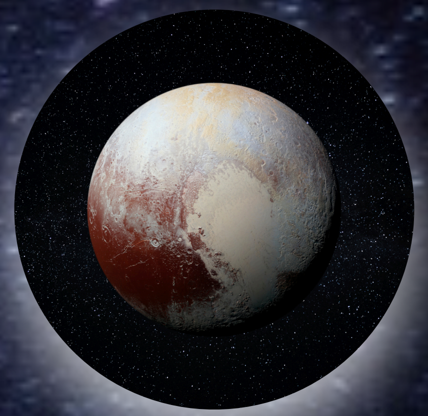
WAVES IN A BOX: RESONANT CAVITIES FOR AXION AND GW DETECTION

Raffaele Tito D’Agnolo - IPhT Saclay

DARK MATTER MASS



Person

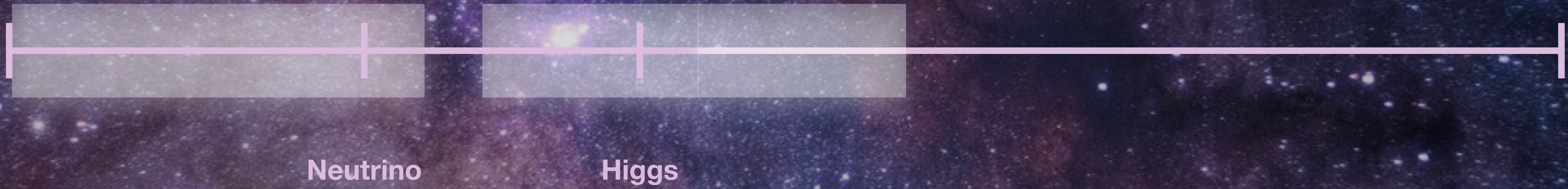
Neutrino

Higgs

Pluto


Self-coupling
50 O.M.
Mass
50 O.M.
80 O.M.
Couplings to ordinary matter

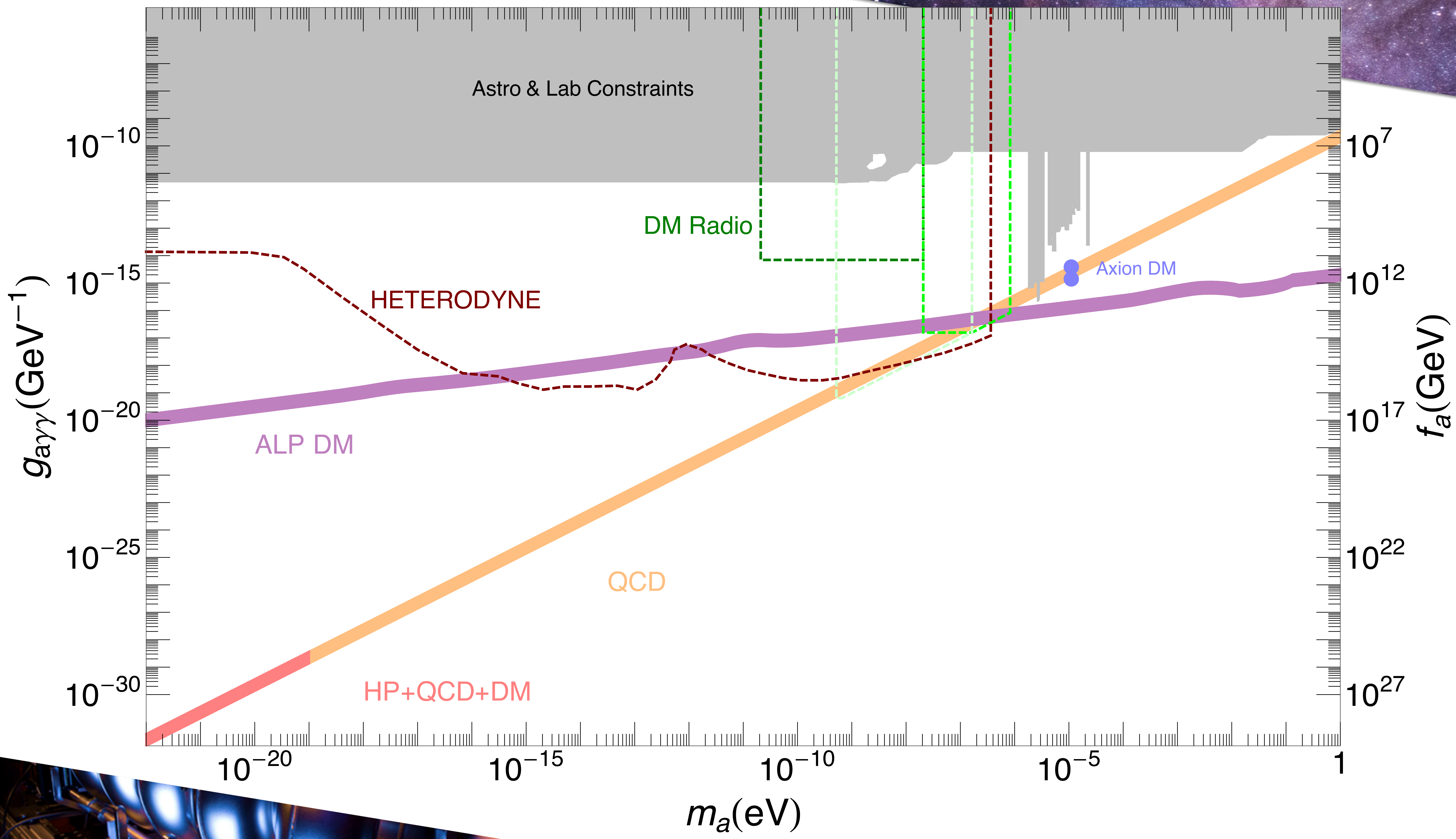
DARK MATTER MASS



Theory Spotlight

DARK MATTER MASS

Theory Spotlight



Astro & Lab Constraints
DM Radio
HETERODYNE
Axion DM
ALP DM
QCD
HP+QCD+DM
$g_{a\gamma\gamma}(\mathrm{GeV}^{-1})$
10^{-10}
10^{-15}
10^{-20}
10^{-25}
10^{-30}
$f_a(\mathrm{GeV})$
10^{7}
10^{12}
10^{17}
10^{22}
10^{27}
$m_a(\mathrm{eV})$
10^{-20}
10^{-15}
10^{-10}
10^{-5}
1

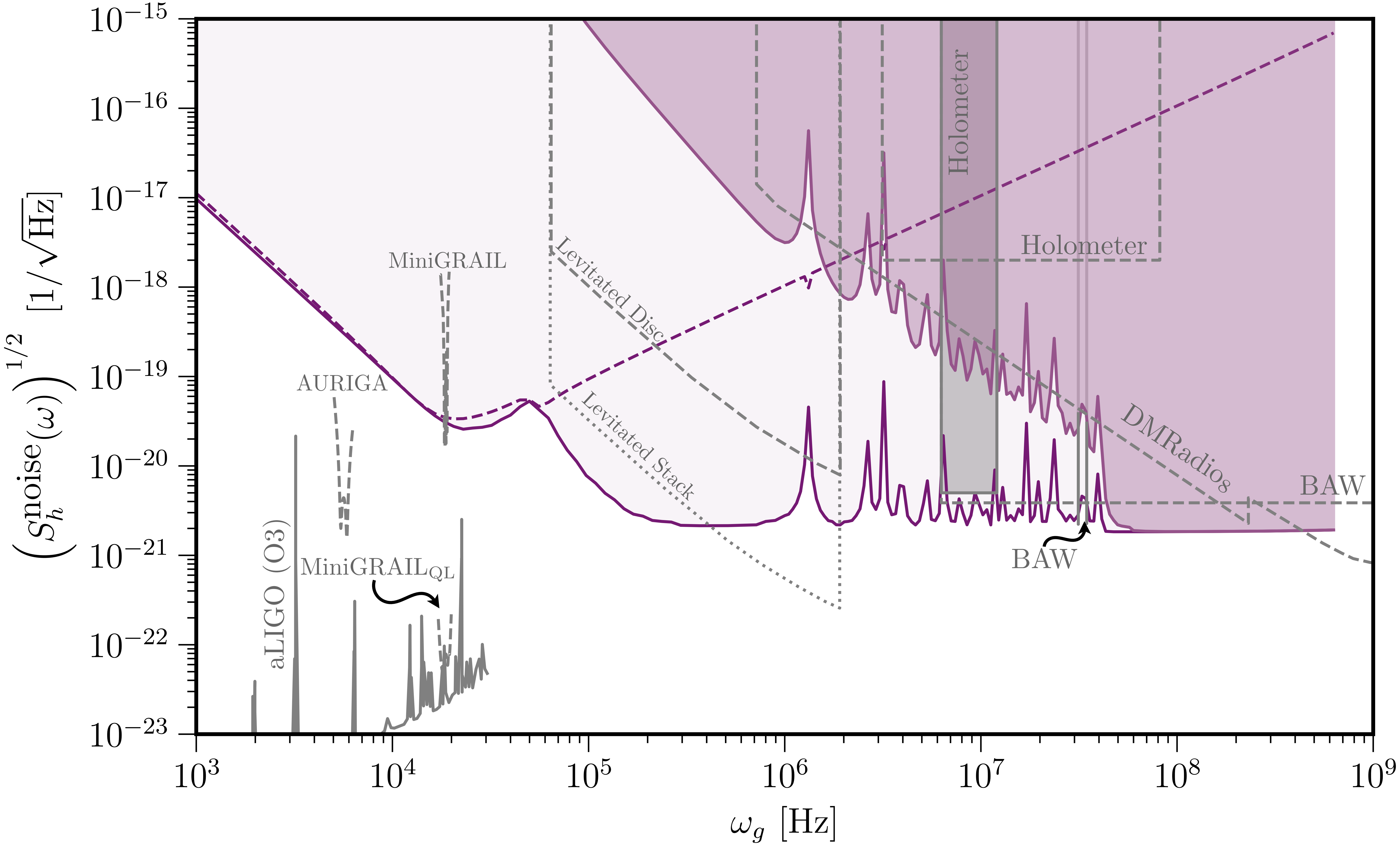
$$\omega_g$$

$$\lambda(T_*) < \frac{1}{H(T_*)}$$ Causality

$$\omega_g$$

$$\lambda(T_*) < \frac{1}{H(T_*)}$$ Causality

$$\omega_0(T_*) = \omega(T_*)\frac{a(T_*)}{a_0} \gtrsim \boxed{100\ \mathrm{MHz}}\left(\frac{T_*}{10^{15}\mathrm{GeV}}\right)\left(\frac{g_*(T_*)}{100}\right)^{1/6}$$



Berlin, Blas, D'Agnolo, Ellis, Harnik, Kahn, Schutte-Engel, Wentzel
arXiv:2303.01518

ALPS DETECTION

Dark Matter Particles in a de Broglie Volume **Today**

Galaxy: $$N_{\mathrm{DM}} \simeq 10^3 \left(\frac{\mathrm{eV}}{m_{\mathrm{DM}}} \right)$$

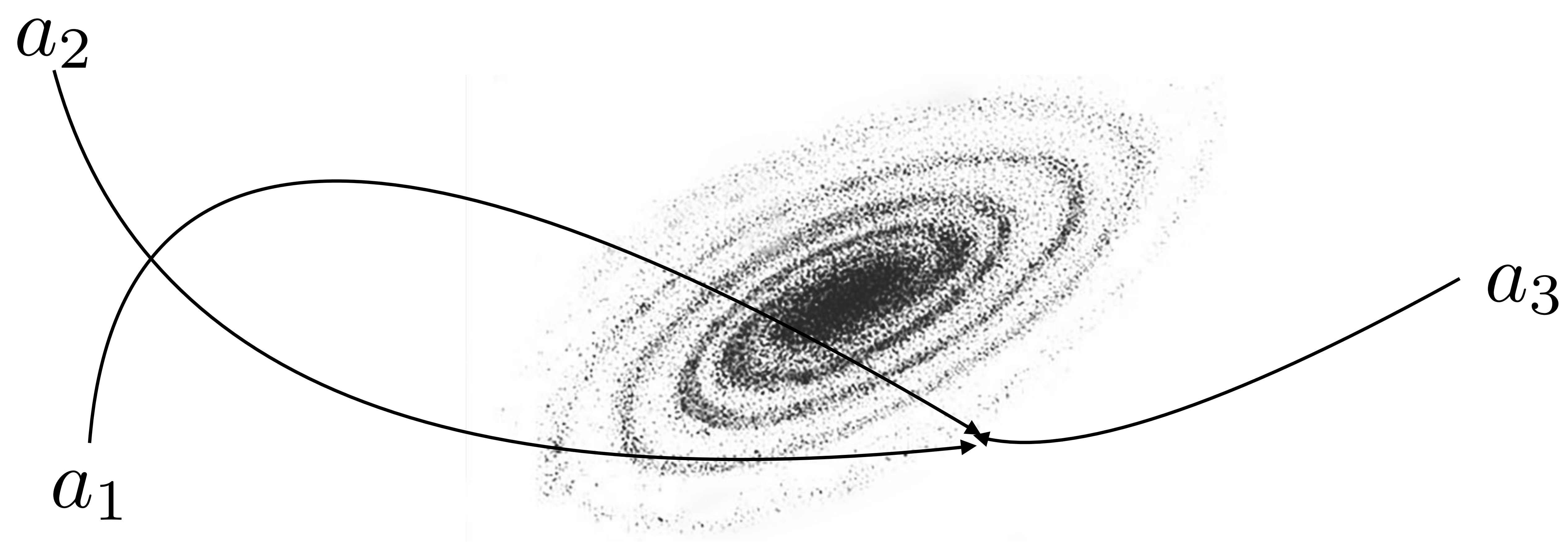
Universe: $$N_{\mathrm{DM}} \simeq 10^{-3} \left(\frac{\mathrm{eV}}{m_{\mathrm{DM}}} \right)$$

