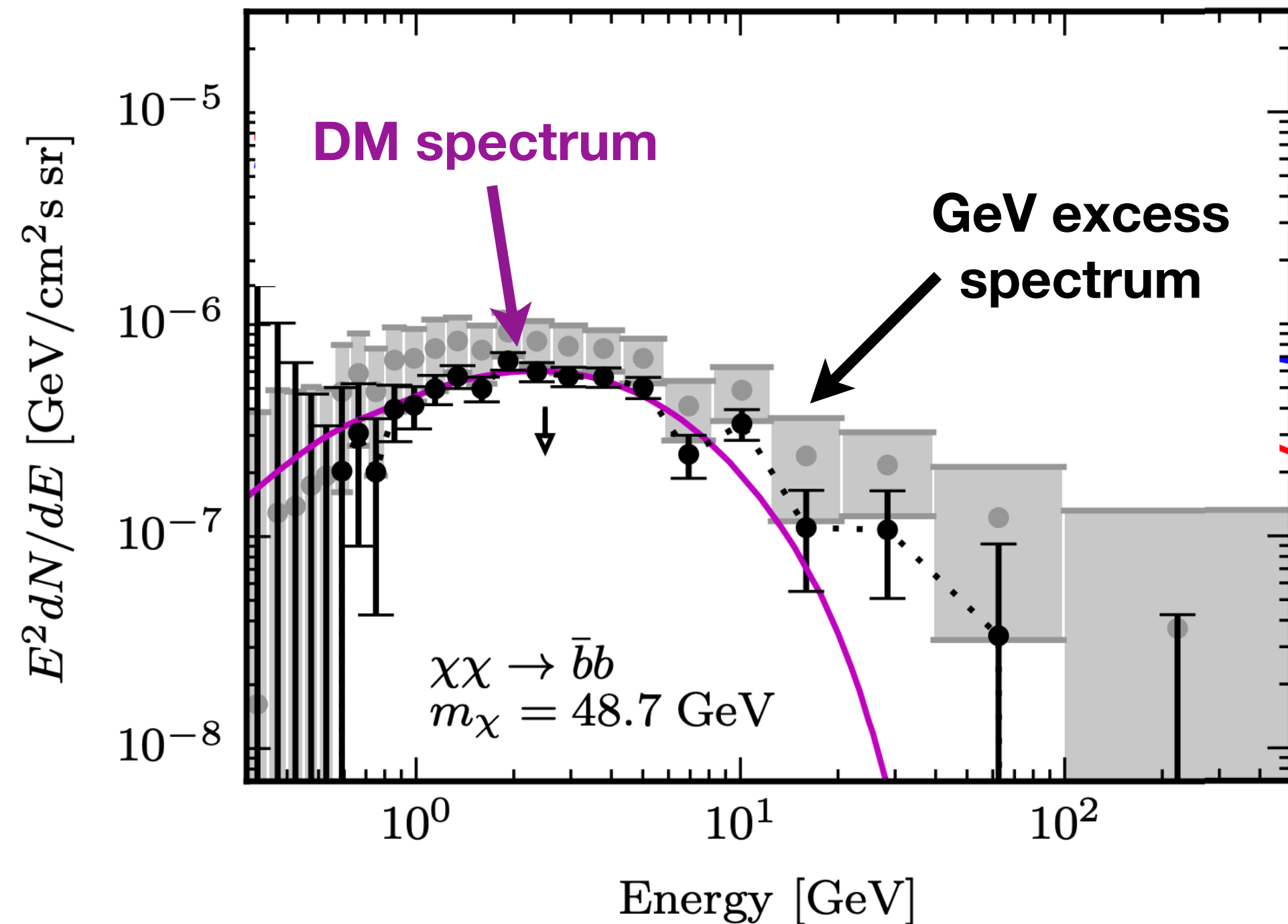
Constraints:

(a) Spectrum & Morphology of the excess? (b) Emission in other wavelengths?

Diffuse processes I

Gamma rays from dark matter (DM) annihilation

- **Decay/Annihilation** of DM particles would lead to the production of **final gamma rays** with specific energy and spatial distribution



Agrawal+JCAP'15; Achterberg+JCAP'15; Bertone, FC+ JCAP'15; Liem, FC+ JCAP'16; O(>100) papers

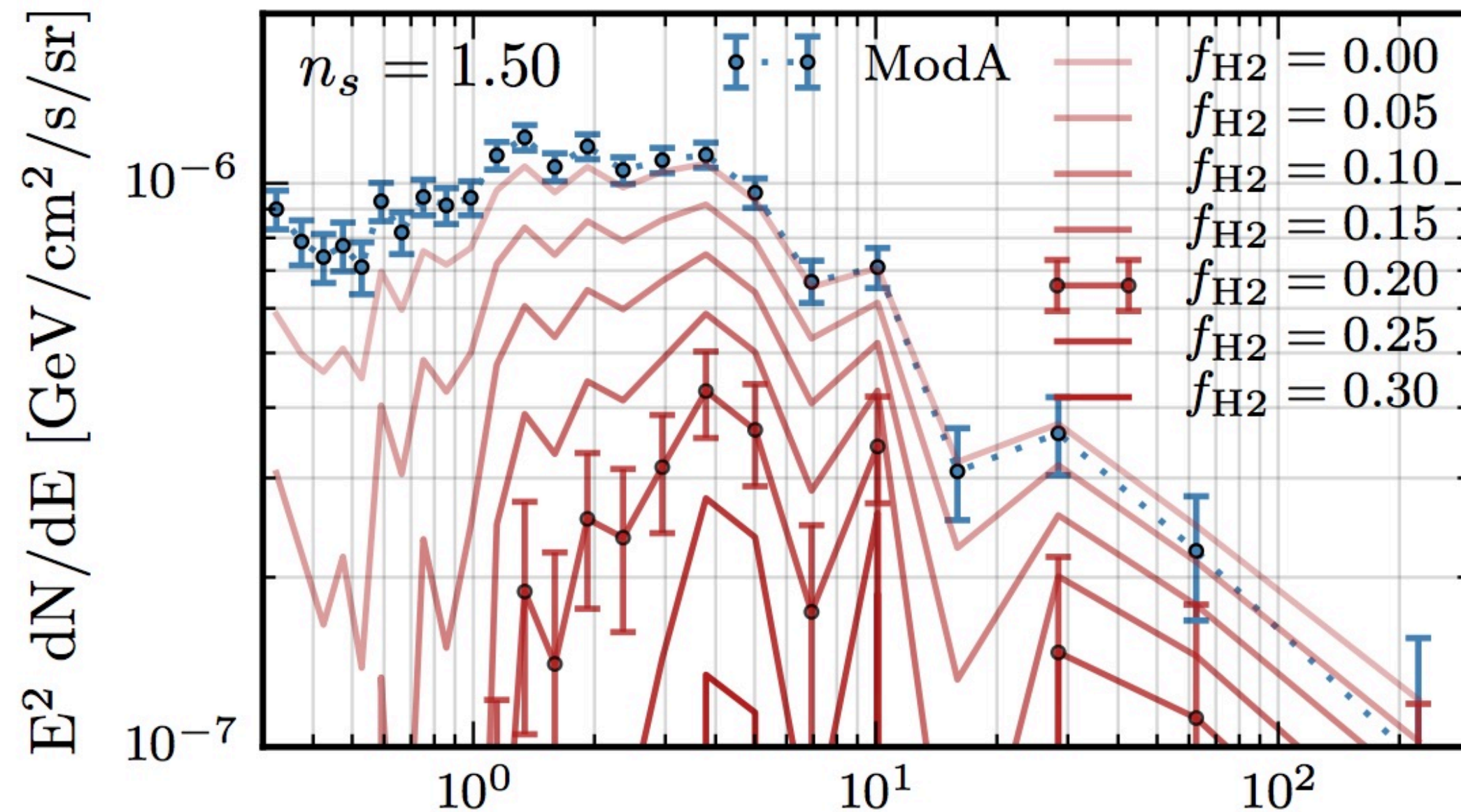
Diffuse processes II

Cosmic rays in the GC

- **New population of cosmic rays** injected at the GC (electrons mostly)
- **Steady state** (from star formation in CMZ) and/or **time-dependent** (from outburst activity of the GC) source term

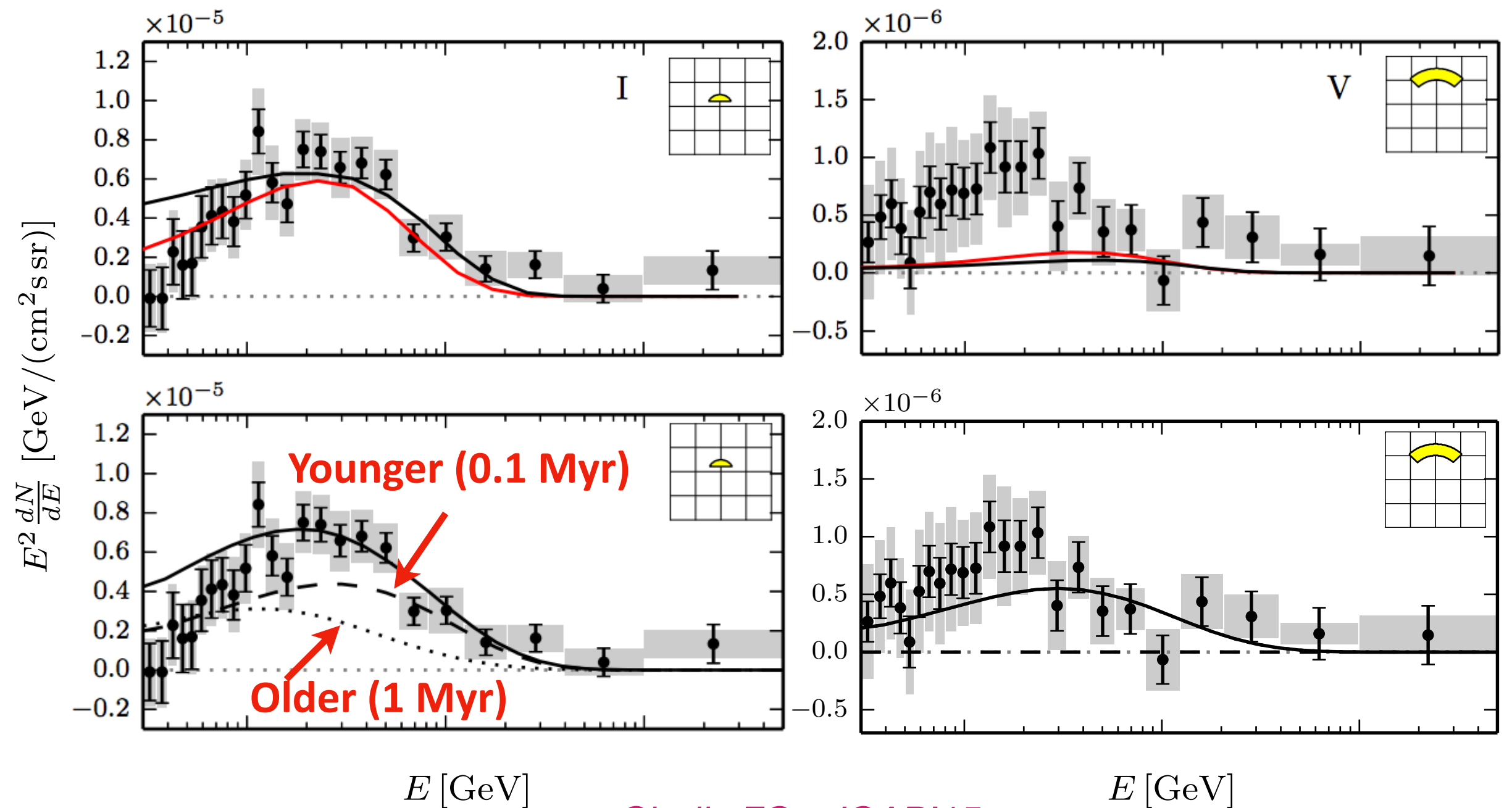
Gaggero+JCAP'15; Carlson+ PRD'16, PRL'16

Petrovic+ JCAP'14; Cholis,FC+ JCAP'15



Carlson+ PRD'16

Additional CR injection at the GC, accounting for enhanced SFR traced by H2 regions (5-10% of total SFR)

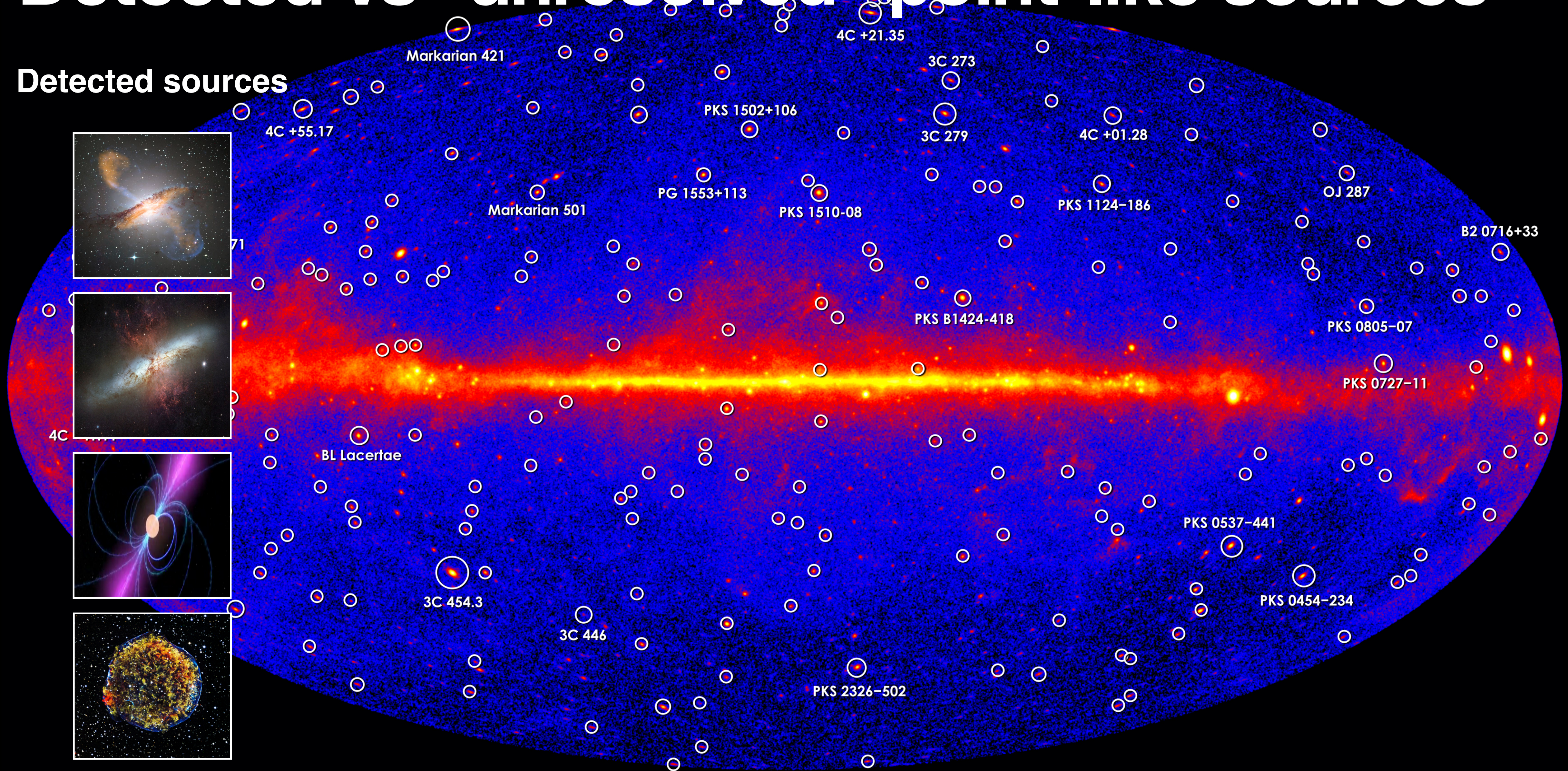


Cholis,FC+ JCAP'15

Time-dependent (burst) injection of leptons at the GC, and tuning of burst parameters (age, duration, injection spectrum, propagation parameters)

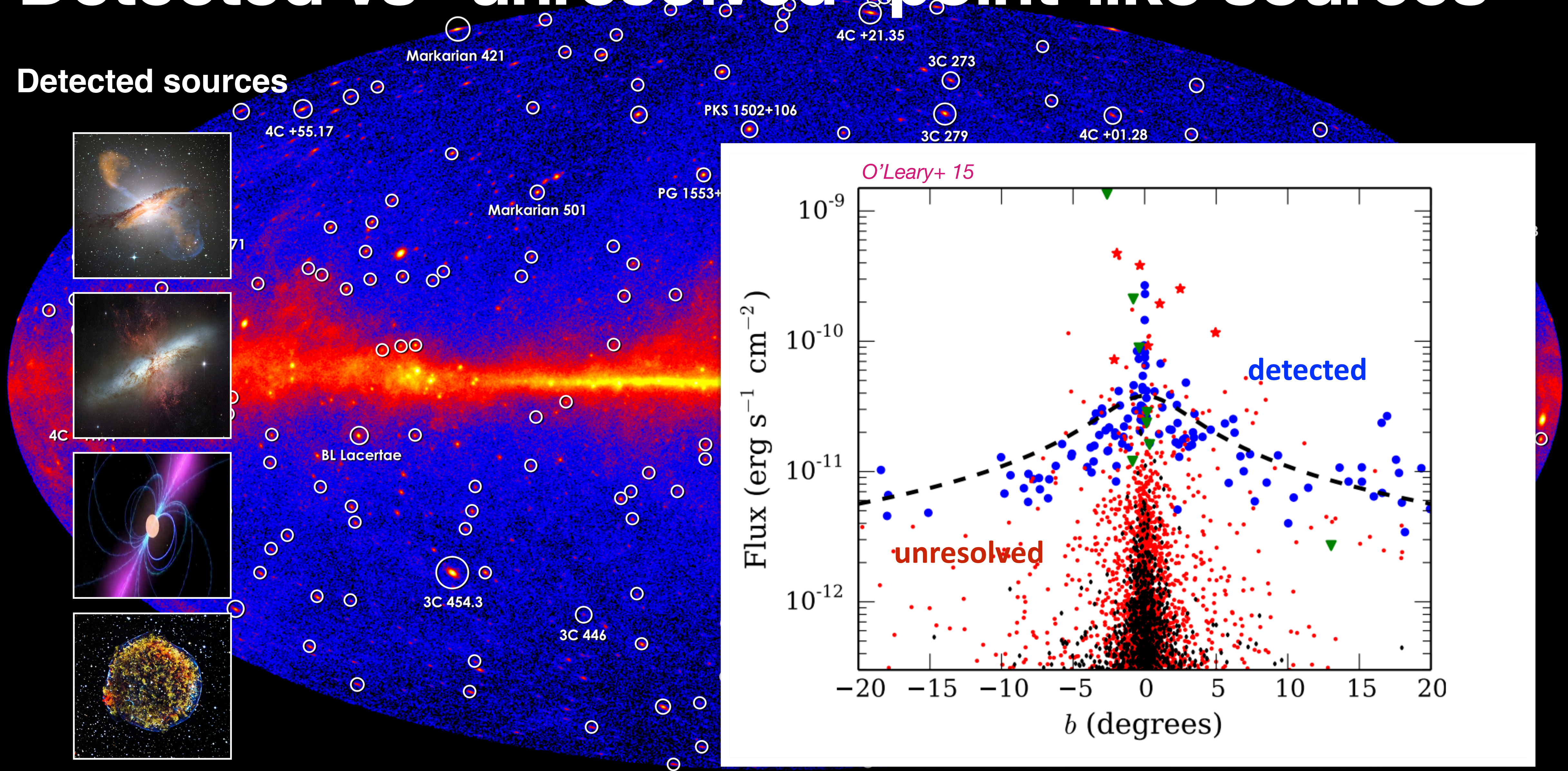
Detected vs “unresolved” point-like sources

Detected sources



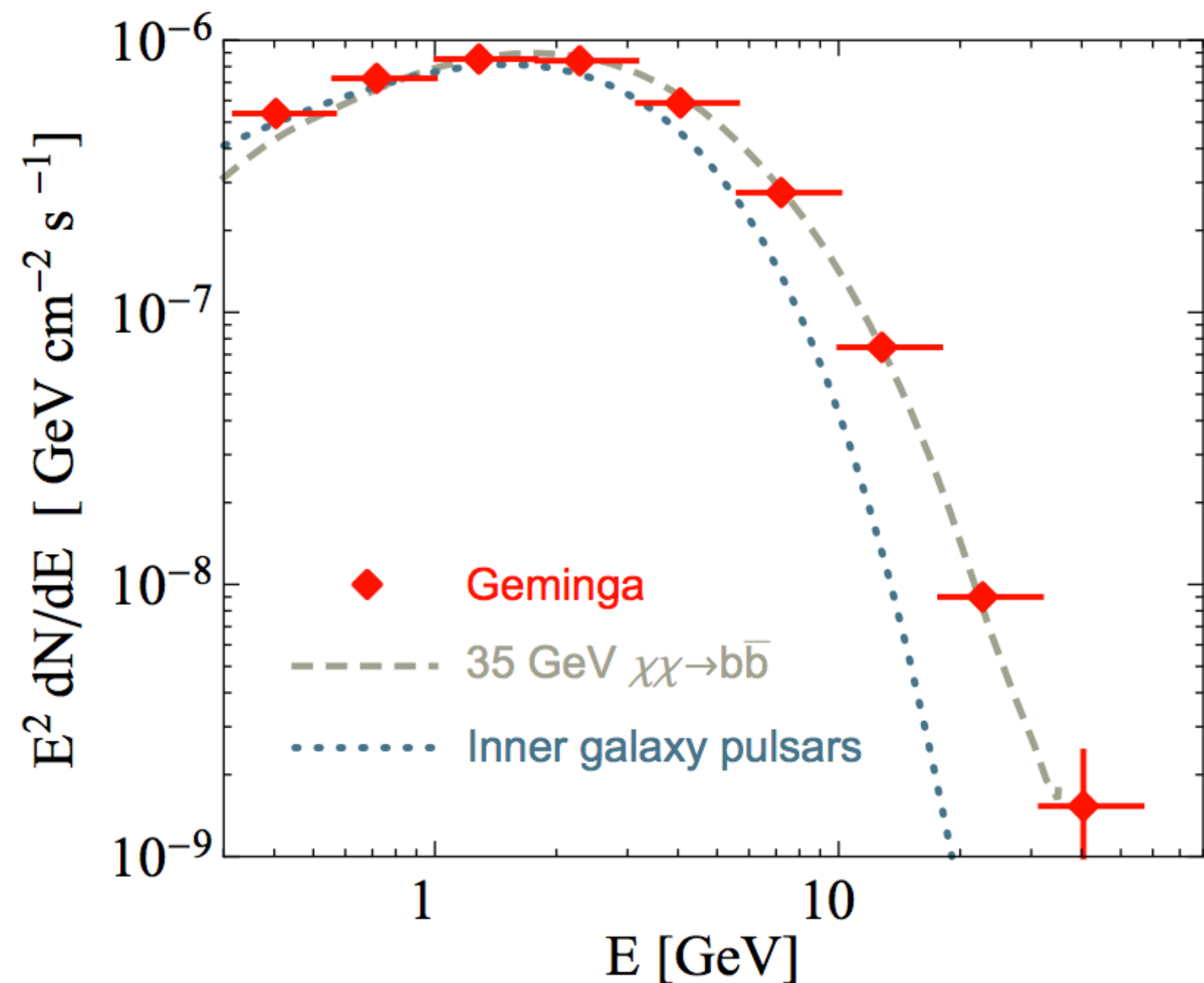
Detected vs “unresolved” point-like sources

