

# Listening for ultra-heavy DM with underwater acoustic detectors

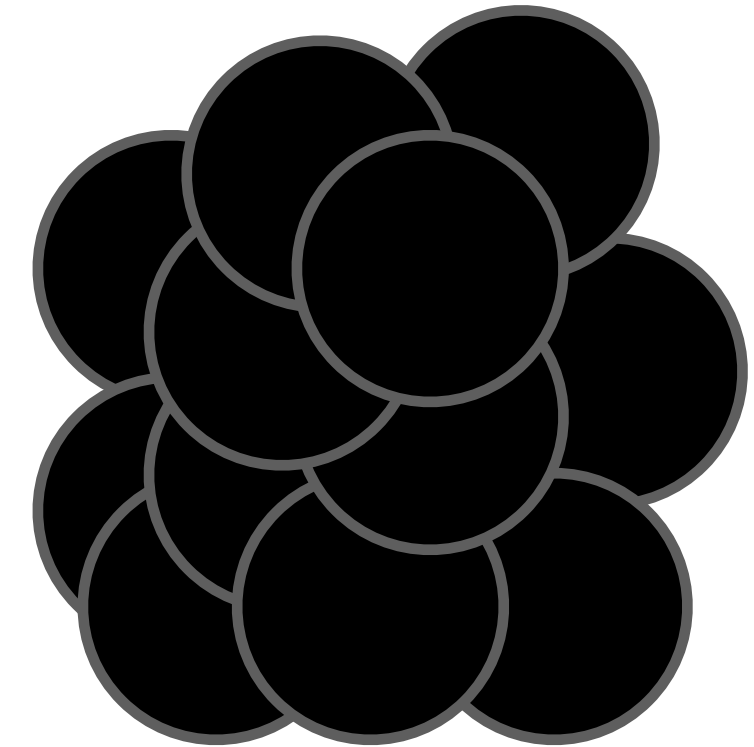
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# Ultra-heavy Dark Matter