ALP DARK MATTER IN THE LAB

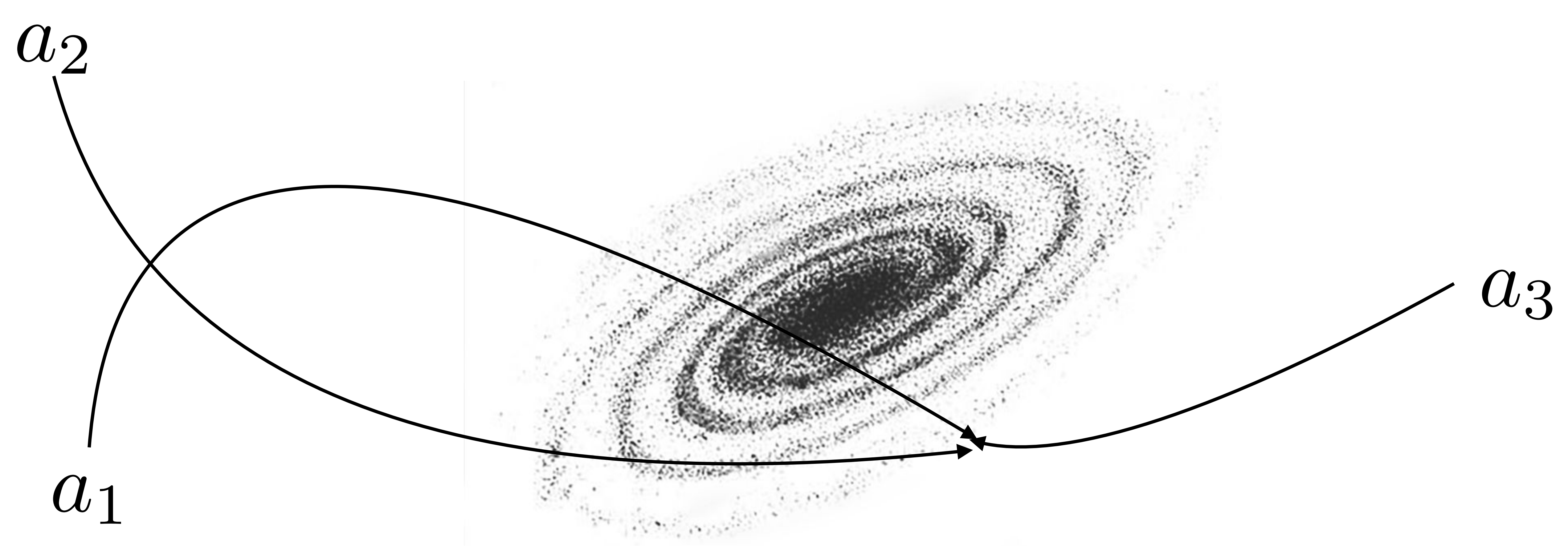
$$a(t) \simeq \frac{\sqrt{2\rho_{\rm DM}}}{m_a}\cos(\omega_a t + \phi)$$

ALP DARK MATTER IN THE LAB

In each experimental bin we are summing over a multitude of plane waves with different phases



In each experimental bin we are summing over a multitude of plane waves with different phases



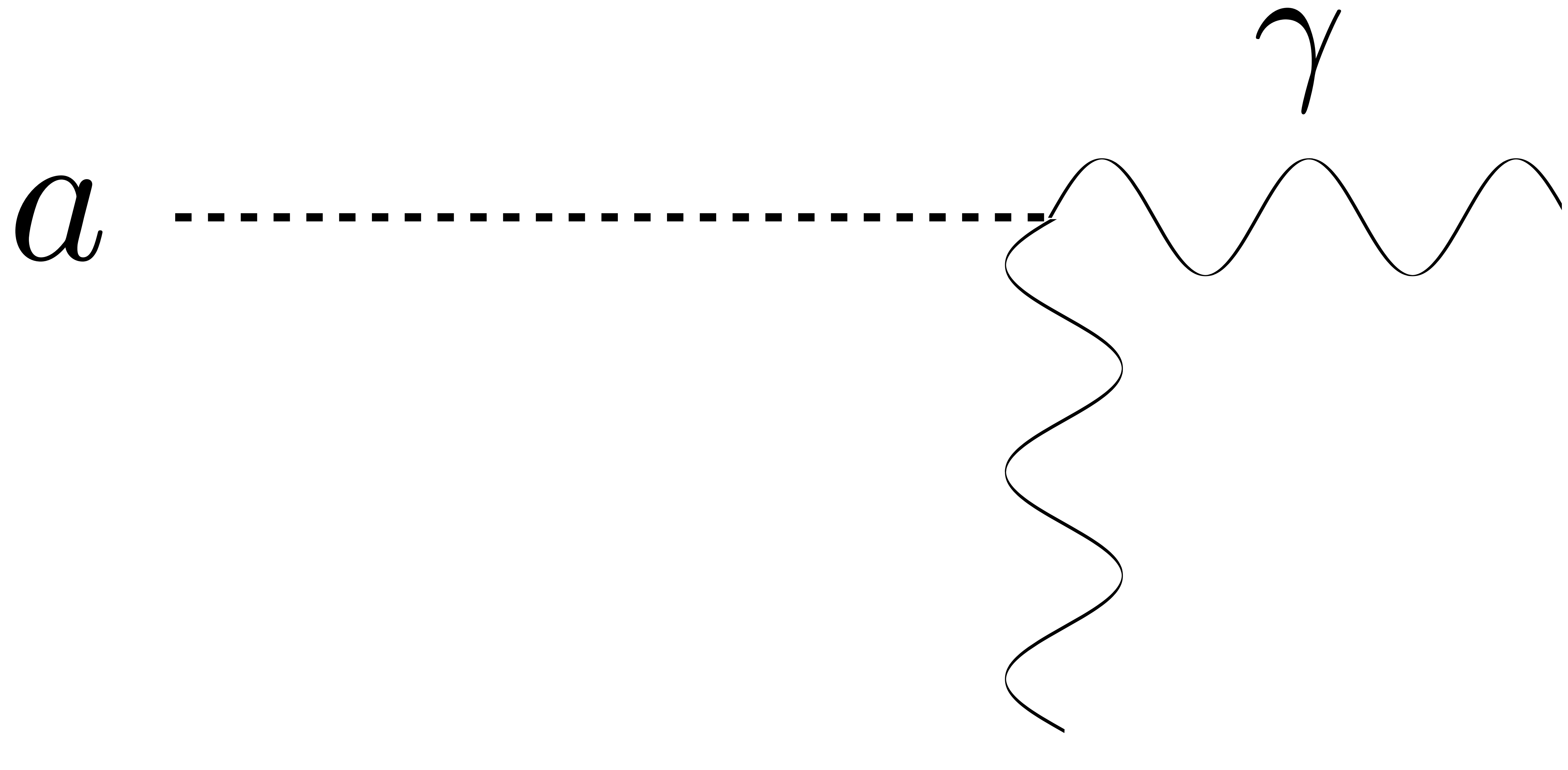
$$a(t) = a_0 \left[\cos\left(m_a \left(1 + \frac{v_1^2}{2}\right) t + \phi_1 \right) + \cos\left(m_a \left(1 + \frac{v_2^2}{2}\right) t + \phi_2 \right) + \ldots \right]$$

$$\simeq a_0 \cos(m_a t + \phi) \left[\cos(\delta\omega_a t + \phi') + \ldots \right]$$

$$\delta\omega_a \simeq \frac{1}{m_a \langle v_{\rm DM}^2 \rangle} \simeq \frac{10^6}{m_a}$$

Effectively: very **slow modulation of** an approximately **monochromatic field**

ALP DARK MATTER DETECTION

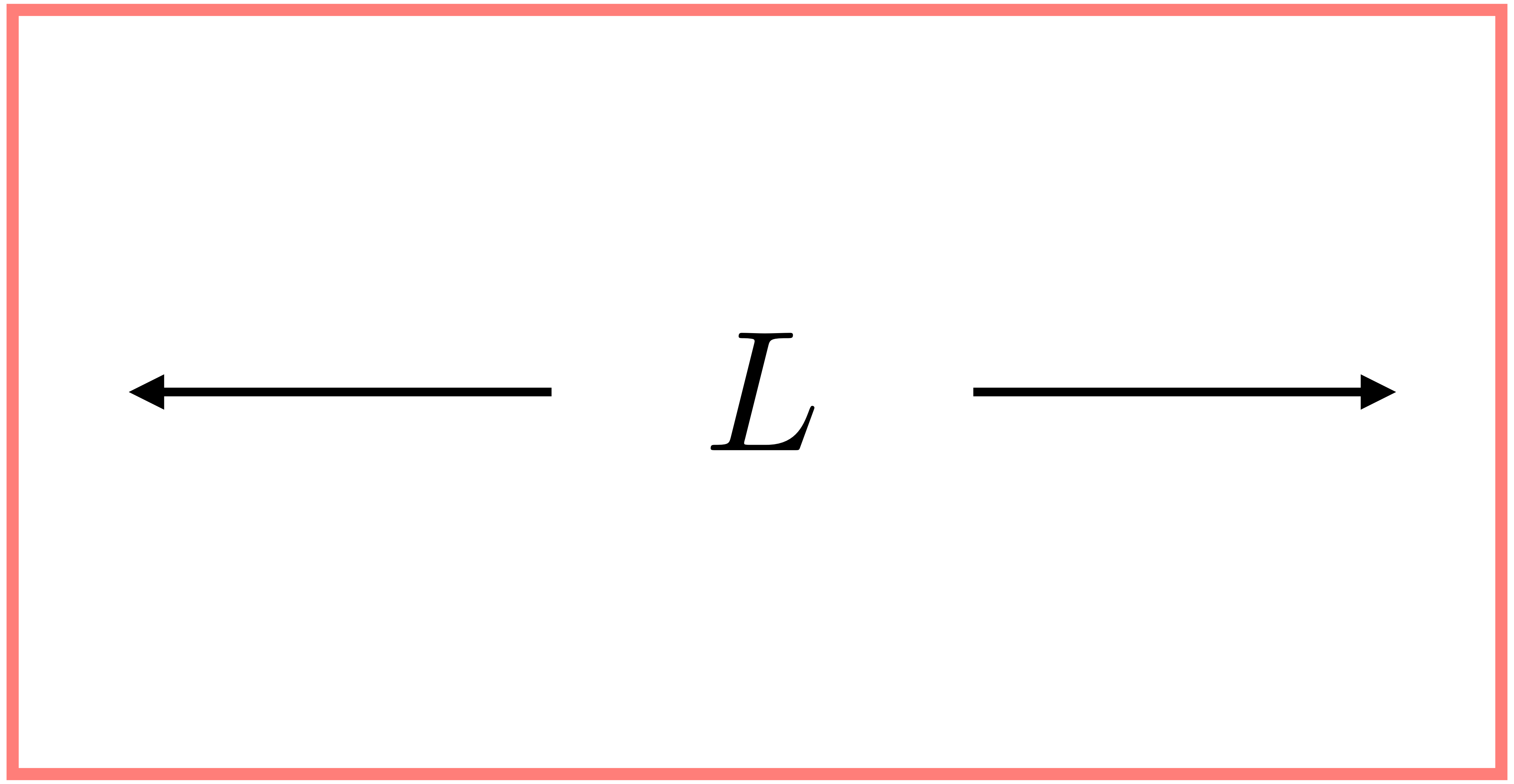


$$\sim \frac{a}{f_a} E_{\rm bkg} \simeq \boxed{10^{-21} E_{\rm bkg}}$$

but you know exactly the waveform
and the signal is always there

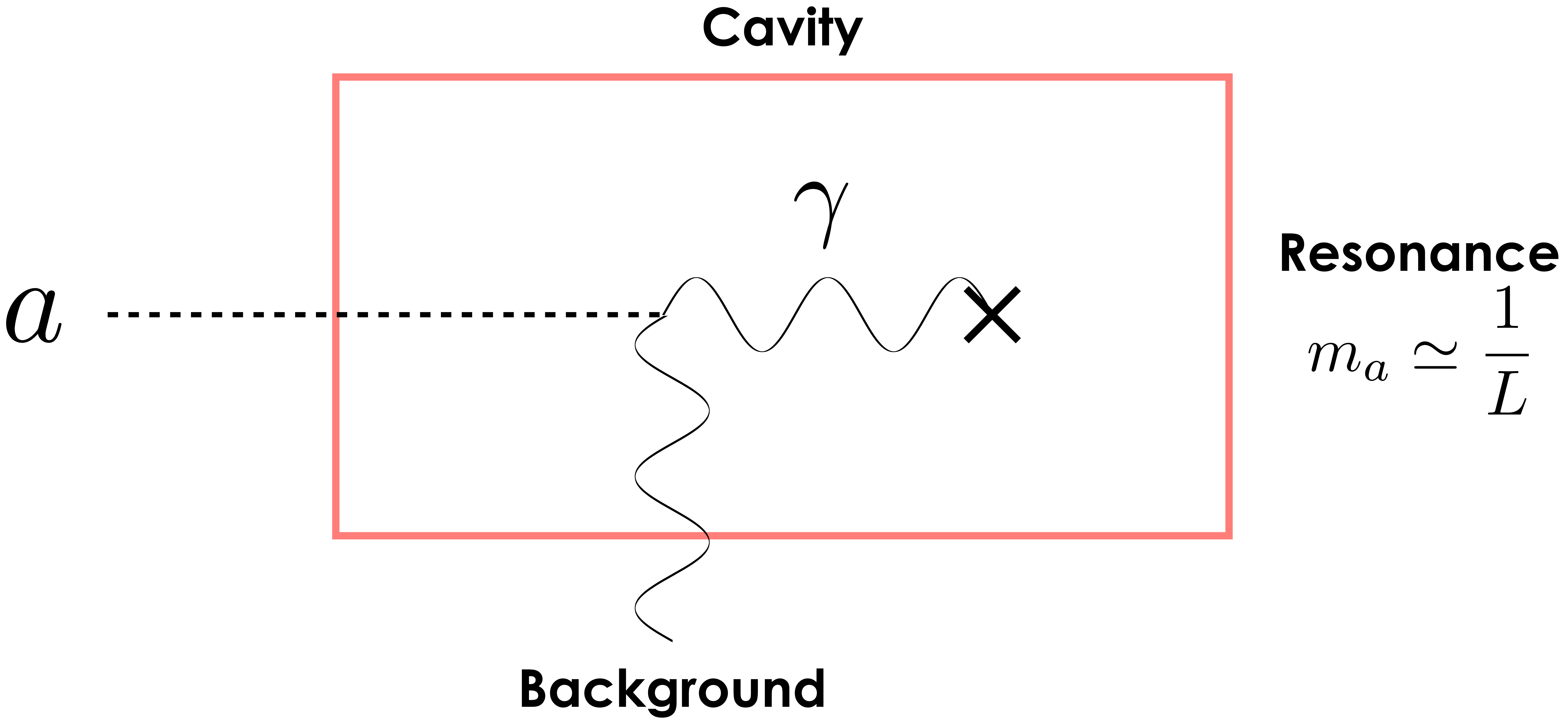
AXION DARK MATTER DETECTION

Cavity



$$m_\gamma \simeq \frac{1}{L}$$

AXION DARK MATTER DETECTION



Astro & Lab Constraints
Impractical
(fine control of
huge cavities)
Axion DM
QCD
HP+QCD+DM
10^{-10}
10^{-15}
10^{-20}
10^{-25}
10^{-30}
$g_{a\gamma\gamma}(\mathrm{GeV}^{-1})$
10^{7}
10^{12}
10^{17}
10^{22}
10^{27}
$f_a(\mathrm{GeV})$
10^{-20}
10^{-15}
10^{-10}
10^{-5}
1
$m_a(\mathrm{eV})$

AXION DETECTION

Cavity:

$$\sum_n \left(\partial_t^2 + \frac{\omega_n}{Q_n} \partial_t + \omega_n^2 \right) \mathbf{E}_n = g_{a\gamma\gamma} \partial_t (\mathbf{B} \partial_t a)$$