Detected sources



Unresolved sources: PSR and MSPs

Spectrum



- ✓ Excess spectrum compatible with observed **millisecond pulsars** (MSPs), and marginally **young pulsars**

Abazajian&Kaplinghat'12

Morphology

$$\epsilon \propto r^{-\Gamma} e^{-r/R_{\text{cut}}}$$

$$\Gamma = 2.5 \quad R_{\text{cut}} = 3 \text{ kpc}$$

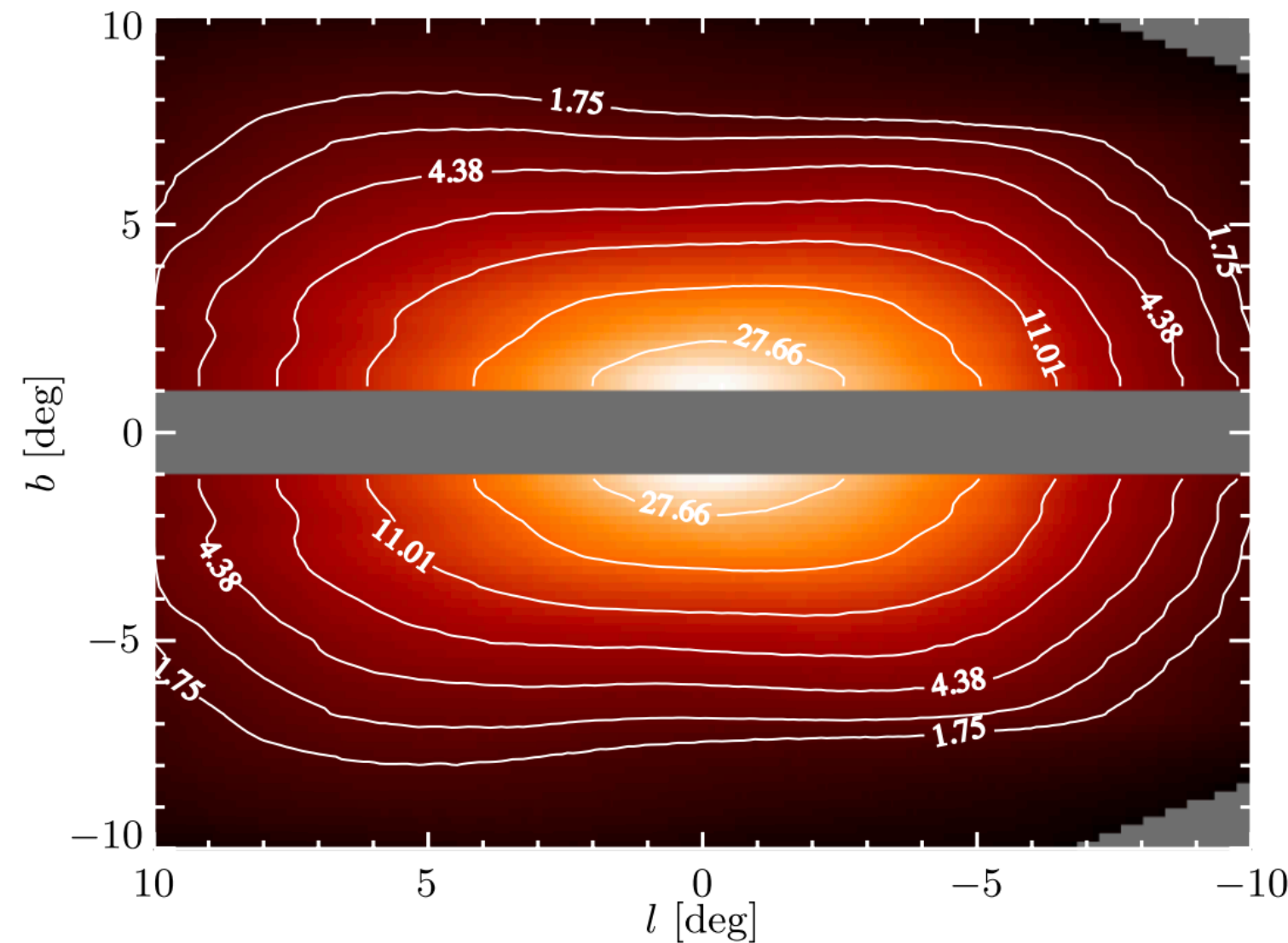
- ✓ Proposed population of **MSPs in the bulge** (vs disc)
Hooper+PRD'14; Petrovic+ JCAP'15; Yuang+ MNRAS'14;
- ✓ **Young pulsars** from SF in the CMZ, but difficult to explain spatial extent and observed bright ones
O'Leary+ '15; Linden PRD'16
- ✓ **Bulge MSPs** from tidally disrupted globular clusters
Brandt&Kocsis ApJ'15; Abbate et al. 2017; Fragione et al. 2017; Arca-Sedda et al. 2017; Macias+JCAP'19
- ✓ Issues in luminosity function of observed MSP and LMXB-to-MSP ratio
Cholis+'14; Hooper+'15; Hooper&Linden JCAP'16; Haggard+ JCAP'17; Ploeg+ JCAP'17

Going beyond dark matter templates

Stellar distribution in the bulge

Boxy bulge

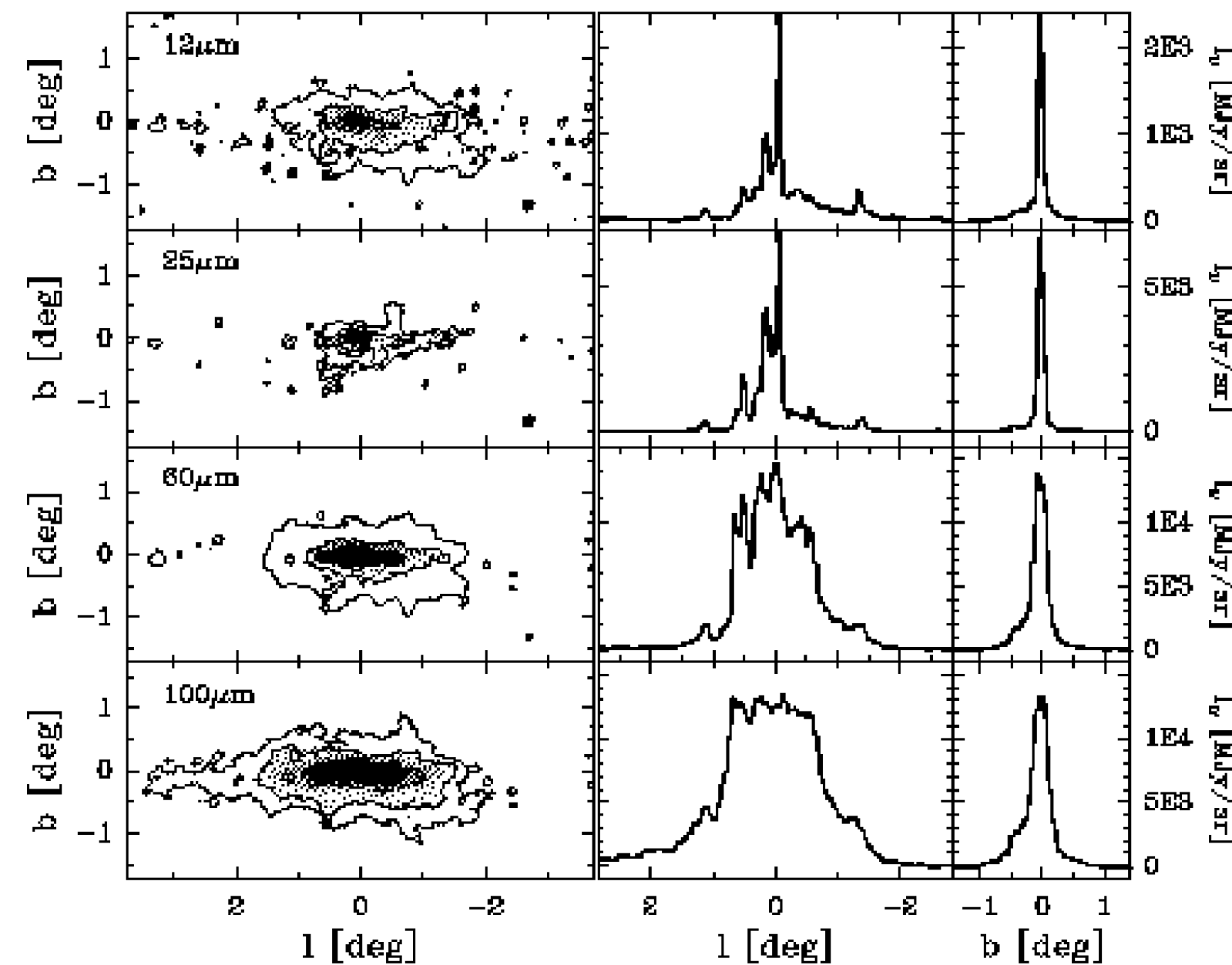
$$0.9 \times 10^{10} M_{\odot}$$



Wegg & Gerhard MNRAS'12

Nuclear bulge

$$1.4 \times 10^9 M_{\odot}$$

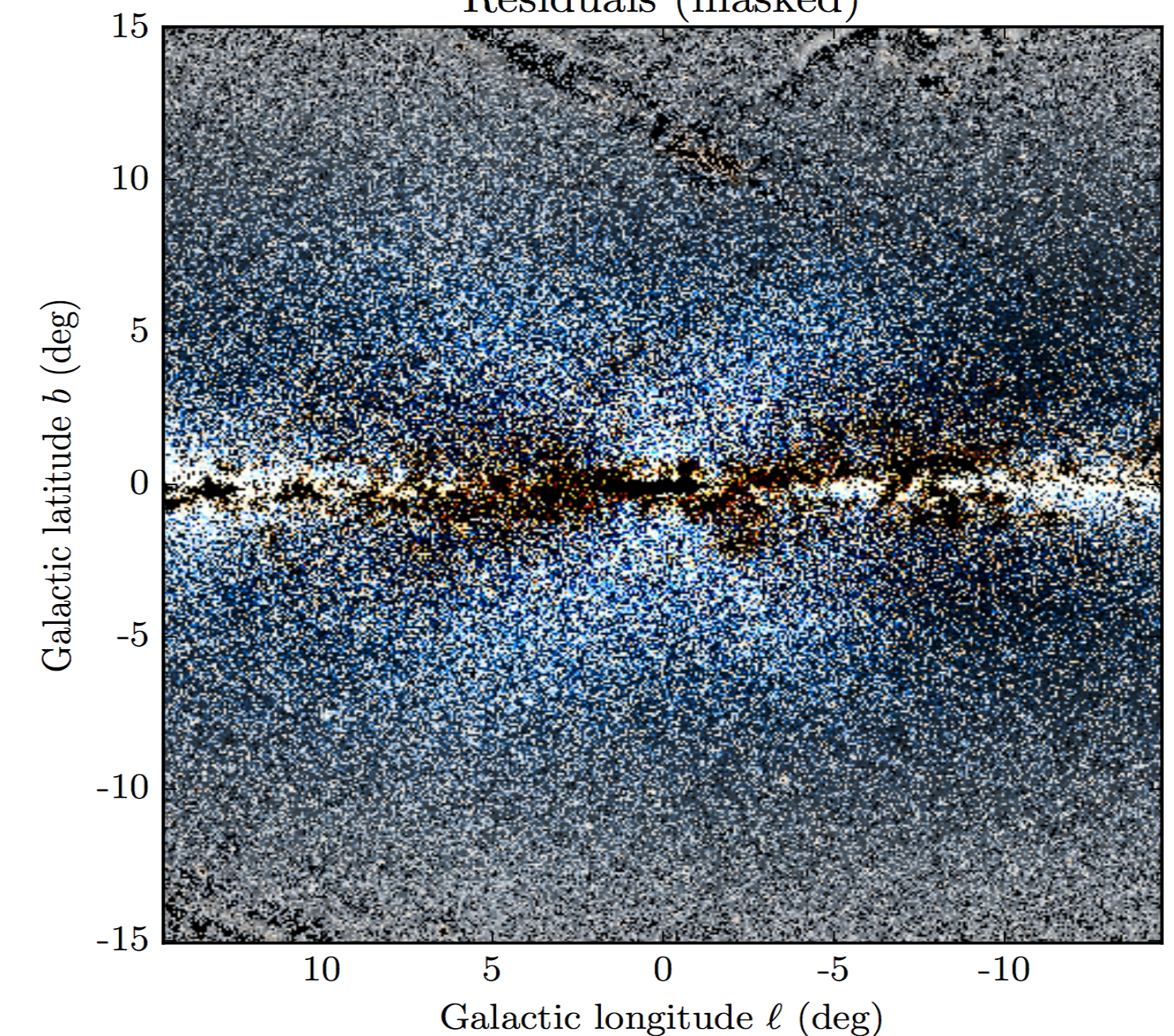


Launhardt+A&A'02

X-shaped bulge

~20% BB mass

Residuals (masked)



Ness&Lang AJ'16

- **Red Clump stars** (near-IR) used to characterise the **three-dimensional density structure** of the BB
- Most recent non-parametrically deconvolved bulge model w/ VISTA Variables in the Via Lactea (VVV) data Coleman+ MNRAS'20
- X-shaped structure characteristic of boxy/peanut like morphology (extragalactic studies of barred galaxies and simulations)

Evidence for the stellar bulge GeV emission

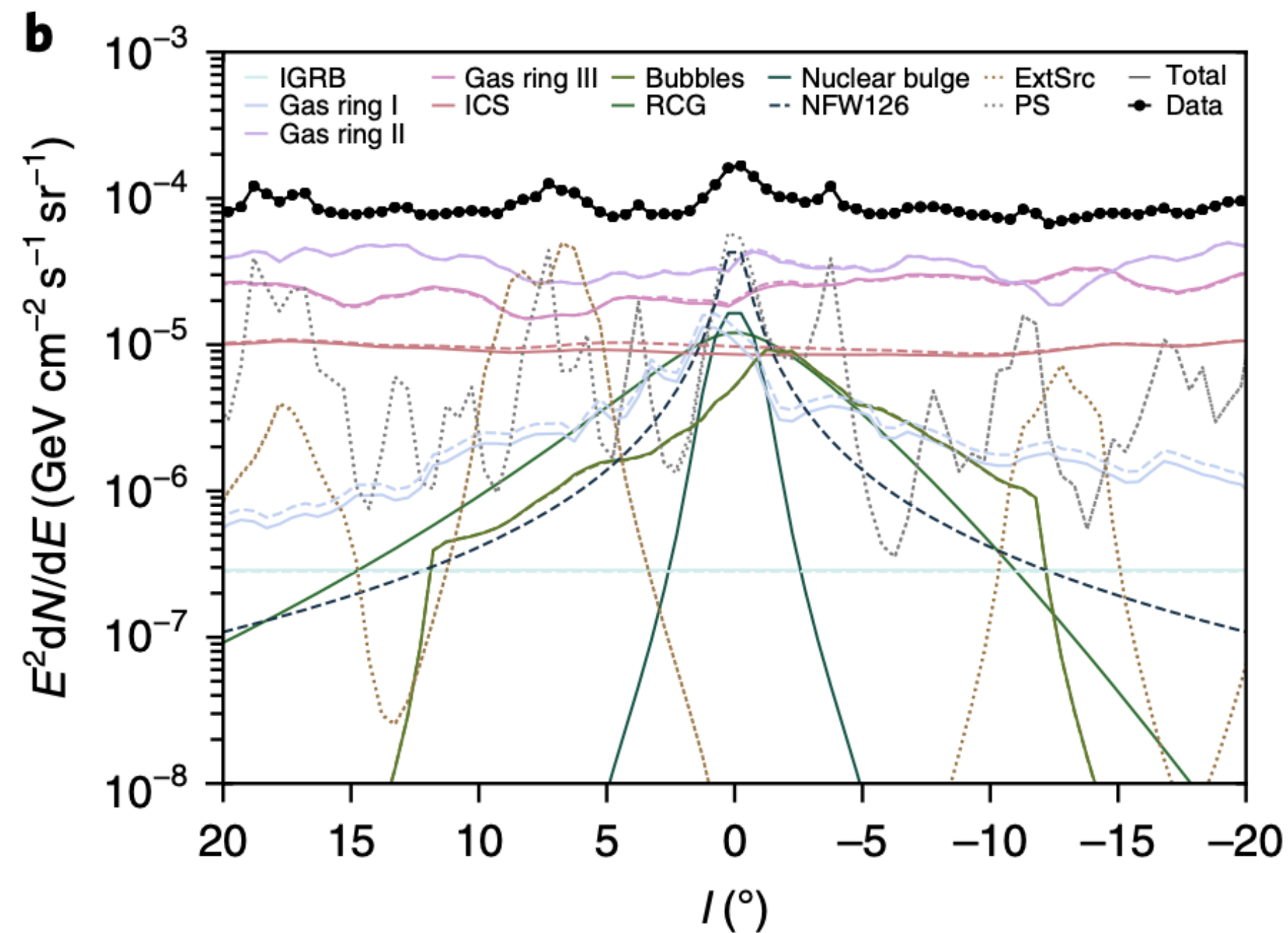
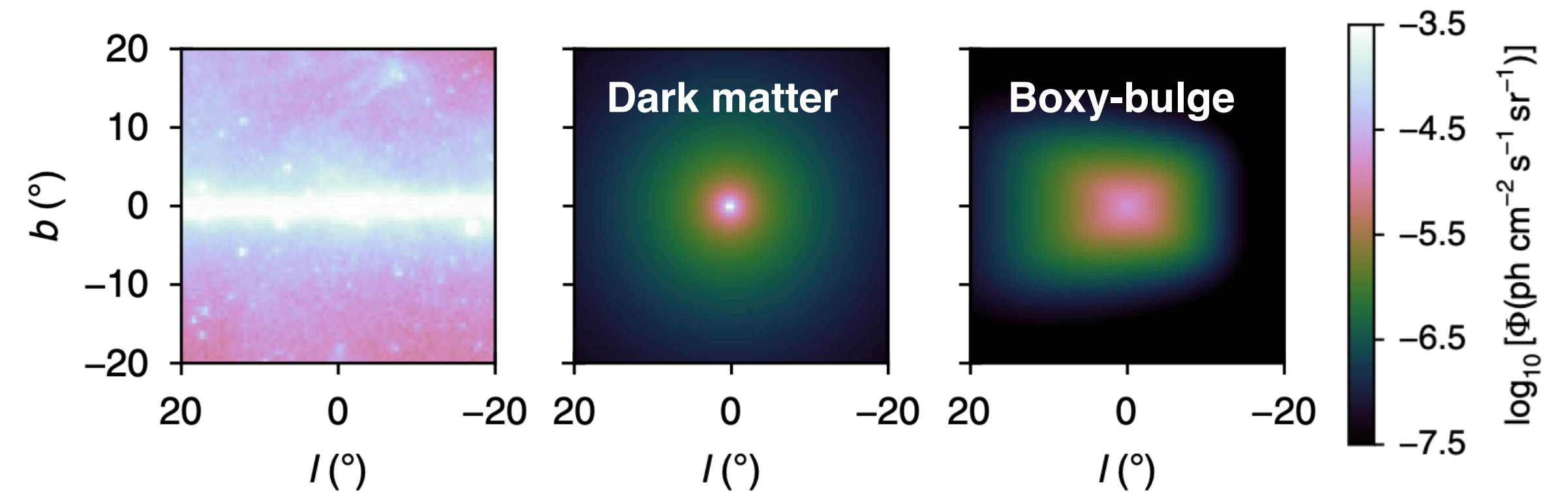
nature
astronomy

ARTICLES

<https://doi.org/10.1038/s41550-018-0531-z>

The Fermi-LAT GeV excess as a tracer of stellar mass in the Galactic bulge

Richard Bartels^{1*}, Emma Storm¹, Christoph Weniger¹ and Francesca Calore²



✓ **Stellar bulge model: Boxy bulge** as traced by red-clump giants + nuclear bulge

Cao+MNRAS'13; Launhardt+ A&A'02

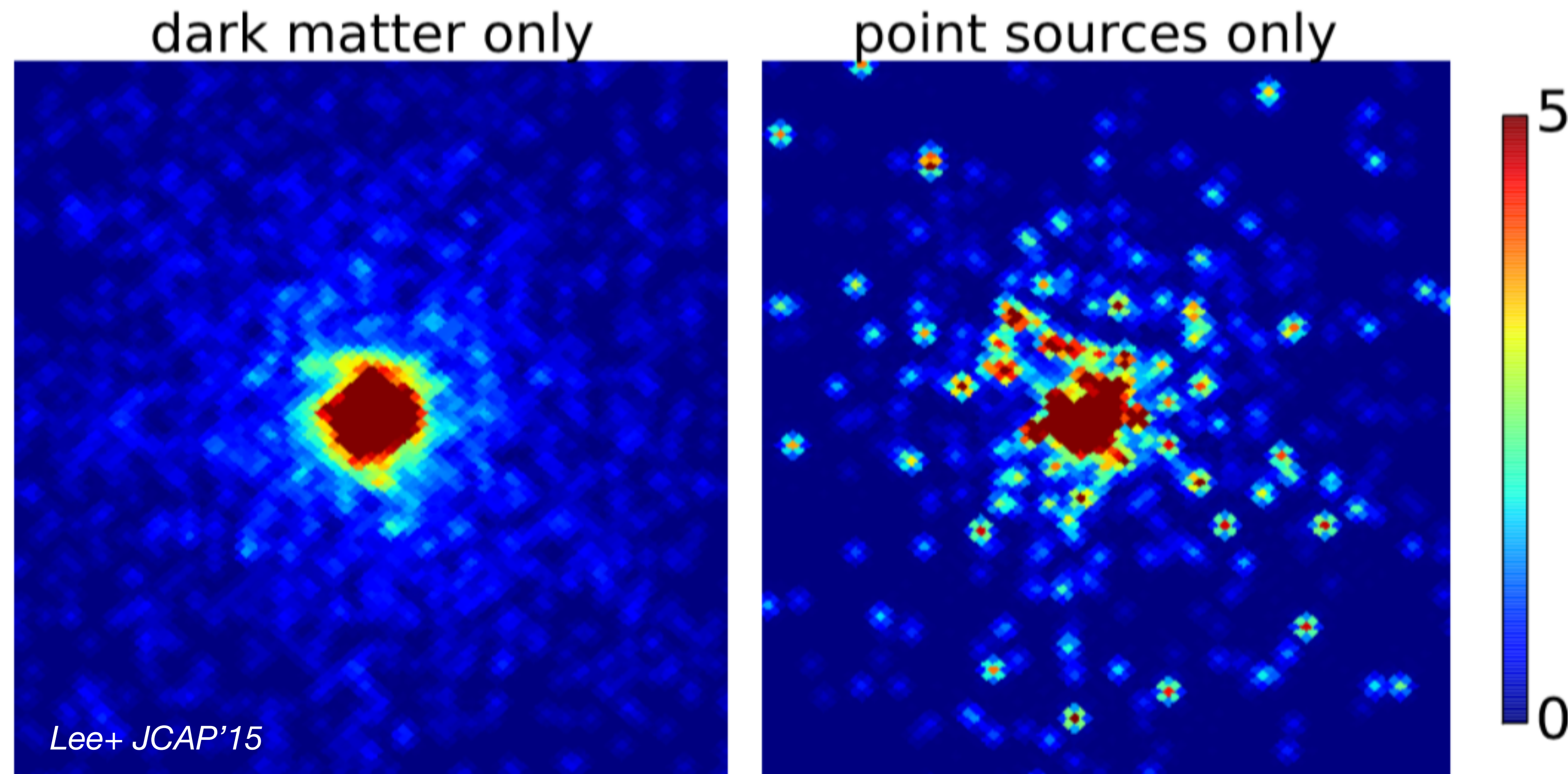
✓ Strong evidence for additional **stellar bulge model** (16σ); no evidence for additional **DM model** ($< 3\sigma$)

