

Analytic computation of 2-loop scattering amplitudes

Ekta Chaubey

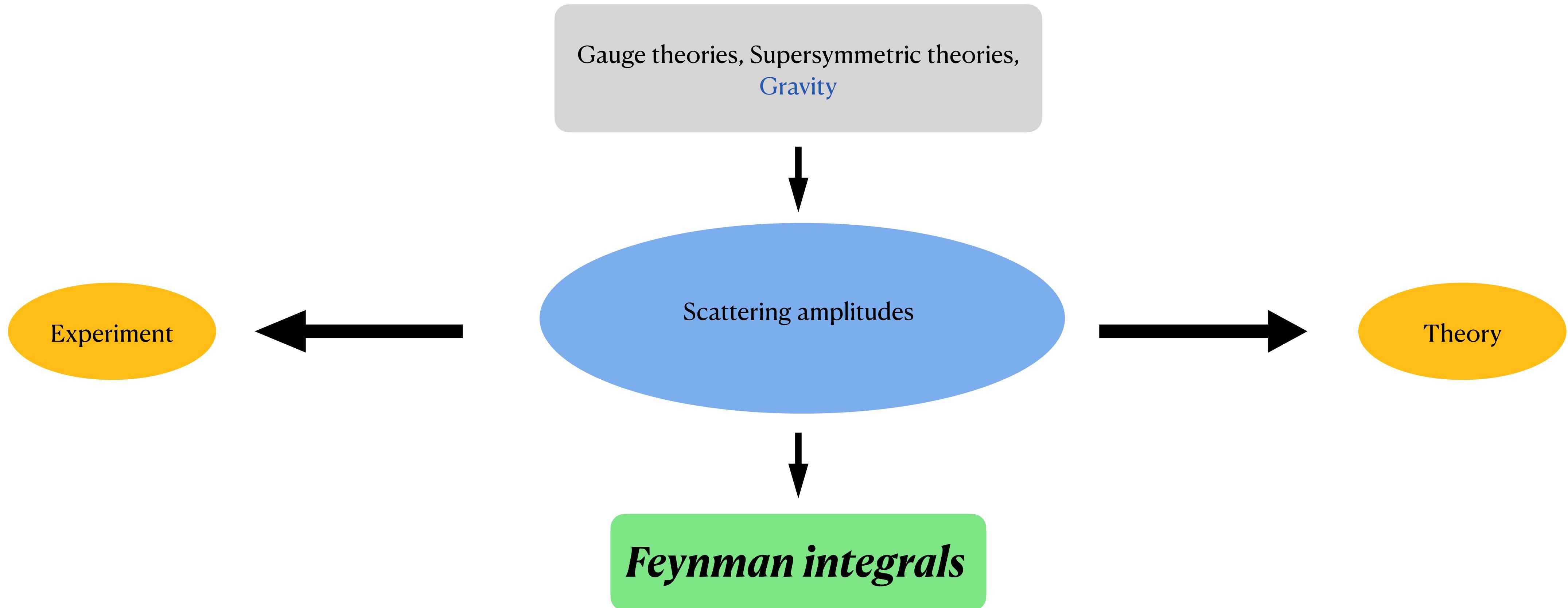
University of Bonn

27th September 2023
Padova, Italy



Bethe Center for
Theoretical Physics





For precision physics, **precise** theoretical predictions are needed
→ computation of **higher-order loop corrections**

Lance
David
Andreas,
Giuseppe,
Simone

Analytic structure

Class of functions?

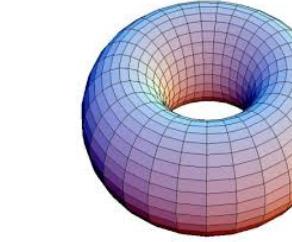
Rich Mathematical structure



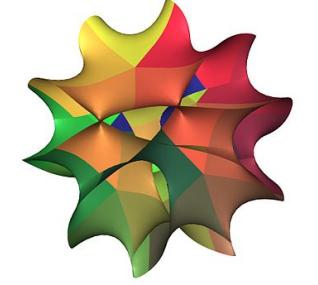
Multiple Polylogarithms



Elliptic Polylogarithms



K₃ surfaces

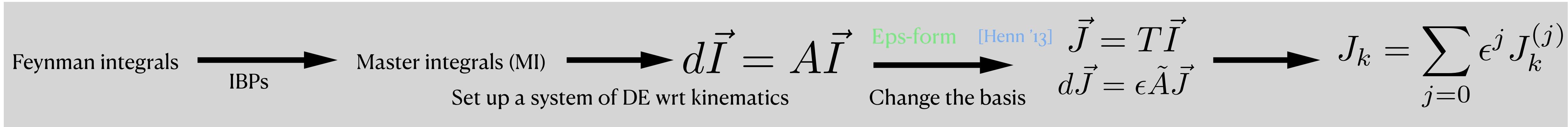


Higher-dimensional Calabi-Yau manifold

Can we identify them?

Can we “use” them?

