

MATHEMATICAL AND PHYSICAL STRUCTURES OF RATIONAL FUNCTIONS IN SCATTERING AMPLITUDES

Title for Mathematicians: Vector Spaces over Fraction Fields of Polynomial Quotient Rings

Giuseppe De Laurentis

Paul Scherrer Institut / University of Edinburgh

arXiv:2203.04269

(GDL, B. Page)

arXiv:2305.17056

(S. Abreu, GDL, H. Ita, M. Klinkert, B. Page, V. Sotnikov)

MathemAmplitudes 2023 - Padova



INTRODUCTION



SCATTERING AMPLITUDES

- Amplitude (integrands) can be written as

$$A(\lambda, \tilde{\lambda}, \ell) = \sum_{\substack{\Gamma, \\ i \in M_\Gamma \cup S_\Gamma}} c_{\Gamma,i}(\lambda, \tilde{\lambda}, \epsilon) \frac{m_{\Gamma,i}(\lambda \tilde{\lambda}, \ell)}{\prod_j \rho_{\Gamma,j}(\lambda \tilde{\lambda}, \ell)} \xrightarrow{\int d^D \ell} \sum_{\substack{\Gamma, \\ i \in M_\Gamma}} \textcolor{red}{c}_{\Gamma,i}(\lambda, \tilde{\lambda}, \epsilon) \textcolor{orange}{I}_{\Gamma,i}(\lambda \tilde{\lambda}, \epsilon)$$

- For a suitable choice of integrands, we get:

$$\textcolor{red}{c}_{\Gamma,i}(\lambda, \tilde{\lambda}, \epsilon) = \frac{\sum_{k=0}^{\text{finite}} \textcolor{red}{c}_{\Gamma,i}^{(k)}(\lambda, \tilde{\lambda}) \epsilon^k}{\prod_j (\epsilon - a_{ij})}, \quad \text{with} \quad a_{ij} \in \mathbb{Q}$$

Some notation:

- Γ : topologies
- M_Γ : masters
- S_Γ : surface terms
- Spinors: $\lambda_i = |i\rangle, \tilde{\lambda}_i = [i|$
- 4-momenta: $\lambda \tilde{\lambda} = \not{p}$
- Loop D -momenta: ℓ



OUTLINE

Disclaimer: I will focus on the $c_{\Gamma,i}^{(k)}(\lambda, \tilde{\lambda})$ – call them $c_i(\lambda, \tilde{\lambda})$ for short, nevertheless some concepts can be extended to whole amplitudes.

We will discuss:

1. Where are the **poles**? What is their order?
2. Relation between pole structure, and **analytic constraints** (e.g. partial fraction decomposition)
3. Reconstructing **sets of functions**: Why is it easier than reconstructing individual ones?
4. Example process @ 2-loop: $pp \rightarrow \gamma\gamma\gamma$ and $pp \rightarrow Wjj$ (preliminary)

Key take away:

What is important is not the number of variables,
but the size of the parametrization (a.k.a. ansatz), and our ability to constrain it.



POLYNOMIAL QUOTIENT RINGS

- Let us start from the polynomial ring of spinor components

$$S_n = \mathbb{F} [|1\rangle, [1|, \dots, |n\rangle, [n|]$$

the field \mathbb{F} can be any of $\mathbb{Q}, \mathbb{R}, \mathbb{C}, \mathbb{F}_p, \mathbb{Q}_p, \dots$

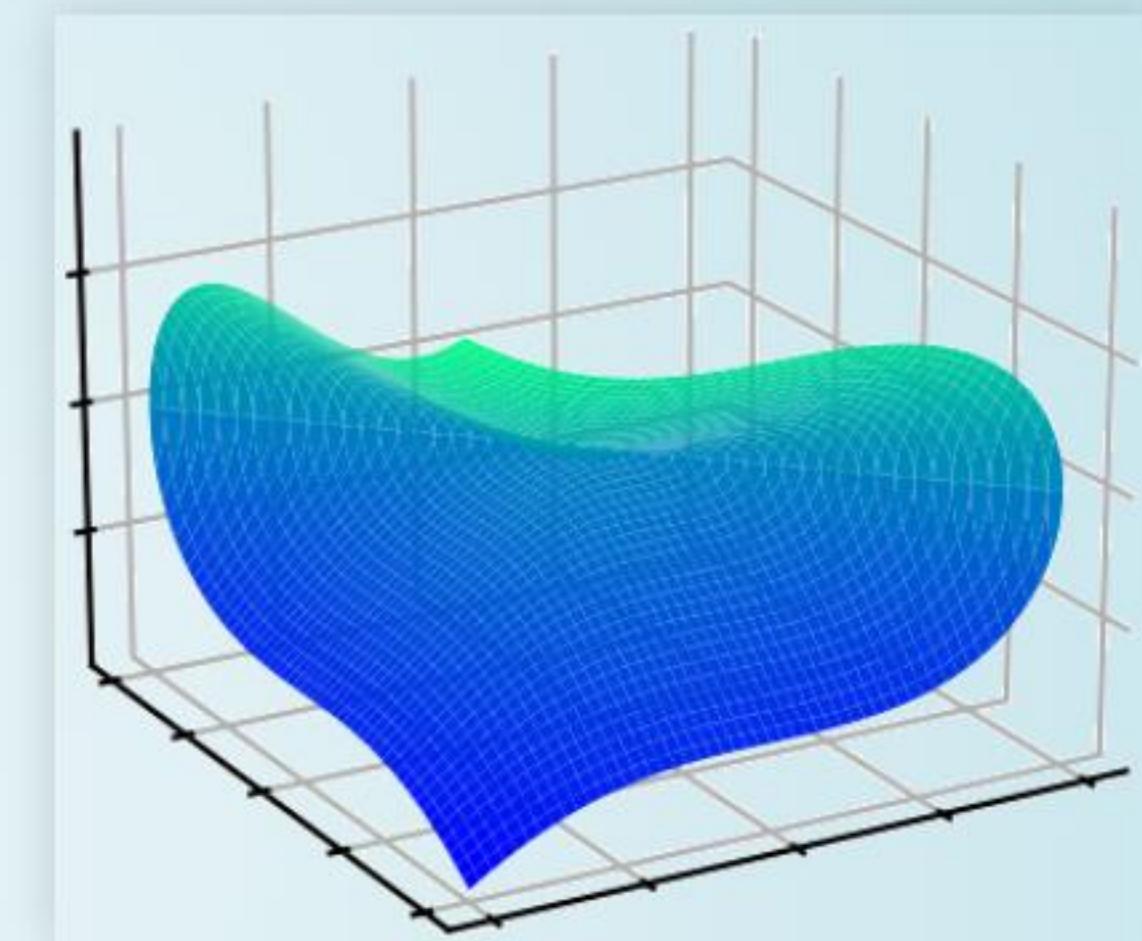
- Define the momentum-conservation ideal as

$$J_{\Lambda_n} = \left\langle \sum_i |i\rangle[i| \right\rangle_{S_n}$$

physically, two polynomials p and q are equivalent if $p - q \in J_{\Lambda_n}$

- This defines the needed polynomial **quotient** ring^{*}: $R_n = S_n / J_{\Lambda_n}$

$c_i(\lambda, \tilde{\lambda})$ at n -point belong to the Field of Fractions[†] of R_n



Artist's Impression of $V(J_{\Lambda_n})$
I can't draw in $4n$ dims!