AXION DETECTION

Cavity:

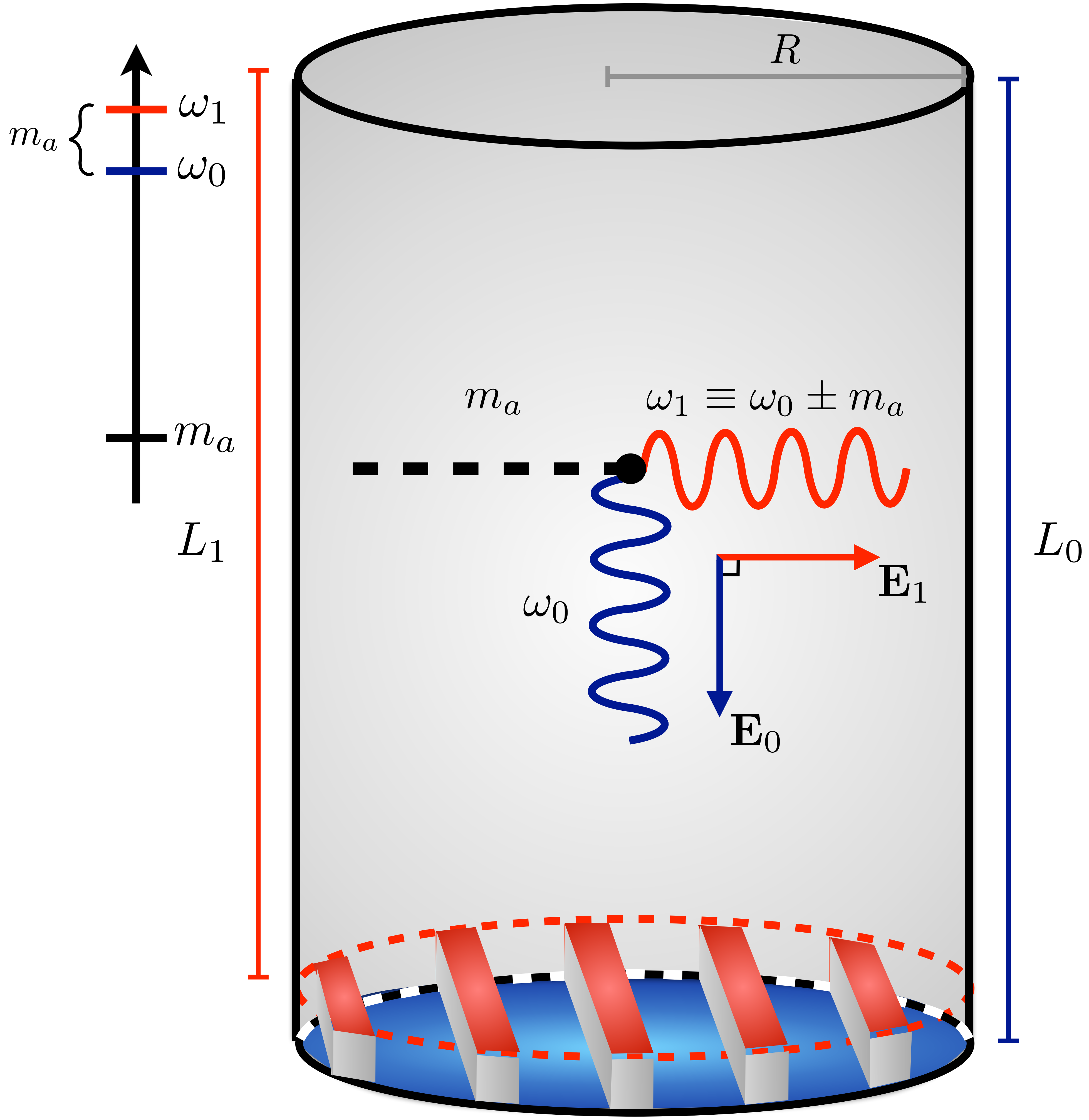
$$\sum_n \left(\partial_t^2 + \frac{\omega_n}{Q_n} \partial_t + \omega_n^2 \right) \mathbf{E}_n = g_{a\gamma\gamma} \partial_t (\mathbf{B} \partial_t a)$$

$$\omega_1 \simeq m_a \qquad \partial_t (\mathbf{B}) \simeq 0$$

$$\left(\partial_t^2 + \frac{m_a}{Q_1} \partial_t + m_a^2 \right) \mathbf{E}_1 = \boxed{g_{a\gamma\gamma} \mathbf{B} \sqrt{\rho_{\mathrm{DM}}} m_a \cos m_a t}$$

HETERODYNE DETECTION

[Berlin, RTD, S. Ellis, C. Nantista, J. Nielson, P. Schuster, S. Tantawi, N. Toro, K. Zhou'19]



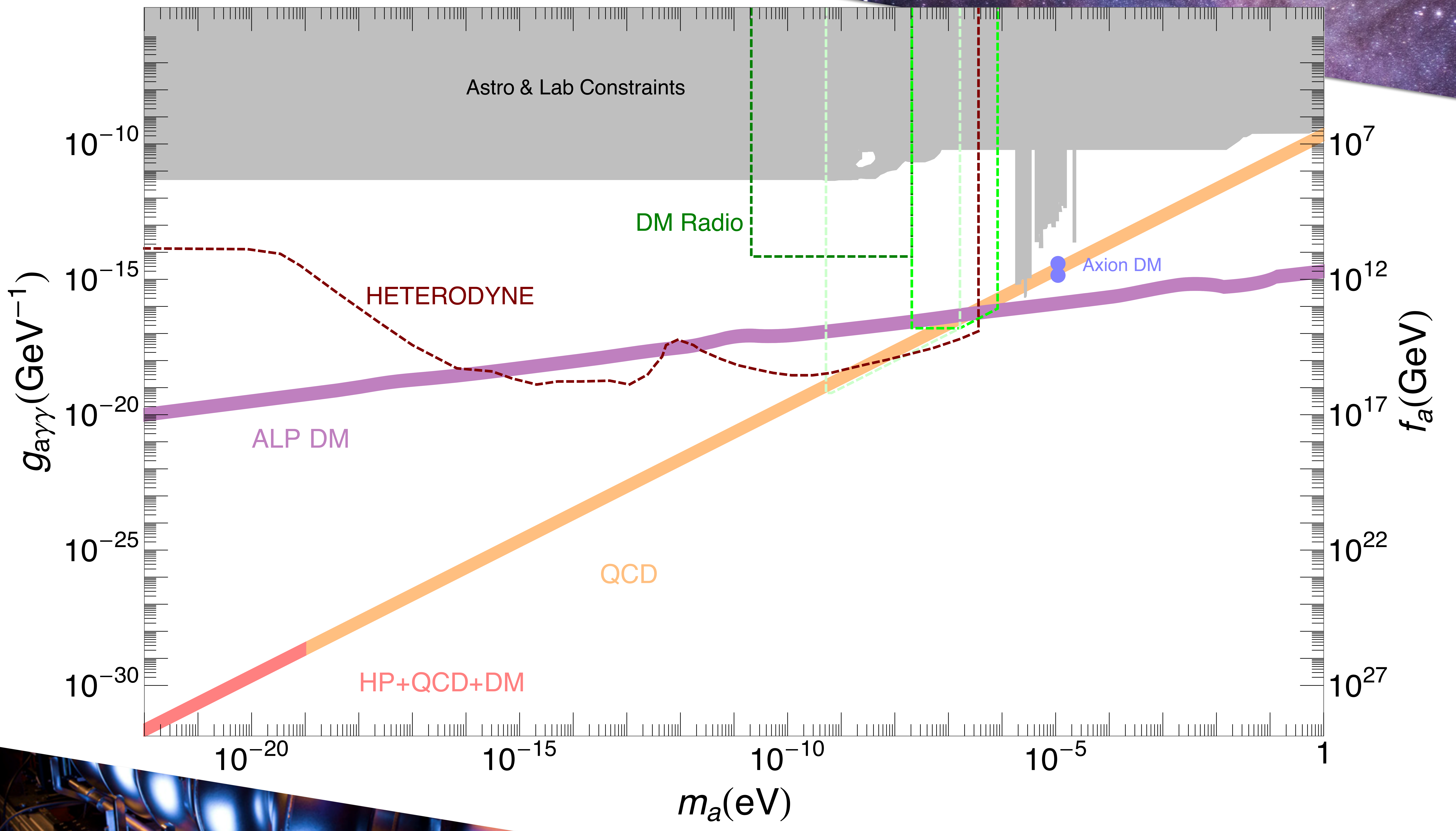
HETERODYNE DETECTION

[Berlin, RTD, S. Ellis, C. Nantista, J. Nielson, P. Schuster, S. Tantawi, N. Toro, K. Zhou'19]

$$\sum_n \left(\partial_t^2 + \frac{\omega_n}{Q_n} \partial_t + \omega_n^2 \right) \mathbf{E}_n = g_{a\gamma\gamma} \partial_t (\mathbf{B} \partial_t a)$$

$$\partial_t(\mathbf{B}) \simeq i\omega_0 \mathbf{B} \qquad \omega_1 \simeq \omega_0 + m_a$$

$$\partial_t J_{\text{eff}} = g_{a\gamma\gamma} \partial_t (\mathbf{B} \partial_t a) \propto \omega_0 m_a \gg m_a^2$$



Astro & Lab Constraints
DM Radio
HETERODYNE
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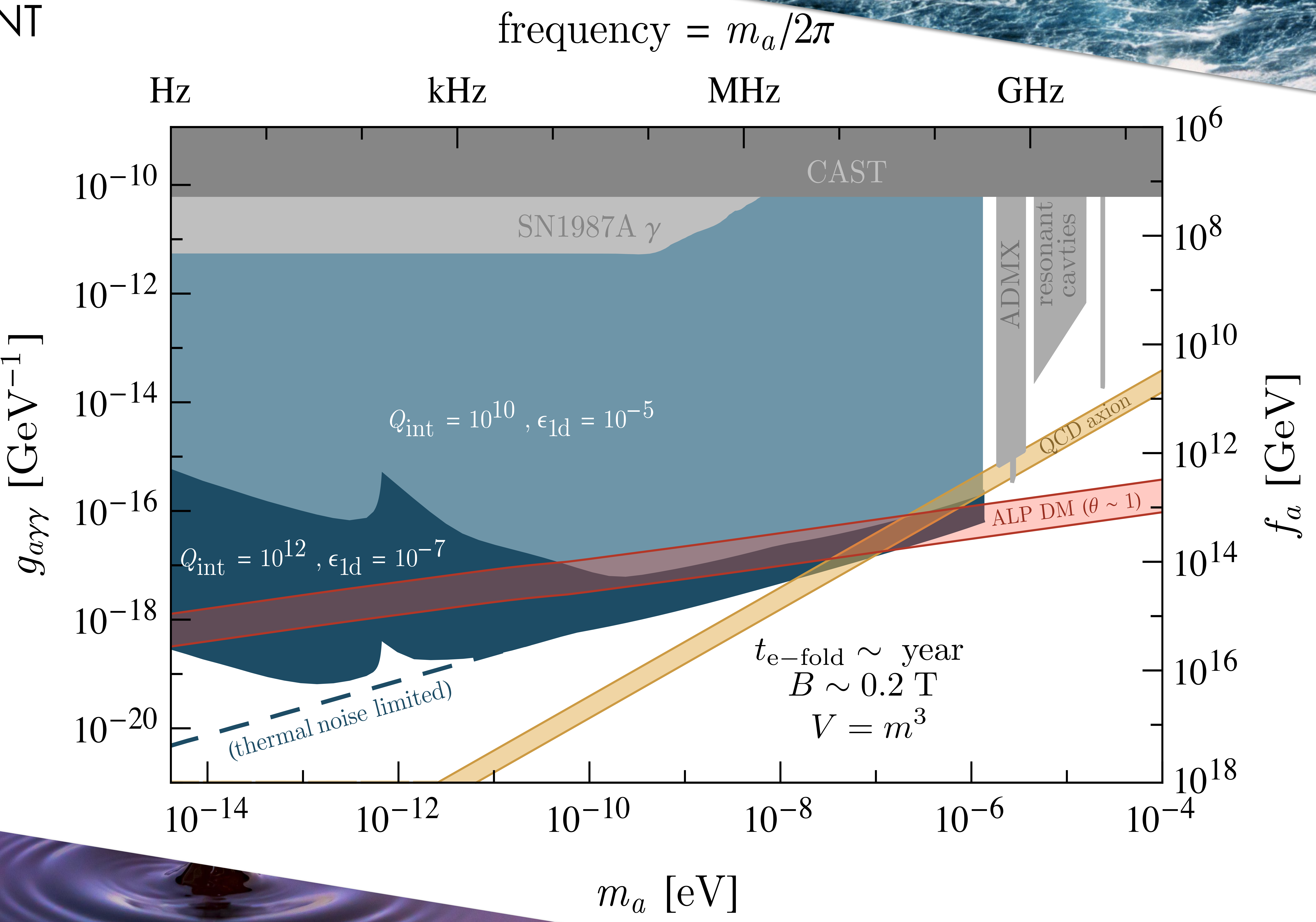
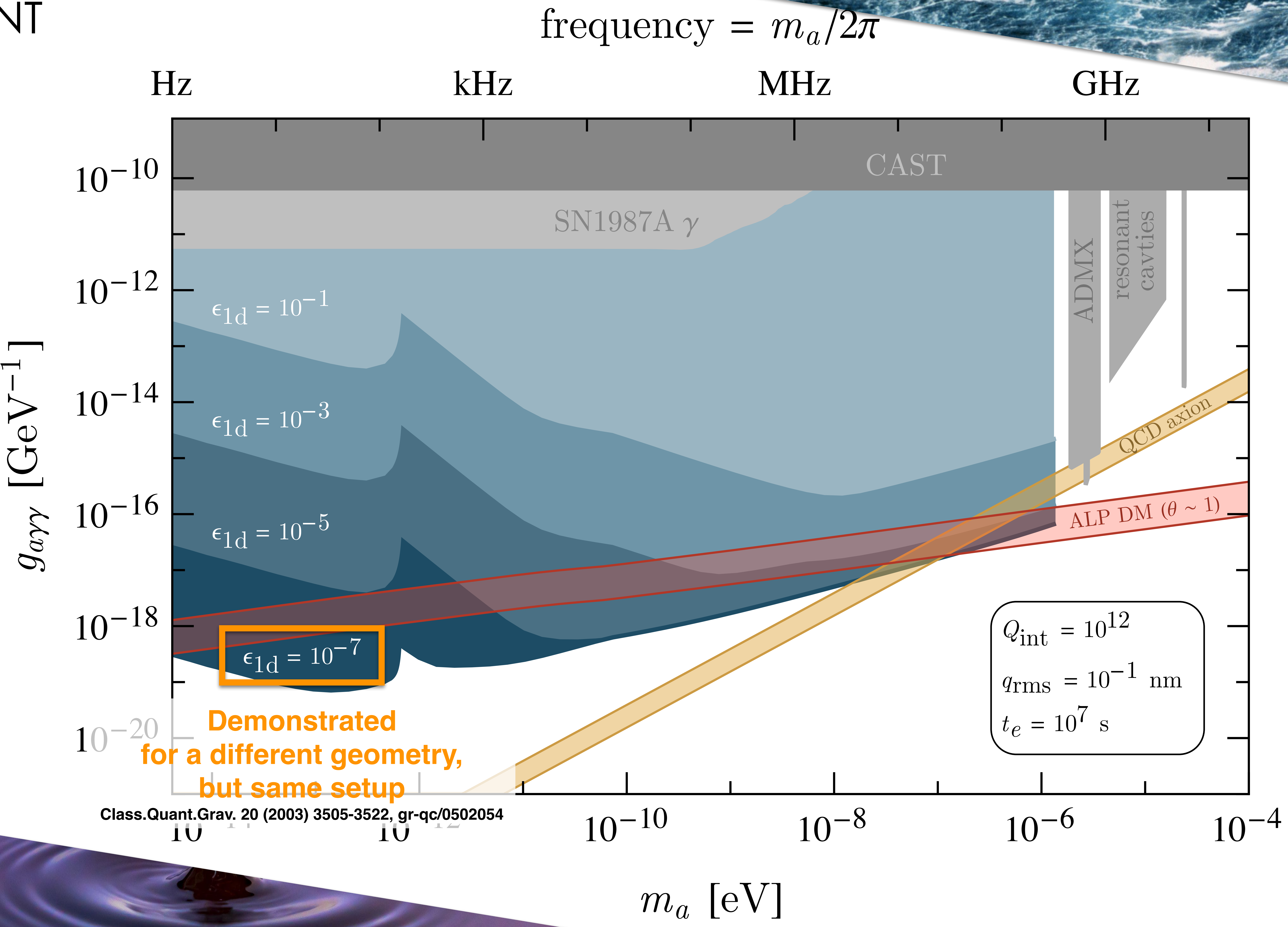
ALP DM
QCD
HP+QCD+DM
10^{-10}
10^{-15}
10^{-20}
10^{-25}
10^{-30}
10^{7}
10^{12}
10^{17}
10^{22}
10^{27}
10^{-20}
10^{-15}
10^{-10}
10^{-5}
1
$g_{a\gamma\gamma}(\mathrm{GeV}^{-1})$
$f_a(\mathrm{GeV})$
$m_a(\mathrm{eV})$