Modern techniques for solving Feynman integrals



The choice of basis not unique; no general methods fits all

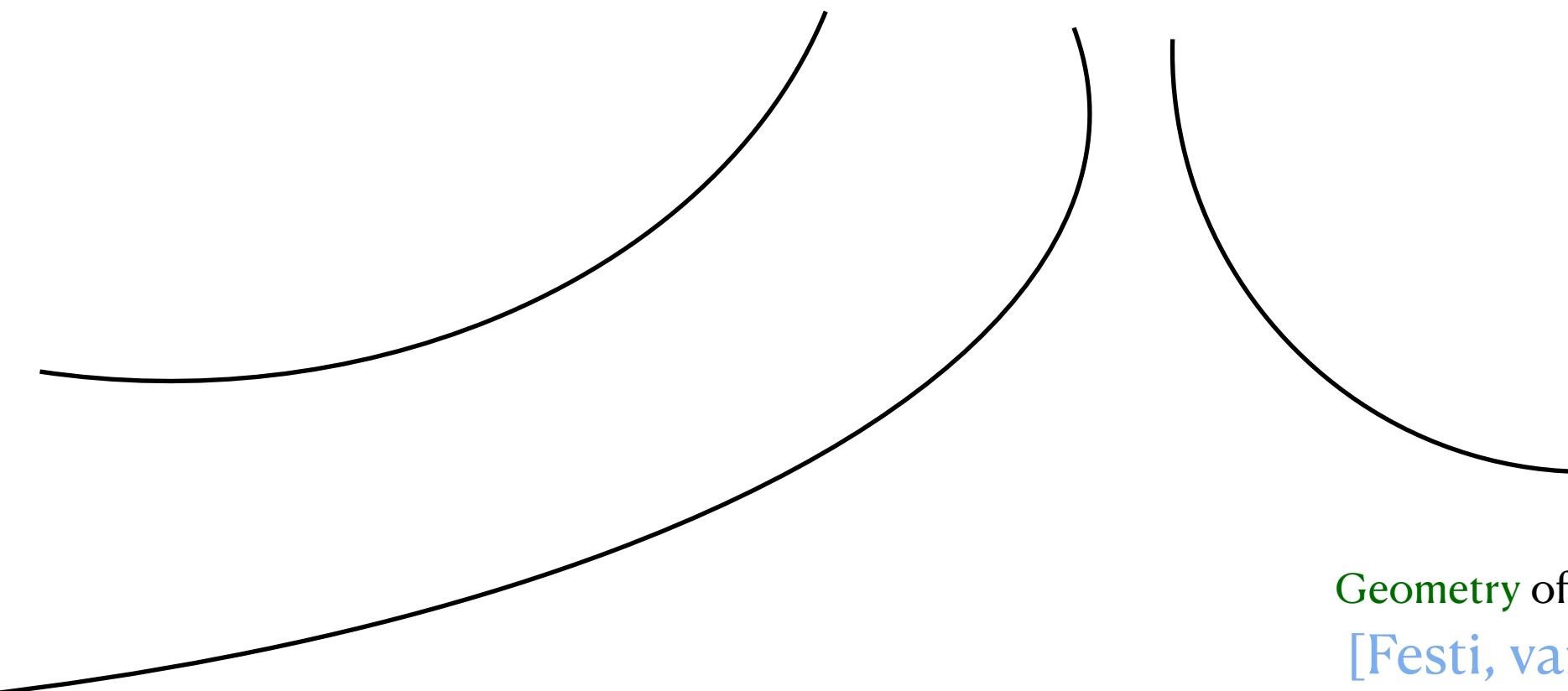
[Lee '14] [Meyer '18] [Wasser '20]
 [Adams, E. Chaubey, Weinzierl, '17]
 [Dlapa, Henn, Wagner, '22]

Maximal cuts

[Baikov, '96, '97]
 [Frellesvig, Papadopoulos, '17]
 [Primo, Tancredi, '17]
 [Adams, E. Chaubey, Weinzierl, '17, '18]

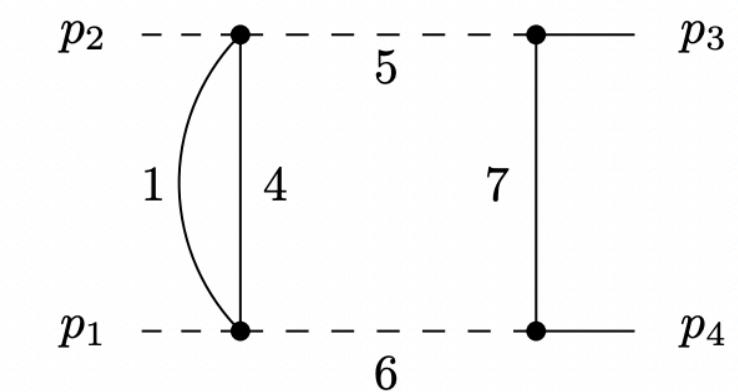
Christoph,
Seva,
William,
Federico,
Jacob

4



Square roots

Geometry of the Feynman integral often manifested by the square root
 [Festi, van Straten, '18] [Adams, E. Chaubey, Weinzierl '18]



$$Maxcut_C I_{2001111}(4 - 2\epsilon) \propto \int_C \frac{dP}{\sqrt{(P-t)(P-t+4m^2)(P^2 + 2m^2 \frac{s+4t}{s-4m^2} P + m^2(m^2 - 4t) \frac{s}{s-4m^2} - \frac{4m^2 t^2}{s-4m^2})}}$$

Functional representation

Chen's definition of iterated integrals: $\gamma : [0, 1] \rightarrow M$ $x_i = \gamma(0)$ $x_f = \gamma(1)$