^{*} R_4 is "weird" (not a UFD), but it proves that polynomial rings are insufficient; [†] The field of fractions of R_3 does not exist.



THE POLE STRUCTURE

PRIME IDEALS & IRREDUCIBLE VARIETIES

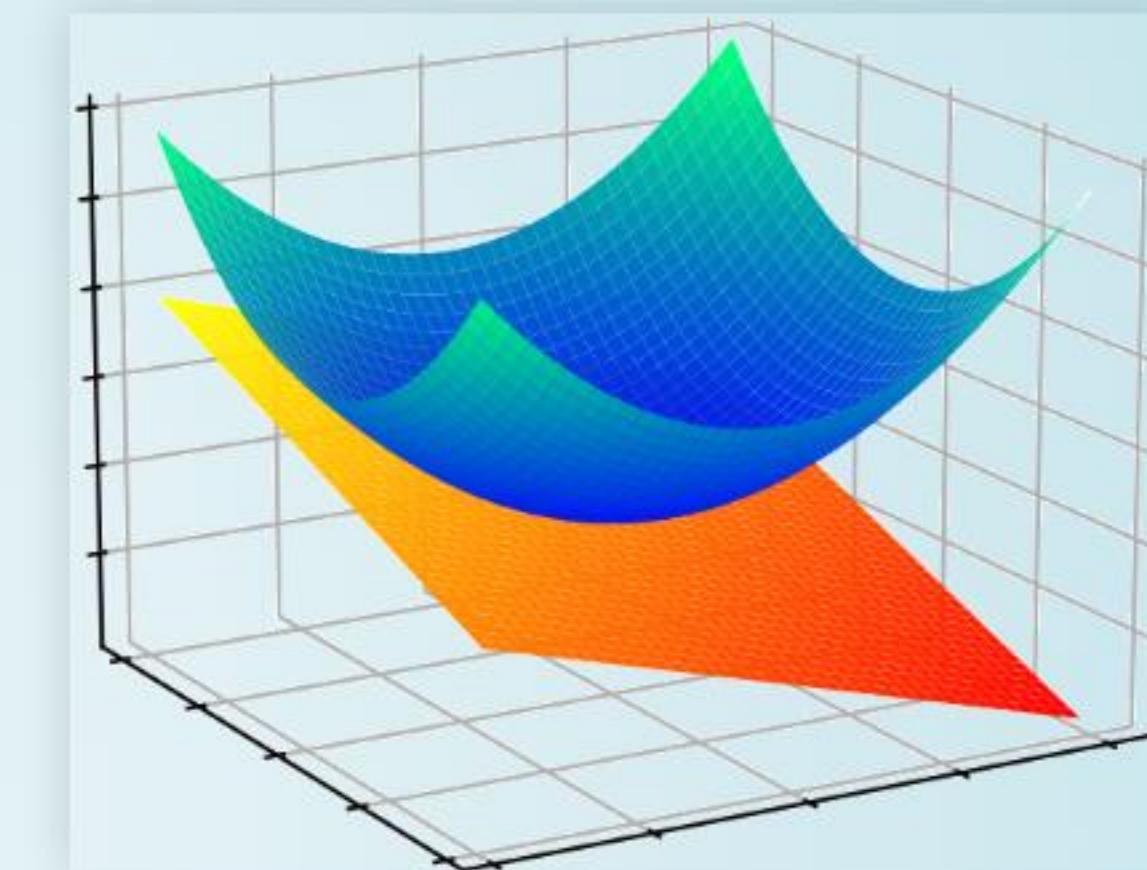
- Let us consider a very simple example

$$iA_{g^-g^-g^+g^+}^{\text{tree}} = \frac{\langle 12 \rangle^3}{\langle 23 \rangle \langle 34 \rangle \langle 41 \rangle} = \frac{[34]^3}{[12][23][41]}$$

is, say, $\langle 23 \rangle$ a pole of this amplitude?

- The question is ill posed!
 $\langle 23 \rangle$ does not identify an irreducible variety in R_4 .
Compute **primary decompositions**, such as

$$\langle\langle 23 \rangle\rangle_{R_4} = \langle\langle 23 \rangle, [14] \rangle_{R_4} \cap \langle\langle 12 \rangle, \langle 13 \rangle, \langle 14 \rangle, \langle 23 \rangle, \langle 24 \rangle, \langle 34 \rangle \rangle_{R_4}$$



Artist's Impression of $V(\langle\langle 23 \rangle\rangle_{R_4})$
as the union of two irreducibles

On the **first branch** there is a simple pole, on the **latter branch** the amplitude is regular.

Poles & Zeros \Leftrightarrow Irreducible Varieties \Leftrightarrow Prime Ideals

Physics

Geometry

Algebra



CONSTRAINTS FROM POLES

BOOTSTRAPPING TREES (?)

- The degree of divergence / vanishing on various surfaces imposes strong constraints, e.g.

$$A_{q^+g^+g^+\bar{q}^-g^-g^-}^{\text{tree}} = \frac{\mathcal{N}(m.d.=6, p.w.=[-1,0,0,1,0,0])}{\langle 12 \rangle \langle 23 \rangle \langle 34 \rangle [45][56][61] s_{345}}$$

- Pretend this is an unknown integral coefficient, \mathcal{N} has 143 free parameters.
- List the various prime ideal, such as

$$\langle \langle 12 \rangle, \langle 23 \rangle, \langle 13 \rangle \rangle, \langle |1\rangle \rangle, \langle \langle 12 \rangle, |1 + 2|3] \rangle, \dots$$

and impose that \mathcal{N} vanishes to the correct order. We determine it up to an overall constant.

GDL, Page (22)

- Likewise, the ansatz for $A_{g^+g^+g^+g^-g^-g^-}^{\text{tree}}$ shrinks $1326 \rightarrow 1$, etc..

*Effectively, we can **compute** trees, just from their poles orders.
Note: compared to BCFW there is no information about residues.*



PARTIAL FRACTION DECOMPOSITIONS

- For true integral coefficients, we can't rely on the Ansatz to shrink to an overall constant.
- Partial fraction decompositions (PFDs) are a popular method to tame algebraic complexity.
- In my opinion, a PFD algorithm needs
 1. to say if two poles W_a and W_b are separable into different fractions;
 2. ideally, to answer (1.) without having access to an analytic expression.
- Hilbert's nullstellensatz: if \mathcal{N} vanishes on all branches of $\langle W_a, W_b \rangle$, then the PFD is possible[†].
- Generalizing to powers > 1 can be done via symbolic powers and the Zariski-Nagata Theorem.
- Similarly, generalizing to non-radical ideals requires ring extensions.GDL, Page ('22)
Campbell, GDL, Ellis ('22)

Issue: evaluations on singular surfaces are expensive, but are not always needed!