RESONANT

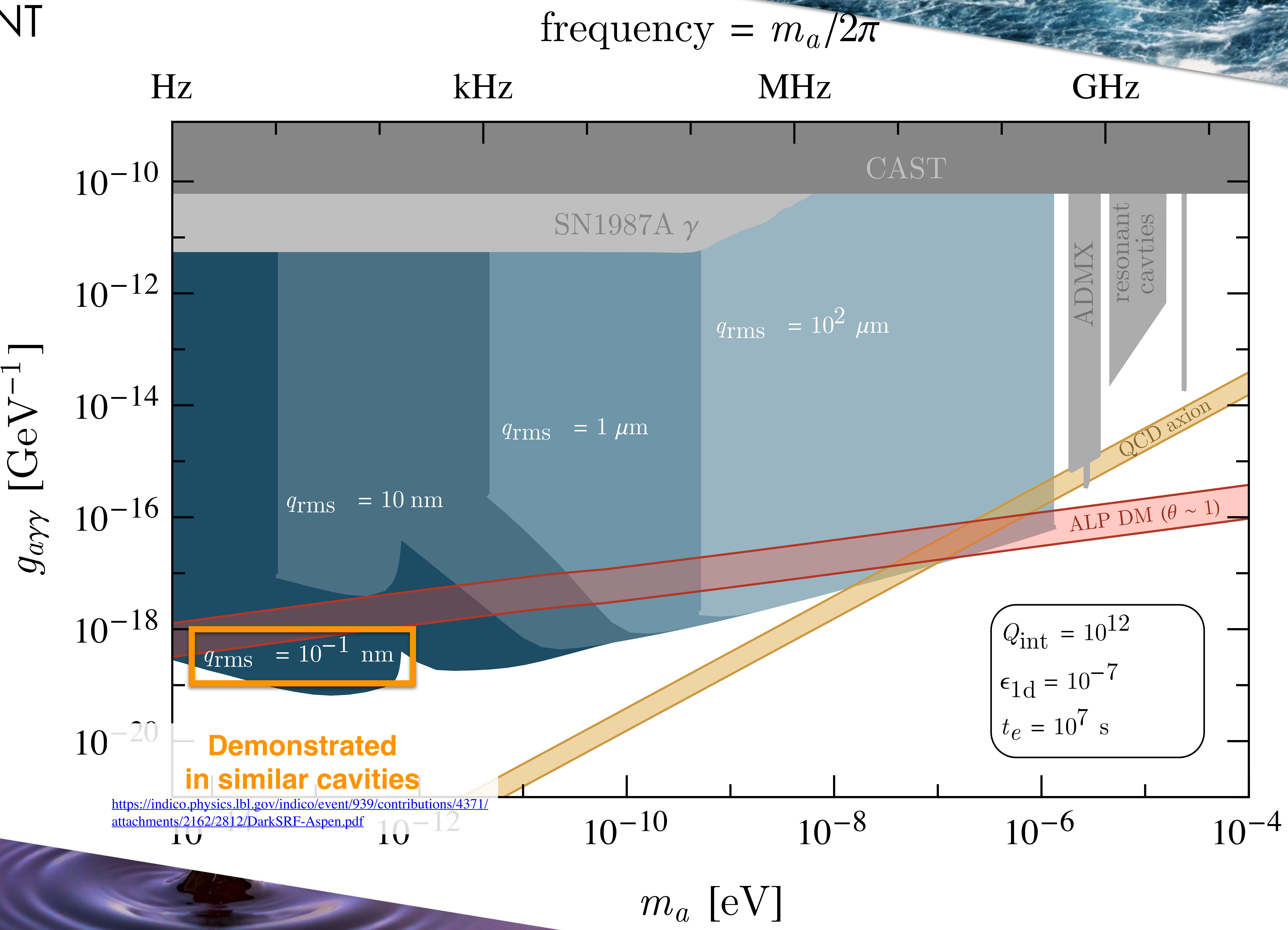


frequency = $m_a/2\pi$
Hz
kHz
MHz
GHz
CAST
SN1987A γ
ADMX
resonant cavties
$Q_{\rm int} = 10^{10}$, $\epsilon_{\rm 1d} = 10^{-5}$
$Q_{\rm int} = 10^{12}$, $\epsilon_{\rm 1d} = 10^{-7}$
QCD axion
ALP DM ($\theta \sim 1$)
(thermal noise limited)
$t_{\rm e-fold} \sim$ year
$B \sim 0.2$ T
$V = m^3$
$g_{a\gamma\gamma}$ [GeV^{-1}]
f_a [GeV]
m_a [eV]
10^{-10}
10^{-12}
10^{-14}
10^{-16}
10^{-18}
10^{-20}
10^{6}
10^{8}
10^{10}
10^{12}
10^{14}
10^{16}
10^{18}
10^{-14}
10^{-12}
10^{-10}
10^{-8}
10^{-6}
10^{-4}

RESONANT

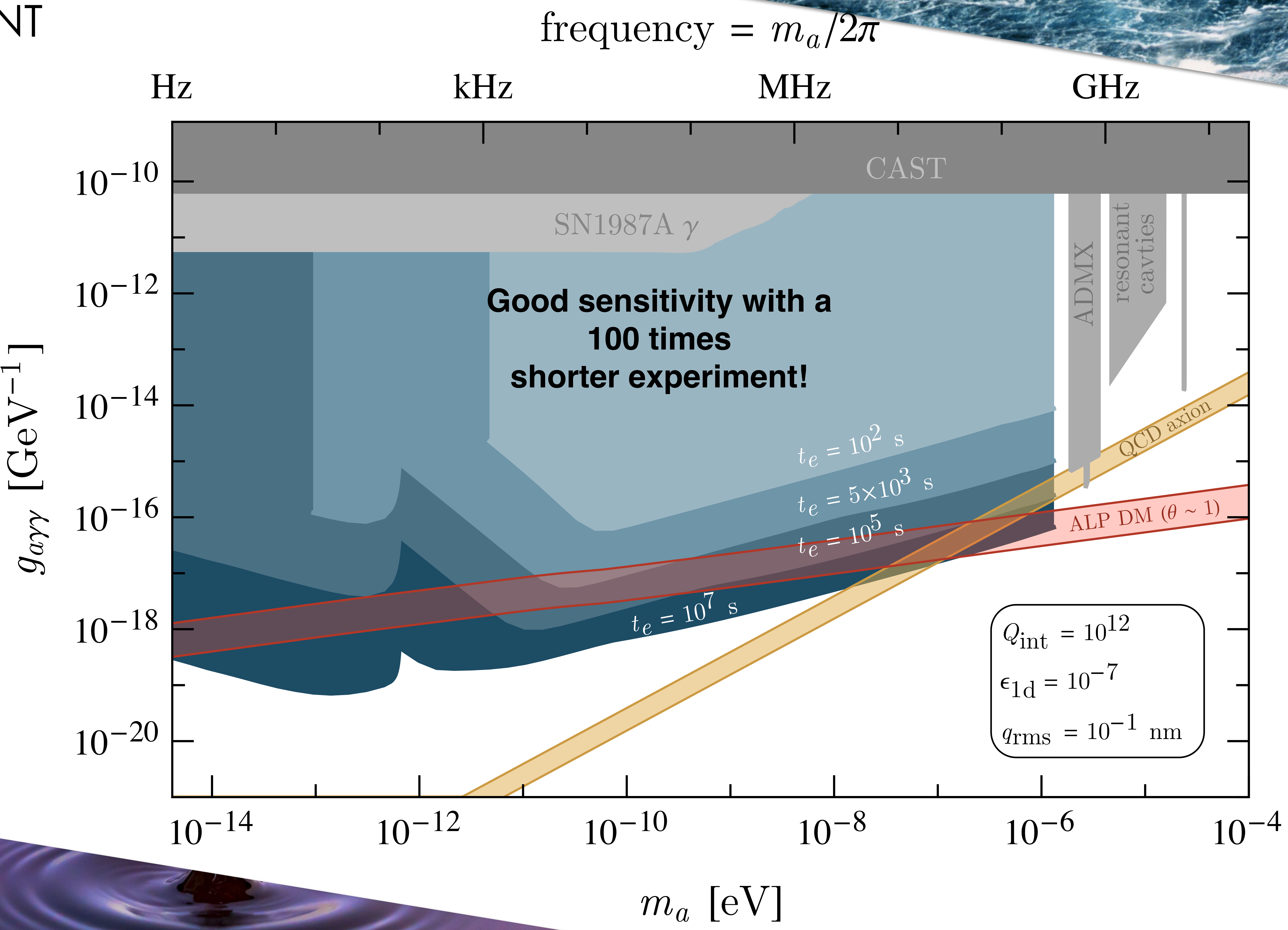

frequency = $m_a/2\pi$
Hz
kHz
MHz
GHz
CAST
SN1987A γ
ADMX
resonant cavties
$\epsilon_{1d} = 10^{-1}$
$\epsilon_{1d} = 10^{-3}$
$\epsilon_{1d} = 10^{-5}$
$\epsilon_{1d} = 10^{-7}$
QCD axion
ALP DM ($\theta \sim 1$)
$Q_{int} = 10^{12}$
$q_{rms} = 10^{-1}$ nm
$t_e = 10^{7}$ s
$g_{a\gamma\gamma}$ [GeV^{-1}]
10^{-10}
10^{-12}
10^{-14}
10^{-16}
10^{-18}
10^{-20}
10^{-10}
10^{-8}
10^{-6}
10^{-4}
m_a [eV]
Demonstrated
for a different geometry,
but same setup
Class.Quant.Grav. 20 (2003) 3505-3522, gr-qc/0502054

RESONANT



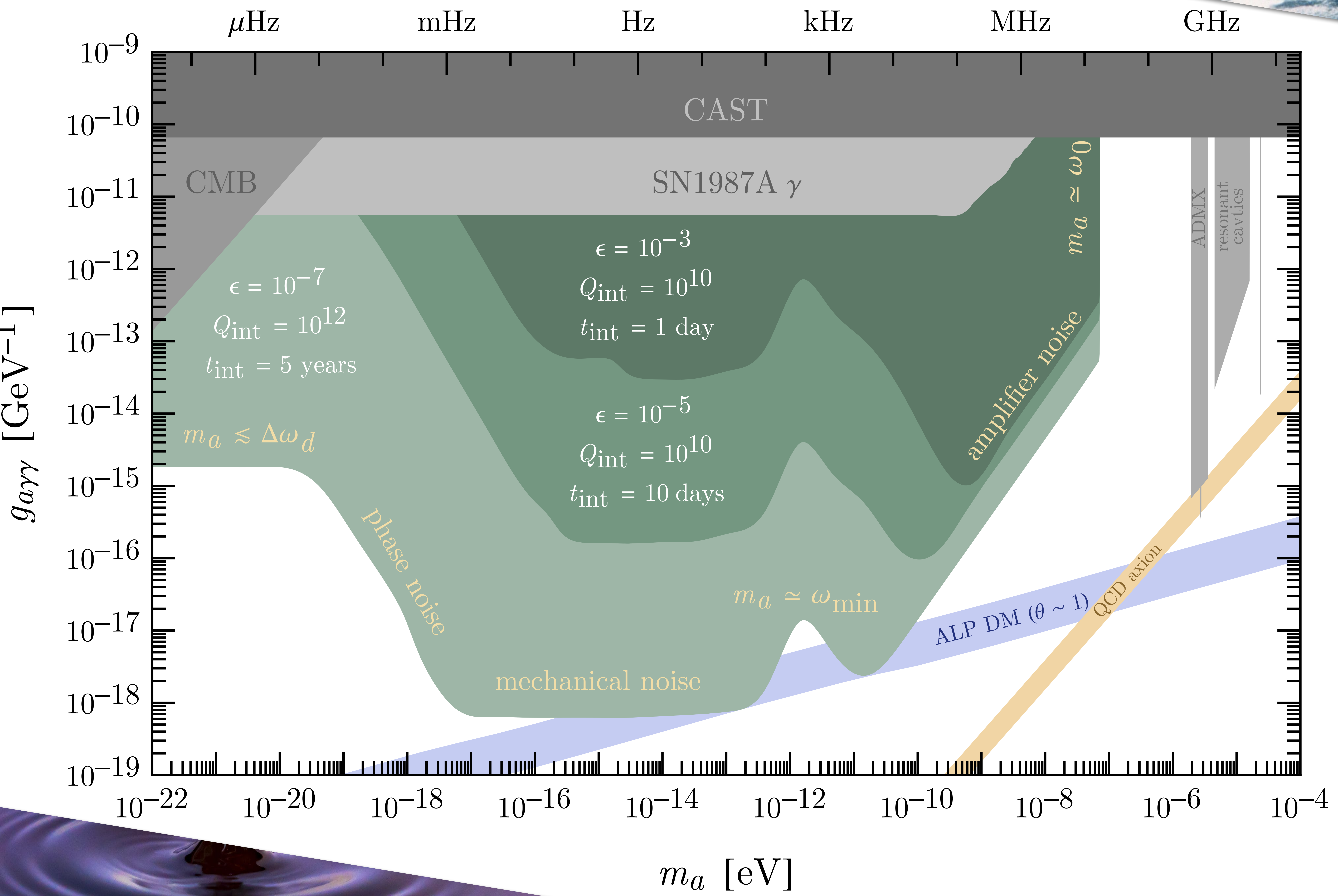
https://indico.physics.lbl.gov/indico/event/939/contributions/4371/attachments/2162/2812/DarkSRF-Aspen.pdf