[Chen '77]

Martijn

$$\begin{aligned} I_\gamma(\omega_1, \dots, \omega_k; \lambda) &= \int_0^\lambda d\lambda_1 f_1(\lambda_1) \int_0^{\lambda_1} d\lambda_2 f_2(\lambda_2) \dots \int_0^{\lambda_{k-1}} d\lambda_k f_k(\lambda_k) \\ &= \int_0^\lambda d\lambda_1 f_1(\lambda_1) I_\gamma(\omega_2, \dots, \omega_k; \lambda_1), \end{aligned}$$

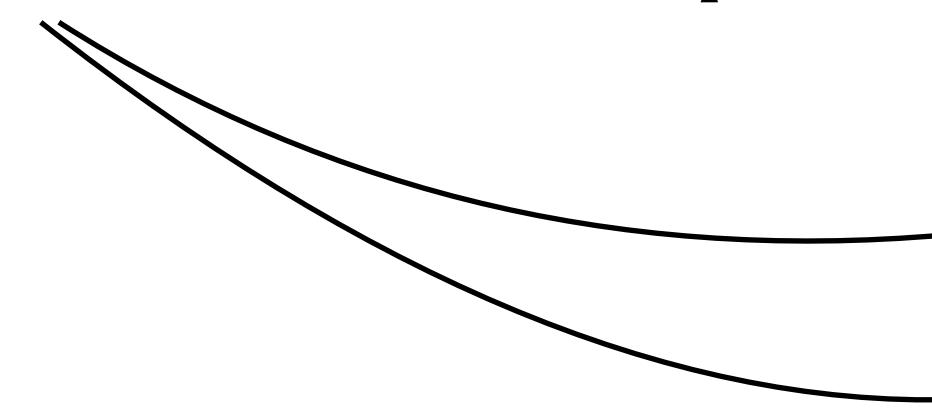


Goncharov's Polylogarithms (MPLs)

Elliptic Polylogarithms (eMPLs)

Numerical evaluation

Numerical evaluation of the functions at various phase-space points



Multiple polylogarithms
Elliptic integrals

GINAC PolyLogTools

[Vollinga, Weinzierl, '05] [Duhr, Dulat, '19]

Local series expansion methods

[Moriello, '19] [Hidding, '21]



Iterated integrals
Canonical DE

Ya-Qing,
Xiao,
Felix

Ofcourse PySecDec, Fiesta, Feyntrhop
Auxiliary mass flow methods

[Liu, Ma, '22]

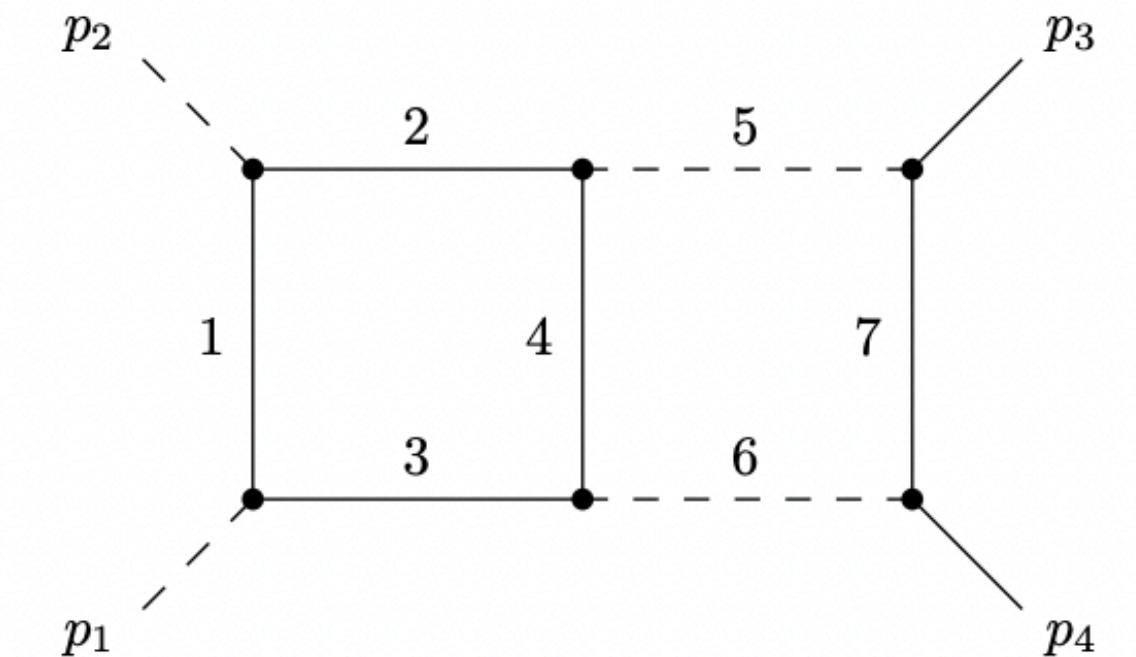
Some selected examples

Scattering amplitudes ($t\bar{t}$)

[Badger, E. Chaubey, Hartanto, Marzucca, '21]

Needed for precise understanding of top quarks

- Computation of helicity amplitudes for top-quark pair production involves elliptic integral contributions
- Topbox evaluates to 3 different elliptic curves
- Faster numerical evaluation of elliptic functions needs optimisation of the basis [E. Chaubey' 21]