Opportunity: we get more than partial fraction decompositions.

[†] $\langle W_a, W_b \rangle$ needs to be radical.



BEYOND PARTIAL FRACTIONS

- **Case 0:** the ideal does **not involve denominator factors**.

E.g. a 6-point function c_i has a pole at $\langle 1|2 + 3|4]$ but not at $\langle 4|2 + 3|1]$, yet it is regular on the irreducible surface $V(\langle\langle 1|2 + 3|4], \langle 4|2 + 3|1]\rangle)$. Then

$$c_i \sim \frac{\langle 4|2 + 3|1]}{\langle 1|2 + 3|4]} + \mathcal{O}(\langle 1|2 + 3|4]^0) \text{ instead of } c_i \sim \frac{1}{\langle 1|2 + 3|4]} + \mathcal{O}(\langle 1|2 + 3|4]^0)$$

- **Case 1:** the **degree of vanishing is non-uniform** across branches, for example:

$$\frac{s_{14} - s_{23}}{\langle 1|3 + 4|2]\langle 3|1 + 2|4]}$$

has a double pole on the first branch, and a simple pole on the second branch of

$$\langle\langle 1|3 + 4|2], \langle 3|1 + 2|4]\rangle_{R_6} = \langle\langle 13], [24]\rangle_{R_6} \cap \langle\langle 1|3 + 4|2], \langle 3|1 + 2|4], (s_{14} - s_{23})\rangle_{R_6}$$

- **Case 2:** ideal is **non-radical** (example on last slide)

$$\sqrt{\langle\langle 3|1 + 4|2], \Delta_{23|14|56}\rangle_{R_6}} = \langle\langle 3|1 + 4|2], s_{124} - s_{134}\rangle_{R_6}$$



ANALYTIC RECONSTRUCTION

CHOOSING INDEPENDENT FUNCTIONS

- The set c_i can be very large, so pick a set of independent ones, and write:

$$c_i = \tilde{c}_j M_{ji} \quad \text{with} \quad M_{ji} \in \mathbb{Q}$$

with \tilde{c} an independent subset of c .

$\implies M_{ji}$ is, up to a permutation of columns, in row reduced echelon form.

- We might as well use a set \tilde{c} which is not a subset of c , at the trivial cost of changing M_{ji} .
- Consider a PFD of one of the \tilde{c} 's:

$$\tilde{c}_i(\lambda, \tilde{\lambda}) = \frac{\mathcal{N}_i(\lambda, \tilde{\lambda})}{\prod_j W_j^{q_{ij}}(\lambda, \tilde{\lambda})} = \sum_k \frac{\mathcal{N}_{ik}(\lambda, \tilde{\lambda})}{\prod_j W_j^{q_{ijk}}(\lambda, \tilde{\lambda})} = \sum_k \tilde{c}_{ik}(\lambda, \tilde{\lambda})$$

We **cannot have** $\tilde{c}_i \in \text{span}(\tilde{c}_{j \neq i})$, but we **can have** $\tilde{c}_{ik} \in \text{span}(\tilde{c}_{j \neq i})$, for some, but not all, k .



LEAST LEAST-COMMON-DENOMINATOR

- In other words, the \tilde{c} span a vector space, and we should **consider one modulo the others**