✓ Discriminating feature: Asymmetry at ~ 10 deg longitude => **Morphology** of the GCE **more oblate** than what found before

Macias+ Nature Astronomy'18; Macias+ JCAP'19

Statistics of photon counts

How to discriminate diffuse vs point-like emission

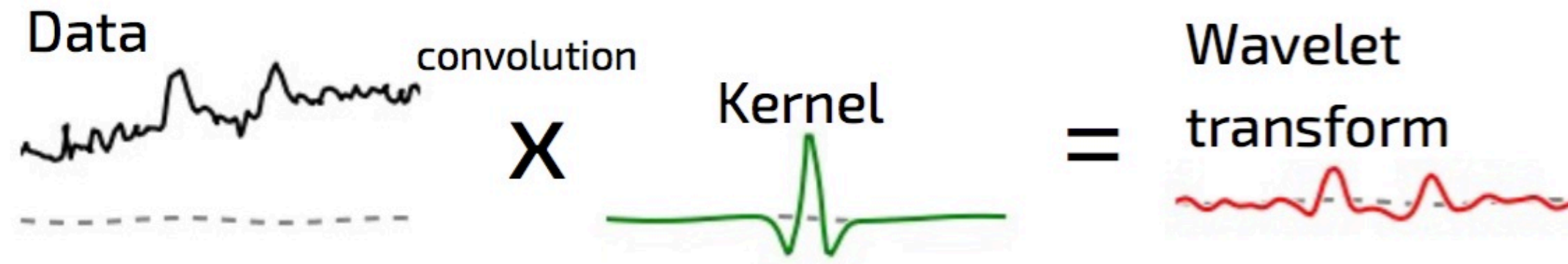


Differences in the **statistics of the photon counts** can be quantified and used for model comparison

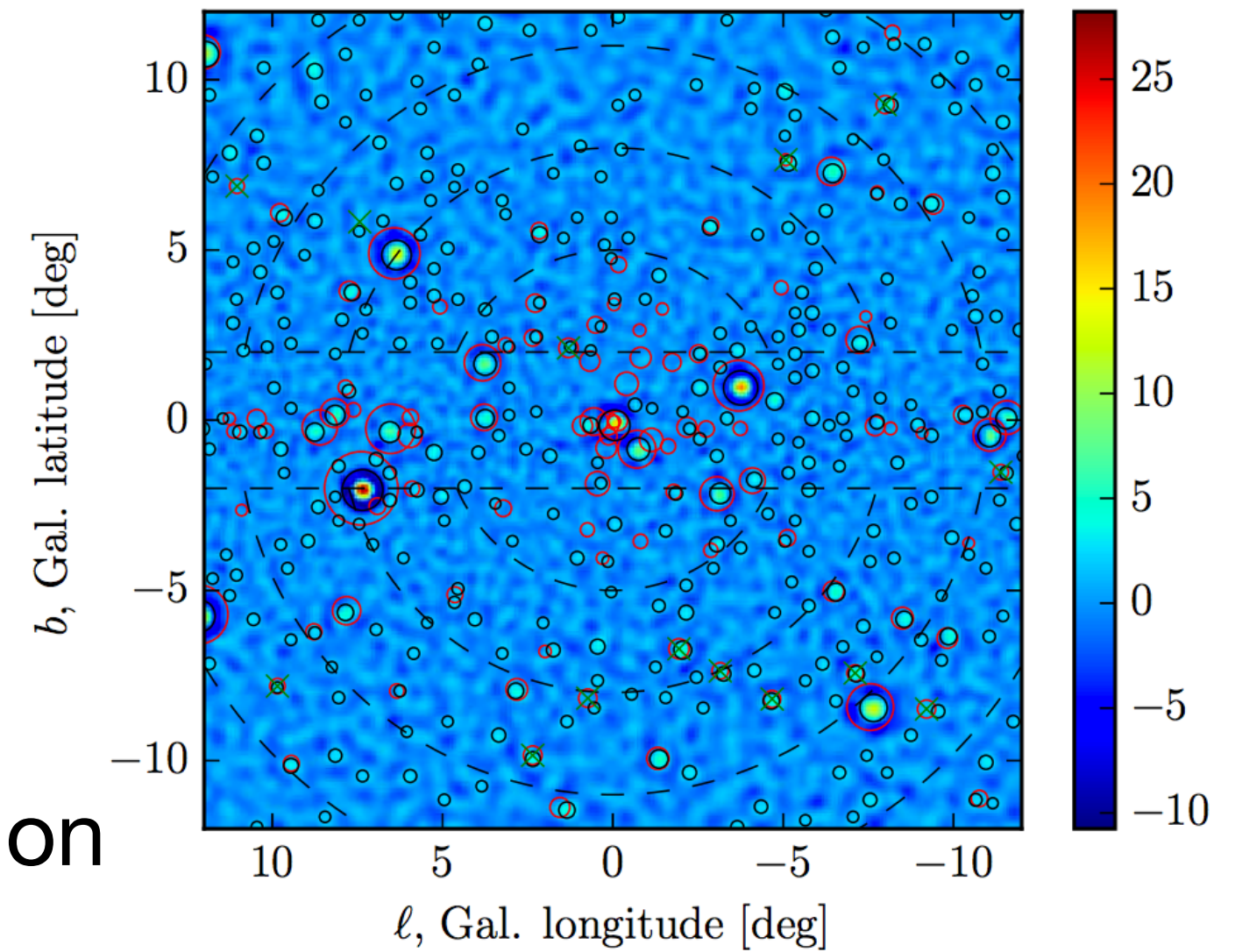
Support for unresolved point sources (PS)

Local maxima of normalised wavelet transform

Bartels+ PRL'16

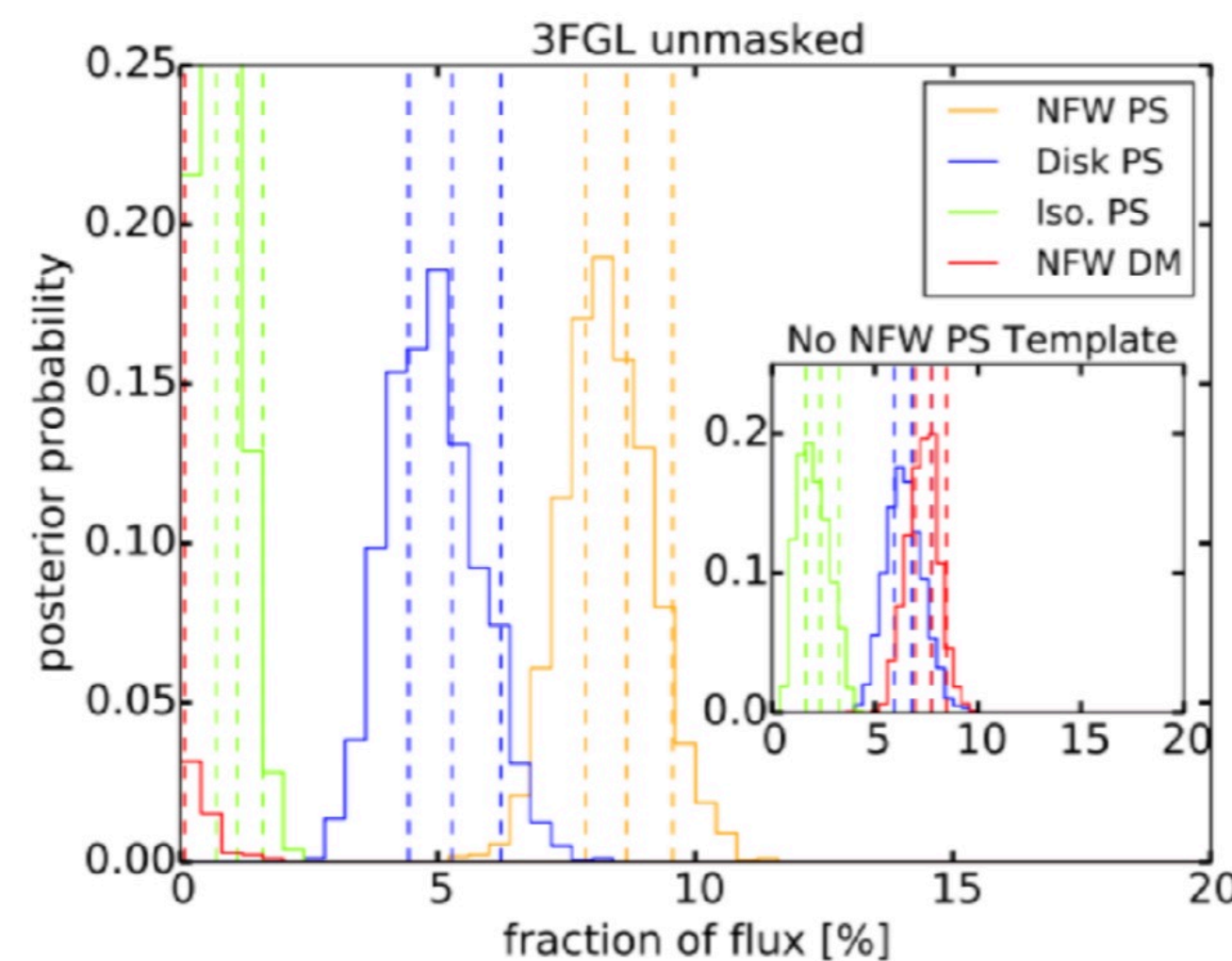
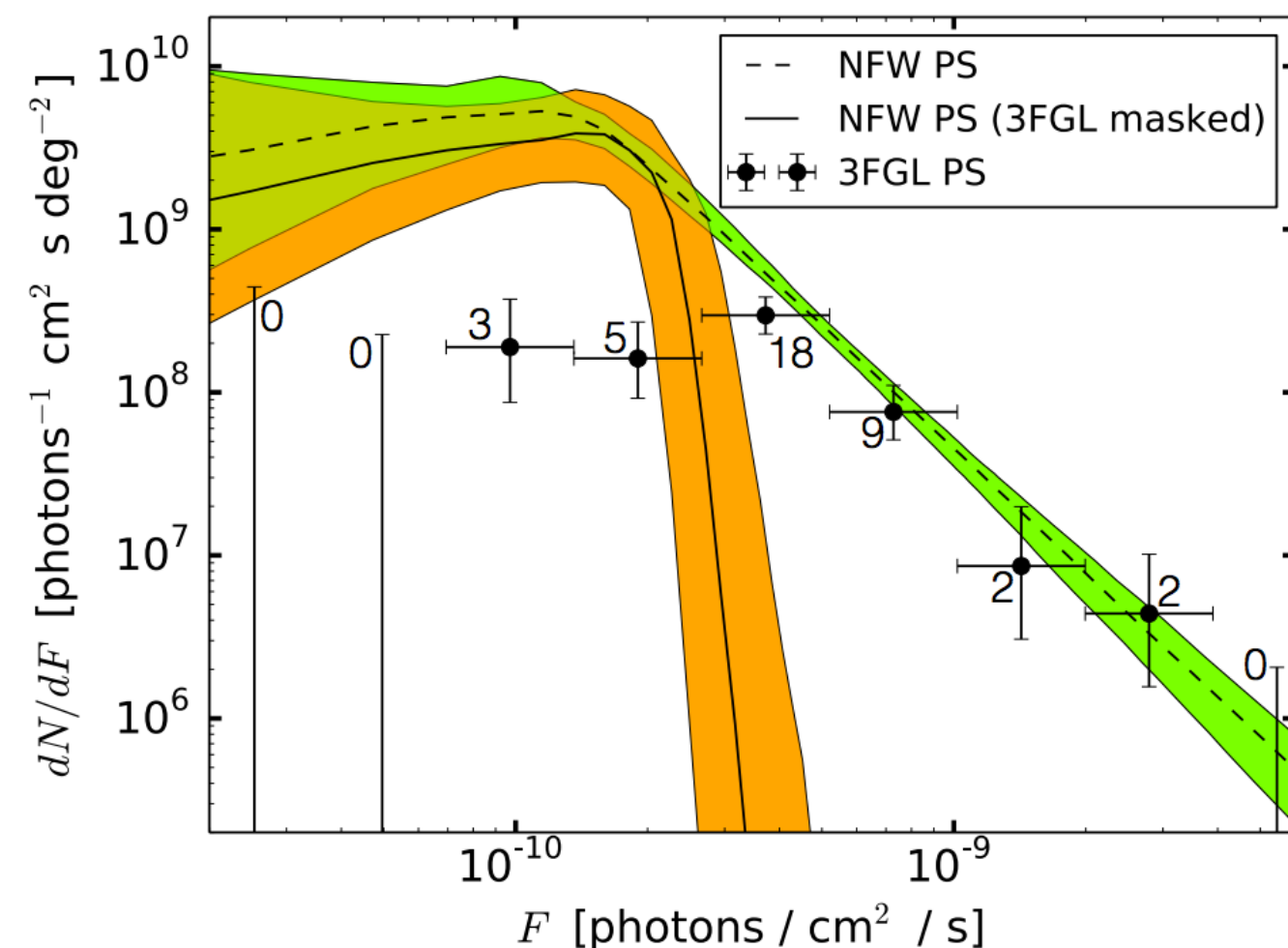


- Wavelet transform to look for **peaks** in data
- Enough peaks were found to explain the cumulative excess emission
- Evidence for unresolved PS population and constraints on luminosity function
- No modelling of diffuse emission required



Non-Poissonian template fitting

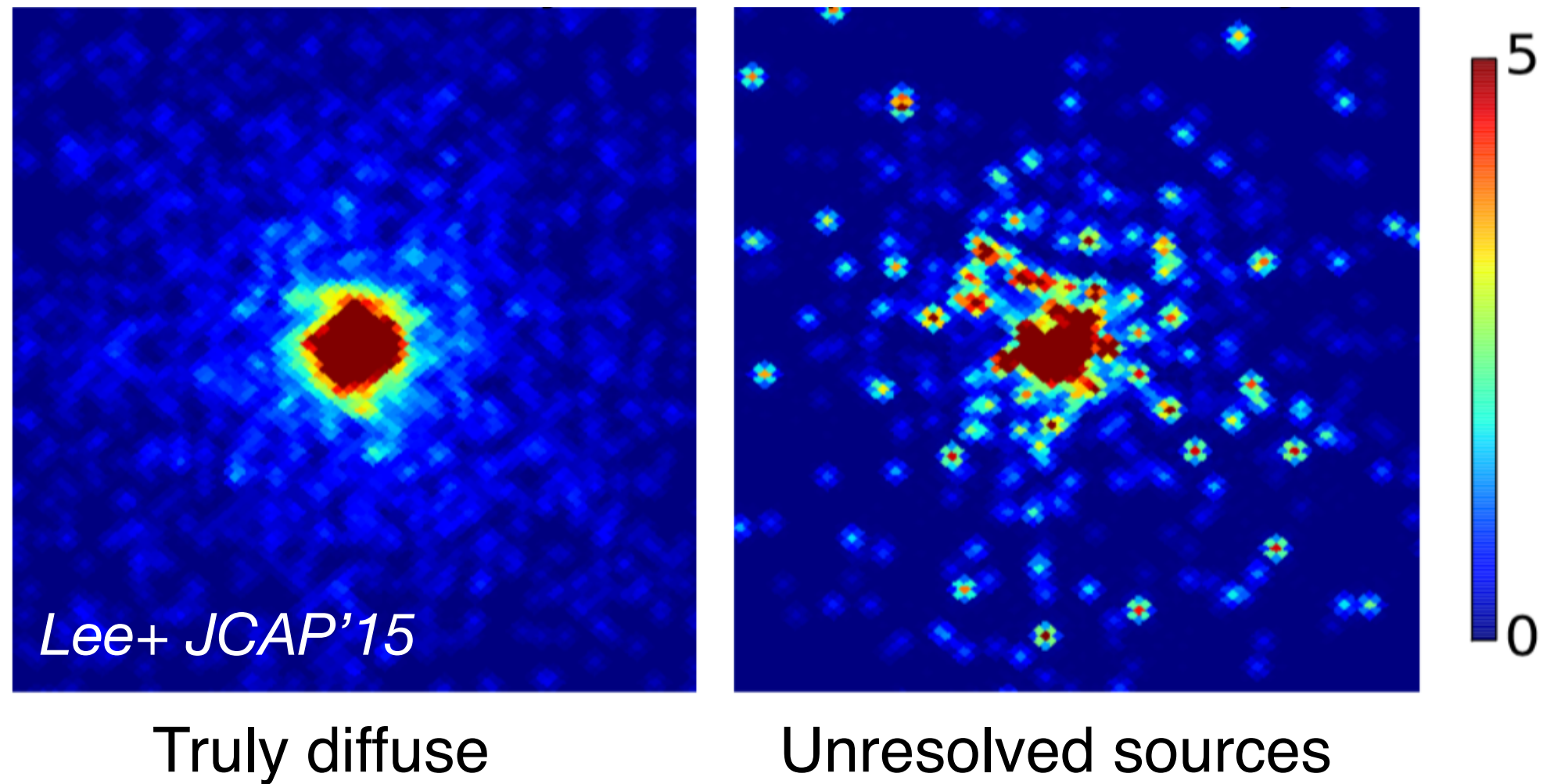
Lee+ PRL'16



- Exploits difference in photon statistics: smooth signal (DM) vs larger variance across pixels (PS)
- PS fluctuations follow non-Poissonian statistics
- Sensitivity to spatial distribution and luminosity function of PS
- Required modelling of diffuse emission

The GeV excess nature

The gamma-ray perspective



- Difference in **statistics of photon counts** can be quantified and used for model comparison

Bartels+ PRL'16; Lee+PRL'16

- **Strong bias** from mis-modelling of foreground diffuse emission and controversial results

Zhong+PRL'19; Leane&Slatyer PRL'20, PRD'20; Chang+ PRD'20, Buschmann+PRD'20

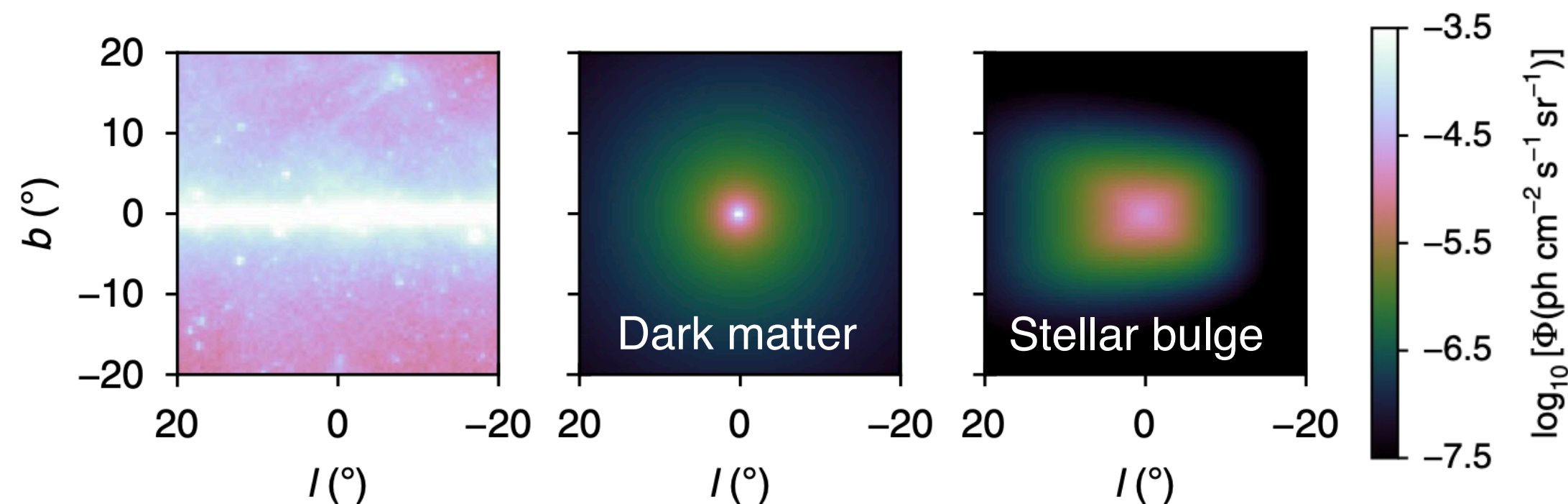
- Nonetheless: **evidence for unresolved point sources** is there with different, independent, methods

Buschmann+PRD'20; FC+ PRL'21; List+ PRL'21

- **Stellar bulge morphology preferred over DM** also when modelling faint point sources