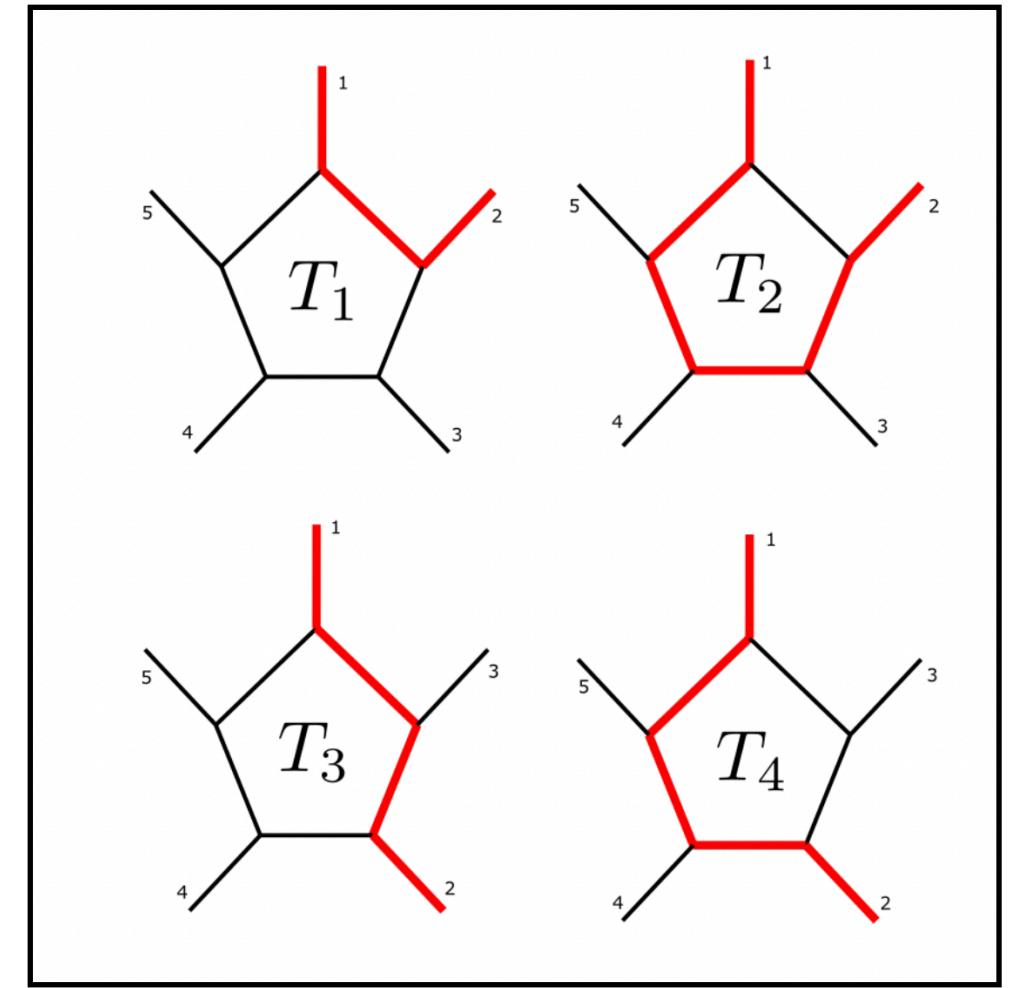


[Adams, E. Chaubey, Weinzierl, '18]

1-loop amplitude for tt-jet

[Badger, Becchetti, E. Chaubey, Marzucca, Sarandrea, '22]

- Analytic helicity amplitudes for 1-loop QCD corrections,
Previously missing ingredient for NNLO, expansion of
the 1-loop helicity amplitudes up to $O(\epsilon^2)$