$$\tilde{c}_i = \sum_{j \neq i} q_j \tilde{c}_j + \tilde{c}'_i$$

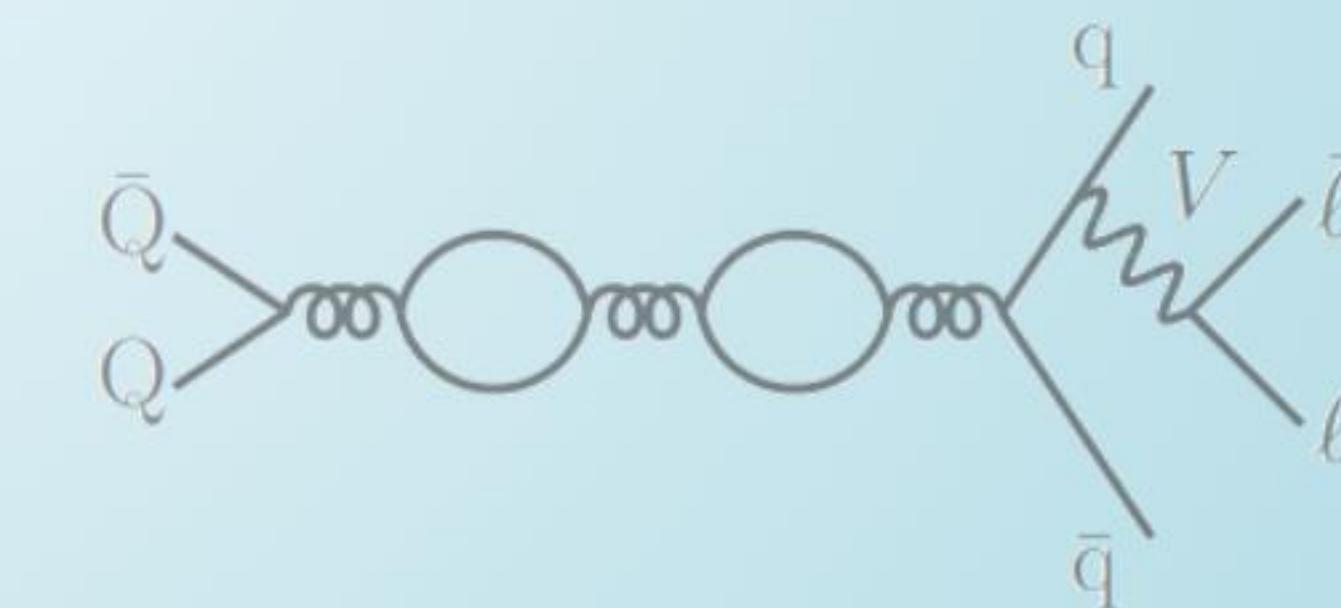
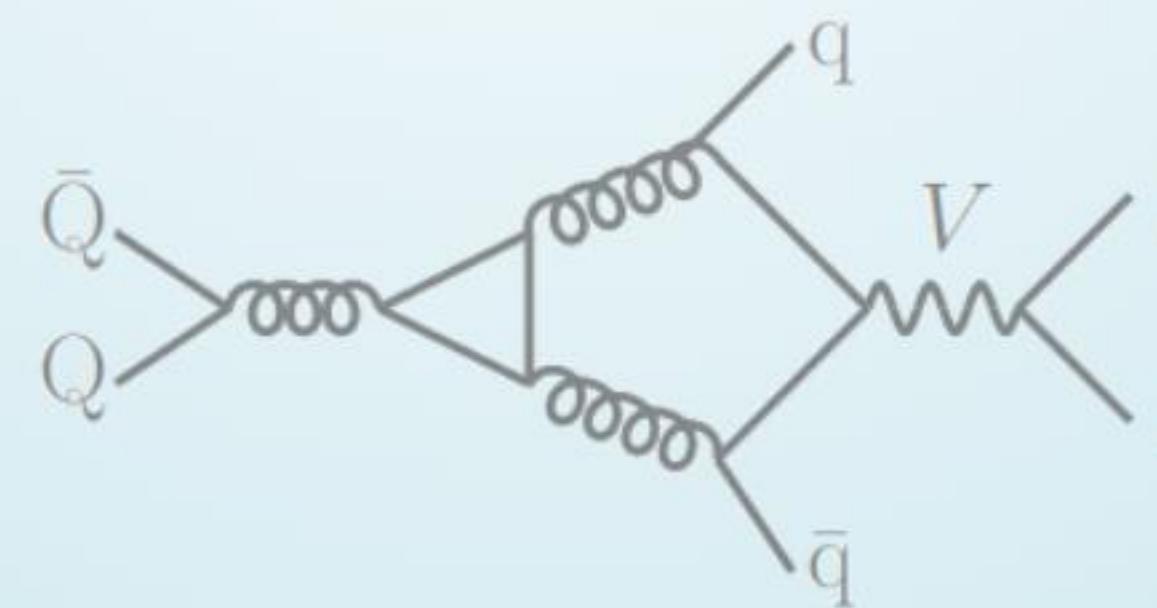