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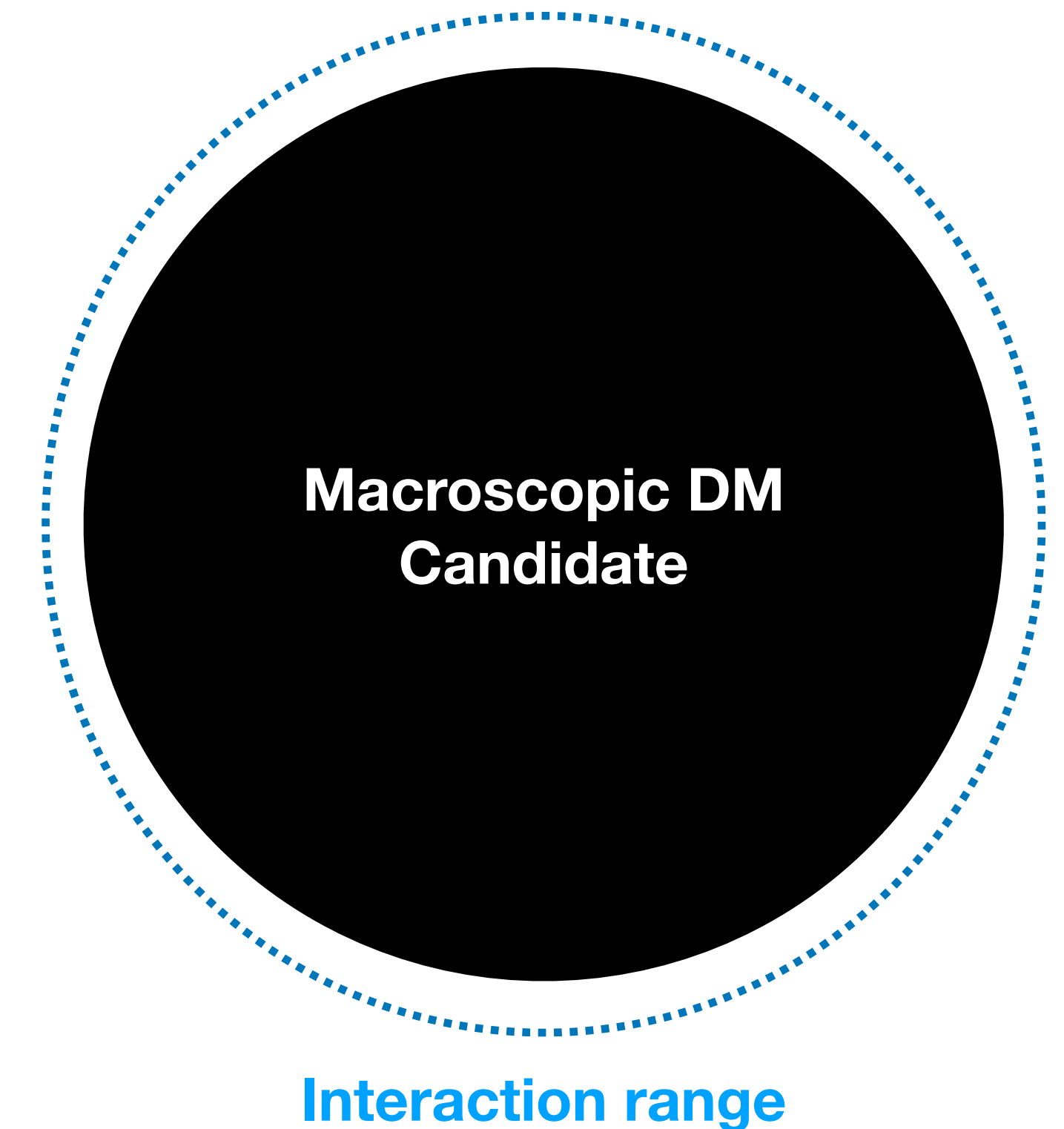
- Ultra-heavy dark matter is necessarily composite (if thermally produced) due to s-wave unitarity
- Many different models for UHDM
  - Nuggets, Blobs, WIMPonium, Q-Balls etc...
- Is there a nice model-independent way to treat them?
- Answer: **Yes** (for some parts of parameter space)



# Macros

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- Consider parameters of models where:
  - The DM is Planck-mass or larger
  - DM Radius  $R_\chi$  much larger than interaction length scale
  - Geometric cross section dominates i.e.  $\sigma_\chi \approx \pi R_\chi^2$
  - Parameterise the interaction in terms of  $R_\chi$  -> set by the theory -> make experimental statements about multiple models!



# Macro Direct Detection

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- Macro's often parameterised in grams (g)

- DM Flux:  $\phi_\chi \approx 6 \left( \frac{1 \text{ g}}{m_\chi} \right) \text{ km}^{-2} \text{ yr}^{-1}$

- Need a *very* large detector (or very long integration time) to have significant number of events.