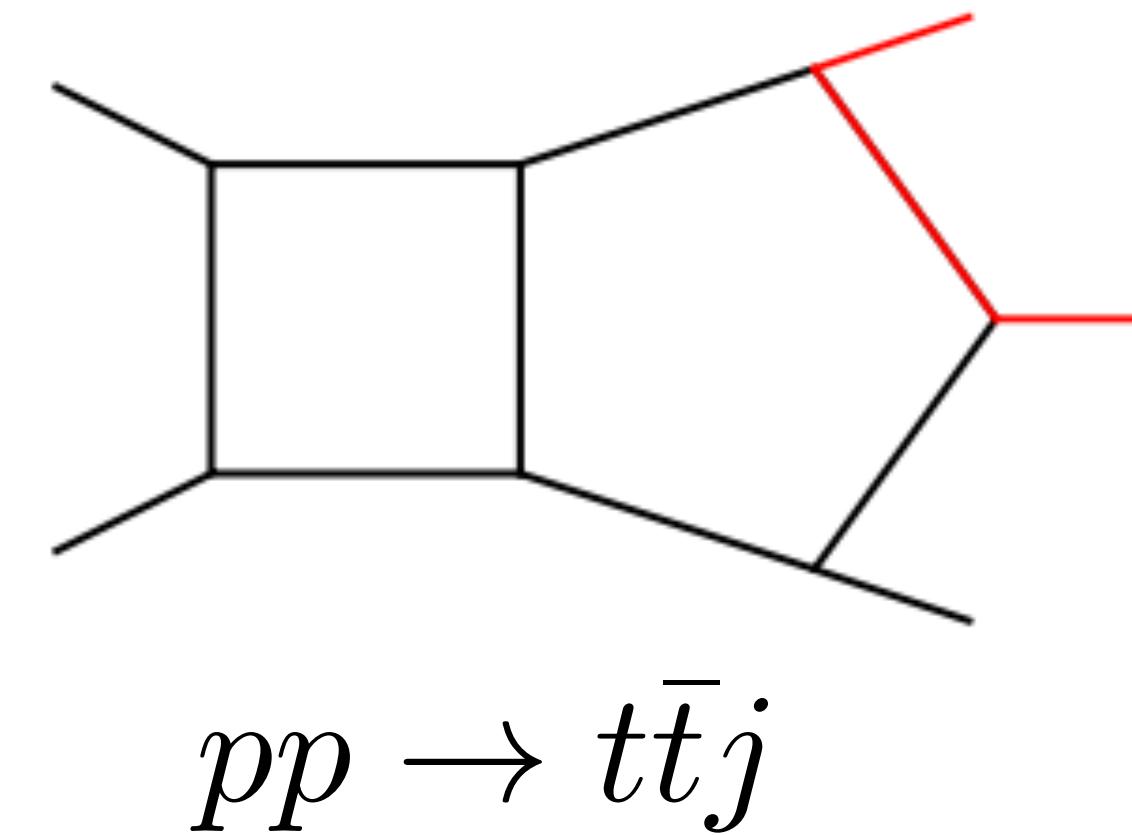
- Canonical form DE for all 130 MIs across 4 pentagon topologies
- Numerical solution using generalised power series expansion in DiffExp
- Analytic result of boundary constants

2-loop 5-point integrals

[Badger, Becchetti, E. Chaubey, Marzucca, '22]

- One of the topologies that appear in the NNLO corrections for tt-jet production
- 5 Mandelstam variables and mass dependence through top-quarks
- 88 master integrals
- Well-thought basis important to reconstruct analytically

$$\mu_{ij} = -k_i^{(-2\epsilon)} \cdot k_j^{(-2\epsilon)}$$



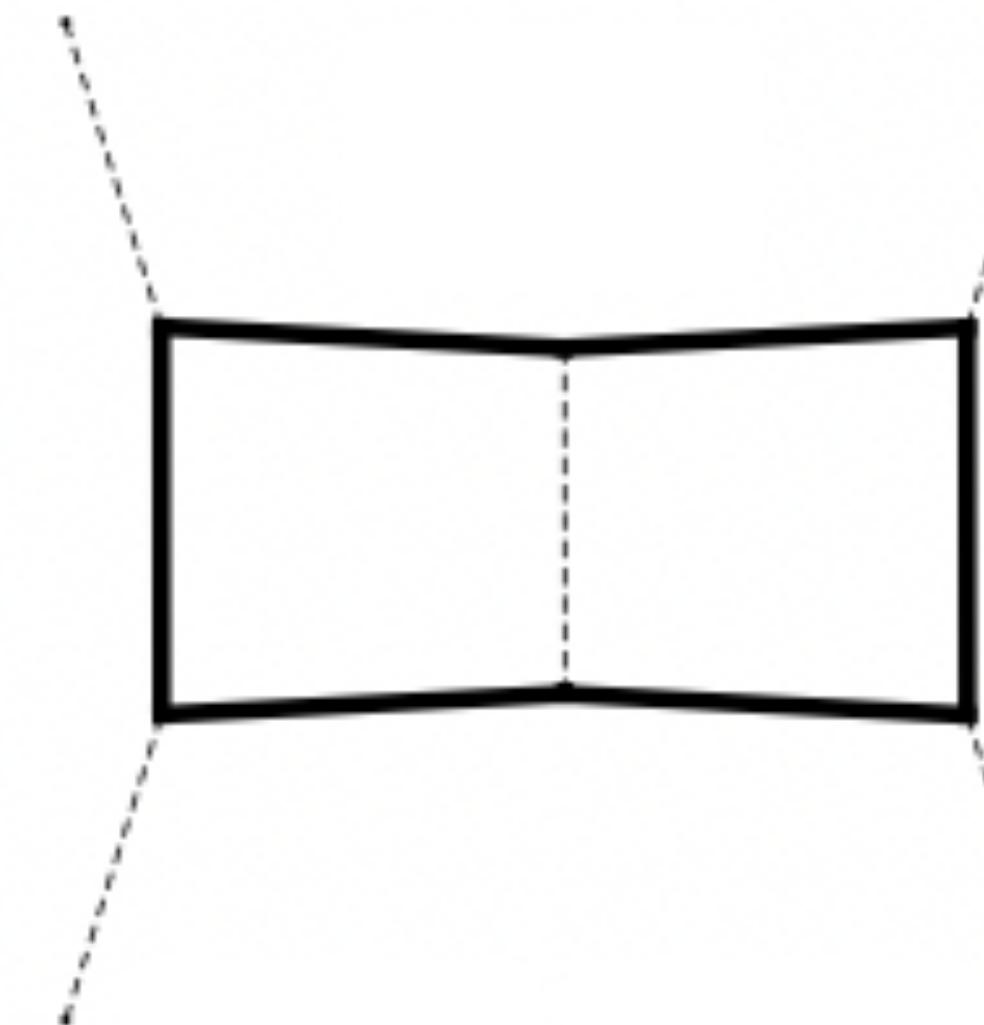
- Evaluation at 1 boundary point with a precision of $O(100)$ using AMFlow
- Integration of analytic differential equations using generalised power series expansion as implemented in DiffExp