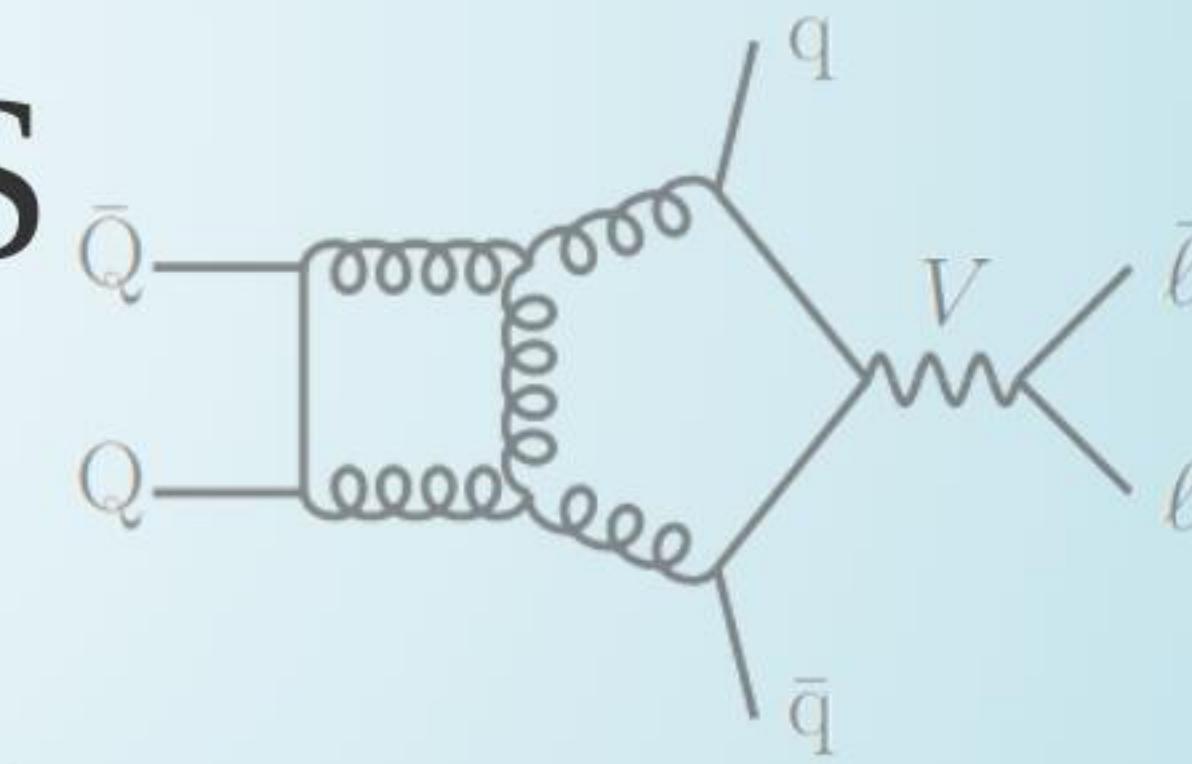
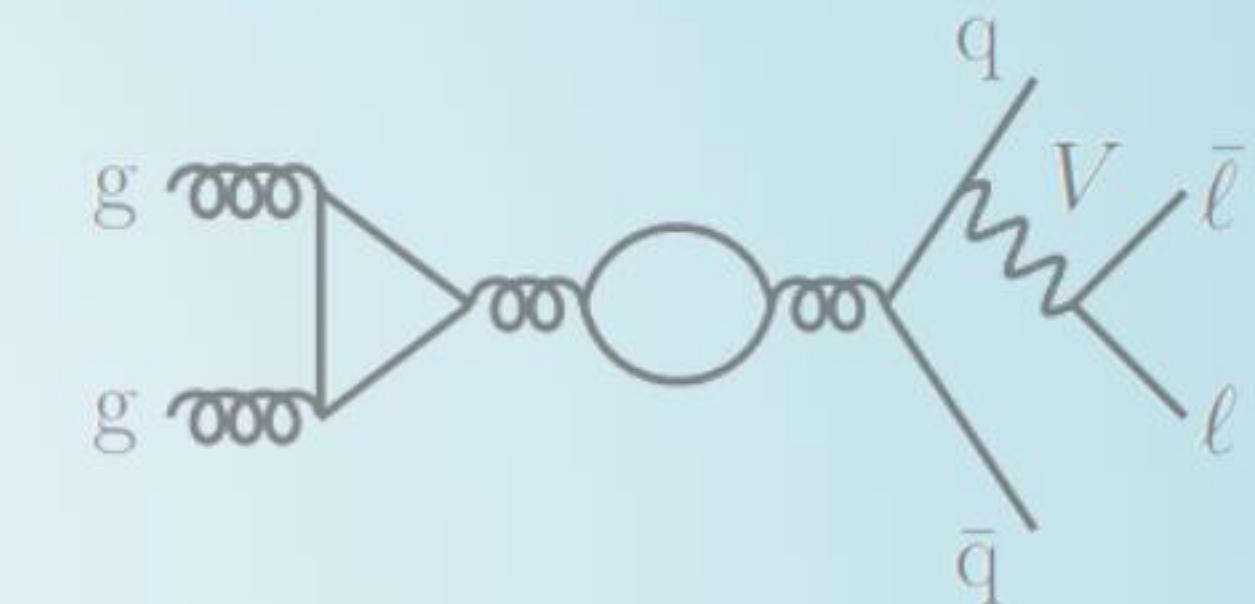
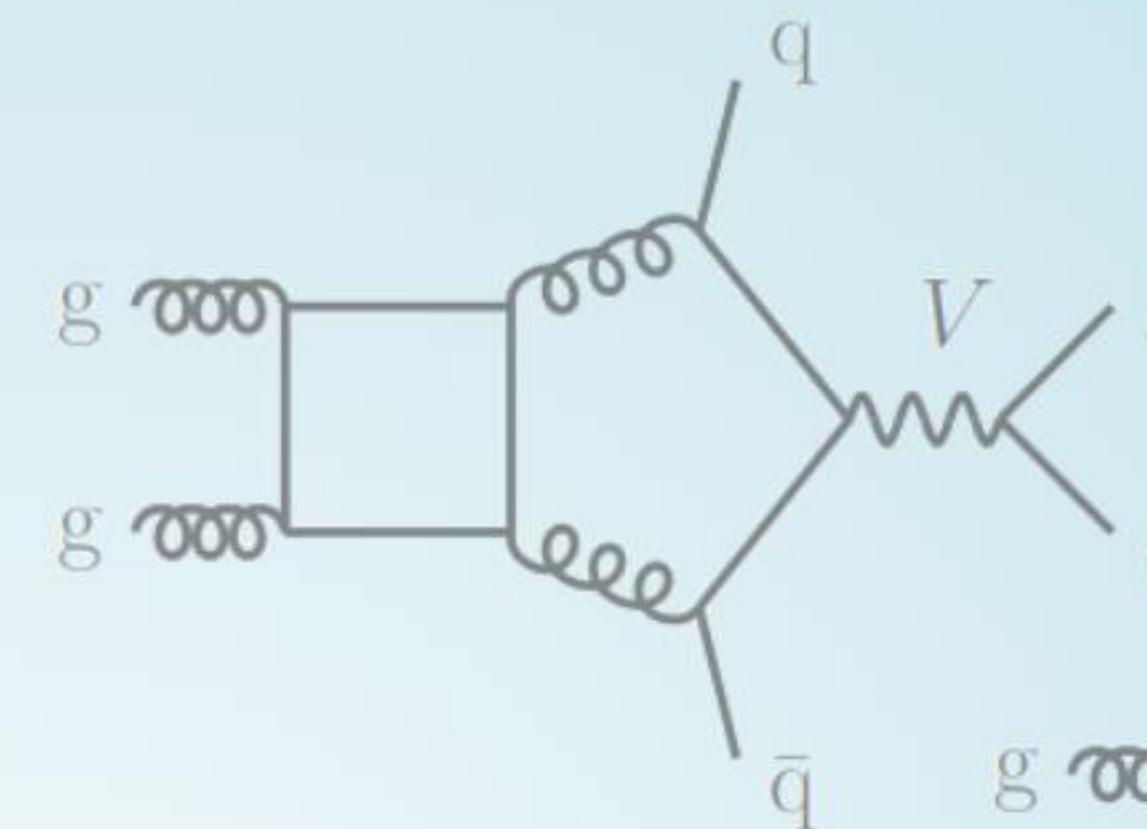
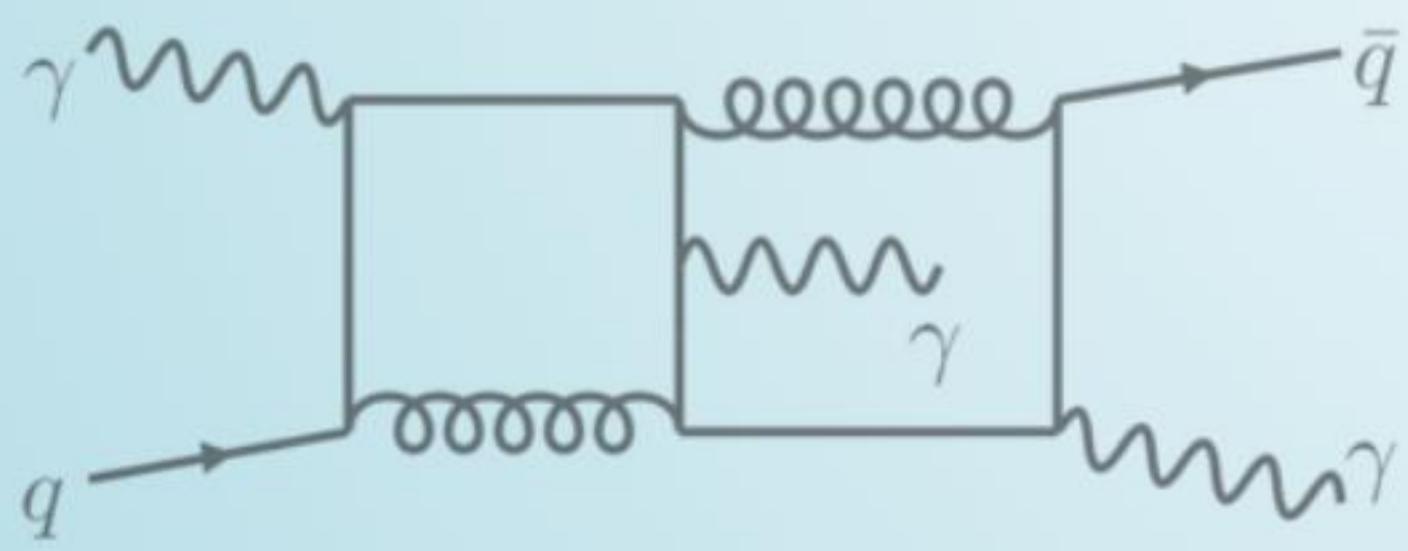