# Macro Kinematics

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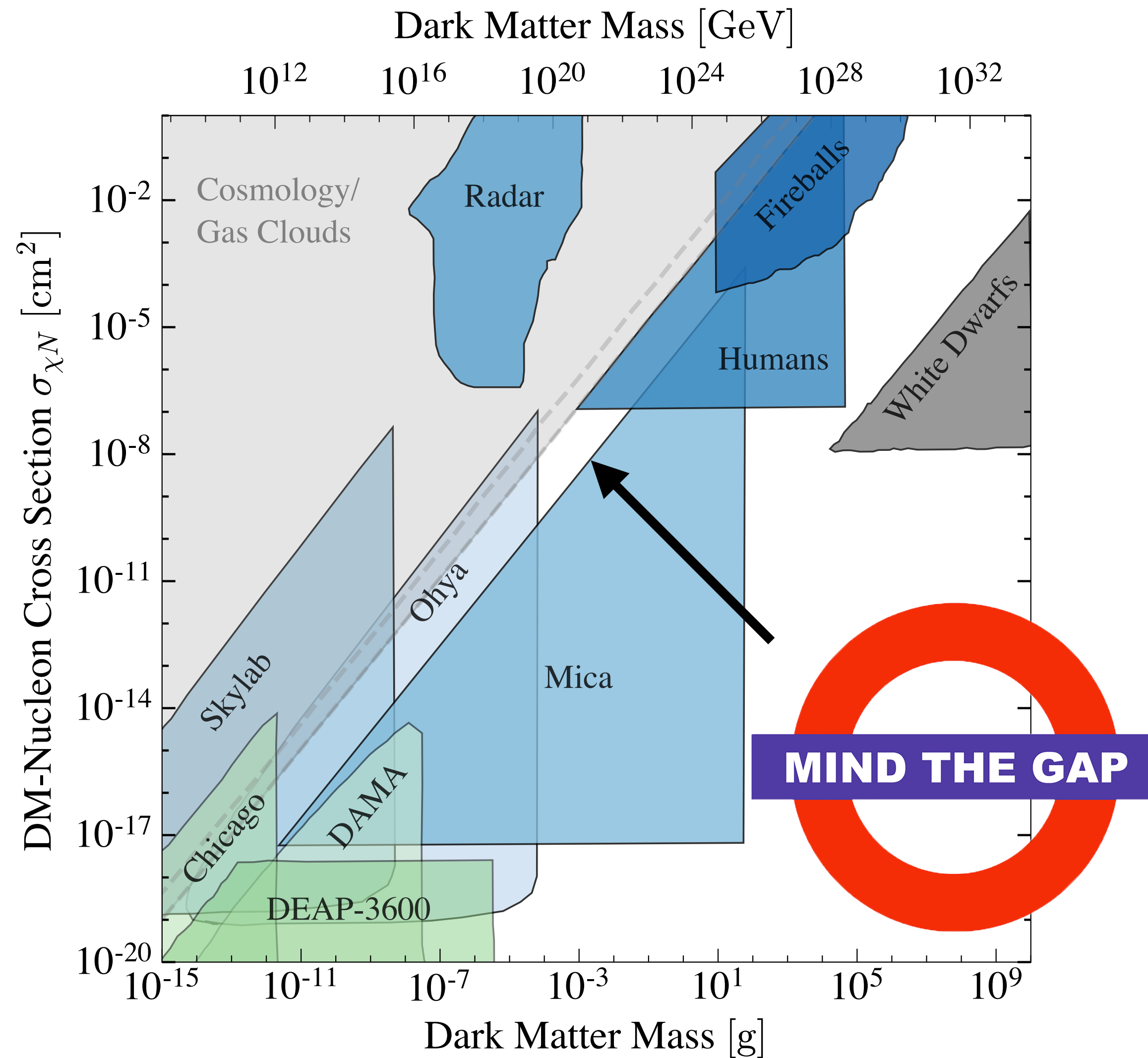
Total energy after traversing path length  $L$  in detector given by (determined by classical scattering):

$$E_{\chi f}(L) = \frac{1}{2} m_{\chi} v_{\chi,SL}^2 \exp \left( -2 \frac{\sigma_{\chi N}}{m_{\chi}} \rho_{\text{med}} L \right)$$

Differentiating, can find the energy deposition rate into the medium:

$$\frac{dE}{dL} = -\frac{dE_{\chi}}{dL} = \rho_w \sigma_{\chi N} v_{\chi,SL}^2 \exp \left( -2 \frac{\sigma_{\chi N}}{m_{\chi}} \rho_{\text{med}} L \right)$$