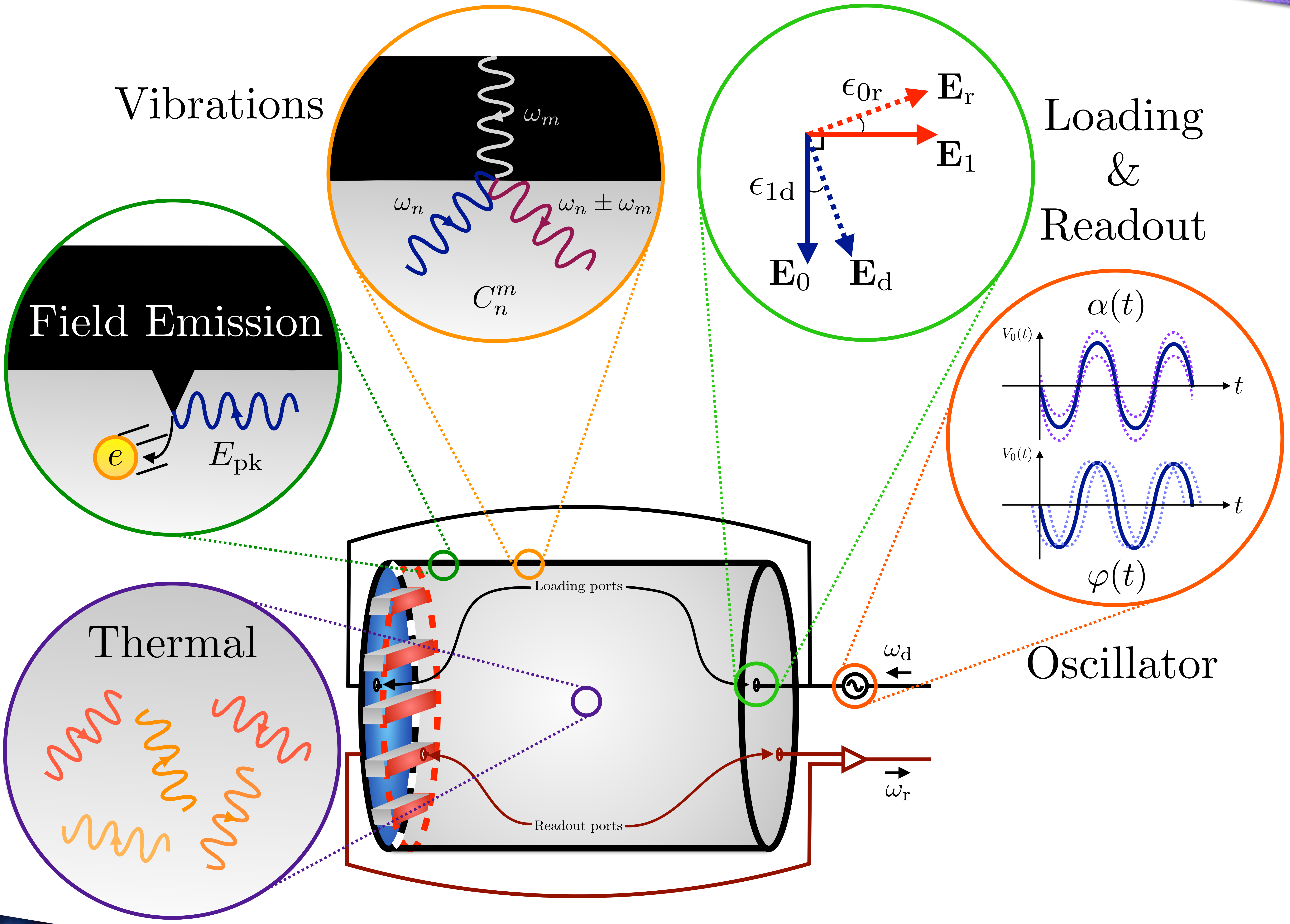
RESONANT



frequency = $m_a/2\pi$
Hz
kHz
MHz
GHz
CAST
SN1987A γ
ADMX
resonant cavties
Good sensitivity with a 100 times shorter experiment!
$t_e = 10^2$ s
$t_e = 5\times10^3$ s
$t_e = 10^5$ s
$t_e = 10^7$ s
QCD axion
ALP DM ($\theta \sim 1$)
$Q_{\rm int} = 10^{12}$
$\epsilon_{\rm 1d} = 10^{-7}$
$q_{\rm rms} = 10^{-1}$ nm
$g_{a\gamma\gamma}$ [GeV^{-1}]
10^{-10}
10^{-12}
10^{-14}
10^{-16}
10^{-18}
10^{-20}
10^{-14}
10^{-12}
10^{-10}
10^{-8}
10^{-6}
10^{-4}
m_a [eV]

A.Berlin, RTD, S. Ellis, K. Zhou 2007.15656





Vibrations
ω_m
ω_n
$\omega_n \pm \omega_m$
C_n^m
Field Emission
e
E_{pk}
Thermal
ϵ_{0r}
$\mathbf{E}_r$
$\mathbf{E}_1$
ϵ_{1d}
$\mathbf{E}_0$
$\mathbf{E}_d$
Loading
&
Readout
$\alpha(t)$
$V_0(t)$
t
$V_0(t)$
t
$\varphi(t)$
Oscillator
ω_d
ω_r
Loading ports
Readout ports

LEAKAGE NOISE

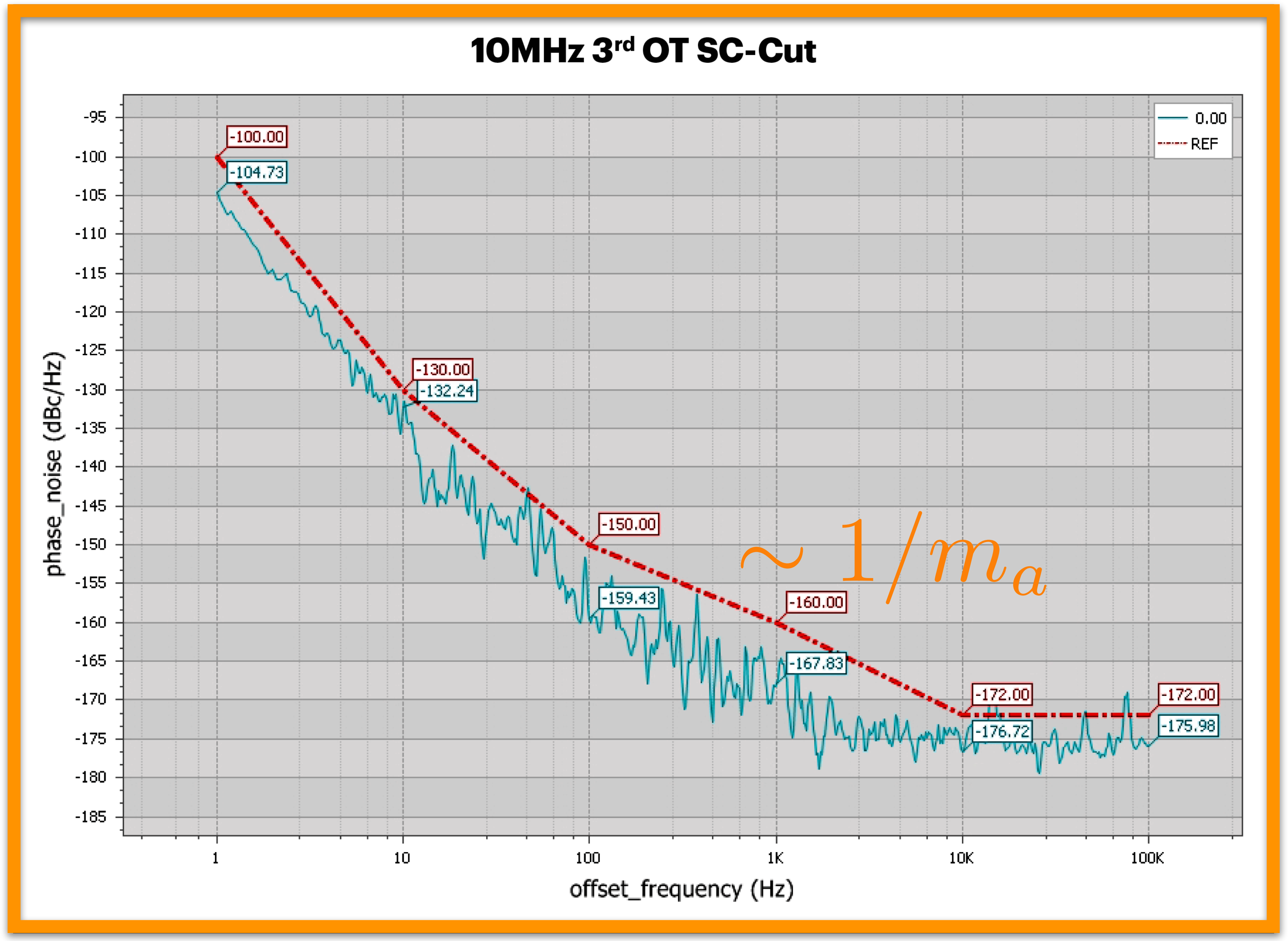
$$S_{\text{phase}}(\omega) \simeq \frac{1}{2}\,\epsilon_{1\text{d}}^2\, S_\phi(\omega-\omega_0)\,\boxed{\frac{(\omega\,\omega_1/Q_1)^2}{(\omega^2-\omega_1^2)^2+(\omega\,\omega_1/Q_1)^2}}\,\frac{\omega_0 Q_1}{\omega_0 Q_0}\,P_{\text{in}}$$

Cavity Response

LEAKAGE NOISE

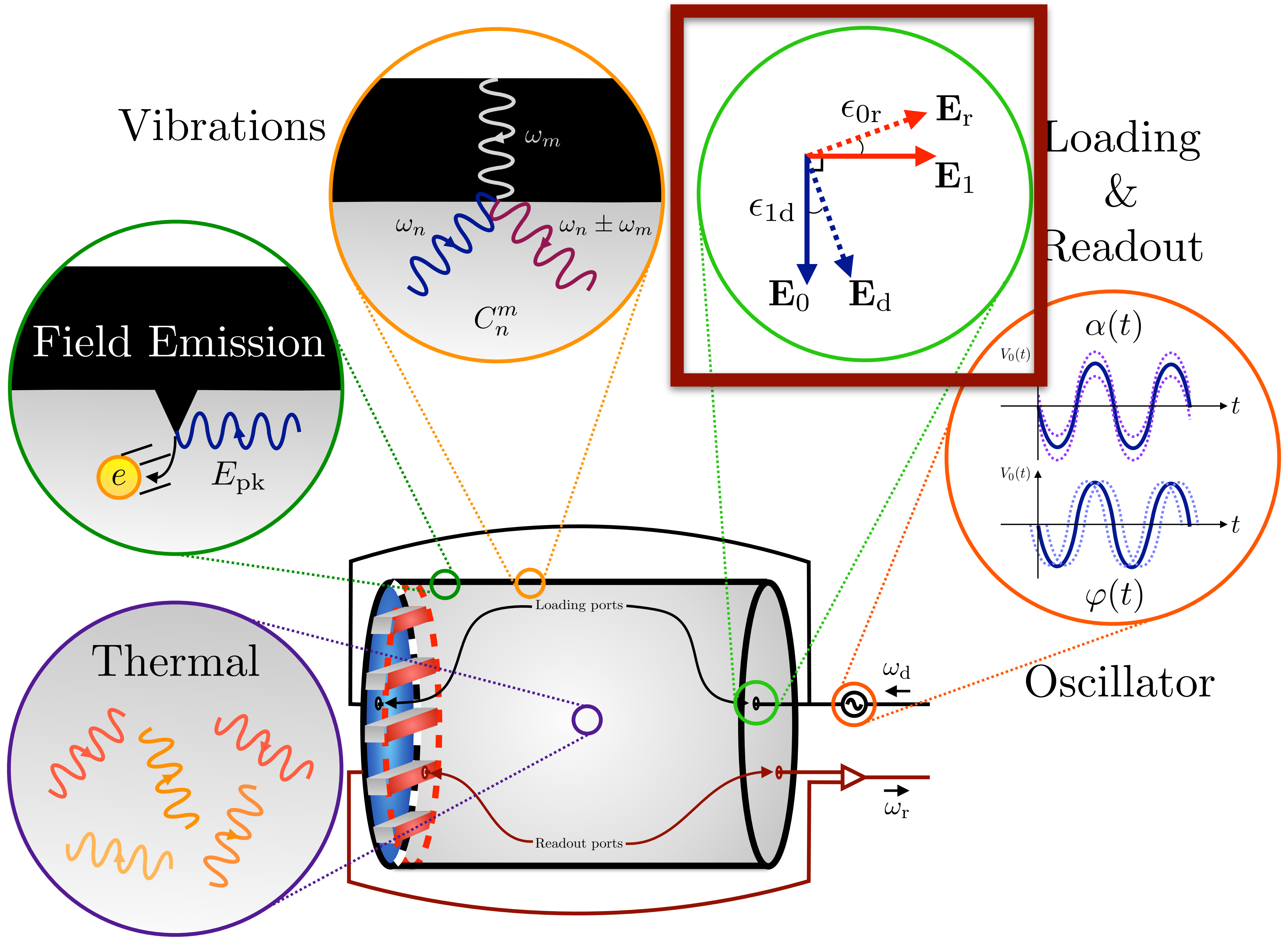
$$S_{\text{phase}}(\omega) \simeq \frac{1}{2}\,\epsilon_{1\text{d}}^2\, S_\phi(\omega-\omega_0)\, \frac{(\omega\,\omega_1/Q_1)^2}{(\omega^2-\omega_1^2)^2+(\omega\,\omega_1/Q_1)^2}\,\frac{\omega_0 Q_1}{\omega_0 Q_0}\,P_{\text{in}}$$

$\sim 1/m_a$



LEAKAGE NOISE

$$S_{\text{phase}}(\omega) \simeq \frac{1}{2}\epsilon_{1\text{d}}^2 S_\phi(\omega-\omega_0)\,\frac{(\omega\,\omega_1/Q_1)^2}{(\omega^2-\omega_1^2)^2+(\omega\,\omega_1/Q_1)^2}\,\frac{\omega_0 Q_1}{\omega_0 Q_0}\,P_{\text{in}}$$