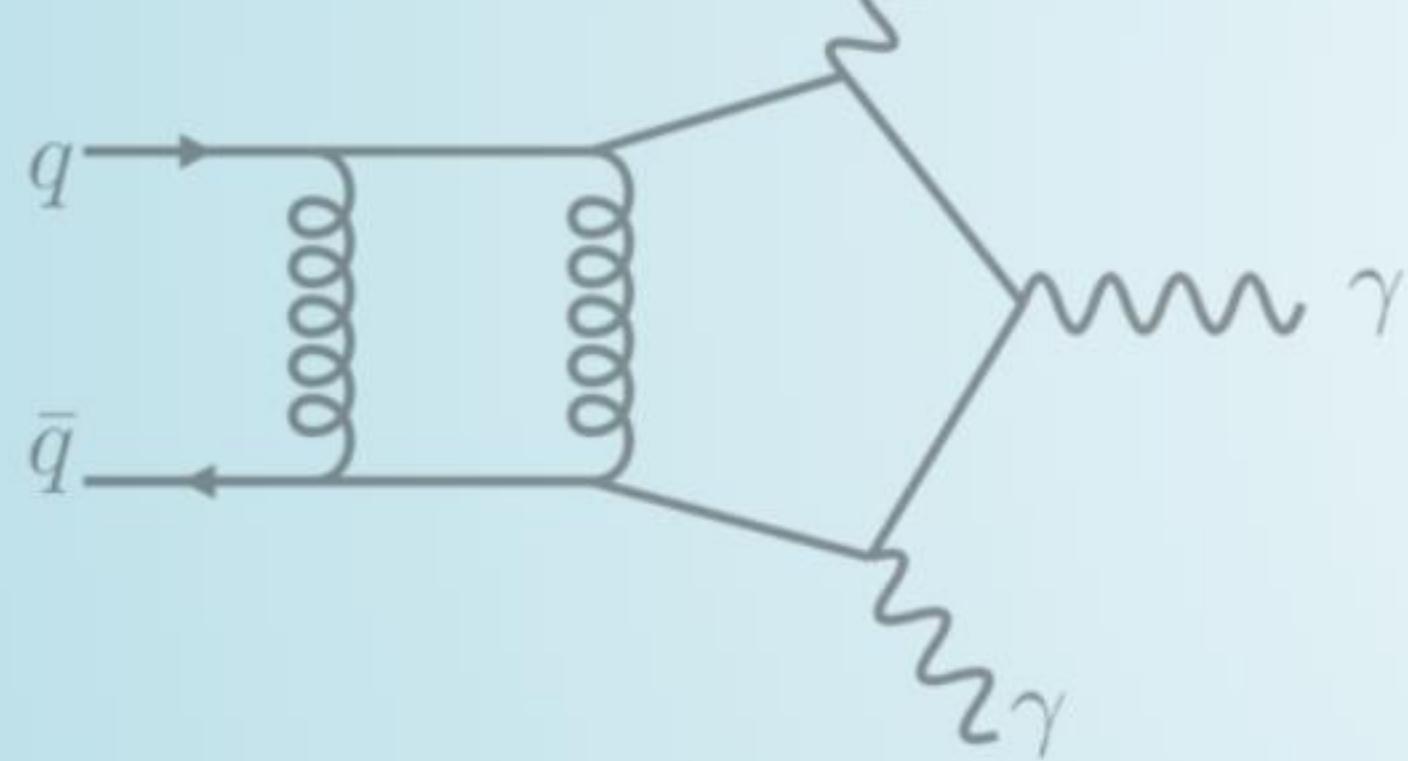
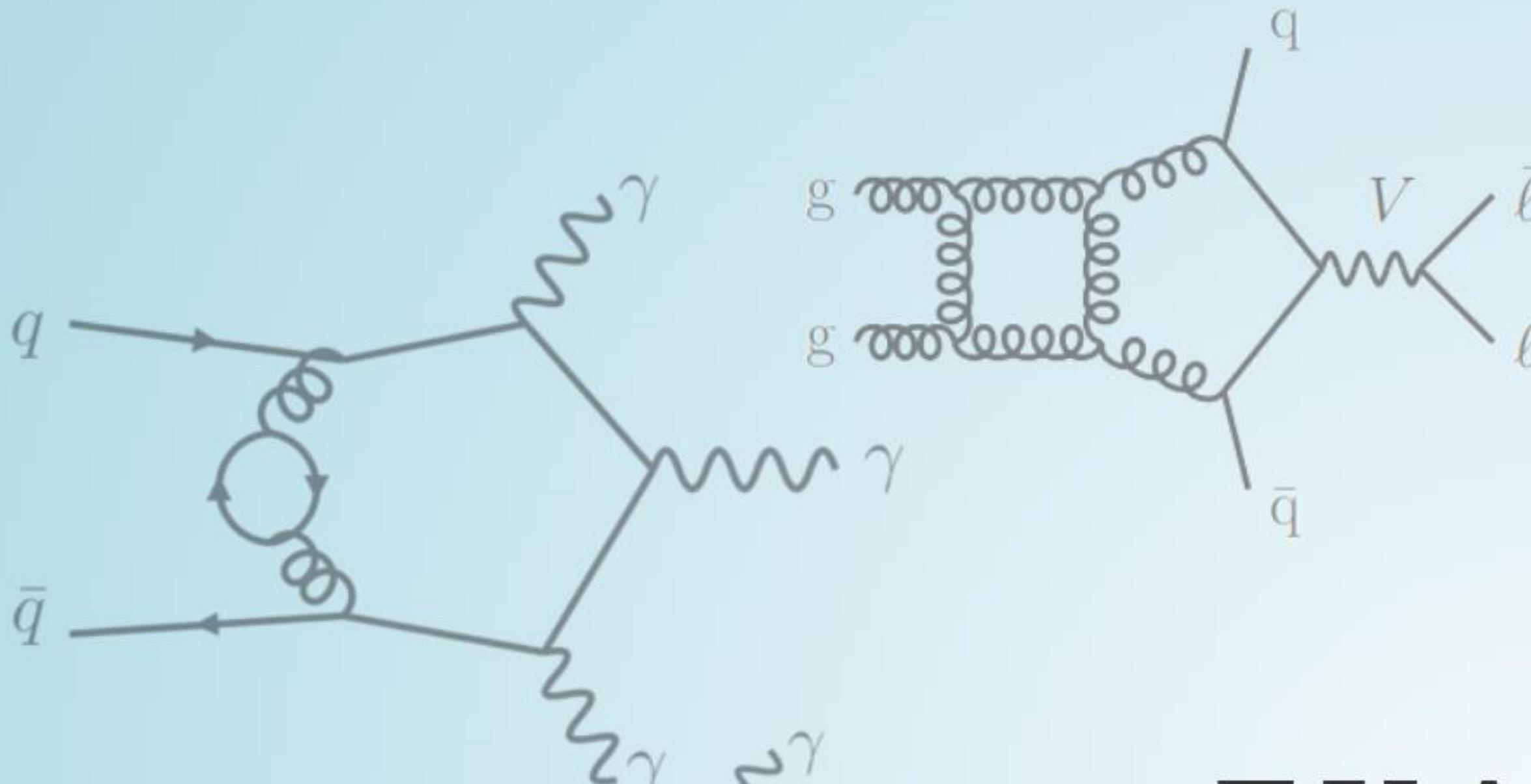
the basis function \tilde{c}_i can be replaced by \tilde{c}'_i without changing the vector space.