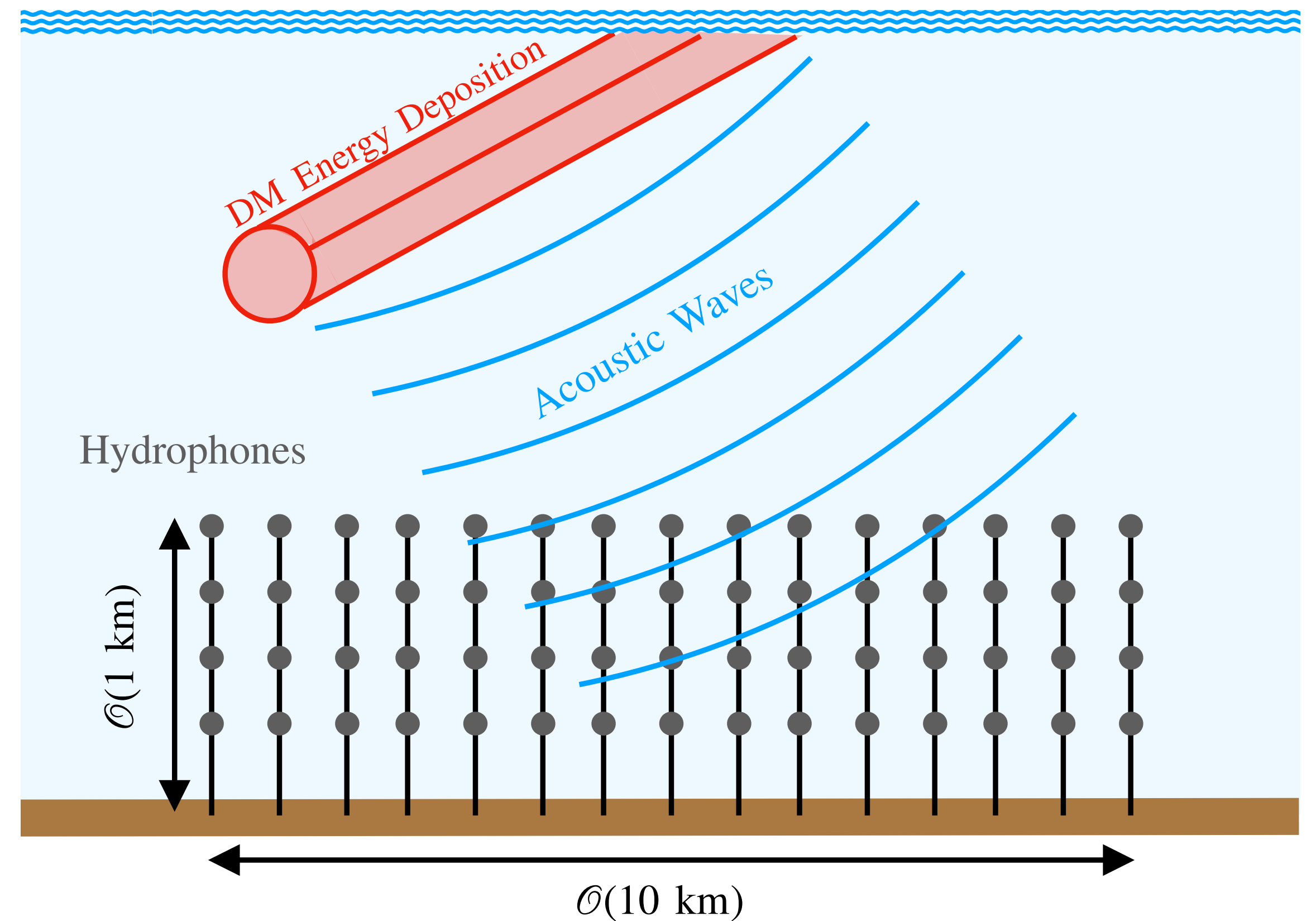
# Current Constraints



- Annoying “gap” in constraints
- Mica underground - too much overburden
- Radar not sensitive enough - not enough ionisation
- **What phenomena could we use to constrain this region?**

# Acoustic Detection

- Idea: DM is weakly interacting enough to make it through the atmosphere
- Reaches much more dense medium: **the ocean**
- DM deposits energy into the ocean creating pressure waves
- Detect pressure waves using **a large hydrophone array**