FC+ PRL' 21; Manconi, FC+ PRD'24

Macias+ Nature Astronomy'18; Macias+ JCAP'19



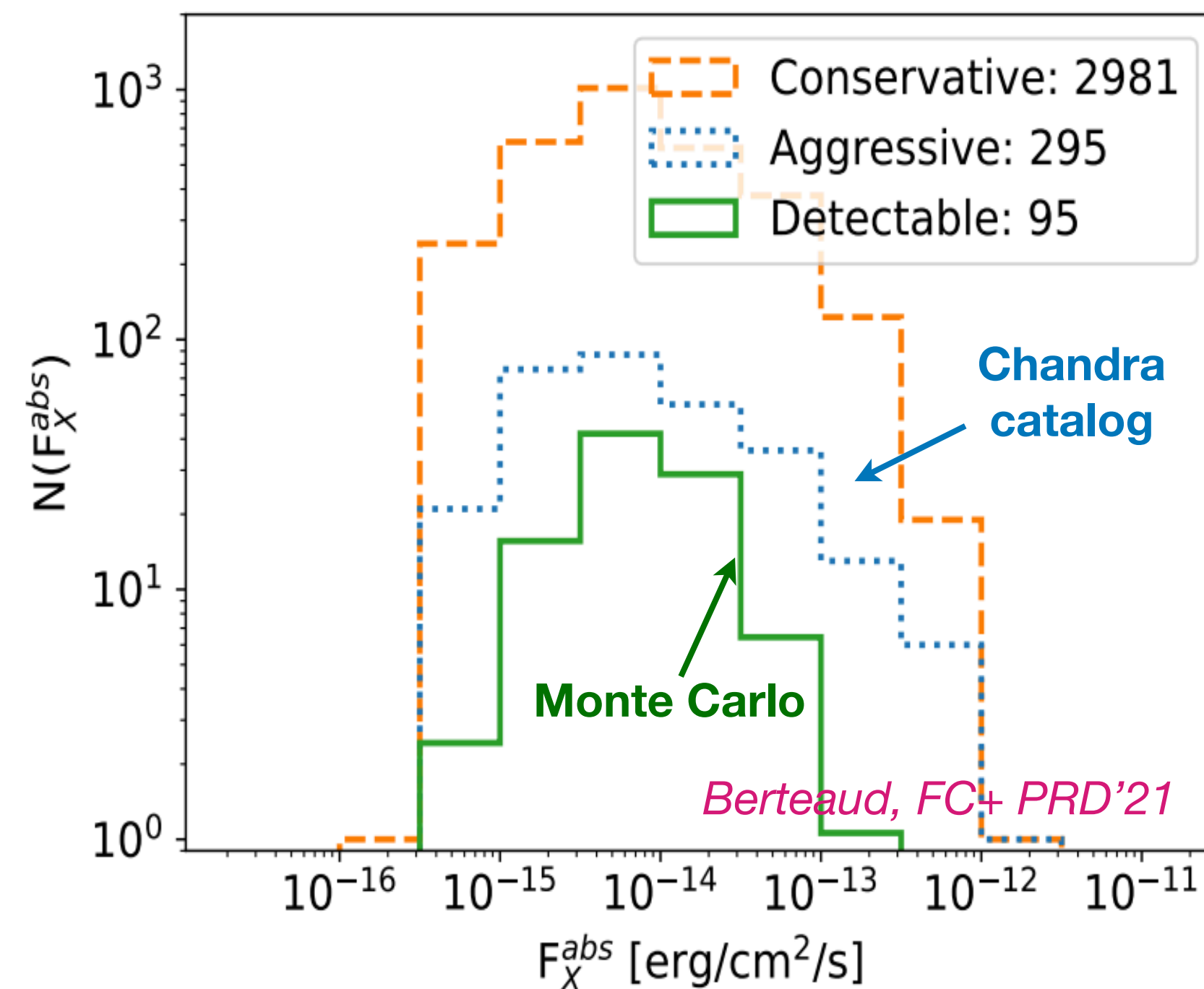
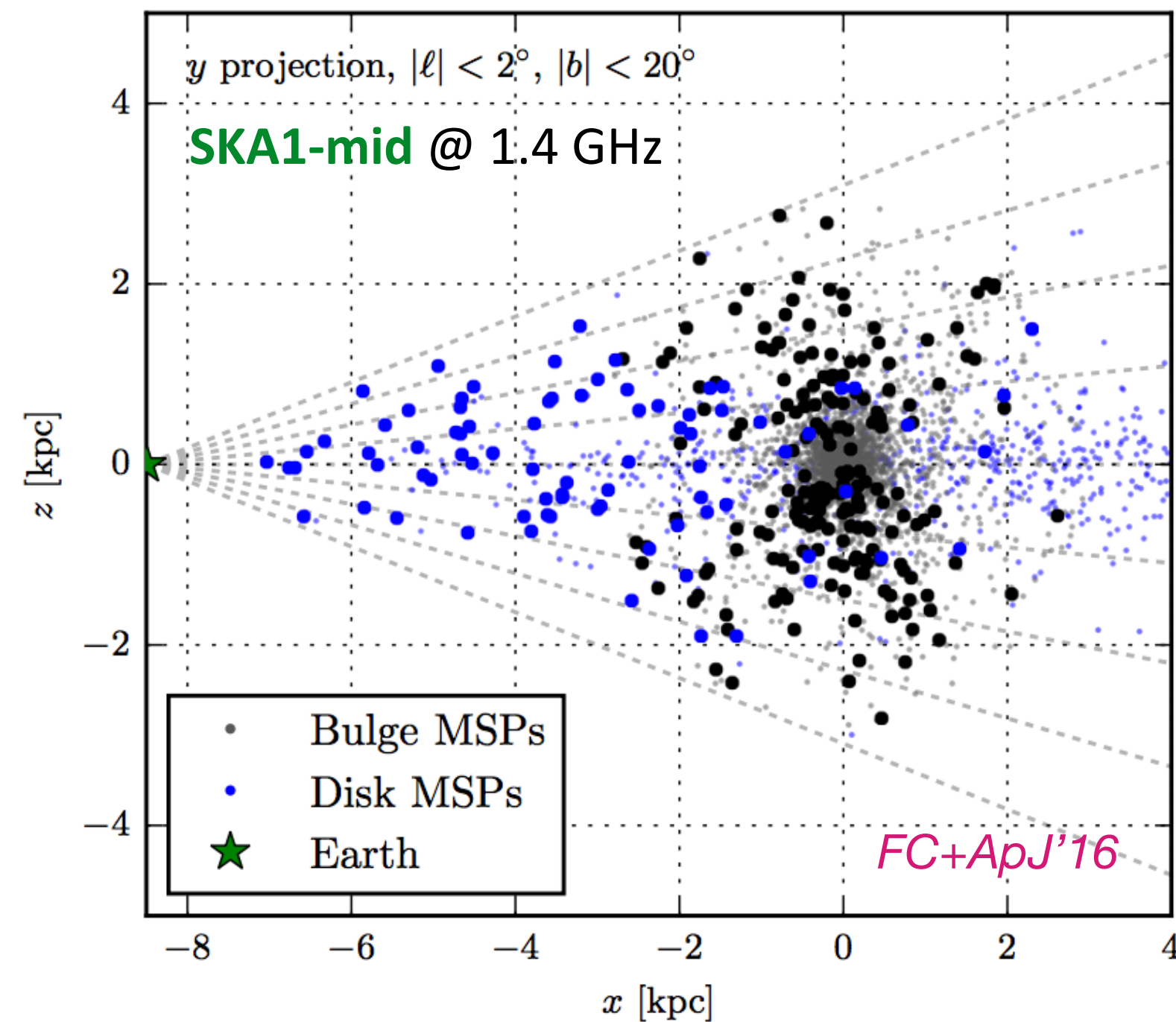
An (at least) partial **stellar origin of the GeV excess** seems to be confirmed

Multi-messenger tests of the GeV excess

Complementary techniques and **multi-wavelength searches** to test the excess nature:

- * Radio, X-ray, and (future) gravitational waves searches

FC+ApJ'16; FC+PRL'19; Berteaud, FC+ PRD'21



- * Very high-energy photons with CTA

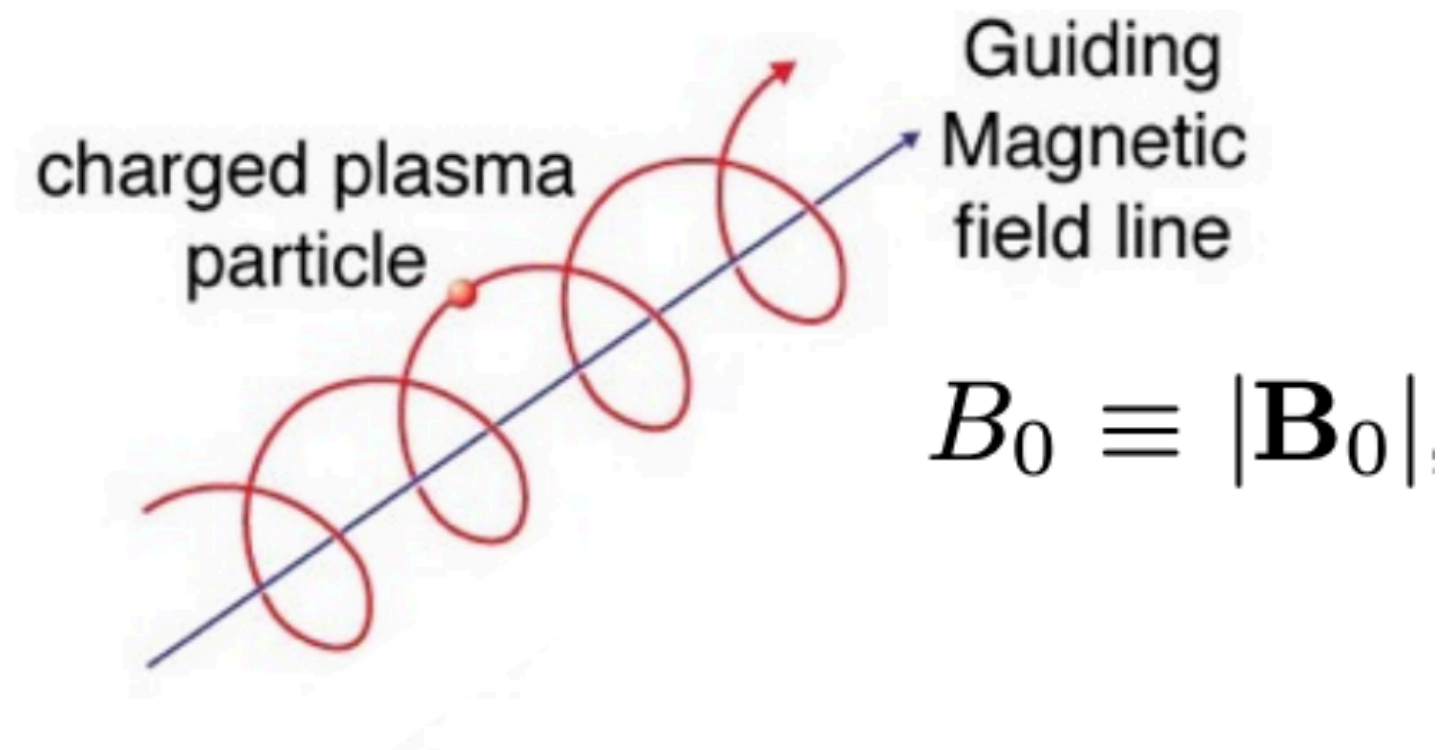
Macias+ MNRAS'21

- * DM constraints from gamma rays (dwarf galaxies) and cosmic-ray antiprotons

Di Mauro & Winkler PRD'21

Cosmic-ray in regular B-field

$$\frac{d\mathbf{p}}{dt} = q\mathbf{E} + q\frac{\mathbf{v}}{c} \times \mathbf{B}$$



Solutions (no electric field):

$$p_z = \text{const}$$

$$v_x = v_0 \cos(\Omega t)$$

$$v_y = v_0 \sin(\Omega t)$$

$$r_L = \gamma r_g = \sqrt{1 - \mu^2} \frac{\mathcal{R}}{B_0} \simeq 10^{-6} \sqrt{1 - \mu^2} \frac{\mathcal{R}}{\text{GV}} \frac{\mu\text{G}}{B_0} \text{pc}$$

Larmor radius

$$\Omega = \frac{qB_0}{E} \simeq 10^{-2} \frac{B_0}{\mu\text{G}} \frac{\text{GeV}}{E} \text{rad/s}$$

Larmor frequency

$$\mu = p_z/p$$

$$\mathcal{R} = p/q [\text{GV}]$$

Cosmic-ray in perturbed B-field

Small-scale stochastic perturbations

$$|\delta\mathbf{B}| \ll |\mathbf{B}_0| \text{ and } \delta\mathbf{B} \perp \mathbf{B}_0$$

$$\frac{d\mathbf{p}}{dt} = q \frac{\mathbf{v}}{c} \times (\mathbf{B} + \delta\mathbf{B})$$

Changes x and y component
of the momentum

Changes only the
direction of p_z

$$\mu = p_z/p$$

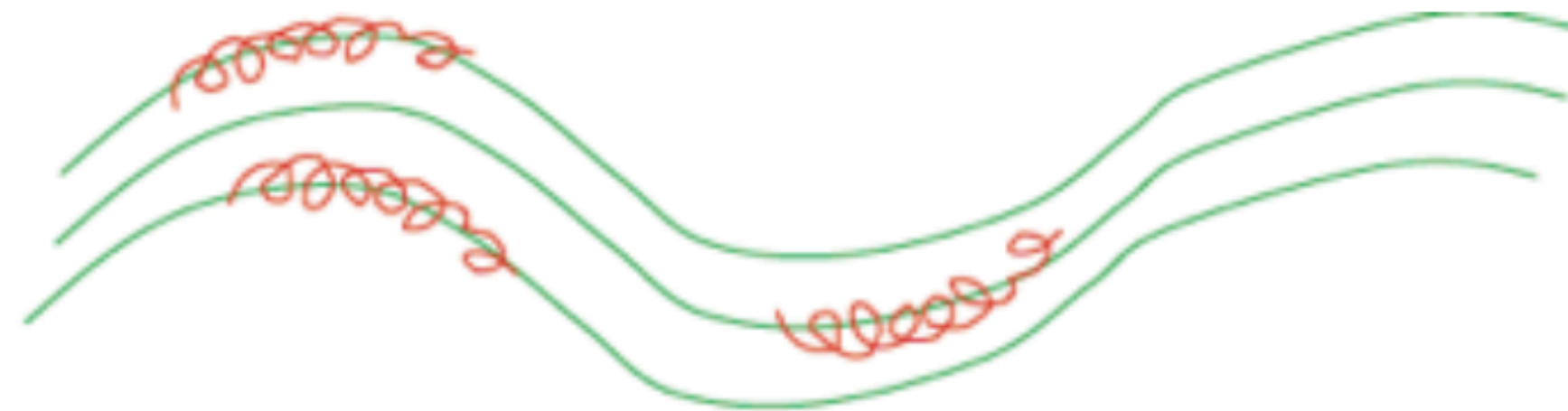
Diffusive process

Cosmic-ray in perturbed B-field

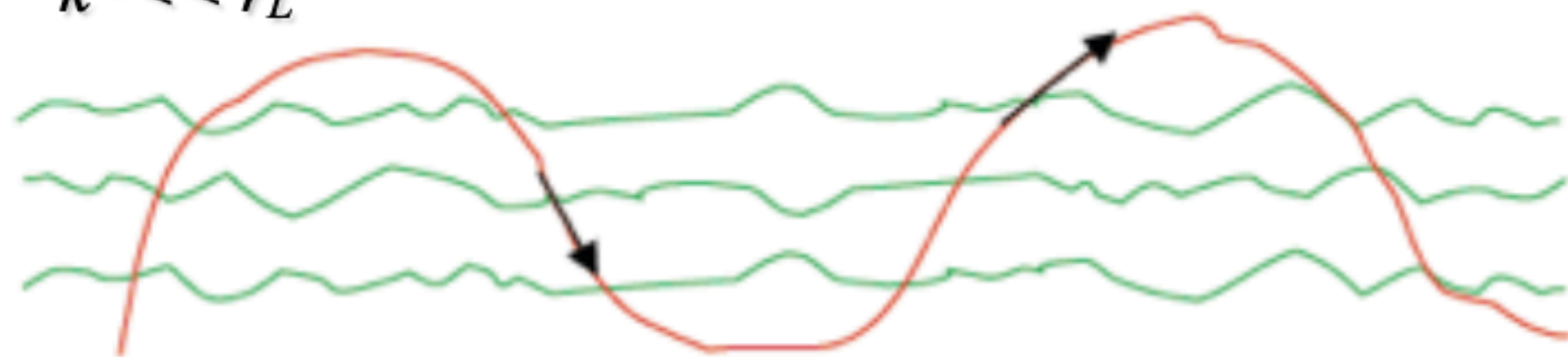
Small-scale (k) stochastic perturbations

$$|\delta\mathbf{B}| \ll |\mathbf{B}_0| \text{ and } \delta\mathbf{B} \perp \mathbf{B}_0$$

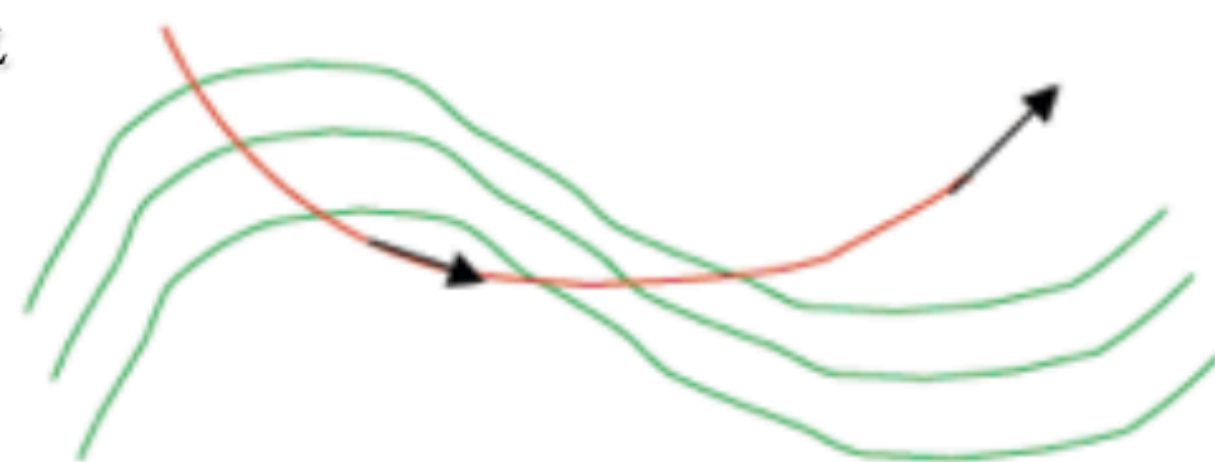
$k^{\parallel} \gg r_L$



$k^{\parallel} \ll r_L$



$k^{\parallel} \sim r_L$



Collisionless Diffusion

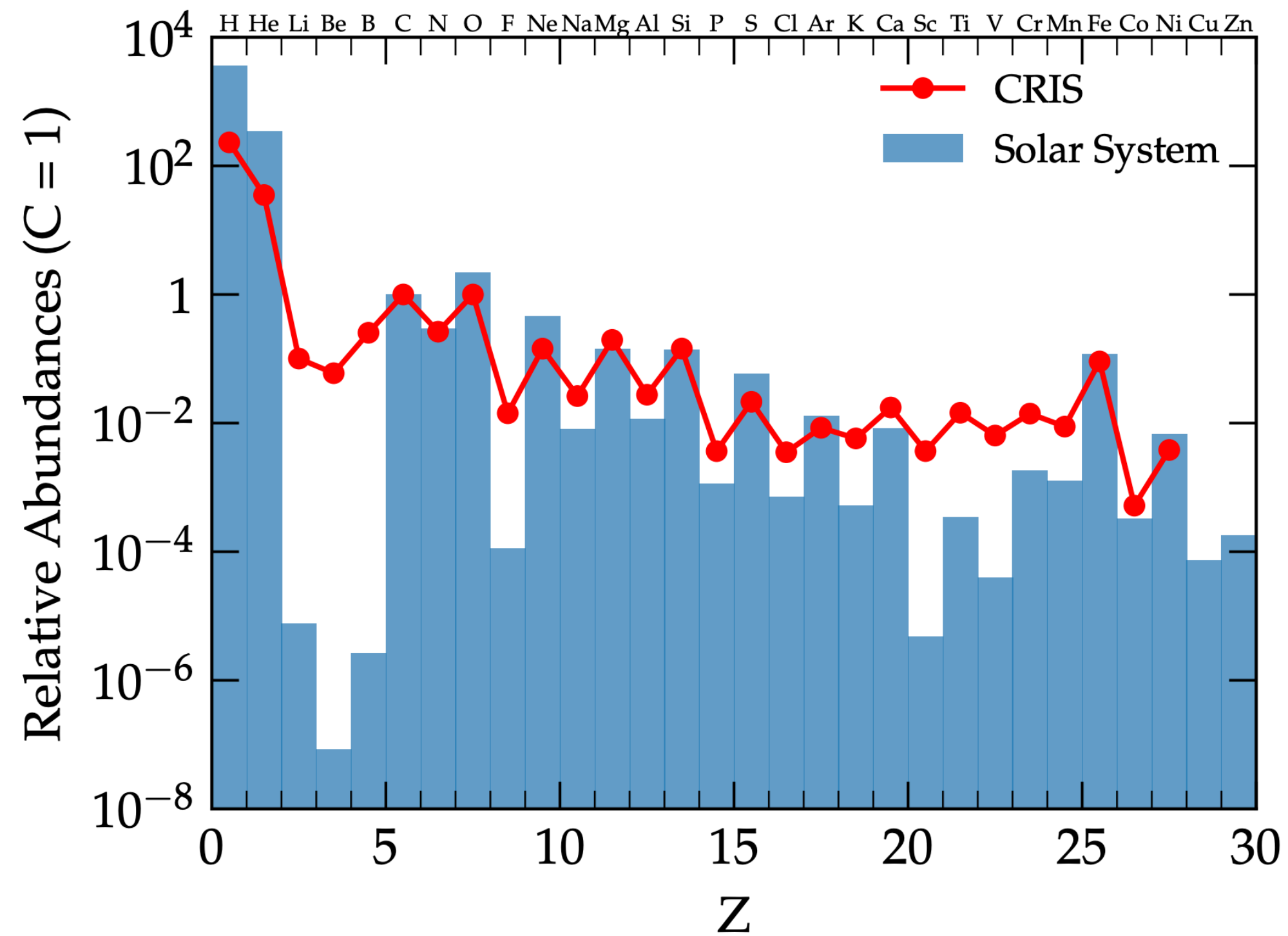
[i.e. scattering on inhomogeneities of the magnetic field]