Light-by-light scattering

[AH, E. Chaubey, Shao, '23]

$$\gamma(p_1, \lambda_1) + \gamma(p_2, \lambda_2) + \gamma(p_3, \lambda_3) + \gamma(p_4, \lambda_4) \rightarrow 0$$

- 2-loop QCD & QED corrections in the ultra-relativistic limit ($s, t, u \gg m^2$, massless internal lines)
[Bern, De Freitas, Dixon, Ghinculov & Wong, 2001]
- Two-loop corrections to light-by-light scattering in supersymmetric QED
[Binoth, Glover, Marquard, & van Der Bij, 2002]



Amplitude computation

$$\mathcal{M} = \varepsilon_{1,\mu_1} \varepsilon_{2,\mu_2} \varepsilon_{3,\mu_3} \varepsilon_{4,\mu_4} \mathcal{M}^{\mu_1\mu_2\mu_3\mu_4}(p_1, p_2, p_3, p_4)$$

$$\begin{aligned}\mathcal{M}^{\mu_1\mu_2\mu_3\mu_4} = & A_1 g^{\mu_1\mu_2} g^{\mu_3\mu_4} + A_2 g^{\mu_1\mu_3} g^{\mu_2\mu_4} + A_3 g^{\mu_1\mu_4} g^{\mu_2\mu_3} + \sum_{j_1,j_2=1}^3 (B_{j_1 j_2}^1 g^{\mu_1\mu_2} p_{j_1}^{\mu_3} p_{j_2}^{\mu_4} \\ & + B^2 j_1 j_2 g^{\mu_1\mu_3} p_{j_1}^{\mu_2} p_{j_2}^{\mu_4} + B^3 j_1 j_2 g^{\mu_1\mu_4} p_{j_1}^{\mu_2} p_{j_2}^{\mu_3} + B^4 j_1 j_2 g^{\mu_2\mu_3} p_{j_1}^{\mu_2} p_{j_2}^{\mu_4} \\ & + B^5 j_1 j_2 g^{\mu_2\mu_4} p_{j_1}^{\mu_1} p_{j_2}^{\mu_3} + B^6 j_1 j_2 g^{\mu_3\mu_4} p_{j_1}^{\mu_1} p_{j_2}^{\mu_2}) \\ & + \sum_{j_1,j_2,j_3,j_4=1}^3 C_{j_1 j_2 j_3 j_4} p_{j_1}^{\mu_1} p_{j_2}^{\mu_2} p_{j_3}^{\mu_3} p_{j_4}^{\mu_4}.\end{aligned}$$