# Neutrino Experiments

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- Propositions for acoustic neutrino experiments with  $\mathcal{O}(100 \text{ km}^3)$  hydrophone arrays in the ocean [Lahmann, 2016]
- Detect UHE neutrinos. Similar number density issues, but similarly high cross section
- Acoustic propagation distance in water much greater than light -> less dense instrumentation required
- Energy deposition comes from particle showers



# What is the signal?

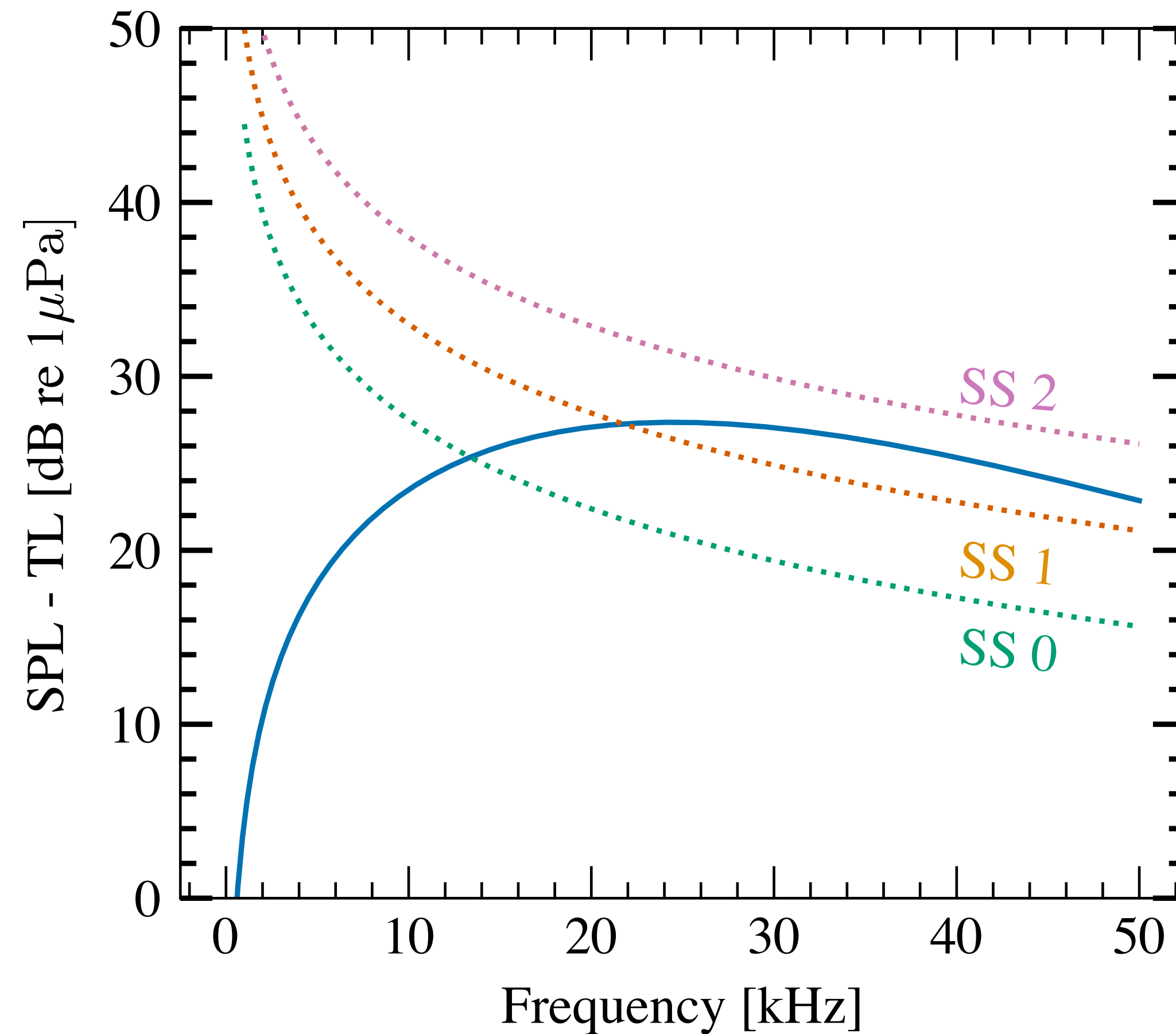
Pressure waves created from **thermo-acoustic heating**.  
[Learned, 1979]