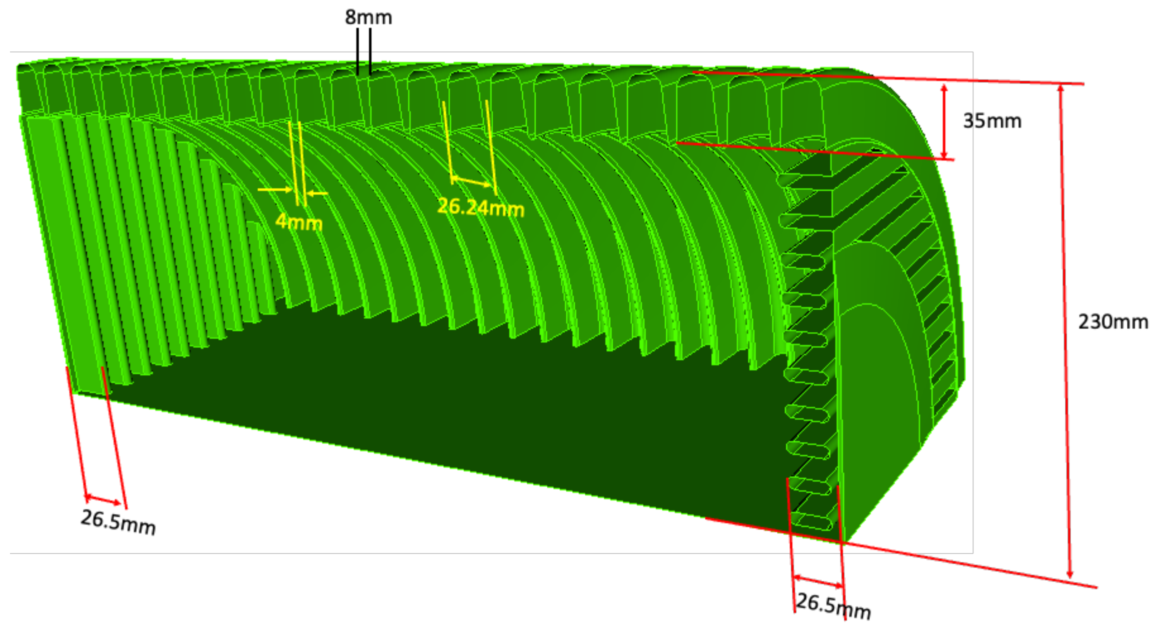


From MAGO and other similar cavities

TWO PROTOTYPES [~ 1 YEAR]

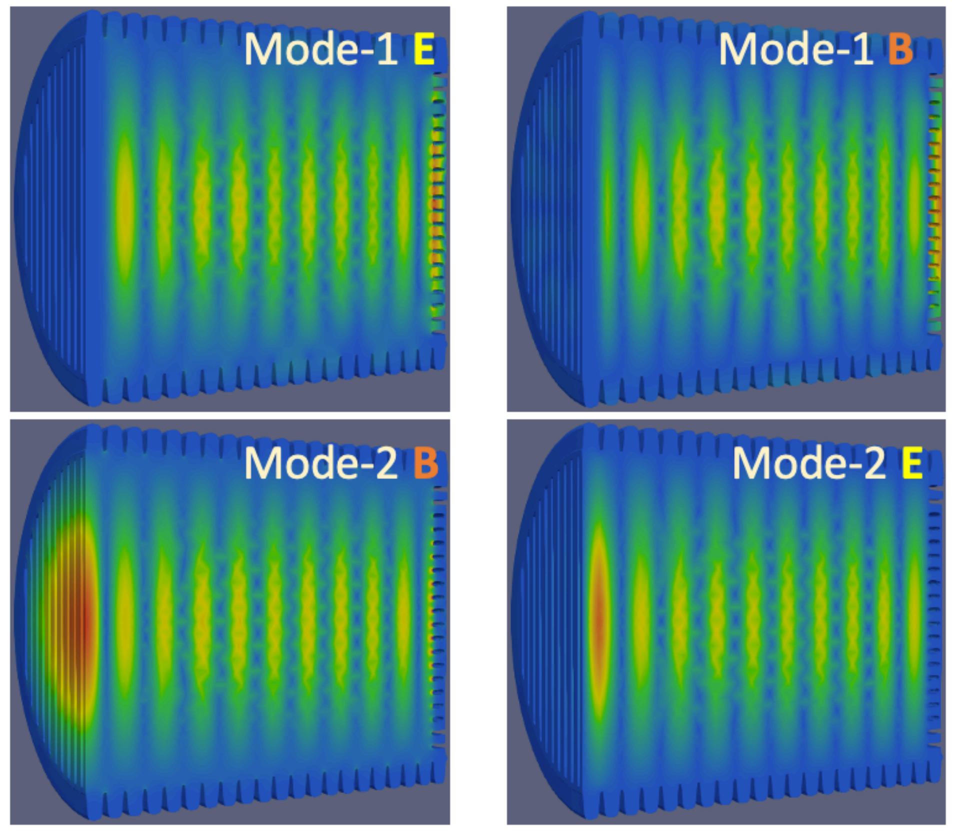
SLAC NATIONAL ACCELERATOR LABORATORY

LDRD [only internal documents]



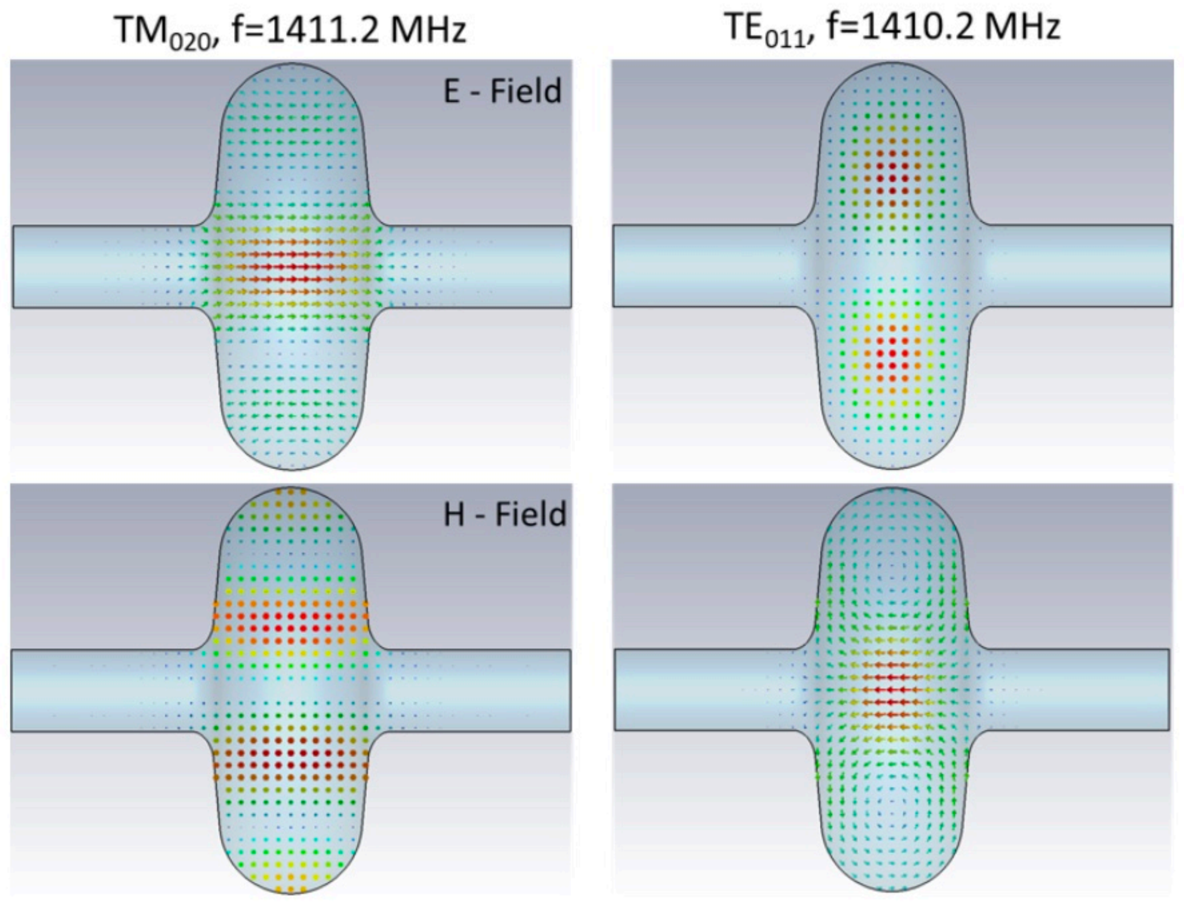
HE11
polatization-1 (E,B)

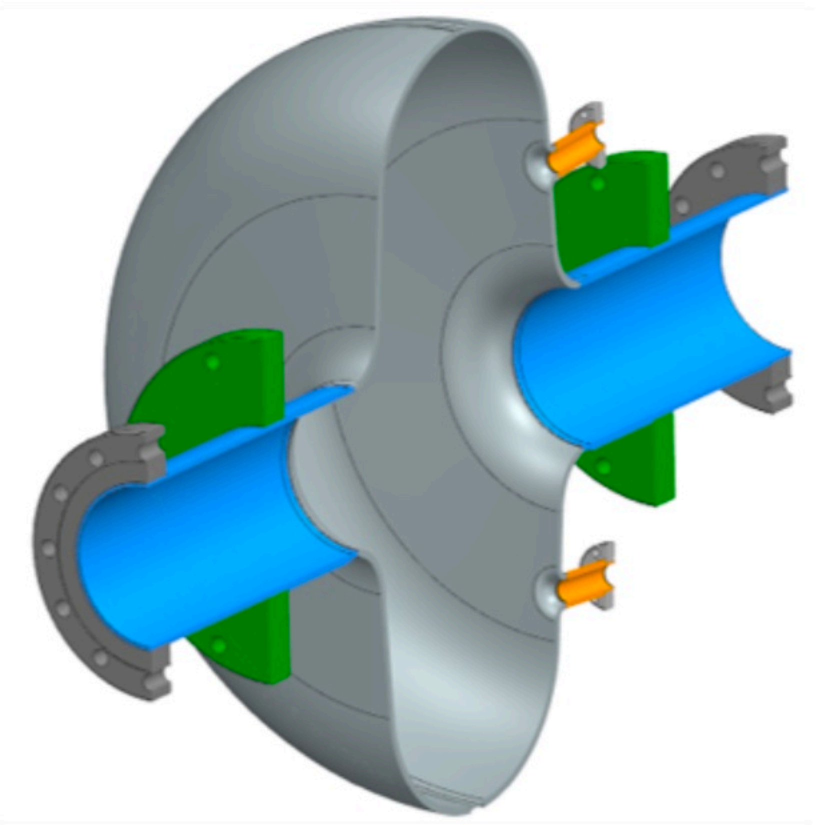
HE11
polatization-2 (B,E)



Fermilab

arXiv:2207.11346






GRAVITATIONAL WAVES

$$S \supset -\frac{1}{2} \int d^4x \; j^\mu_{\text{eff}} \, A_\mu$$

GW

$$j^\mu_{\text{eff}} = \partial_\nu \left(\frac{1}{2} \, h \, \underline{F^{\mu\nu}} + h^\nu{}_\alpha \, \underline{F^{\alpha\mu}} - h^\mu{}_\alpha \, \underline{F^{\alpha\nu}} \right)$$

Axion

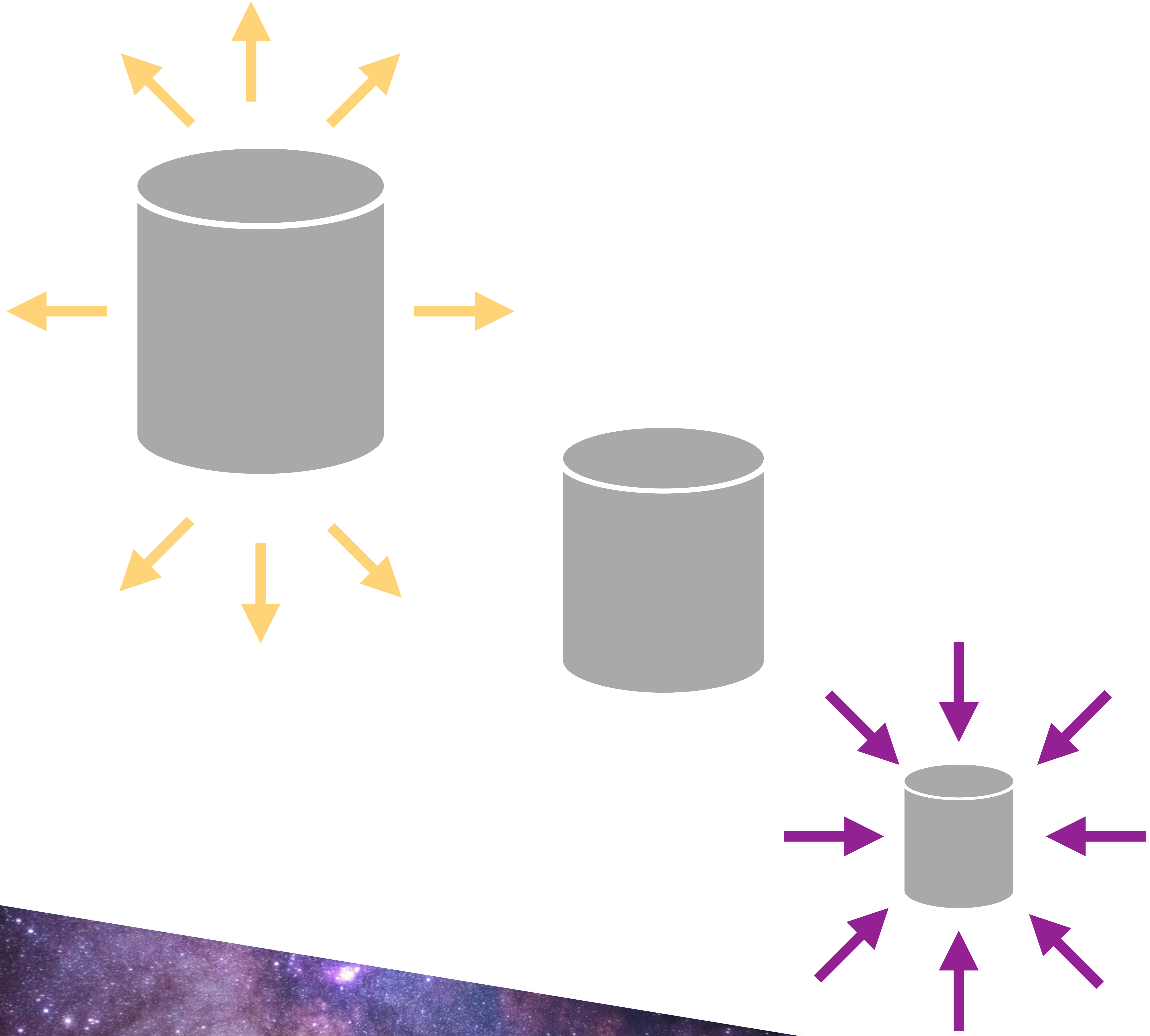
$$j^\mu_{\text{eff}} = \epsilon^{\mu\nu\rho\sigma} \partial_\nu (a \underline{F_{\rho\sigma}})$$