Number of independent functions

$$A_1(s, t, u)$$

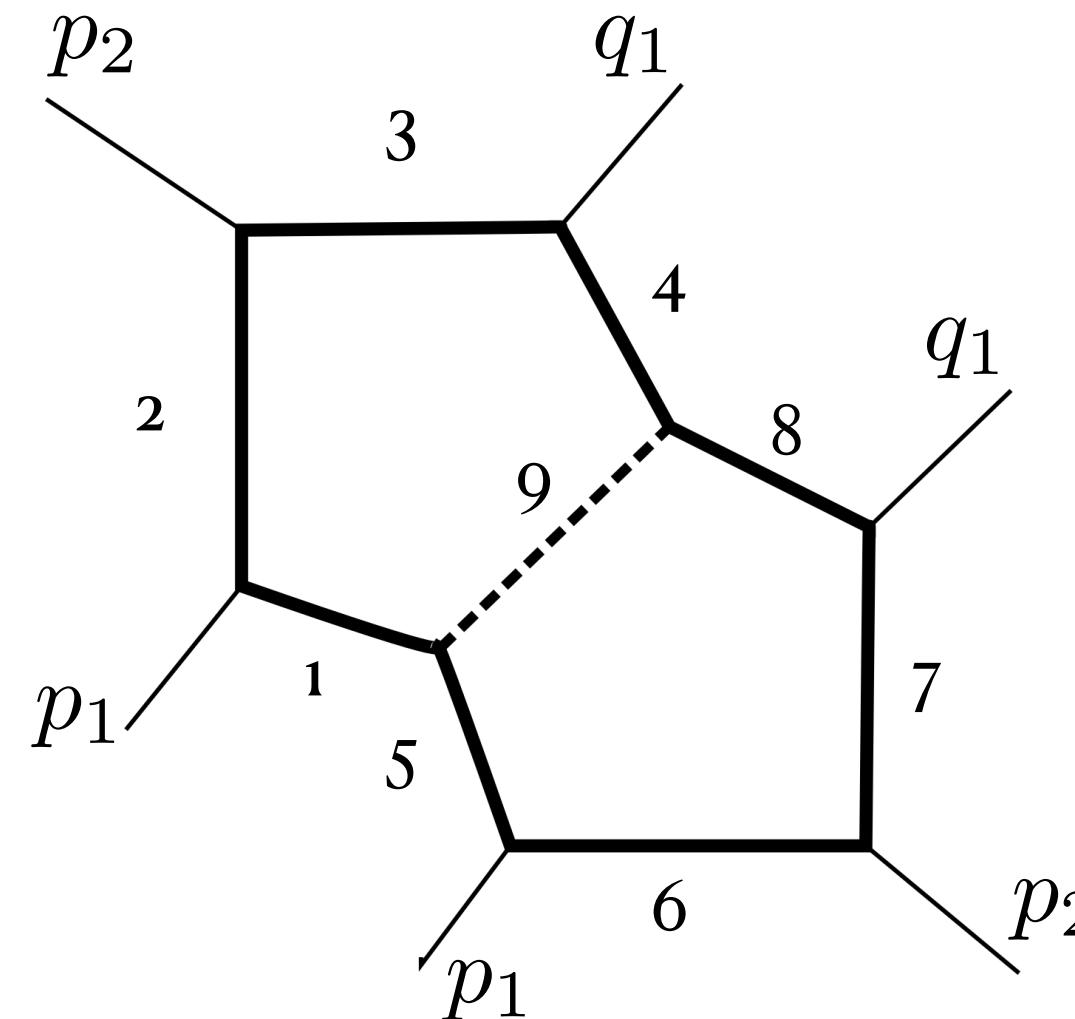
$$B_{11}^1(s, t, u)$$

$$C_{2111}(s, t, u)$$

$\varepsilon_j \cdot p_j = 0$
Bose symmetry
Gauge symmetry

[Binoth, Glover, Marquard, & van Der Bij, 2002]

Reduction to MIs & Simplification



60 diagrams in total

7798 integrals before IBP

18 top-level sectors

Can be mapped into the 2-loop diagram shown on the right

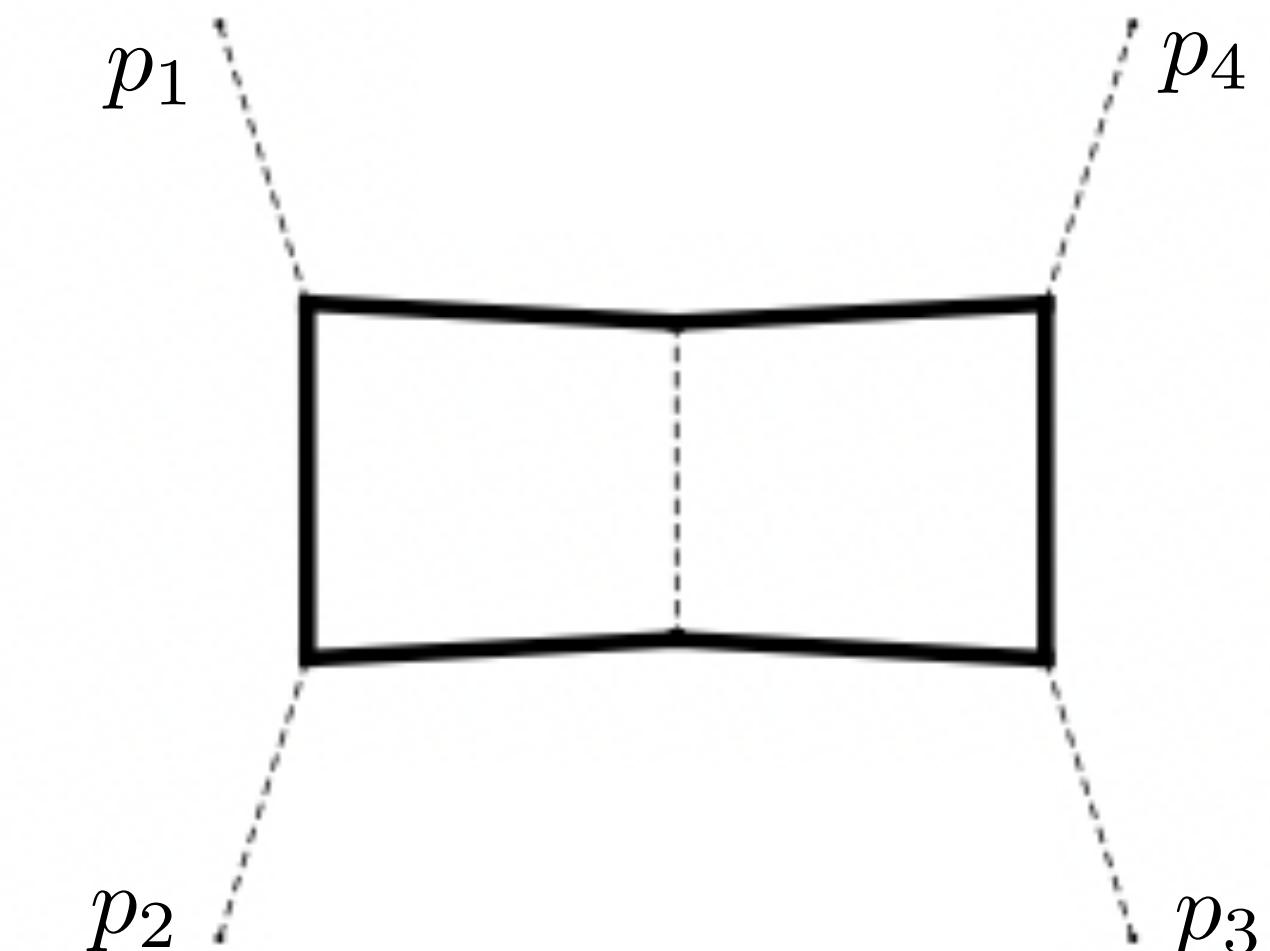
$$I_{a_1, \dots, a_9} = \left(\frac{e^{\epsilon \gamma_E}}{i\pi^{\frac{d}{2}}} \right)^2 \int \prod_{i=1}^2 d^d k_i \frac{D_4^{a_4} D_6^{a_6}}{D_1^{a_1} D_2^{a_2} D_3^{a_3} D_5^{a_5} D_7^{a_7} D_8^{a_8} D_9^{a_9}},$$