$$\nabla^2 \underbrace{P(\mathbf{r}, t)}_{\text{Acoustic pressure}} - \frac{1}{c_s^2} \frac{\partial^2 P(\mathbf{r}, t)}{\partial t^2} = - \frac{\alpha}{C_p} \frac{\partial^2 \underbrace{q(\mathbf{r}, t)}_{\text{Energy Deposition Density}}}{\partial t^2}$$

General solution to this equation given by:

$$P(\vec{r}, t) = \frac{\alpha}{4\pi C_p} \int_V \frac{dV'}{|\vec{r} - \vec{r}'|} \frac{\partial^2 q(\vec{r}', t')}{\partial t'^2}, \quad t' = t - \frac{|\vec{r} - \vec{r}'|}{c_s};$$

# Signal and Noise Characteristics



- The signal is **broadband**
- After transmission losses, **signal power greatest in 10-30kHz band**
- Dominant ambient noise in this band is **sea-state noise (surface agitation from wind)**

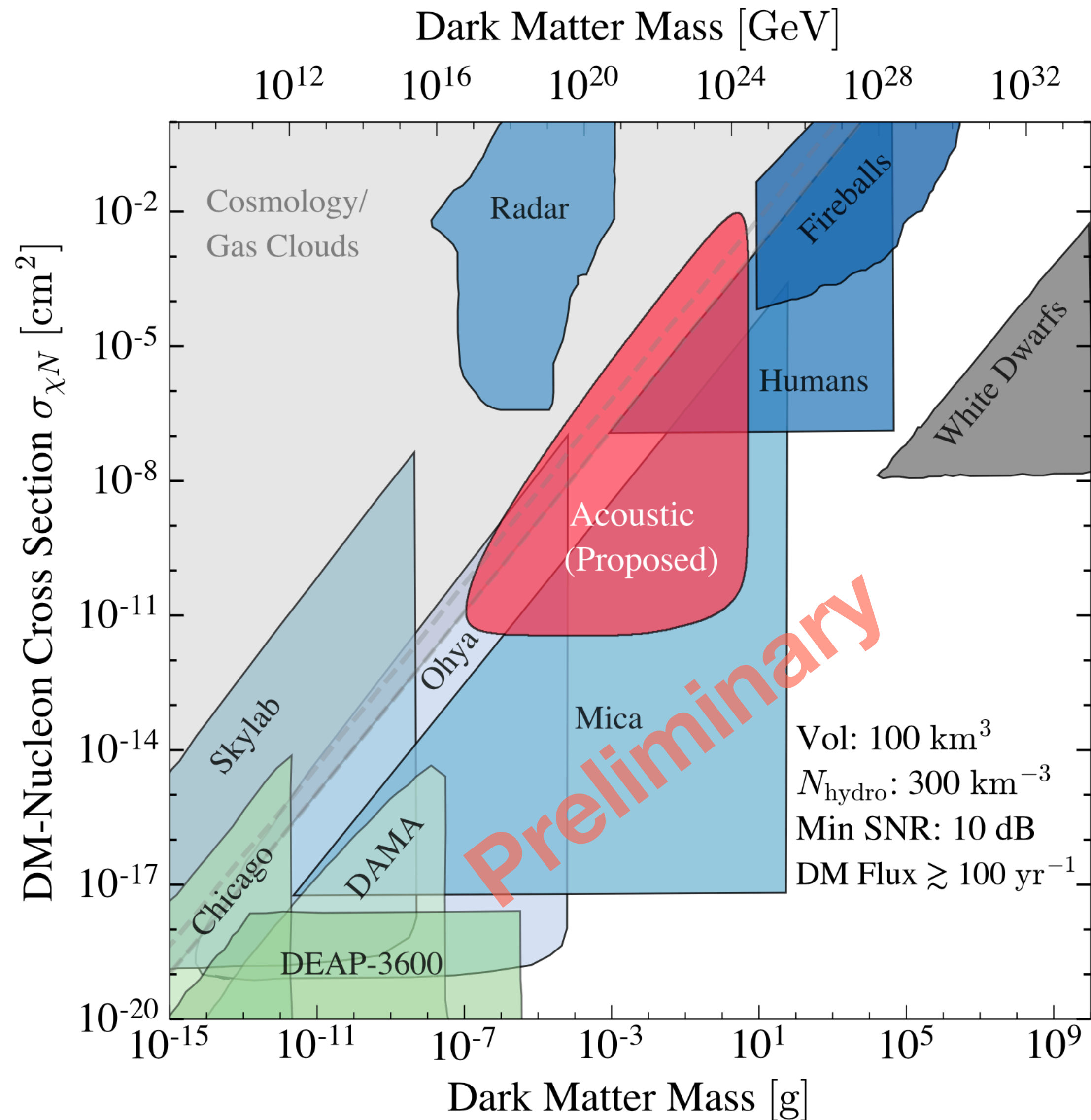
# What we require

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DM signal is:

- Larger than the typical hydrophone threshold (approx  $1 \mu\text{Pa}$  @ 15kHz) after propagation losses.
- Carries a minimum SNR to sea-state noise (conservatively defaulted to 10).
- Each hydrophone signal can be coherently summed to integrate the signal i.e. a  $\sqrt{N_{\text{hydro}}}$  enhancement to the signal.

# Preliminary Sensitivities



- Assuming proposed acoustic neutrino experiment parameters, **could constrain the gap!**
- Complementary to Humans, Mica, Ohya and Cosmological Bounds

# Punchline

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**Future acoustic neutrino experiments could detect (or constrain) ultra-heavy dark matter**