- In particular, \tilde{c}'_i needs not have all the poles of \tilde{c}_i , thus it can be much simpler.
In other words, **the LCD of \tilde{c}'_i can be smaller than that of \tilde{c}_i** .
- Brute-force search works well when an analytic expressions is available.
- In a future publication we will provide an algorithm based on finite field evaluations.

Reconstructing a set of c_i is not as bad as reconstructing the most complex function in the set.



EXAMPLE PROCESSES



THREE-PHOTON PRODUCTION AT TWO LOOPS

- The denominator factors W_j are conjectured to be restricted to the letters of the symbol alphabet

Abreu, Dormans, Febres Cordero, Ita, Page ('18)

$$\{W_j\} = \bigcup_{\sigma \in \text{Aut}(R_5)} \sigma \circ \{\langle 12 \rangle, \langle 1|2 + 3|1] \} \quad \text{Identical to 1-loop!}$$

- Advantage of spinor variables due to:

1. little group covariant LCD (no spurious poles); 2. avoiding parity even/odd split;

\Rightarrow in LCD form we would need **29 059** evaluations instead of **117 810** (with s_{ij}) for $\mathcal{R}_{2q3\gamma}^{(2)}$.

- To **avoid evaluations on singular surfaces**, use insights from physics (locality).
E.g. conjecture that no 5-point denominator has pairs of $\langle i|j + k|i]$, like at 1 loop.
- Remove some overlap with other \tilde{c} 's, obtain the \tilde{c}_i with higher degree LCD with **4 003** points.
- A posteriori, we find that for the c_i with highest degree LCD the following would have sufficed

$$\tilde{c}_i = \frac{\langle 13 \rangle [14]^2 \langle 24 \rangle \langle 34 \rangle [45]}{\langle 45 \rangle \langle 4|1 + 3|4]^3} - \frac{[14] \langle 25 \rangle \langle 34 \rangle^2 [45]}{\langle 45 \rangle^2 \langle 4|1 + 3|4]^2} - \frac{[14] \langle 24 \rangle \langle 34 \rangle \langle 35 \rangle}{\langle 45 \rangle^3 \langle 4|1 + 3|4]}$$



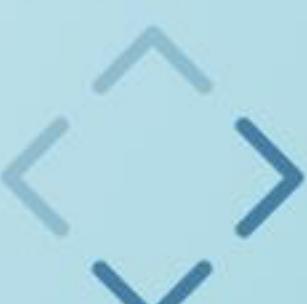
W+2-JETS: SIMPLIFICATION STRATEGY

0. Start from analytics of Abreu, Febres Cordero, Ita, Klinkert, Page, Sotnikov ('21) - 1.2GB of C++ source code.
1. Script to split up the expressions, and compile them ($\sim 20\text{GB}$ binaries) for evaluation over \mathbb{F}_p ;
2. Recombine the 3 projections $p_V \parallel p_1, p_V \parallel p_2, p_V \parallel p_3$ and reintroduce the little group factors to build 6-point spinor-helicity amplitudes (subject to degree bounds on $|5\rangle, [5|, |6\rangle, [6|]$);
3. Perform (rough) PFDs based on expected structures and fit the Ansatze.

Comparison of $q\bar{q} \rightarrow \gamma\gamma\gamma$ (in full color) to $pp \rightarrow Wjj$ (at leading color):

| Kinematics | # Poles (W) | LCD Ansatz | Partial-Fraction Ansatz | Rational Functions |
|------------------|-----------------|------------|-------------------------|----------------------|
| 5-point massless | 30 | 29k | 4k | $\sim 300\text{ KB}$ |
| 5-point 1-mass | >200 | >5M | $\sim 40\text{k}$ | $\sim 25\text{ MB}$ |

$$\{W_j\} = \bigcup_{\sigma \in \text{Aut}(R_6)} \sigma \circ \{\langle 12 \rangle, \langle 1|2 + 3|1], \langle 1|2 + 3|4], s_{123}, \Delta_{12|34|56}, \langle 3|2|5 + 6|4|3] - \langle 2|1|5 + 6|4|2] \}$$



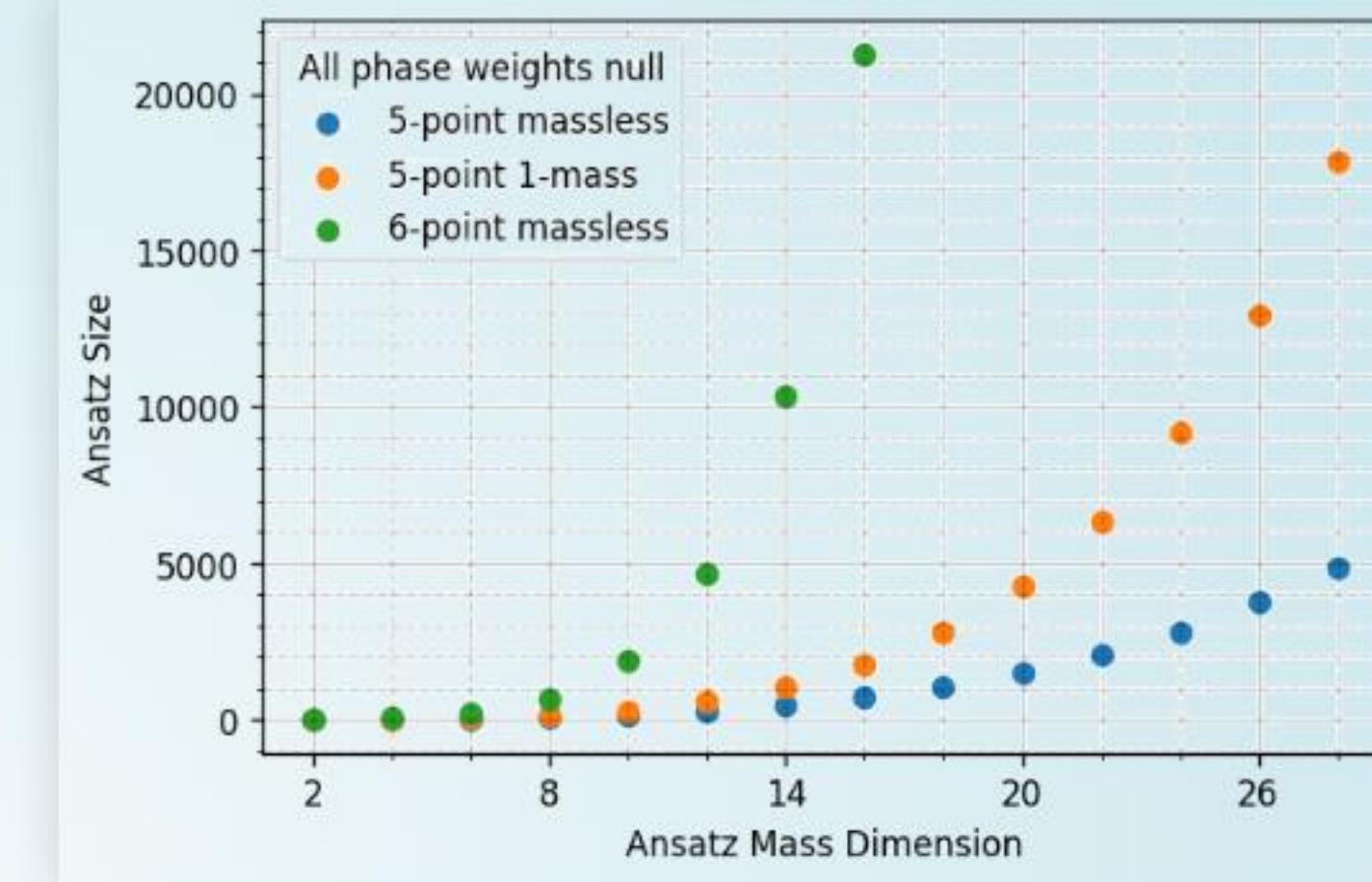
ANALYTIC STRUCTURES OF 2-LOOP 5-POINT 1-MASS AMPLITUDES

- The Ansatz size grows quickly with multiplicity (m) and mass dimension (d):

$$\binom{m(m-3)/2}{d/2}$$

is a lower bound.