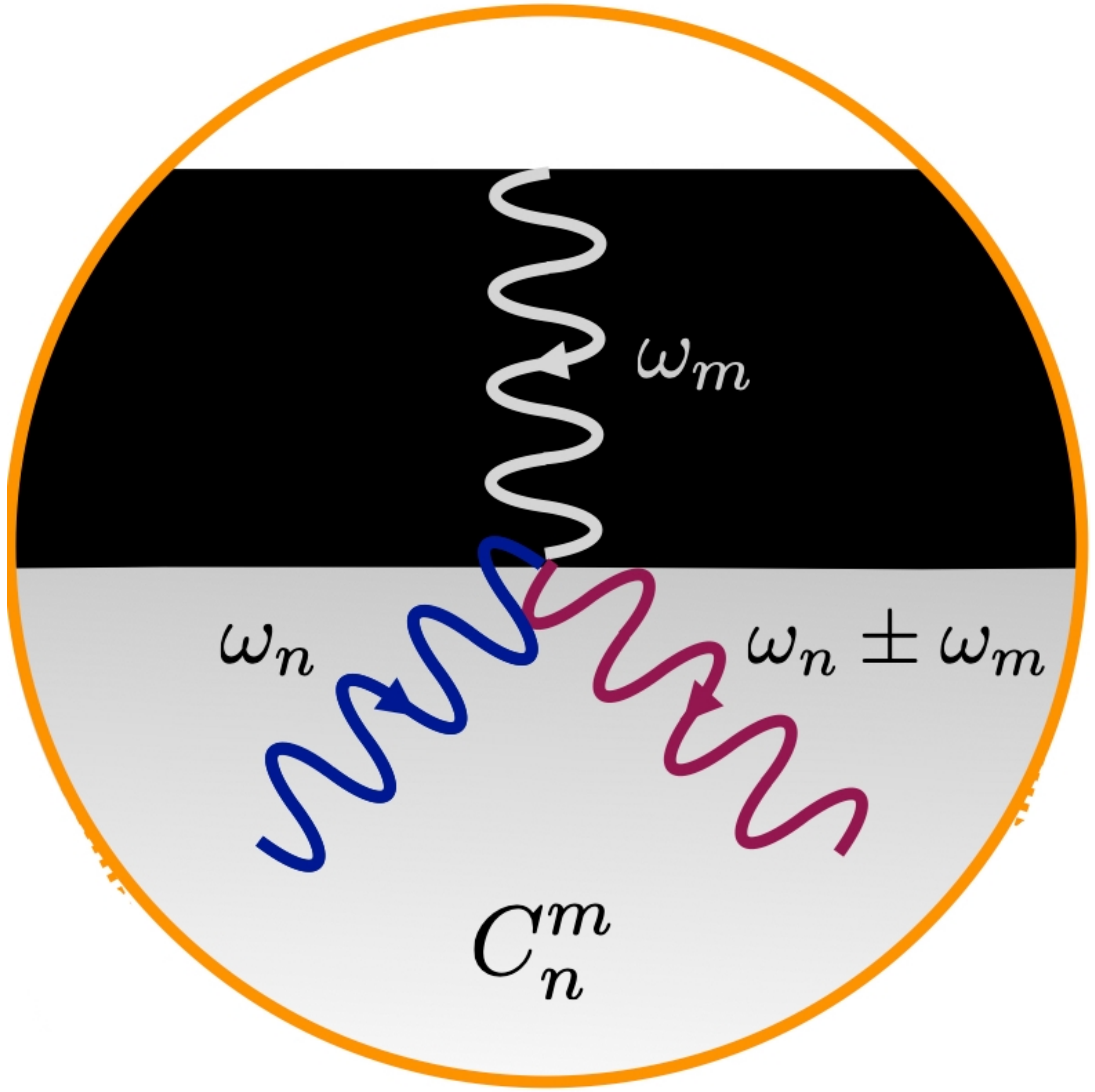
depends on the background field in the laboratory

LOWER FREQUENCIES

$$E_{\text{sig}} \sim \frac{\partial_t j_{\text{eff}}}{\omega_1^2}$$

$$\partial_t j_{\text{eff}} \sim \partial_t^2 (R B_0) \sim \omega_g^2 (h B_0)(\omega_g L_{\text{cav}})^2$$

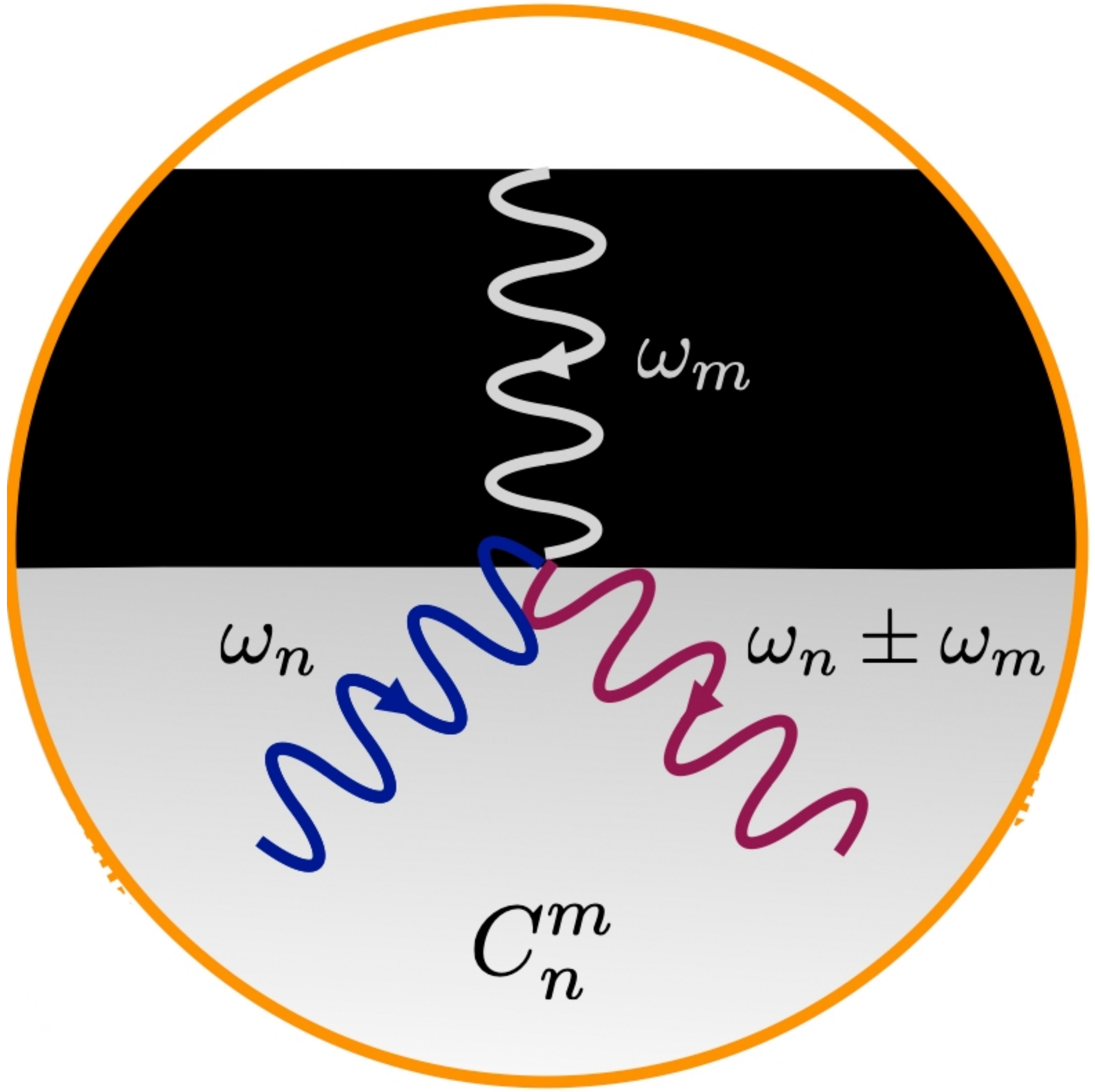



ω_m
ω_n
$\omega_n \pm \omega_m$
C_n^m

Mechanical Mode

$$\left(\partial_t^2 + \frac{\omega_m}{Q}\partial_t + \omega_m^2\right) \underline{u_m} = \frac{F_m}{M_{\rm cav}}$$

MECHANICAL



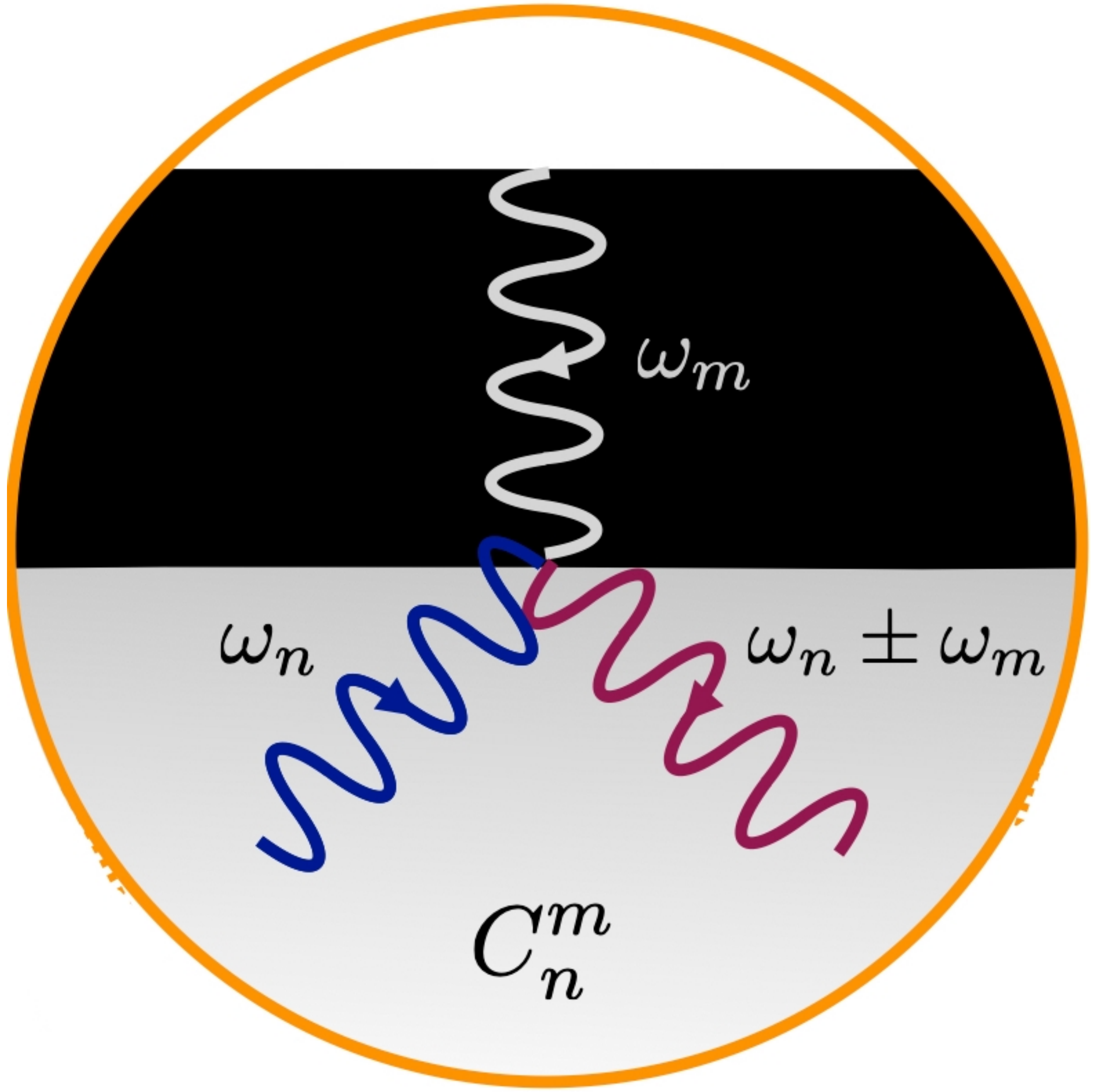
[Caves '78], [Radicati, Pegoraro, Picasso '78]

Mechanical Mode

$$\left(\partial_t^2 + \frac{\omega_m}{Q}\partial_t + \omega_m^2\right) u_m = \frac{F_m}{M_{\rm cav}}$$

$$u_m \sim Q \frac{\omega_g^2}{\omega_m^2} (h L_{\rm cav})$$

MECHANICAL

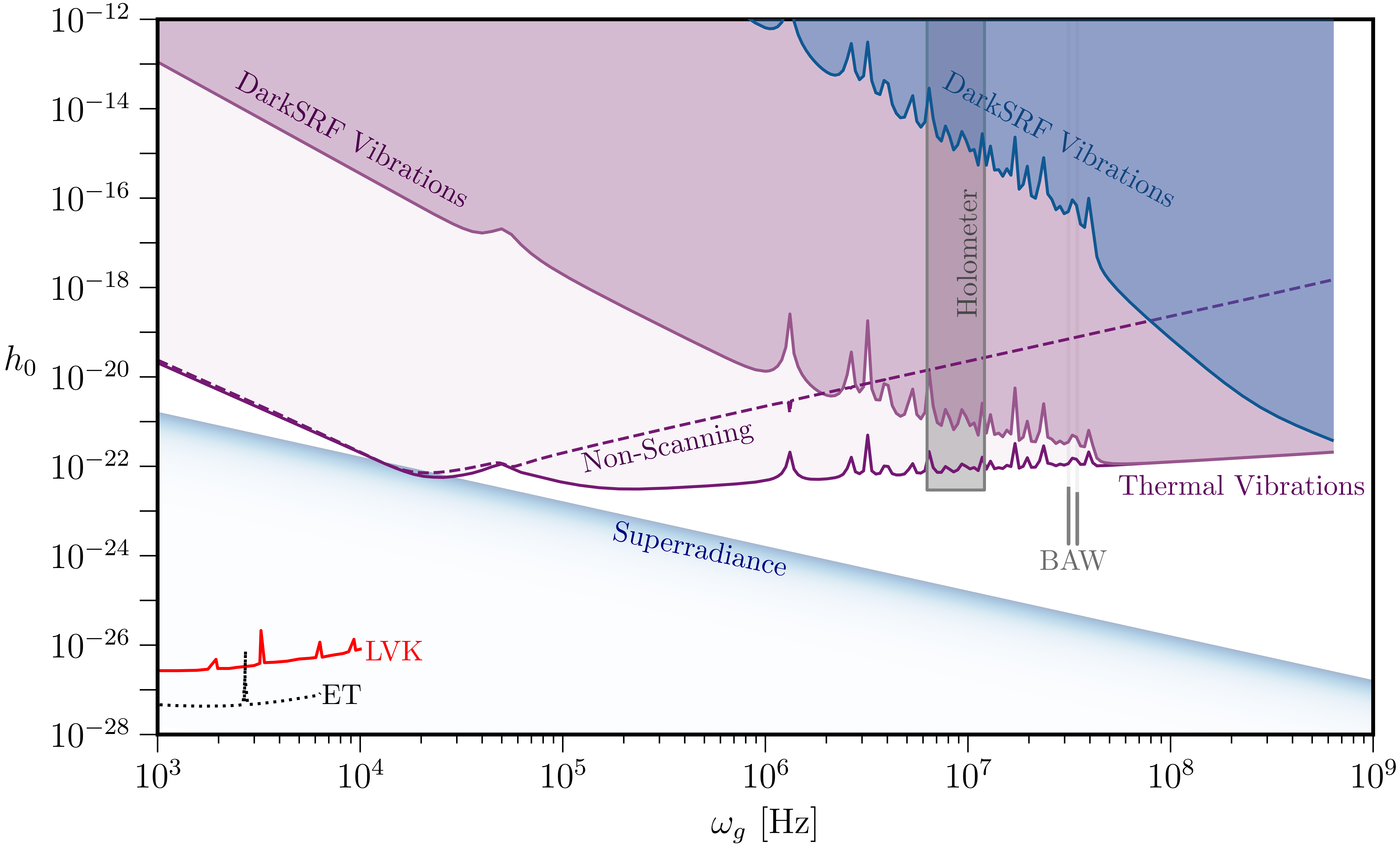


[Caves '78], [Radicati, Pegoraro, Picasso '78]