Diffusion coefficient K

Describes the random change of the pitch angle

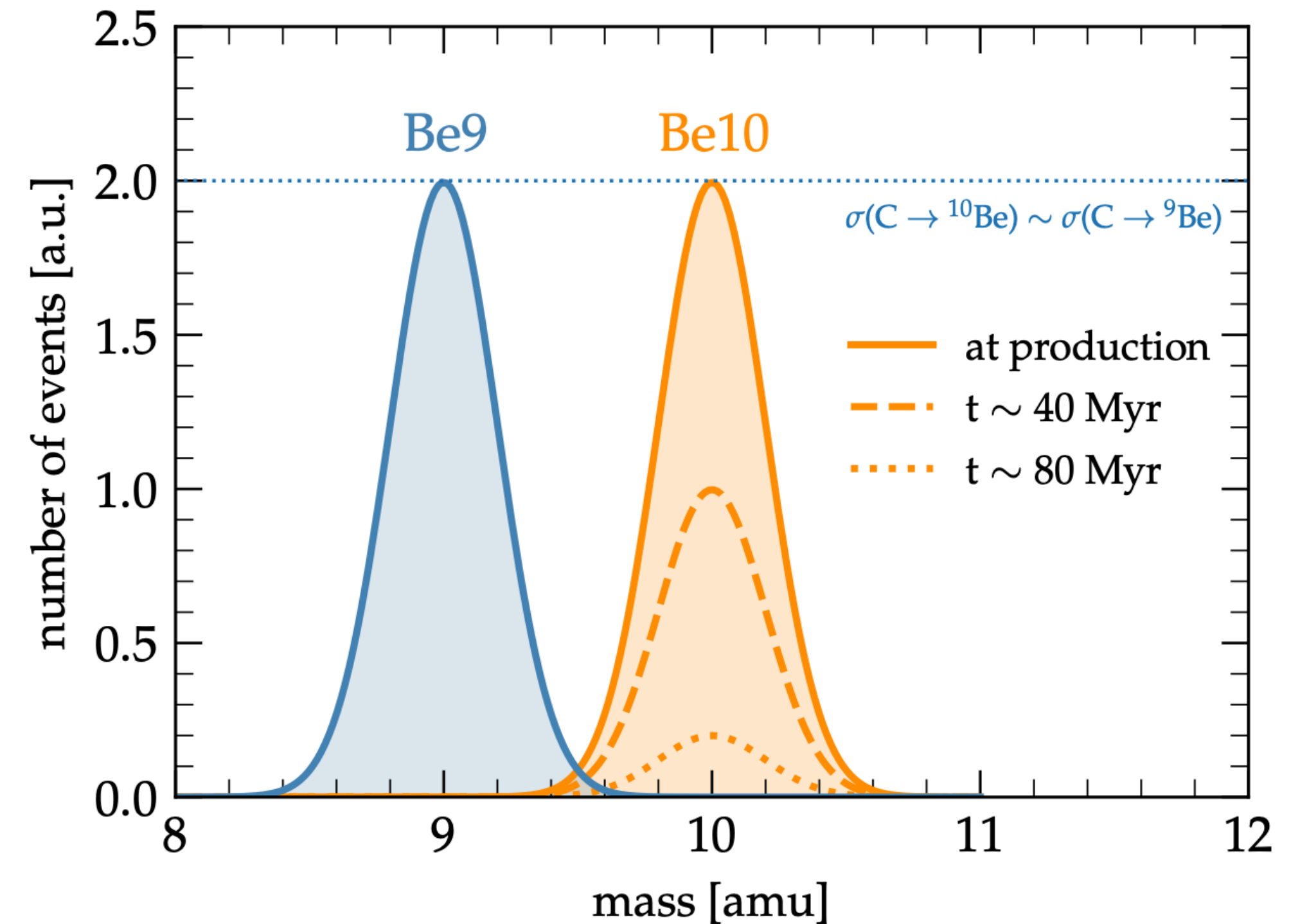
Composition of Galactic CR

Basic indicators of diffuse transport



STABLE ELEMENTS

Besides primary species, produced in stellar nucleosynthesis, the average interstellar medium hosts a population of secondary particles produced by **primary fragmentation** during propagation (spallation)



UNSTABLE ELEMENTS

${}^{10}\text{B}$ is beta unstable with half-life time of 1.5 Myr
 Production rate similar to other stable nuclei (${}^9\text{Be}$)
 The ratio ${}^9\text{Be}/{}^{10}\text{Be}$ can be used as a CR clock => residence time in the Galaxy of $O(100)$ Myr => **DIFFUSIVE** propagation

Particle DM: v-dependent x-section

$$\langle \sigma v \rangle = a + bv^2 + \mathcal{O}(v^4), \quad v/c \sim 10^{-3}$$

Non-relativistic regime:
if present, **s-wave** is dominant

- S-wave can be suppressed (e.g. helicity suppression) or models may allow for v-dependent cross sections (long-range interactions for TeV scale dark matter)
- The connection between Early Universe and today annihilation is altered in a non-trivial way

$$\langle \sigma v \rangle \equiv S(v/c) \times \langle \sigma v \rangle_0$$

$$S(v/c) = (v/c)^n$$

- n=-1: **Sommerfeld**-enhanced annihilation in the Coulomb limit
- n=0: **s-wave** velocity-independent annihilation
- n=2: **p-wave** annihilation. This scenario is relevant if DM is a Majorana fermion, which annihilates to Standard Model fermion/antifermion pairs

Particle DM: v-dependent x-section

Gamma-ray flux

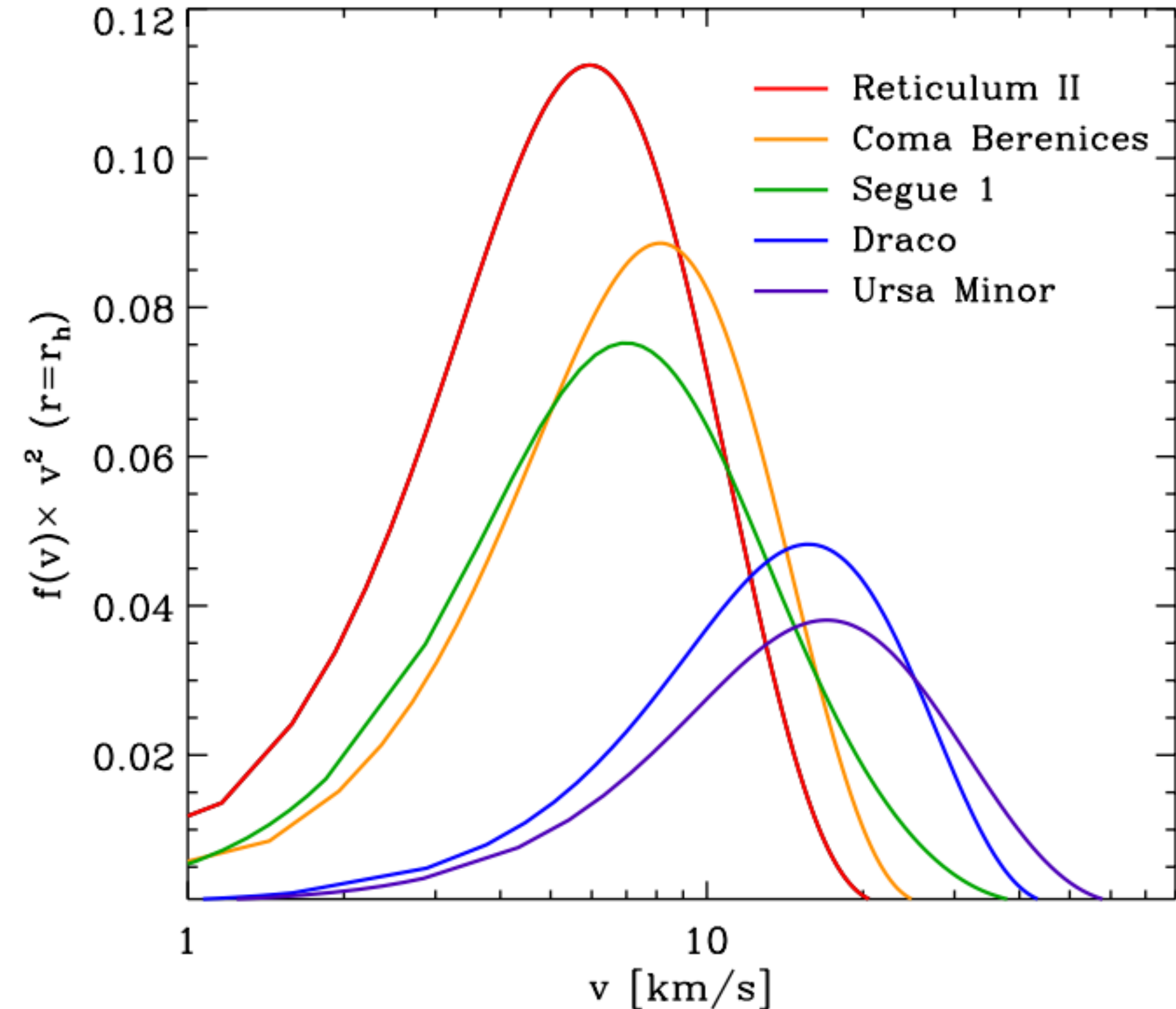
$$\langle \sigma v \rangle \equiv S(v/c) \times \langle \sigma v \rangle_0$$

$$\frac{d\Phi_\gamma}{dE_\gamma} = \frac{\langle \sigma v \rangle_0}{2m_{\text{DM}}^2} \sum_i B_i \frac{dN_\gamma^i}{dE_\gamma} \times \frac{1}{4\pi} \int_0^{\Delta\Omega} d\Omega \int_{\text{l.o.s}} ds \int d^3v_1 f(r(s, \Omega), \mathbf{v}_1) \int d^3v_2 f(r(s, \Omega), \mathbf{v}_2) S(|\mathbf{v}_1 - \mathbf{v}_2|/c)$$

DM phase-space distribution

$$\rho(r) \equiv \int f(r, \mathbf{v}) d^3v$$

$$f(\mathbf{v}) \equiv \int f(r, \mathbf{v}) dr$$



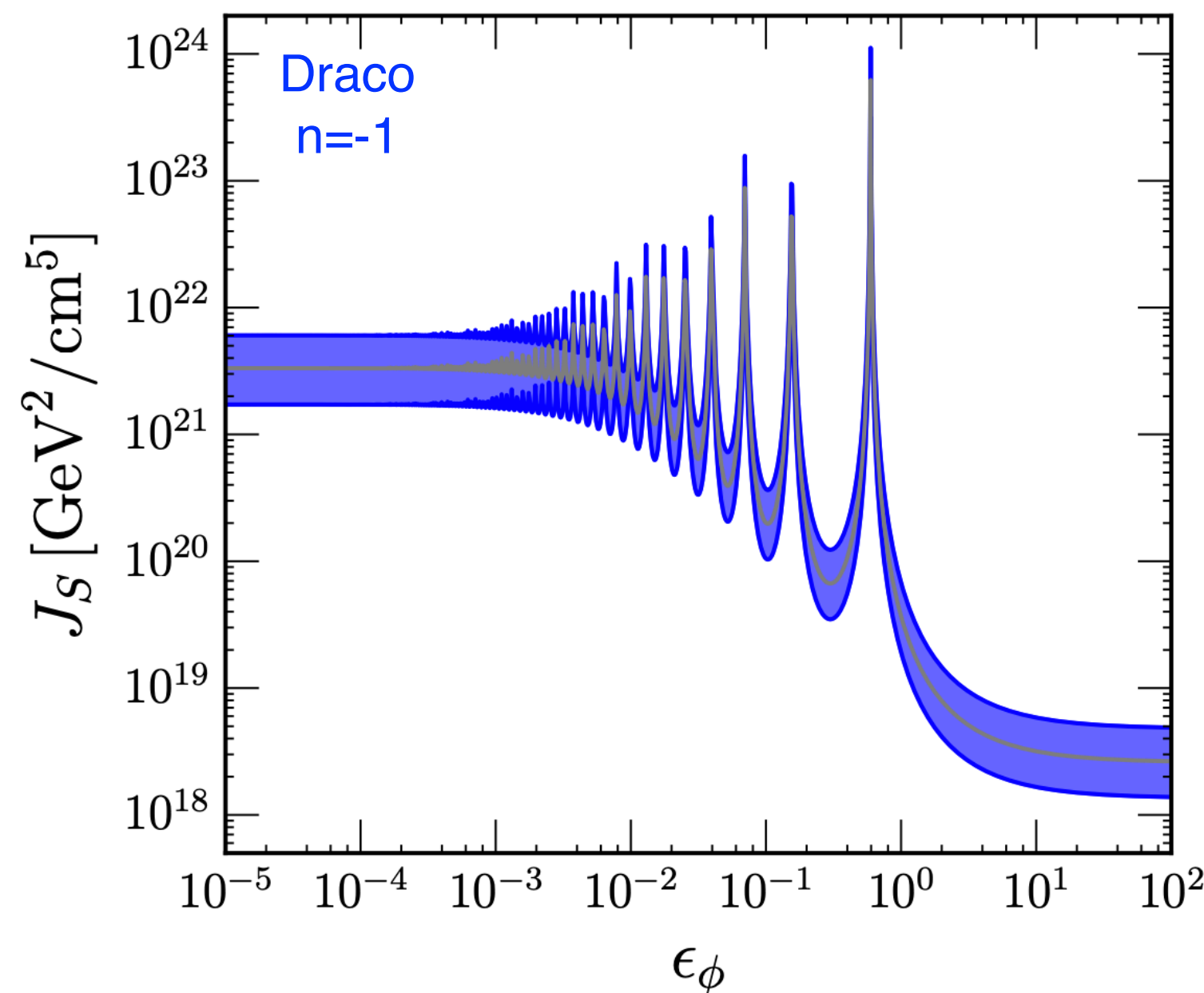
Particle DM: v-dependent x-section

Gamma-ray flux

$$\langle \sigma v \rangle \equiv S(v/c) \times \langle \sigma v \rangle_0$$

$$\frac{d\Phi_\gamma}{dE_\gamma} = \frac{\langle \sigma v \rangle_0}{2m_{\text{DM}}^2} \sum_i B_i \frac{dN_\gamma^i}{dE_\gamma} \times \frac{1}{4\pi} \int_0^{\Delta\Omega} d\Omega \int_{\text{l.o.s}} ds \int d^3v_1 f(r(s, \Omega), \mathbf{v}_1) \int d^3v_2 f(r(s, \Omega), \mathbf{v}_2) S(|\mathbf{v}_1 - \mathbf{v}_2|/c)$$

$$J_S(\Delta\Omega) \xrightarrow{S=1} J(\Delta\Omega) \int_0^{\Delta\Omega} d\Omega \int_{\text{l.o.s}} \rho(r(s, \Omega)) \times \rho(r(s, \Omega)) ds$$



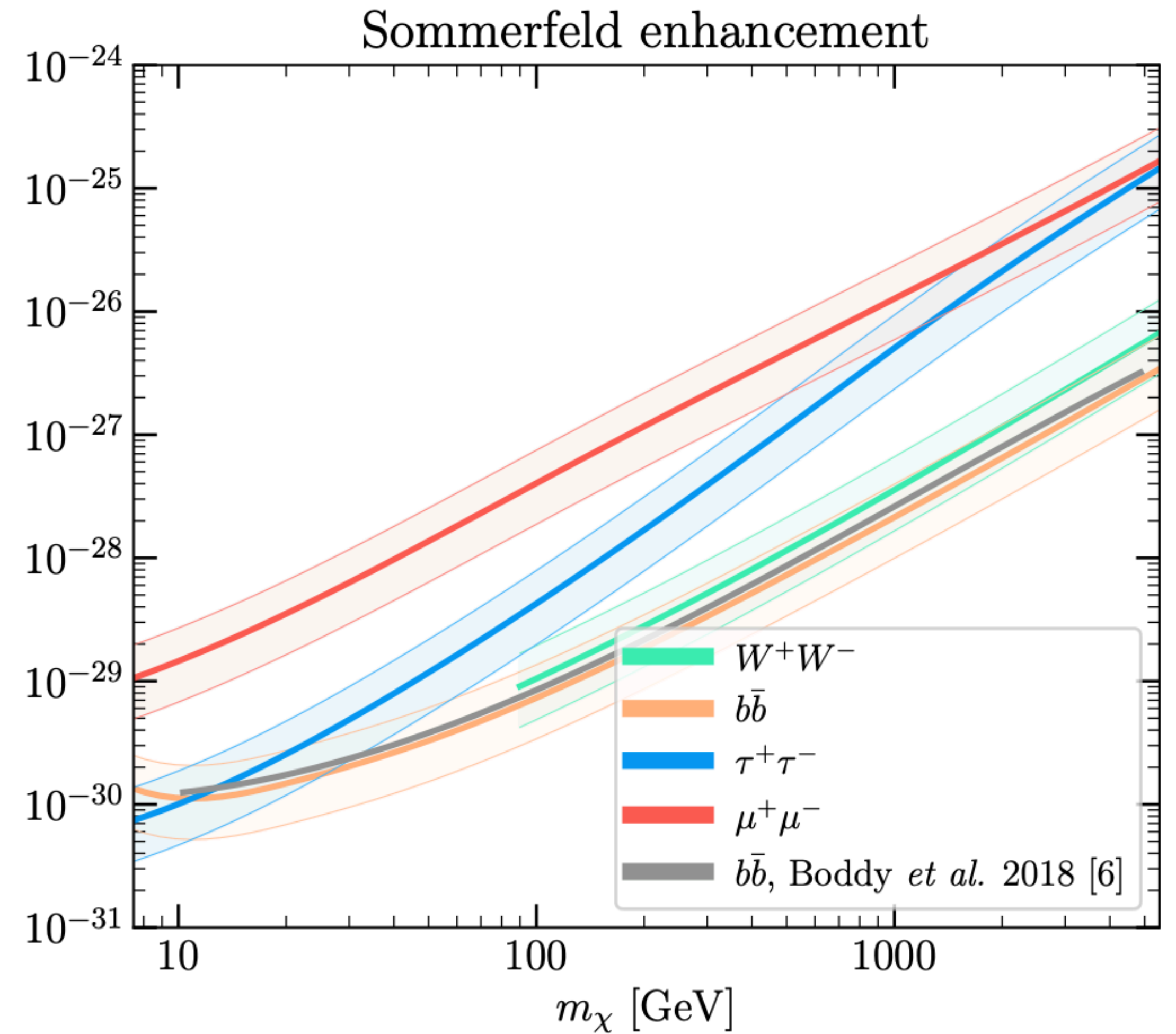
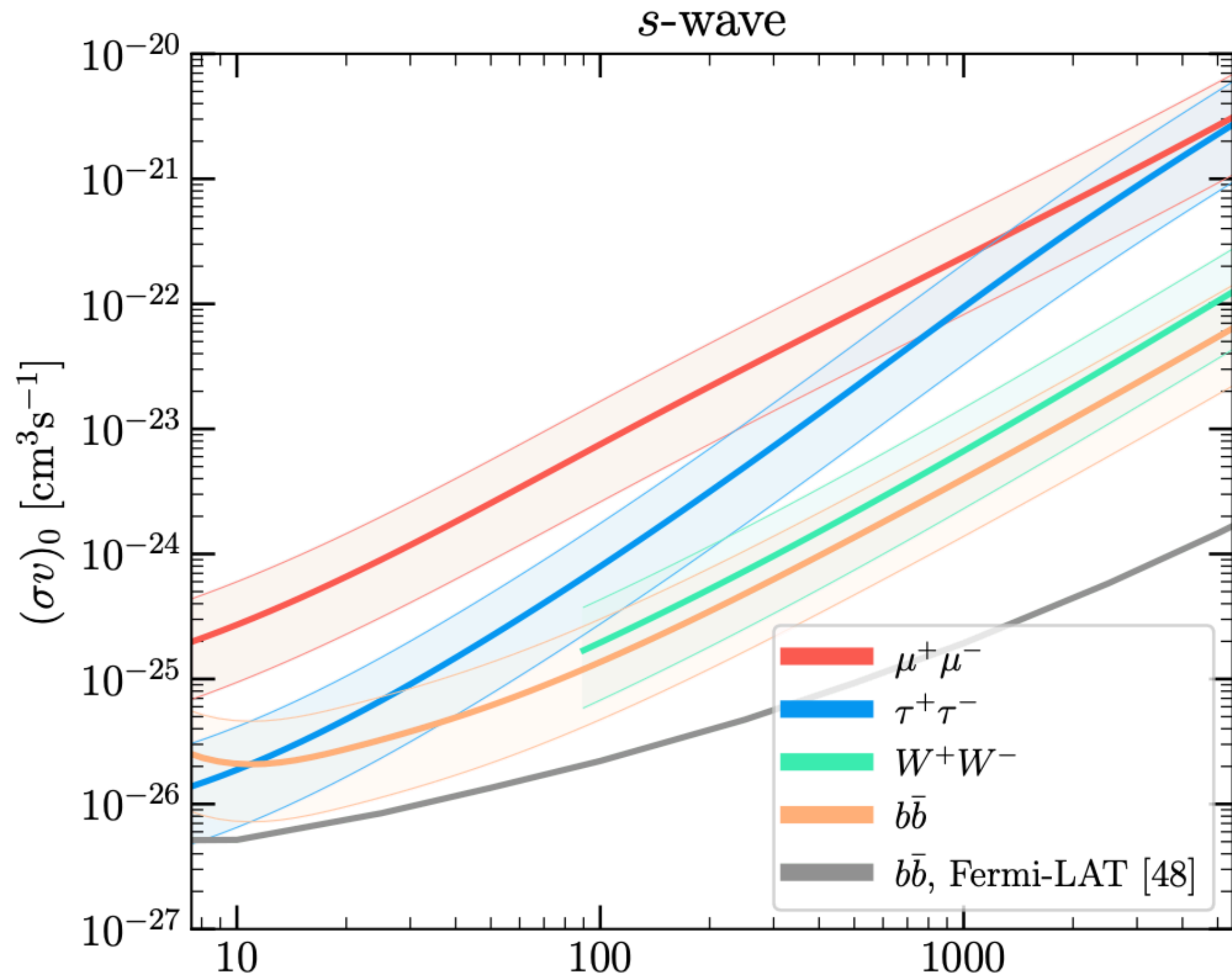
v-dependence of xsec can be translated in v-dependence of J-factors, allowing an easier **recasting** of limits under s-wave assumptions

Boddy et al., Phys. Rev. D 95, 123008 (2017) [1702.00408]

Boddy et al., Phys. Rev. D 102, 023029 (2020) [1909.13197]