GDL, Maître ('20)



- Compact residues for the new 2-loop (spurious?) pole, $\langle k|j|\not{p}_V|l|k\rangle - \langle j|i|\not{p}_V|l|j\rangle$, e.g.:

$$r_{\bar{u}^+ g^+ g^+ d^- (V \rightarrow \ell^+ \ell^-)}^{(5 \text{ of } 54)} = \frac{[12][23]\langle 24\rangle\langle 46\rangle^2\langle 1|2+3|4\rangle\langle 2|1+3|4\rangle}{\langle 12\rangle\langle 23\rangle\langle 56\rangle(\langle 3|2|5+6|4|3] - \langle 2|1|5+6|4|2])^2}$$

- The three mass Grams, $\Delta_{12|34|p_V}$, $\Delta_{14|23|p_V}$, behave analogously to one-loop amplitudes, e.g.:

$$r_{\bar{u}^+ g^- g^+ d^- (V \rightarrow \ell^+ \ell^-)}^{(73 \text{ of } 120)} = \frac{105}{128} \frac{\langle 2|1+4|3\rangle\langle 4|2+3|1\rangle\langle 6|1+4|5\rangle s_{14}s_{23}s_{56} (\mathbf{s}_{124} - \mathbf{s}_{134})(s_{123} - s_{234})(s_{25} + s_{26} + s_{35} + s_{36})}{\langle 3|1+4|2\rangle \Delta_{23|14|56}^4} + \\ \left[-6 \frac{[12]^2\langle 13\rangle[25]\langle 34\rangle\langle 36\rangle\langle 56\rangle[56](\mathbf{s}_{124} - \mathbf{s}_{134})}{\langle 3|1+4|2\rangle^5} \right] + \left[\right]_{1234 \rightarrow \overline{4321}} + \mathcal{O}\left(\frac{1}{\langle 3|1+4|2\rangle^4 \Delta_{23|14|56}^3}\right)$$



**THANK YOU
FOR YOUR ATTENTION!**

These slides are powered by:

[markdown](#), [html](#), [revealjs](#), [hugo](#), [mathjax](#), [github](#)

BACKUP SLIDES



FINITE REMAINDERS & THE RATIONAL / TRANSCENDENTAL SPLIT

- In general, in $D = 4 - 2\epsilon$, with *pure* master integrals $I_{\Gamma,i}$ we have

$$A_n^{\ell-loop}(\lambda, \tilde{\lambda}) = \sum_{\Gamma} \sum_{i \in M_{\Gamma}} \frac{c_{\Gamma,i}(\lambda, \tilde{\lambda}, \epsilon) I_{\Gamma,i}(\lambda \tilde{\lambda}, \epsilon)}{\prod_j (\epsilon - a_{ij})}, \quad \text{with } a_{ij} \in \mathbb{Q}$$

- For NNLO applications, we are interested in the *finite remainder*

$$\mathcal{A}_R^{(2)} = \underbrace{\mathcal{R}}_{\text{finite remainder}} + \underbrace{I^{(1)} \mathcal{A}_R^{(1)} + I^{(2)} \mathcal{A}_R^{(0)}}_{\text{divergent + convention-dependent finite part}} + \mathcal{O}(\epsilon)$$

- Finite remainder as a weighted sum of *pentagon functions*

Chicherin, Sotnikov ('20);

$$\mathcal{R}(\lambda, \tilde{\lambda}) = \sum_i r_i(\lambda, \tilde{\lambda}) h_i(\lambda \tilde{\lambda})$$

Reconstruct $r_i(\lambda, \tilde{\lambda})$ from \mathbb{F}_p samples

von Manteuffel, Schabinger ('14)
Peraro ('16)



THE NUMERATOR ANSATZ

- The numerator Ansatz takes the form

GDL, Maître ('19)

$$\text{Num. poly}(\lambda, \tilde{\lambda}) = \sum_{\vec{\alpha}, \vec{\beta}} c_{(\vec{\alpha}, \vec{\beta})} \prod_{j=1}^n \prod_{i=1}^{j-1} \langle ij \rangle^{\alpha_{ij}} [ij]^{\beta_{ij}}$$

subject to constraints on $\vec{\alpha}, \vec{\beta}$ due to: 1) mass dimension; 2) little group; 3) linear independence.

- Construct the Ansatz via the algorithm from Section 2.2 of GDL, Page ('22)

Linear independence = irreducibility by the Gröbner basis of a specific ideal.

- Efficient implementation using open-source software only



Gröbner bases → constrain $\vec{\alpha}, \vec{\beta}$

Decker, Greuel, Pfister, Schönemann



Google OR-Tools

Integer programming → enumerate sols. $\vec{\alpha}, \vec{\beta}$

Perron and Furnon (Google optimization team)

- All linear systems solved with CUDA over $\mathbb{F}_{p \leq 2^{31}-1}$ on a laptop ($t_{\text{solving}} \ll t_{\text{sampling}}$)



POLYNOMIAL QUOTIENT RINGS

VS.

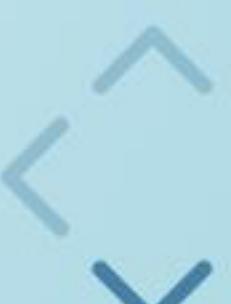
POLYNOMIAL RINGS

- Can we get rid of equivalence relations (redundancies) by changing variables?
In other words, can we solve the redundancies and turn the quotient ring in a ring?

With four-point massless kinematics, you cannot. In R_4 we have

$$\langle 12 \rangle [12] = \langle 34 \rangle [34]$$

this means that R_4 is not a unique factorization domain (UFD).
All polynomial rings are UFDs, so R_4 cannot be isomorphic to one.



THE FIELD OF FRACTIONS OF R_3 DOES NOT EXISTS

- R_3 is not an integral domain, i.e. there exists products of non-zero elements which are zero.

$$\langle 12 \rangle [23] = \langle 1|2|3 \rangle = -\langle 1|1+3|3 \rangle = 0 \text{ but } \langle 12 \rangle \neq 0 \text{ and } [23] \neq 0$$

Hence, one cannot define a field of fractions, as this would not be closed under multiplication.