$$k_1^2 - m_t^2, (k_1 + p_1)^2 - m_t^2, (k_1 + p_1 + p_2)^2 - m_t^2, (k_1 + p_1 + p_2 + p_3)^2 - m_t^2,$$

$$(k_2)^2 - m_t^2, (k_2 + p_1)^2 - m_t^2, (k_2 + p_1 + p_2)^2 - m_t^2, (k_2 + p_1 + p_2 + p_3)^2 - m_t^2,$$

$$(k_2 - k_1)^2$$

LiteRed (FiniteFlow), KIRA
[\[Lee, '13\]](#) [\[Peraro, '19\]](#) [\[Klappert, Lange, Maierhöfer, Usovitsch,](#)



Analytic computation of the MIs

- 29 MIs; use of differential equations;
- Choice of a canonical basis [Caron-huot, Henn, 14]

$$\begin{aligned}\partial_s \vec{f} &= \epsilon A_s \vec{f} \\ \partial_t \vec{f} &= \epsilon A_t \vec{f}\end{aligned}$$

- Square roots:

$$\sqrt{s(s - 4m^2)} \quad \sqrt{t(t - 4m^2)} \quad \sqrt{st(st - 4m^2(s + t))} \quad \sqrt{s(m^4s - 2m^2t(s + 2t) + st^2)}$$

- Choice of variables

$$s = -\frac{4(w - z)^2}{(1 - w^2)(1 - z^2)} \quad t = -\frac{(w - z)^2}{wz} \quad m^2 = 1$$

$$\sqrt{-2wz + z^2 + w^4z^2 - 2w^3z^3 + w^2(1 + z^2 + z^4)} \quad \text{Non-rationalizable??..}$$

The alphabet

$$sq = \sqrt{-2wz + z^2 + w^4z^2 - 2w^3z^3 + w^2(1 + z^2 + z^4)} \quad [\text{Caron-huot, Henn, 14}]$$

$$\begin{aligned} & 1-w, 1+w, 1-wz, w-z, w, w+z, 1+w-z+wz, 1-w+z+wz, 1+wz, 1-z, 1+z, \\ & z, \frac{-sq + w - z - wz + w^2z - wz^2}{sq + w - z - wz + w^2z - wz^2}, \frac{-sq + w^2 - 3wz + z^2}{sq + w^2 - 3wz + z^2}, \frac{-1 - sq + w^2z - wz^2}{-1 + sq + w^2z - wz^2}, \\ & \frac{1 - sq + w^2z - wz^2}{1 + sq + w^2z - wz^2} \end{aligned}$$

6 master integrals containing all 4 square roots in the integrand at weight-4.

Solving the canonical master integrals

We keep the analytic result in terms of iterated integrals with dog one-forms constructed using

[Heller, von Manteuffel, Schabinger, '20]

We convert all the integral in terms of 1-dimensional integrations

[Caron-huot, Henn, 14] [Chicherin, Sotnikov '21][Chicherin, Sotnikov, Zoia '22]

We also express the first two orders in terms of logs and classical polylogs by matching symbols

[Duhr, Gangl, Rhodes, '11]

$$\int_{\gamma} I(\omega_1 \dots \omega_4; \lambda) = \int_0^{\lambda} d\lambda_1 f_1(\lambda_1) \int_0^{\lambda_1} d\lambda_2 f_2(\lambda_2) \underbrace{\int_0^{\lambda_2} d\lambda_3 f_3(\lambda_3) \int_0^{\lambda_3} d\lambda_4 f_4(\lambda_4)}_{\{\text{Log}^2(z), \text{Li}_2(z)\}}$$

Numerical evaluation

Physical phase-space regions of interest

$$0 < w < 1 \quad \& \quad (0 < z < w \mid w < z < 1)$$

$$0 < w < 1 \quad \& \quad (1 < z < \frac{1}{w} \mid z > \frac{1}{w})$$

$$0 < w < 1 \quad \& \quad (-1 < z < -w \mid -w < z < 0)$$

We obtain different analytic representations of the results valid in different regions

Outputs



- *We obtained a completely analytic representation for the squared matrix element at 2-loop*
- *NLO cross section with massive contributions for light-by-light scattering at UPC within reach!*

Take away

- Computation of higher order perturbative corrections important
- Analytic solution of multi-loop integrals requires understanding the **mathematical structure**
- With the inclusion of masses in the loop, the analytic structure starts becoming **more complicated**
- With the inclusion of more legs, often one needs go beyond current mathematical understand; often also desirable to **combine analytic as well as numerical** techniques