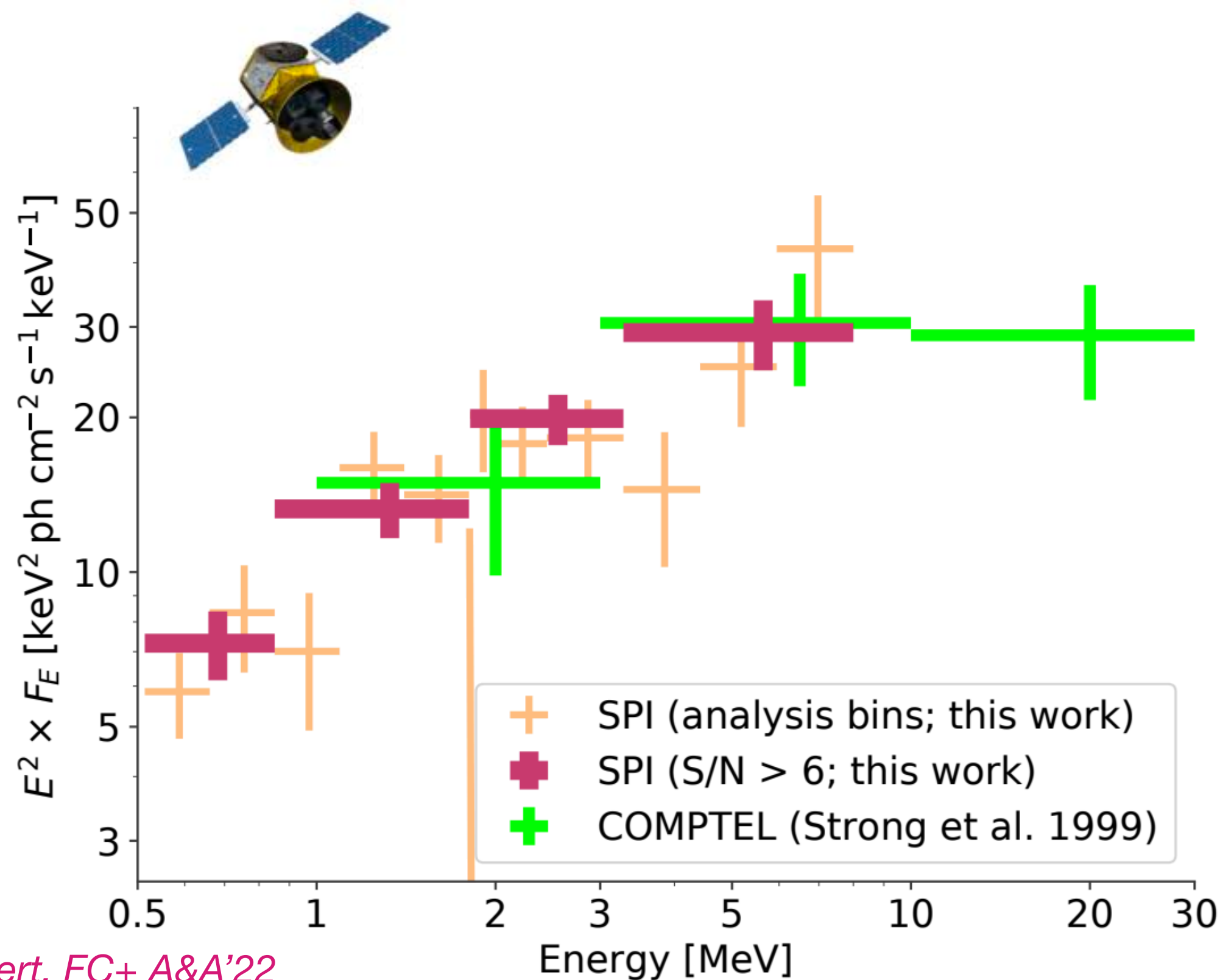
Particle DM: v -dependent x -section

Impact on DM limits



Boddy *et al.*, *Phys. Rev. D* 102, 023029 (2020) [1909.13197]

The MeV diffuse Galactic emission



Siegert, FC+ A&A'22

Constraints on cosmic-ray transport at MeV energy but also on exotic emission mechanisms for the first-time in a self-consistent framework

Berteaud, FC+ PRD'22 ; FC+ MNRAS'23

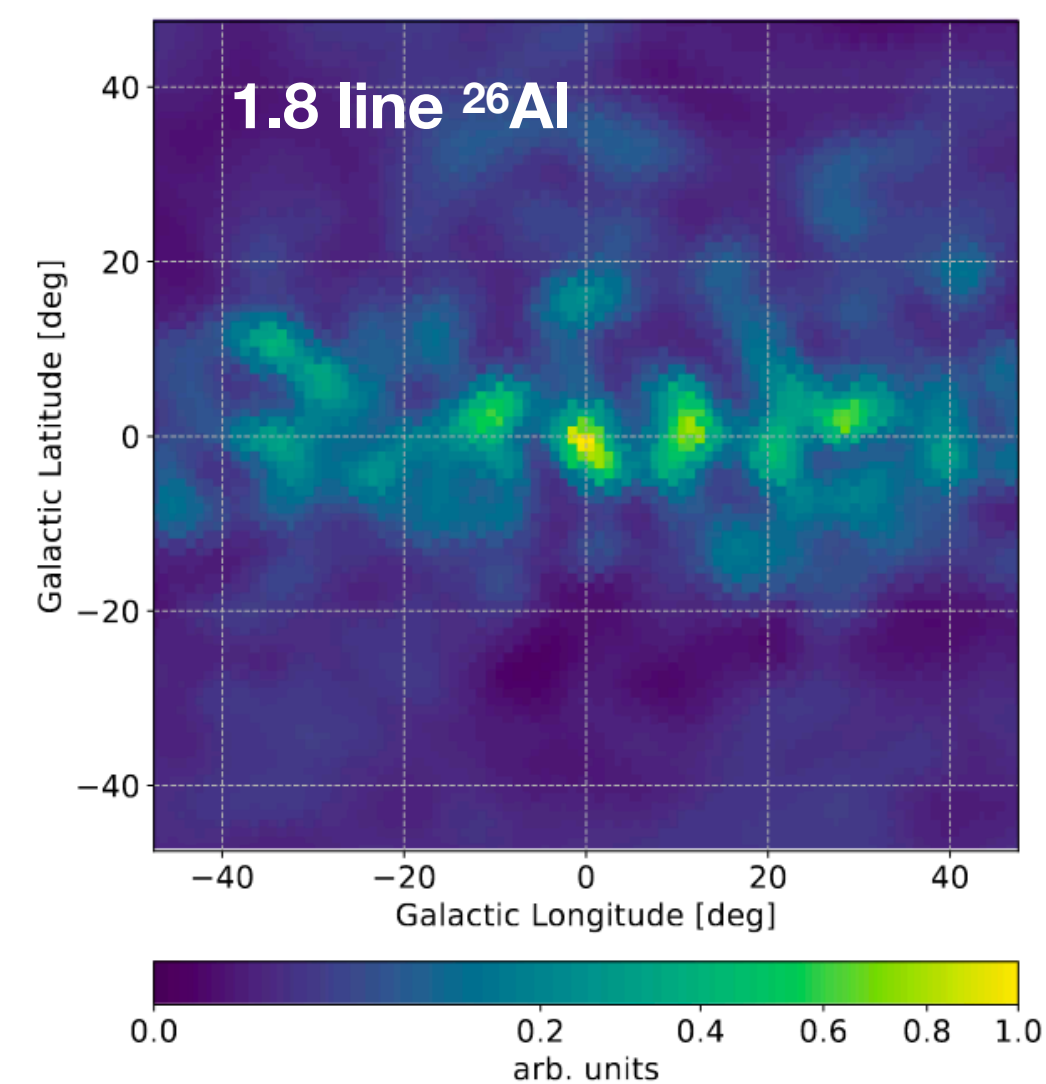
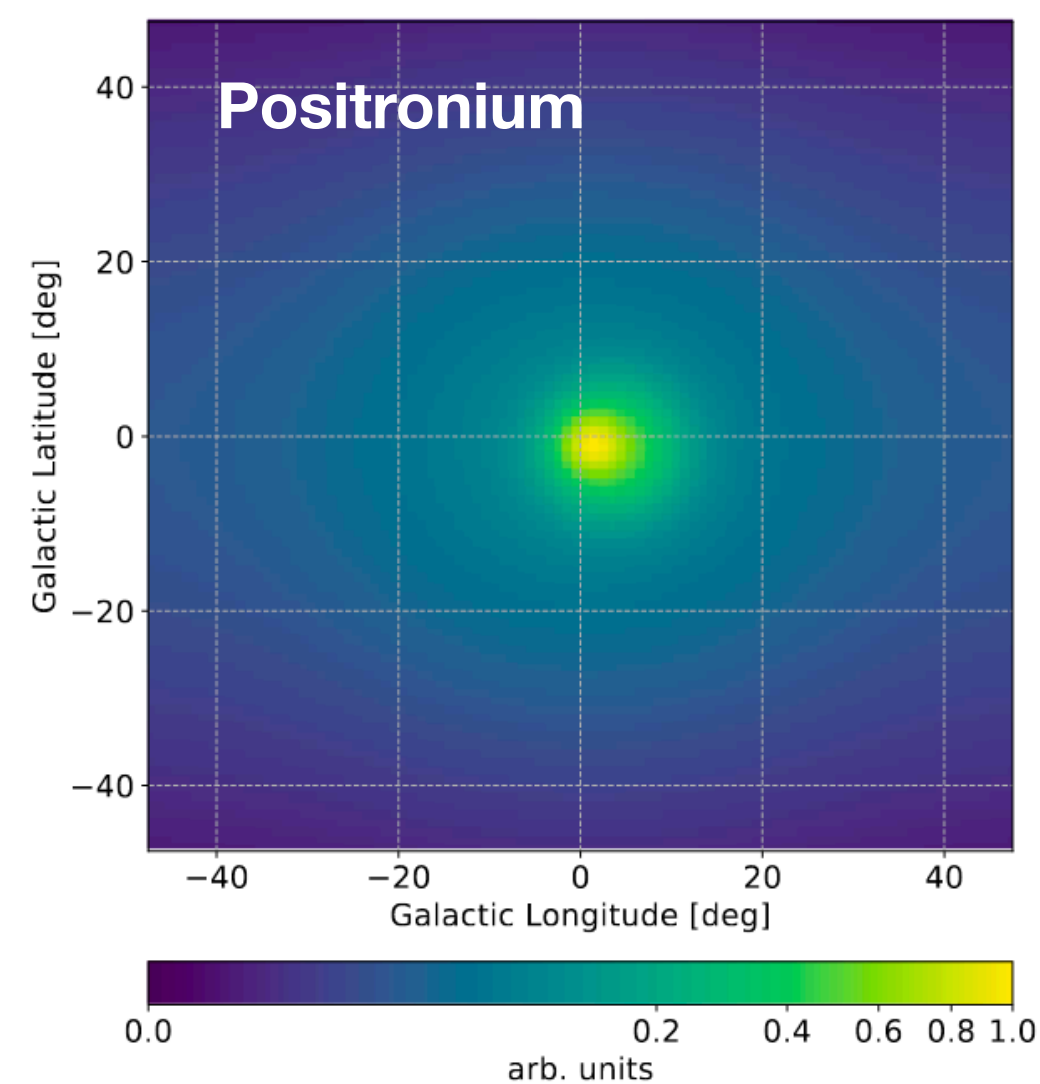
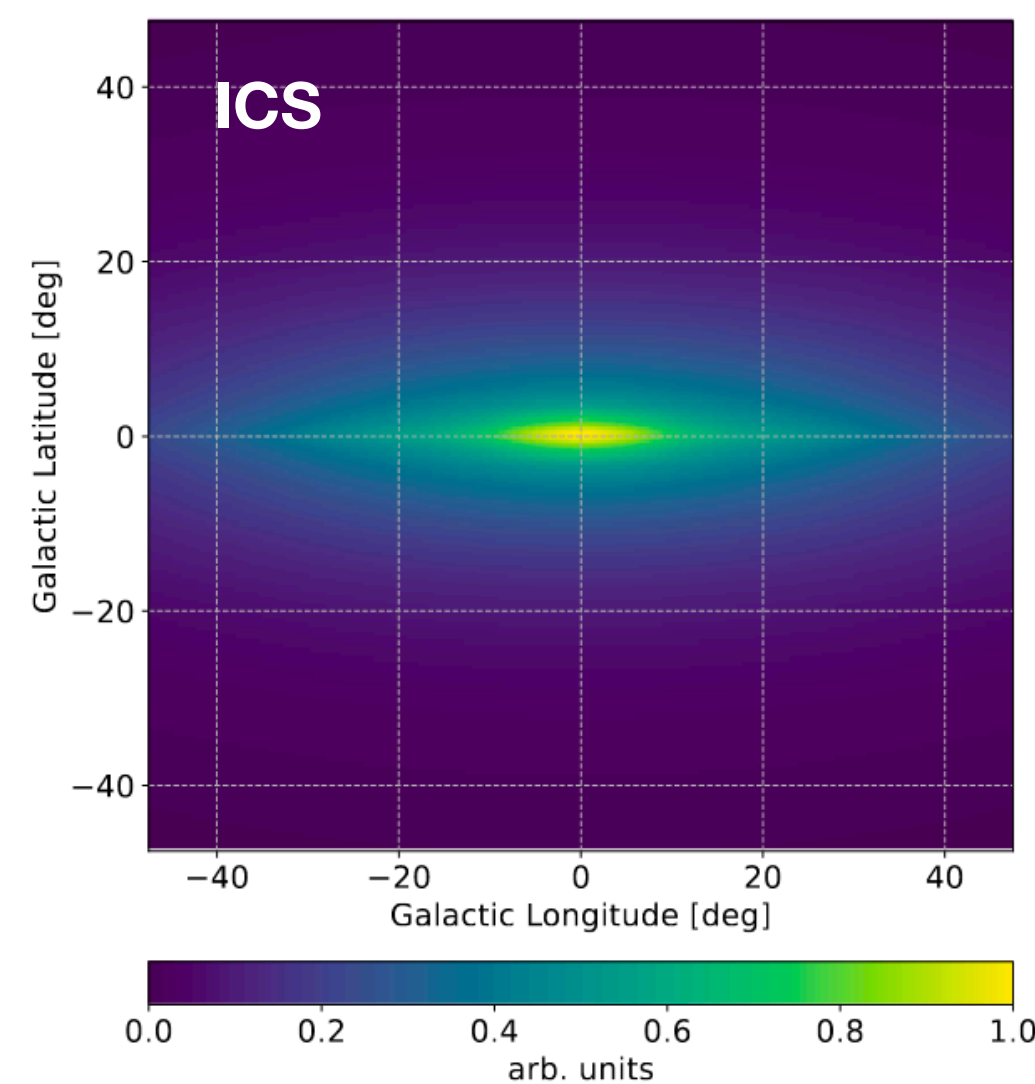
New analysis of 16yr-data from SPI
30 keV — 8 MeV

Extraction of the Galactic diffuse emission

Astrophysical contributions

Modelled **spatial templates** (30 keV — 8 MeV)

- **Inverse Compton scattering** of electrons off the interstellar radiation field $e_{\text{CR}}^{\pm} + \gamma \longrightarrow e^{\pm} + \gamma_{\text{MeV}}$
- Unresolved sources (<100 keV)
- Nuclear lines
- Positronium annihilation line+continuum



Extraction of the Galactic diffuse emission

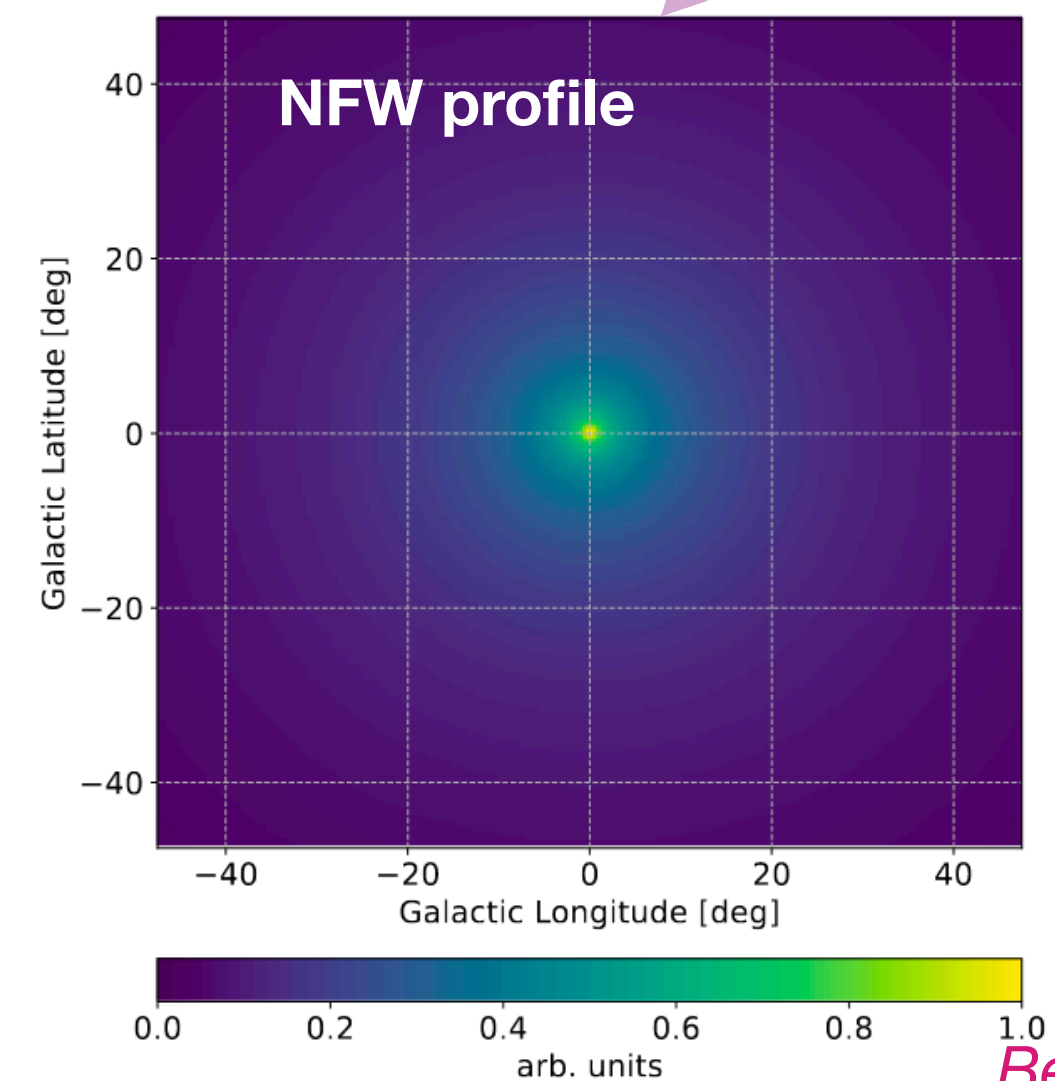
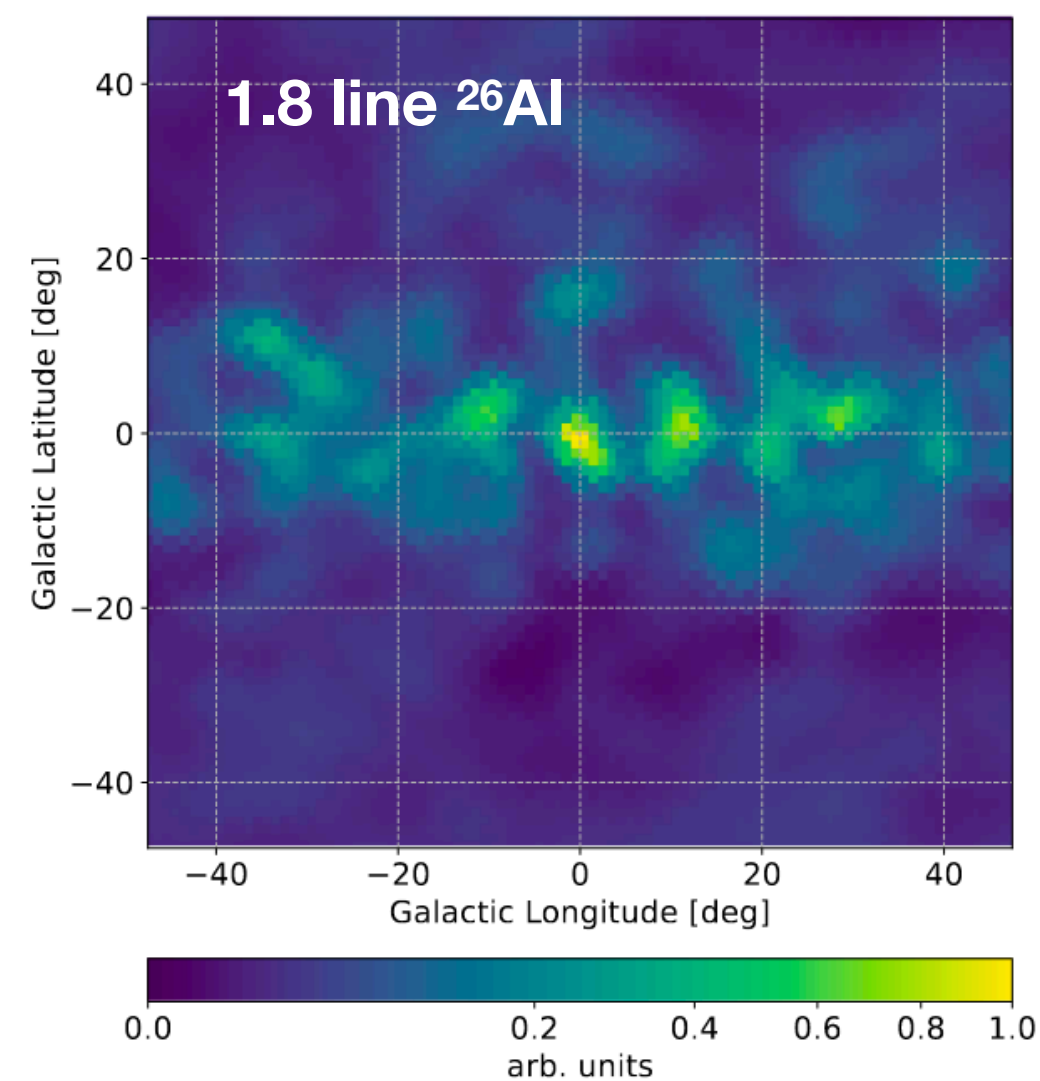
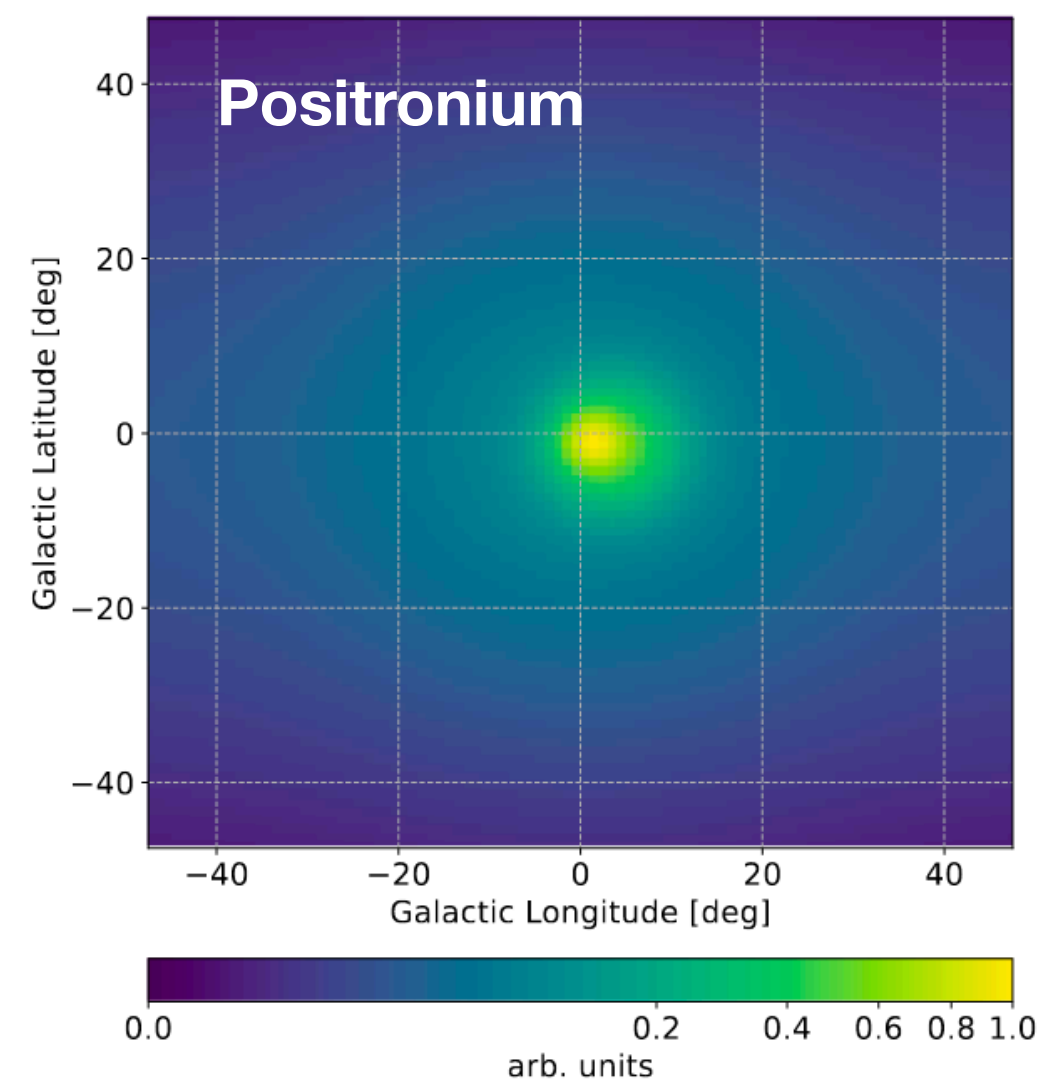
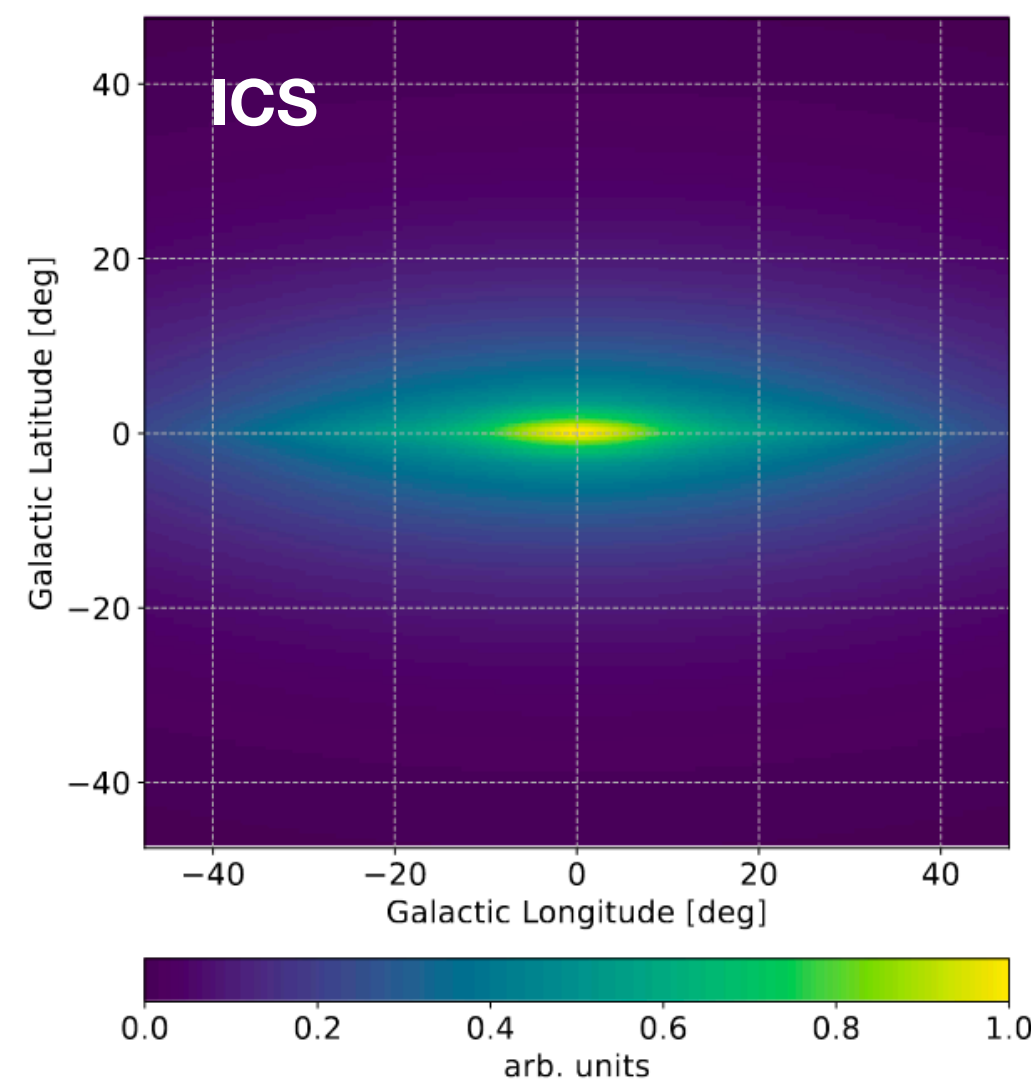
ALPs decay signal?