**Any Questions**

# Backup Slides

# Decibel Formalism

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The pressure wave will have some **frequency dependent absorption** due to the chemical content of the sea water

Convenient to use a **decibel formalism**, as this is most often used in acoustics.  
Sound pressure level:

$$SPL = 20 \log_{10} \left( \frac{P}{P_{\text{ref}}} \right)$$

$$P_{\text{ref}} = 1 \mu\text{Pa}$$



# Transmission Loss Model

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The Transmission loss over a propagation distance  $r$ , is given by a frequency dependent parameter  $\alpha$

$$TL = \alpha r$$
$$\alpha = 0.106 \frac{f_1 f^2}{f^2 + f_1^2} e^{(pH-8)/0.56}$$
$$+ 0.52 \left(1 + \frac{T}{43}\right) \left(\frac{S}{35}\right) \frac{f_2 f^2}{f^2 + f_2^2} e^{-D/6}$$
$$+ 0.00049 f^2 e^{-(T/27+D/17)},$$

[Ainslie, McColm 1998]

Dependent on the temperature  $T$ , salinity  $S$ , depth  $D$  and  $pH$ . Frequency here in kHz. **After transmission losses, DM signal peaks at ~15kHz**

# Sea State Noise

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The dominant background noise source 10-100 kHz band is **sea state noise (surface agitation due to wind etc)**. Parameterised by Knudsen curves:

$$NL(f, n_s) = n_s - 10 \log_{10} \left( f^{-5/3} \right)$$

There will also be transient noise sources from ocean wildlife e.g. dolphins and sperm whales. **Assume differentiable from DM signal** using algorithms being developed for neutrino detection.

# Yearly and Daily Modulations