Modelled **spatial templates** (30 keV – 8 MeV)

- **Inverse Compton scattering** of electrons off the interstellar radiation field $e_{\text{CR}}^{\pm} + \gamma \longrightarrow e^{\pm} + \gamma_{\text{MeV}}$
- Unresolved sources (<100 keV)
- Nuclear lines
- Positronium annihilation line+continuum

- Additional **ALPs decay signal?**

$$\left(\frac{d\Phi_{\gamma}}{dE}\right)_{\text{decay}} = \frac{\Gamma(a \rightarrow \gamma\gamma)}{4\pi m_a} \left(\frac{dN_{\gamma}}{dE}\right)_{\text{decay}} \times \int_{\text{l.o.s.}} \rho_a(\ell) d\ell$$



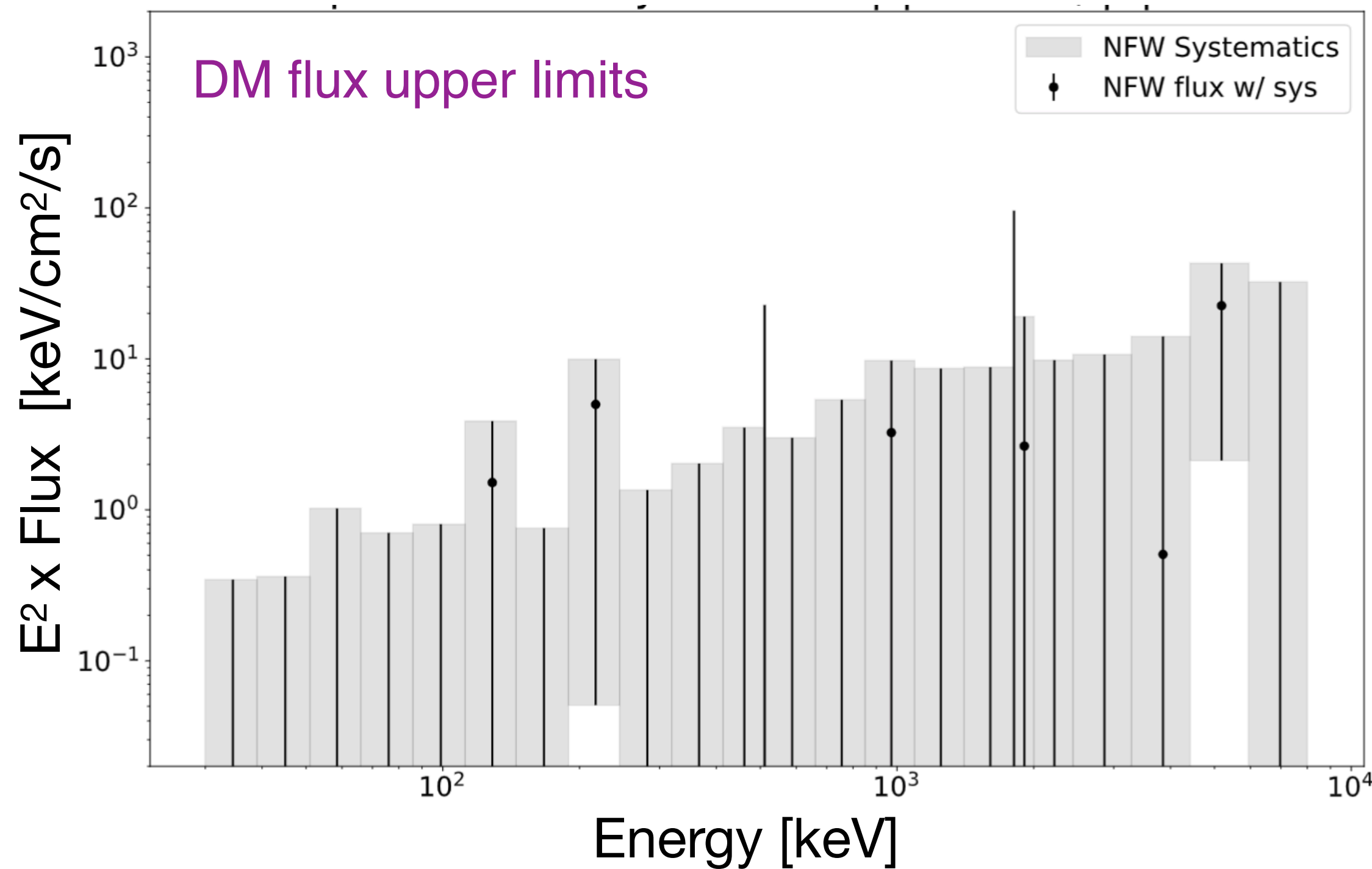
Berteaud, FC+ PRD'22

Constraints on decaying dark matter

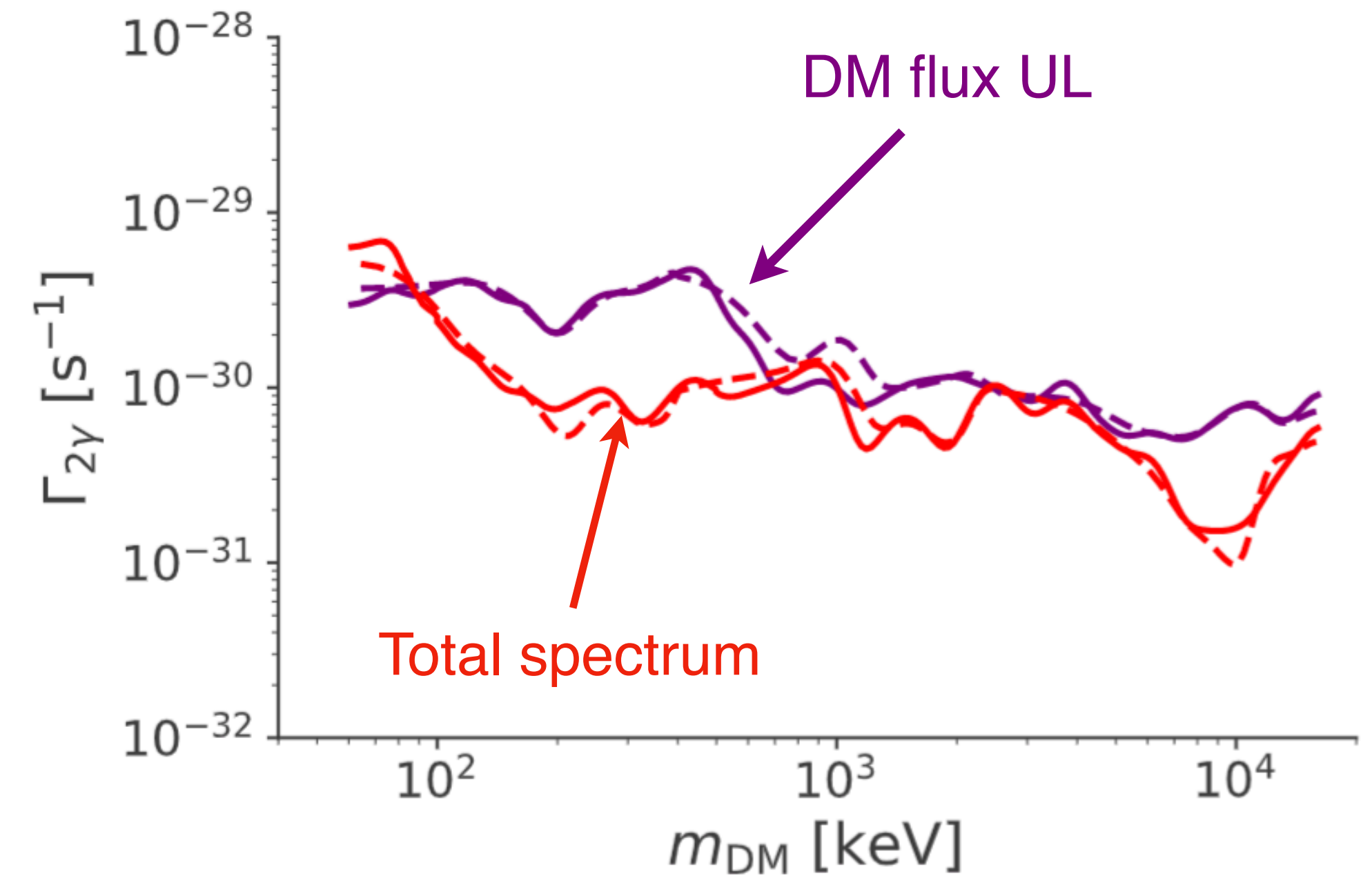
Decay into two photons (line-like spectrum)

FC+ MNRAS'23

No signal detected



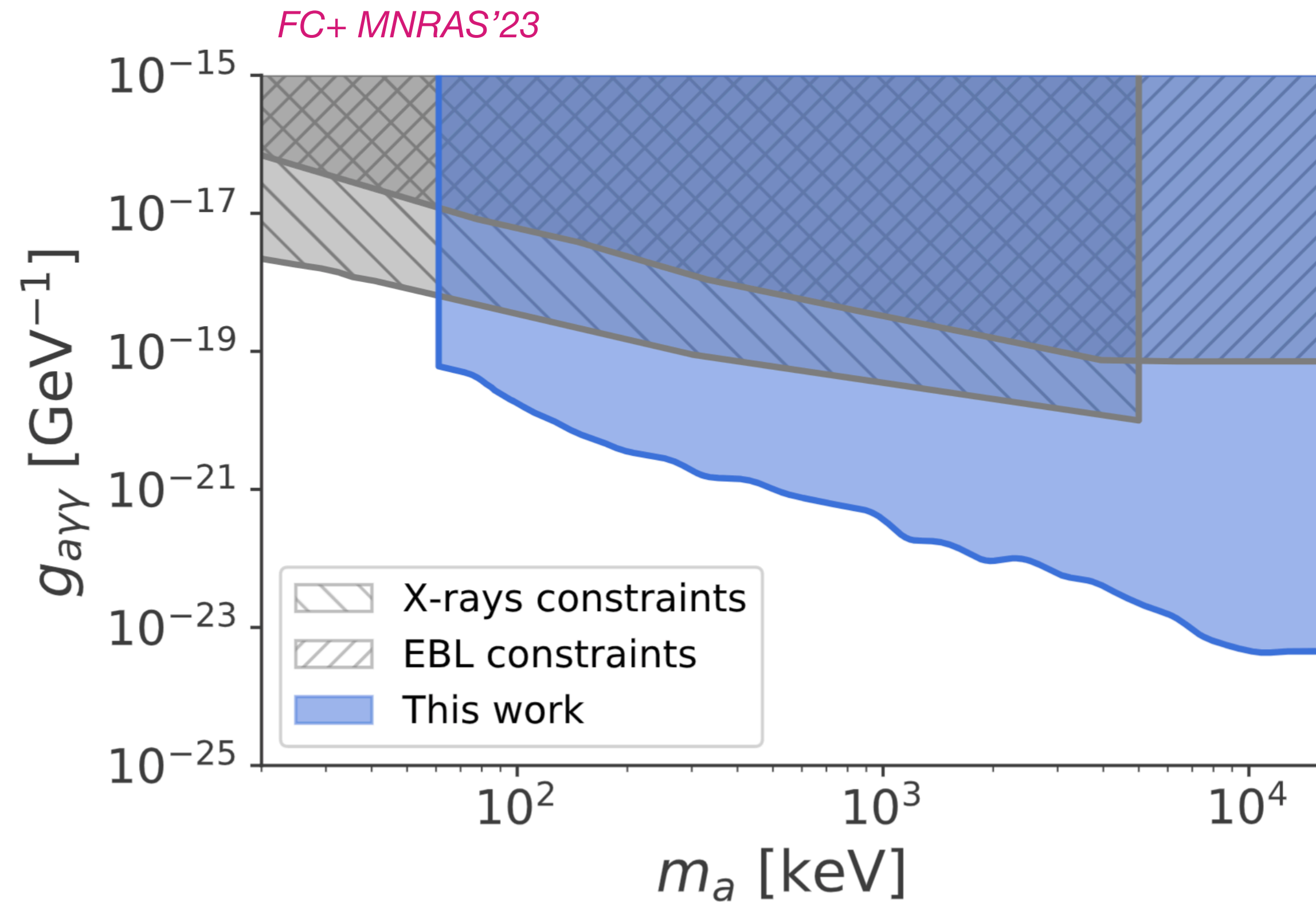
→ Upper limits on **decay rate** into 2 photons, $\Gamma_{2\gamma}$



$$\left(\frac{d\Phi_\gamma}{dE}\right)_{\text{decay}} = \frac{\Gamma(a \rightarrow \gamma\gamma)}{4\pi m_a} \left(\frac{dN_\gamma}{dE}\right)_{\text{decay}} \times \int_{\text{l.o.s.}} \rho_a(\ell) d\ell \quad \frac{dN_\gamma}{dE} = 2\delta\left(E - \frac{Nm_{\text{DM}}}{2}\right)$$

<https://zenodo.org/record/7984451>

Constraints on ALPs dark matter



$$\tau_a = \frac{64\pi}{m_a^3 g^2}$$

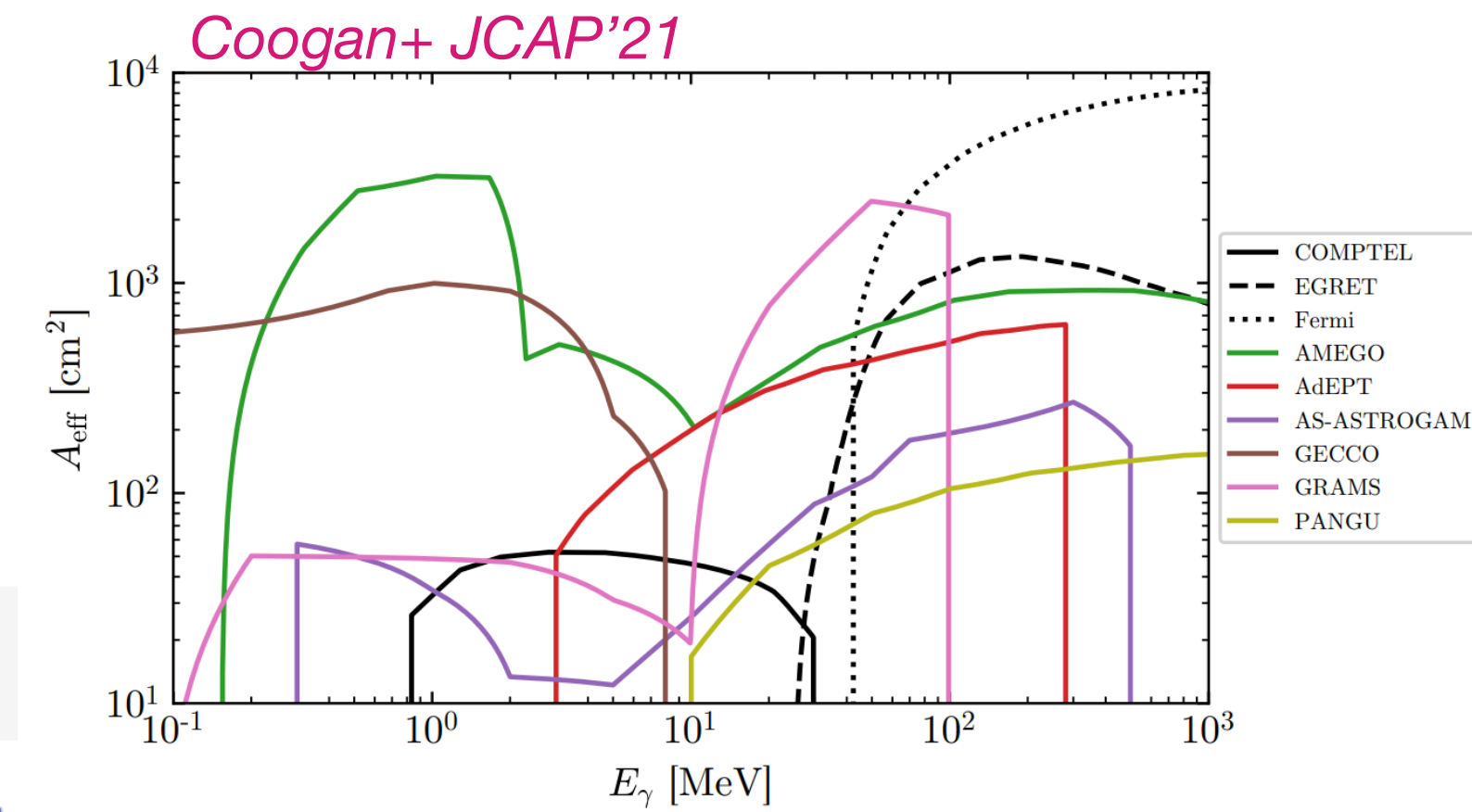
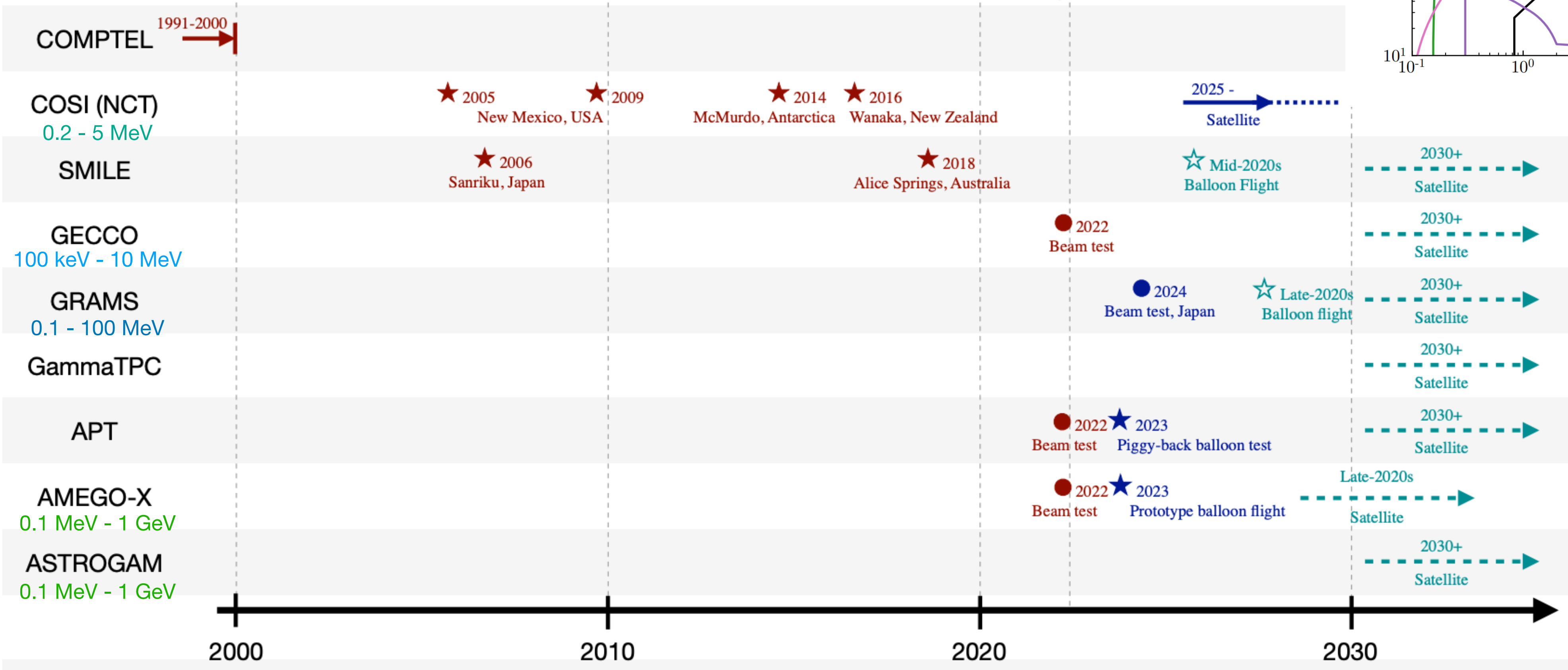
$$\Gamma_{2\gamma} = \frac{g_{a\gamma\gamma}^2 m_a^3}{64\pi} = 0.755 \times 10^{-30} \left(\frac{g_{a\gamma\gamma}}{10^{-20} \text{ GeV}^{-1}} \right)^2 \left(\frac{m_a}{100 \text{ keV}} \right)^3 \text{ s}^{-1}$$

Re-analysis of Integral/SPI data provides the **strongest constraints** on (light) particle and non-particle DM

Future: MeV Galactic diffuse emission

Covering the MeV sensitivity gap

MeV Gamma-ray missions



$$N_{\gamma} = T_{\text{obs}} \int_{E_{\text{min}}}^{E_{\text{max}}} dE A_{\text{eff}} \frac{d\Phi}{dE_{\gamma}}$$

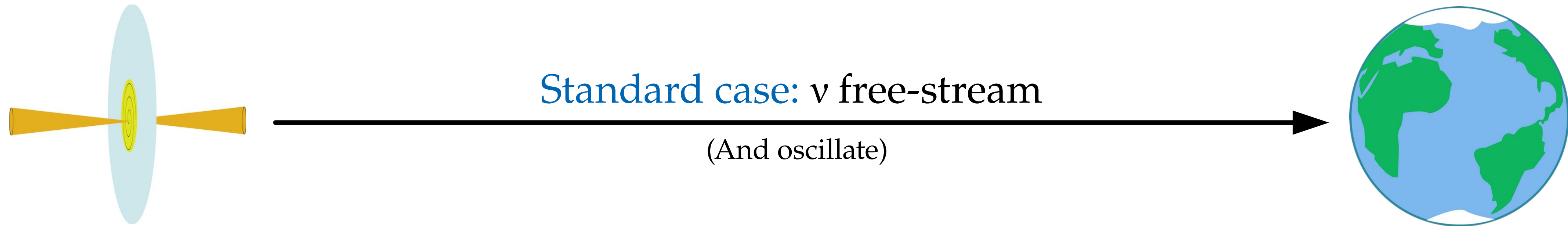
Aramaki+ Snowmass'21 CF

Anomalous propagation effects

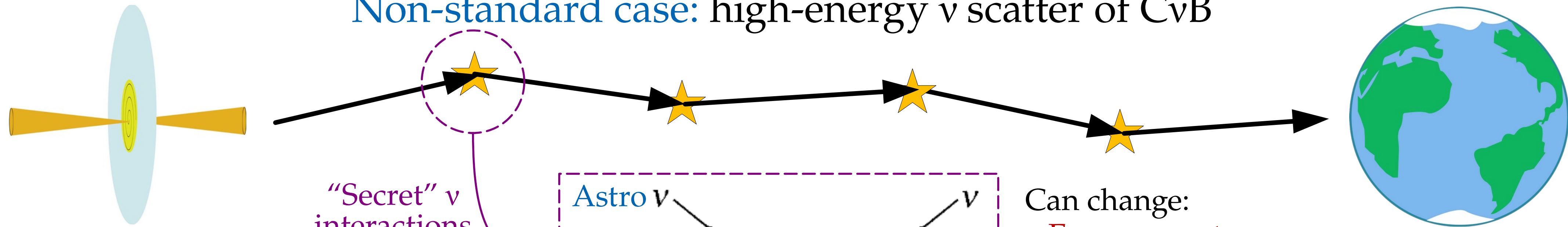
Astrophysical neutrino sources

Earth

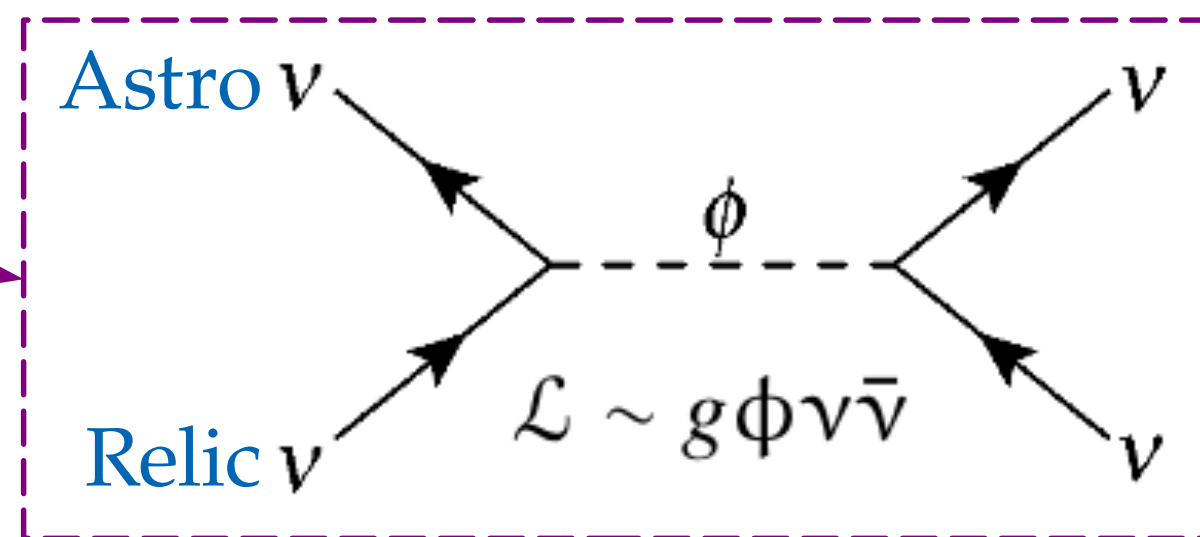
Galactic (kpc) or extragalactic (Mpc – Gpc) distance



Non-standard case: high-energy ν scatter of CvB



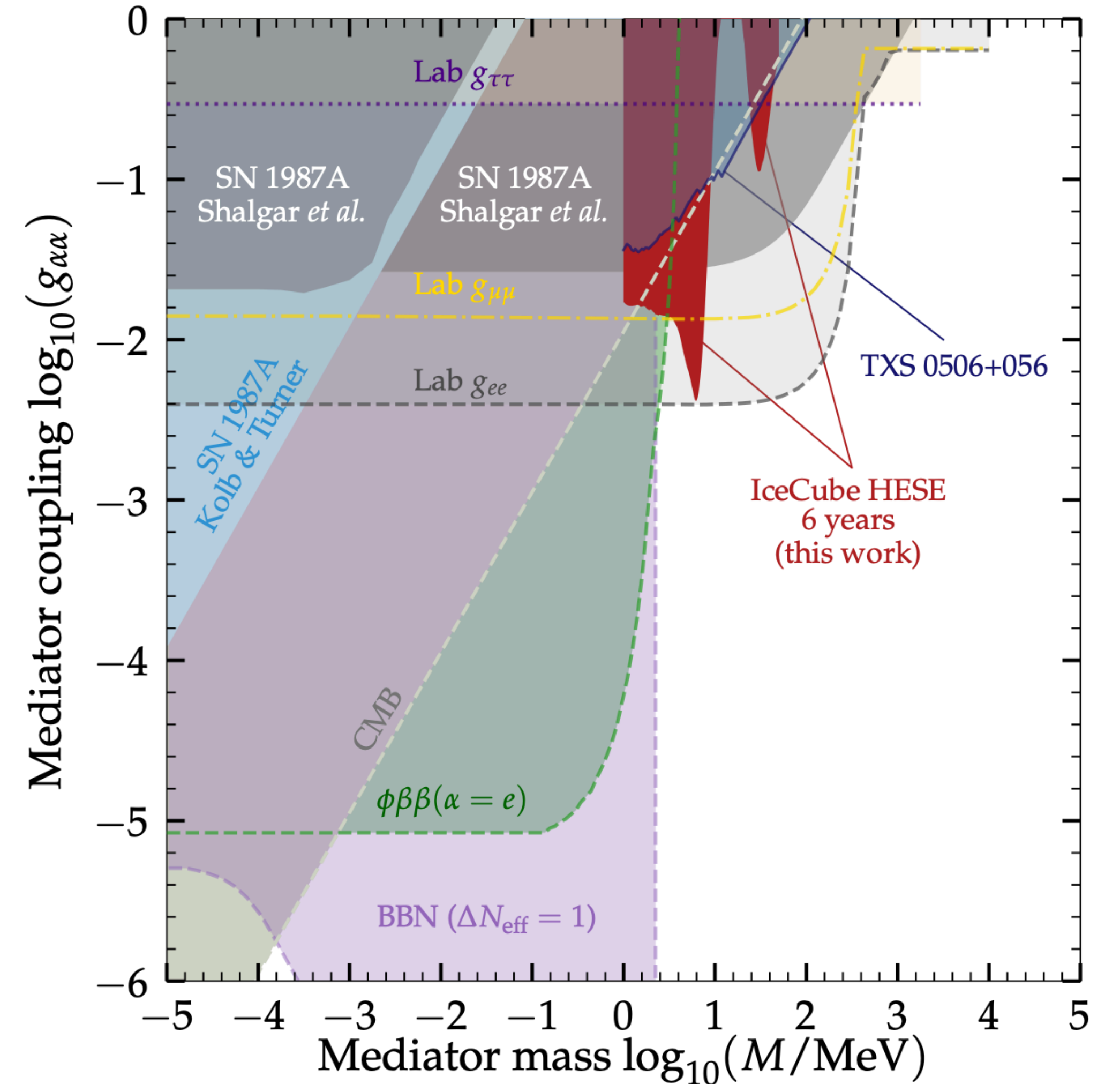
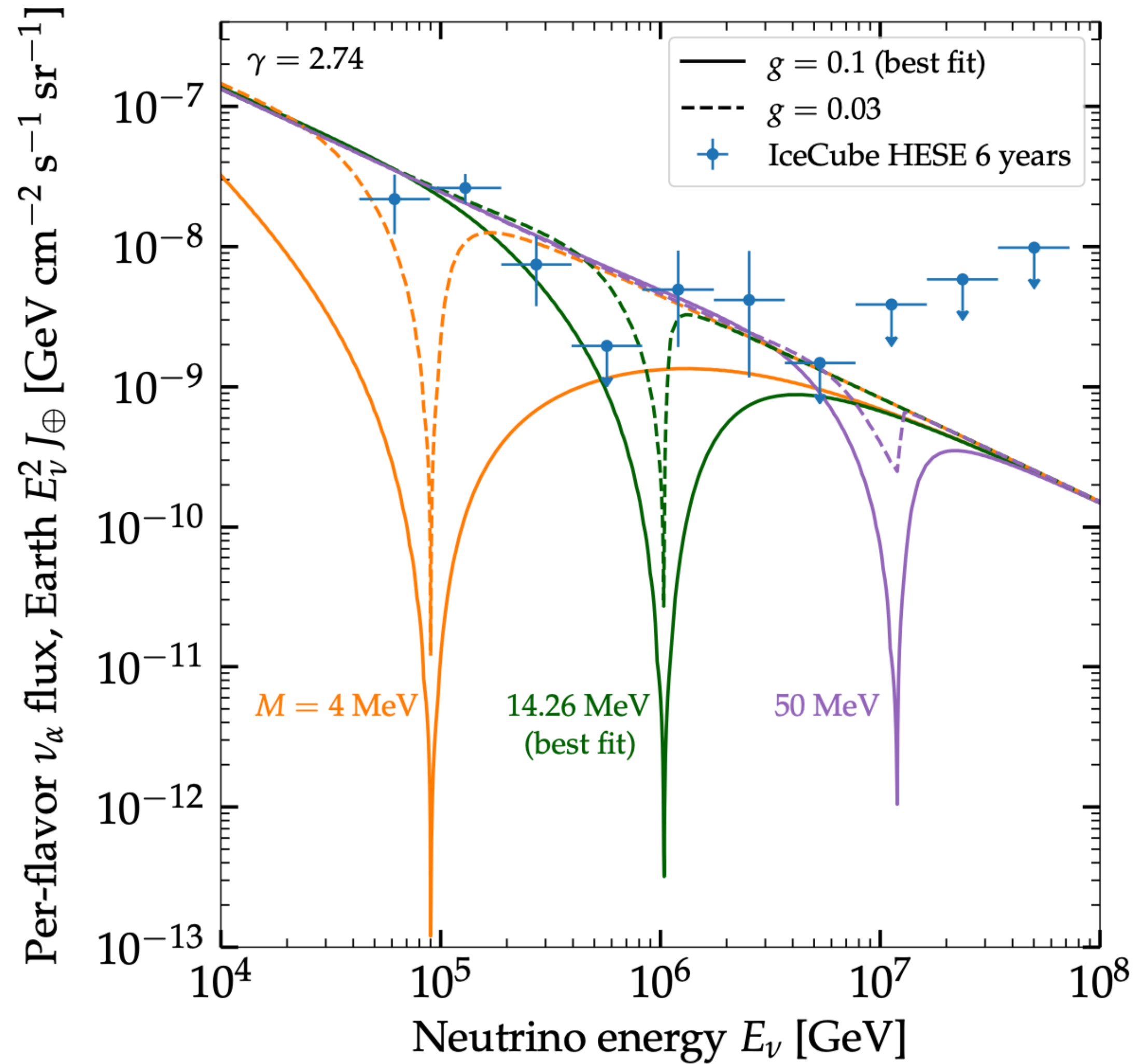
“Secret” ν interactions
≡
BSM ν self-interactions



- Can change:
- ▶ Energy spectrum
 - ▶ Flavor composition
 - ▶ Direction
 - ▶ Arrival times

Neutrino spectrum and bounds

Bustamante+ PRD'20



No statistically significant evidence for νSI in the 6-year HESE sample

Tests of LIV

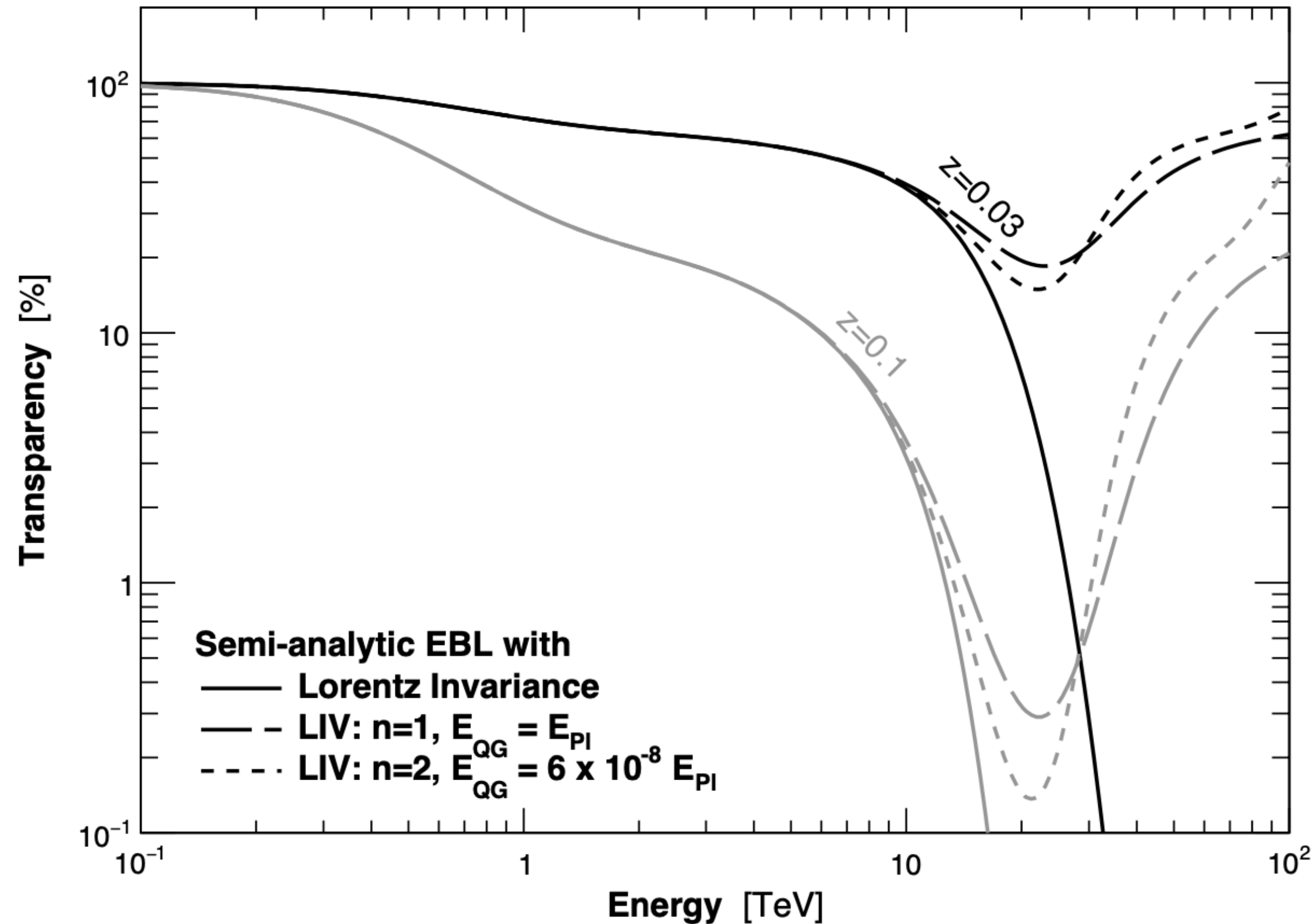


Figure 7.8: Attenuation factor, in percent, as a function of gamma-ray energy on Earth, for sources located at $z = 0.03$ and $z = 0.1$, following the semi-analytical EBL model of [Gilmore et al. \(2012\)](#) and the LIV formalism of [Jacob and Piran \(2008\)](#). Both linear ($n = 1$) and quadratic ($n = 2$) modifications of the pair-creation threshold are shown, at the Planck scale for $n = 1$, and at an ad-hoc energy scale of $6 \times 10^{-8} E_{\text{Pl}}$ for $n = 2$, resulting in a similar effect.

Tests of LIV

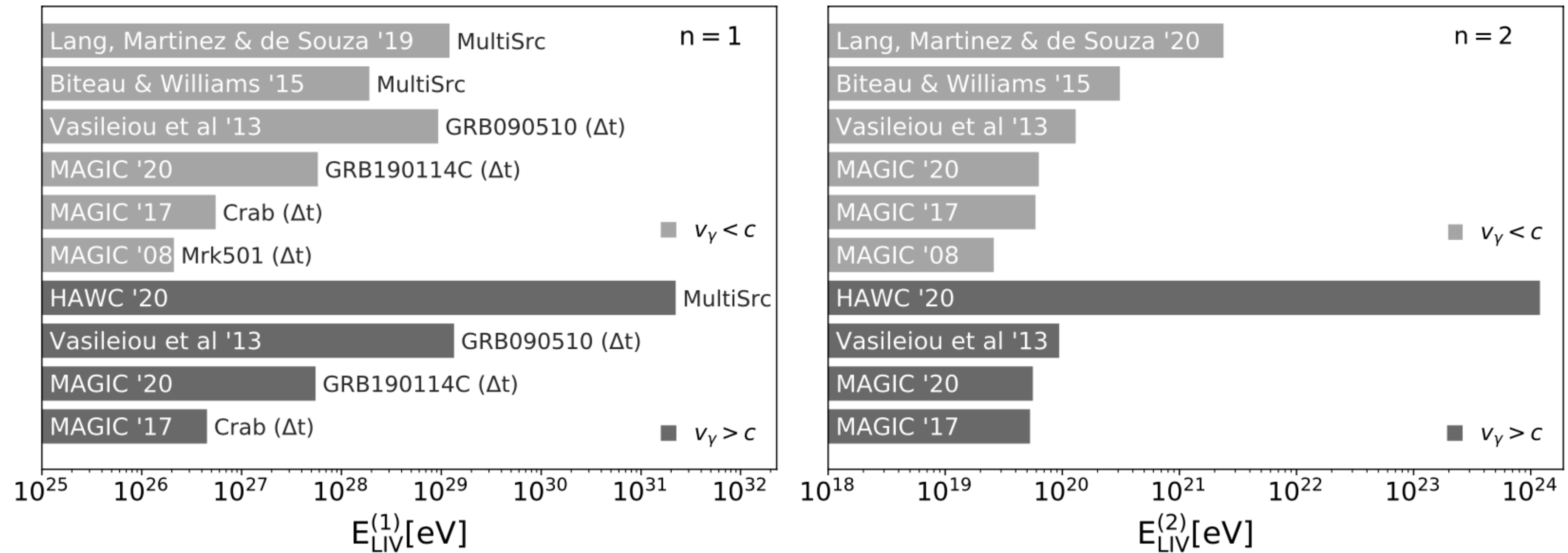


Figure 7.10: Summary of the lower limits on the energy scale of LIV, E_{LIV} . The left panel shows the limits on the linear term, the right panel on the quadratic term. Light grey bars assume a subluminal modification, while dark grey bars assume a superluminal modification. Credit: Humberto Martínez-Huerta, adapted from [Martínez-Huerta et al. \(2020\)](#).