MECHANICAL VS ELECTROMAGNETIC RESONANCE

$$\omega_m \simeq \frac{c_s}{L} \ll \omega_{\rm em} \simeq \frac{c}{L}$$

MONOCHROMATIC SIGNAL



Berlin, Blas, D'Agnolo, Ellis, Harnik, Kahn, Schutte-Engel,Wentzel '23

arXiv:2303.01518

EXISTING PROTOTYPE

MAGO '05

R. Ballantini, A. Chincarini, S. Cuneo, G. Gemme,* R. Parodi, A. Podest`a, and R. Vaccarone

INFN and Universita` degli Studi di Genova, Genova, Italy

Ph. Bernard, S. Calatroni, E. Chiaveri, and R. Losito

CERN, Geneva, Switzerland

R.P. Croce, V. Galdi, V. Pierro, and I.M. Pinto

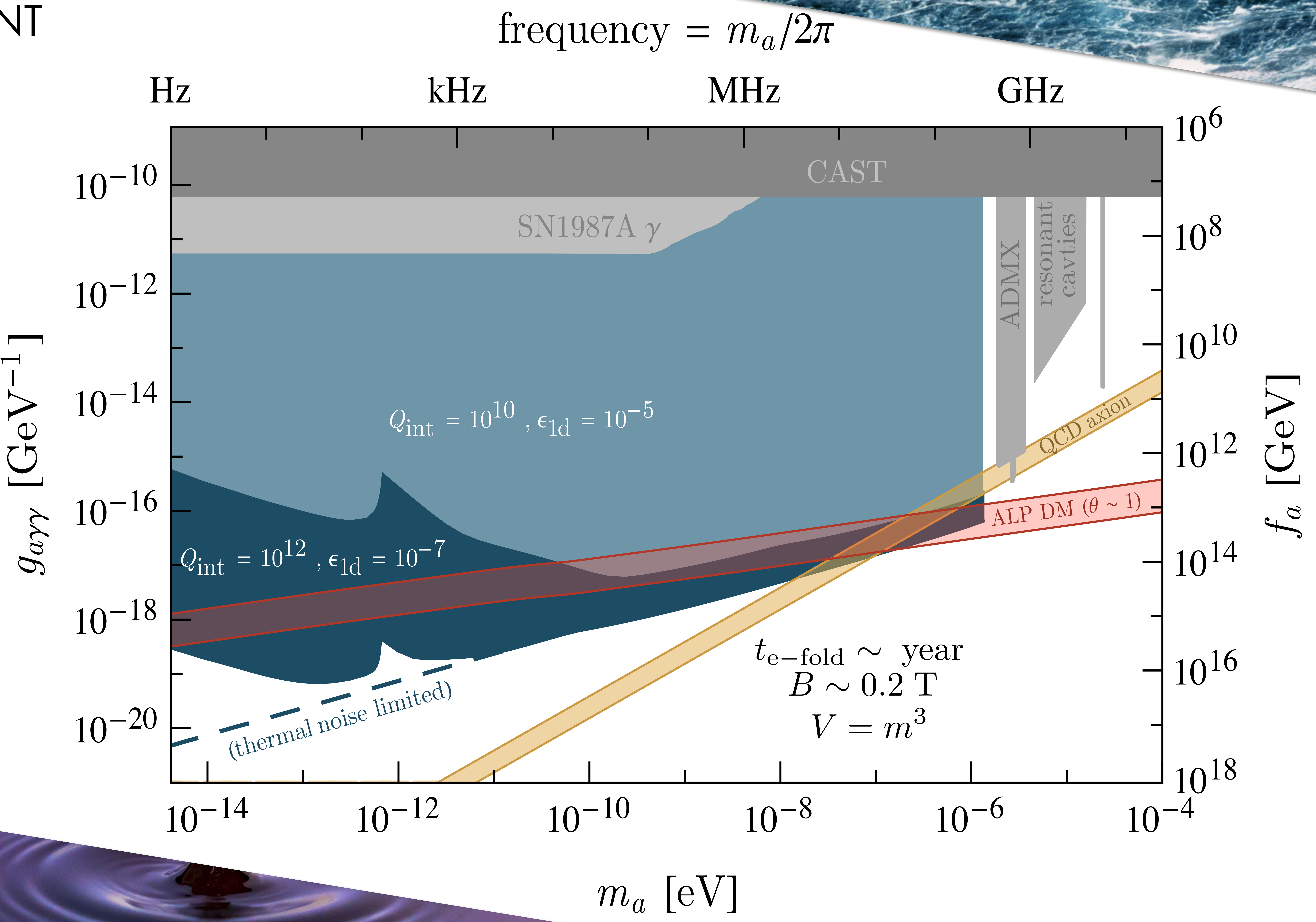
INFN, Napoli, and Universita` degli Studi del Sannio, Benevento, Italy

E. Picasso

INFN and Scuola Normale Superiore, Pisa, Italy and CERN, Geneva, Switzerland

BACKUP

RESONANT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Hz
kHz
MHz
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CAST
SN1987A γ
ADMX
resonant cavties
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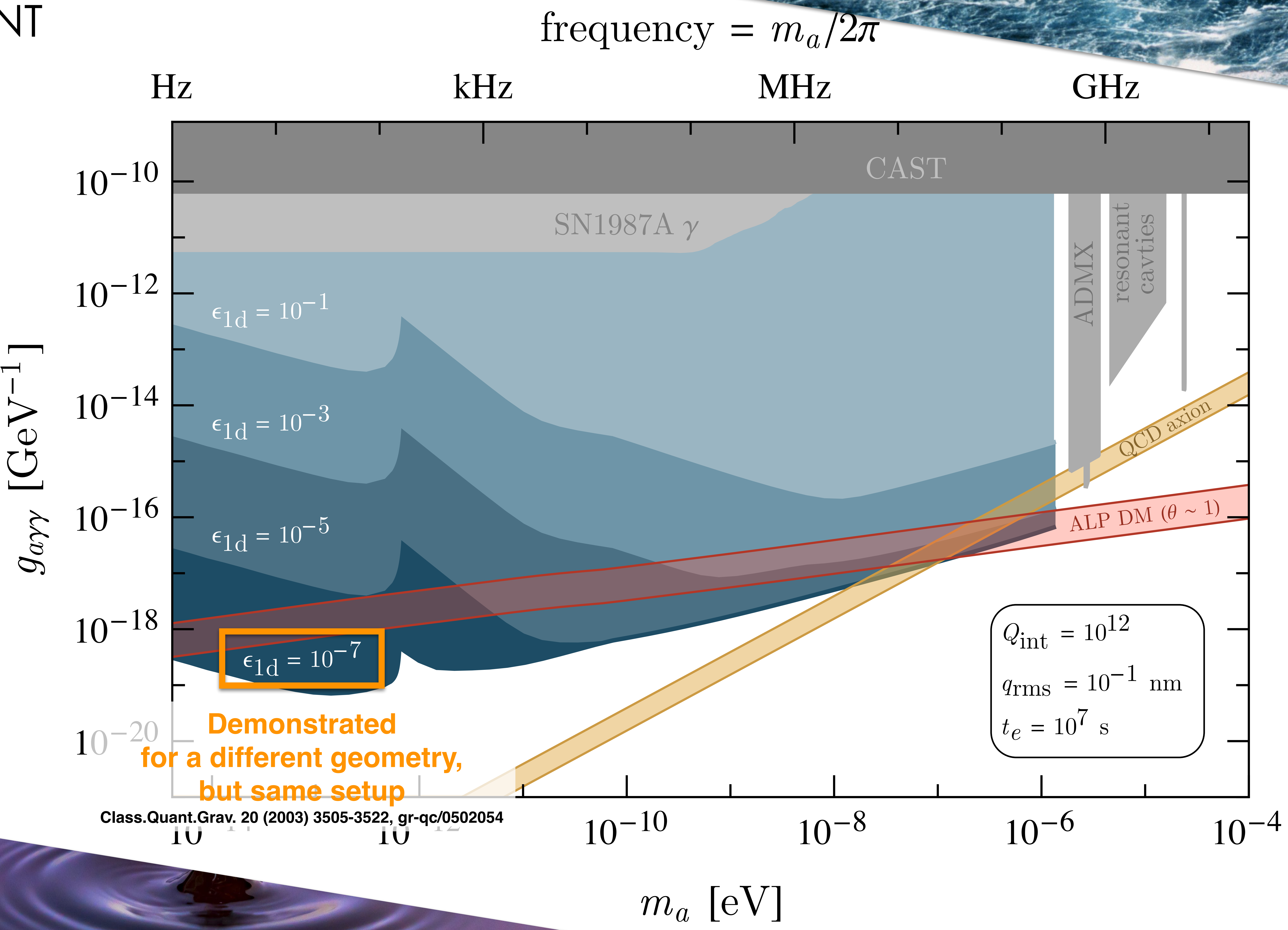
10^{12}
10^{14}
10^{16}
10^{18}
m_a [eV]
10^{-14}
10^{-12}
10^{-10}
10^{-8}
10^{-6}
10^{-4}

RESONANT



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Demonstrated
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Class.Quant.Grav. 20 (2003) 3505-3522, gr-qc/0502054
$g_{a\gamma\gamma}$ [GeV^{-1}]
10^{-10}
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10^{-20}
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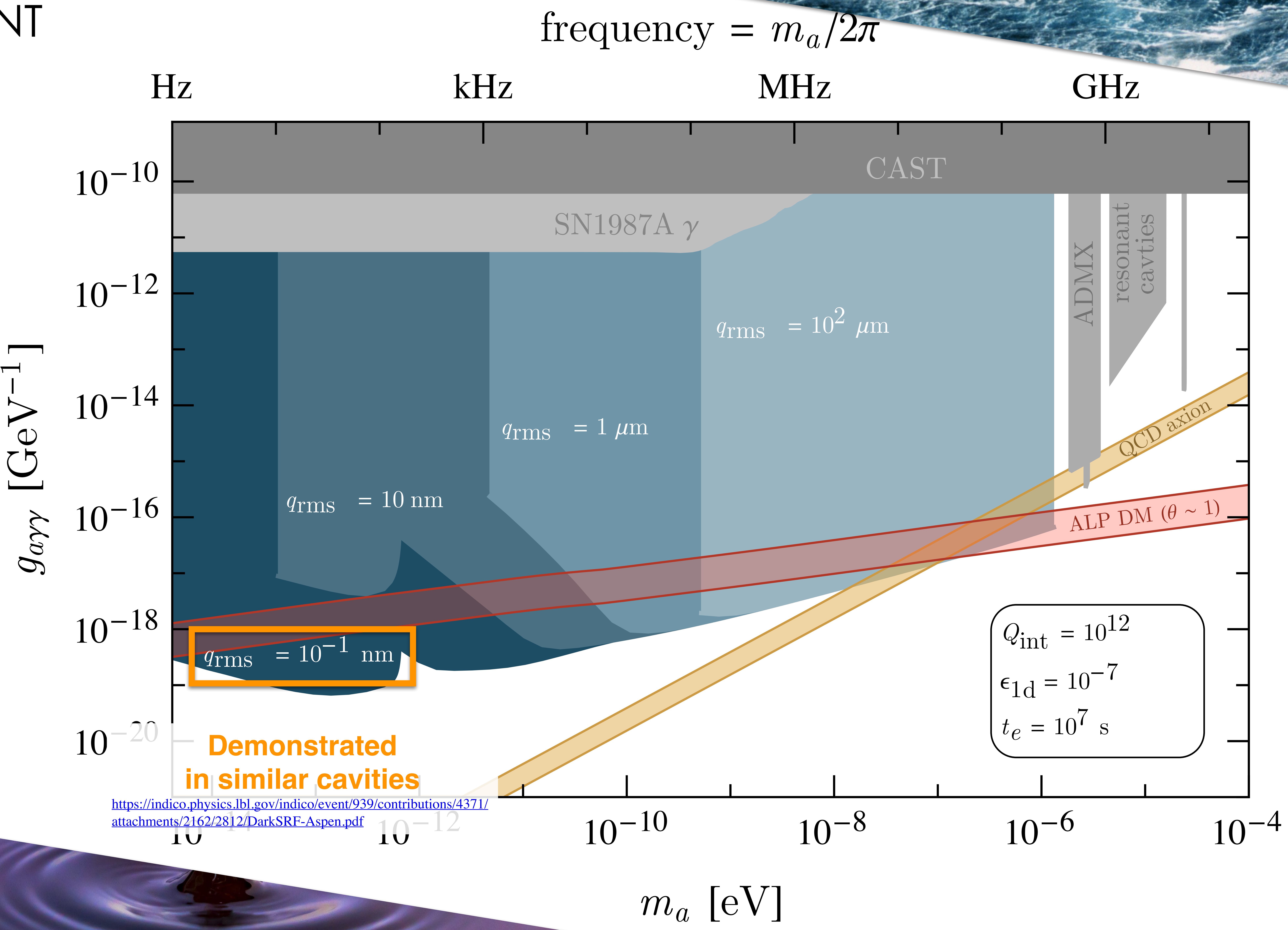
10^{-8}
10^{-6}
10^{-4}
m_a [eV]

RESONANT



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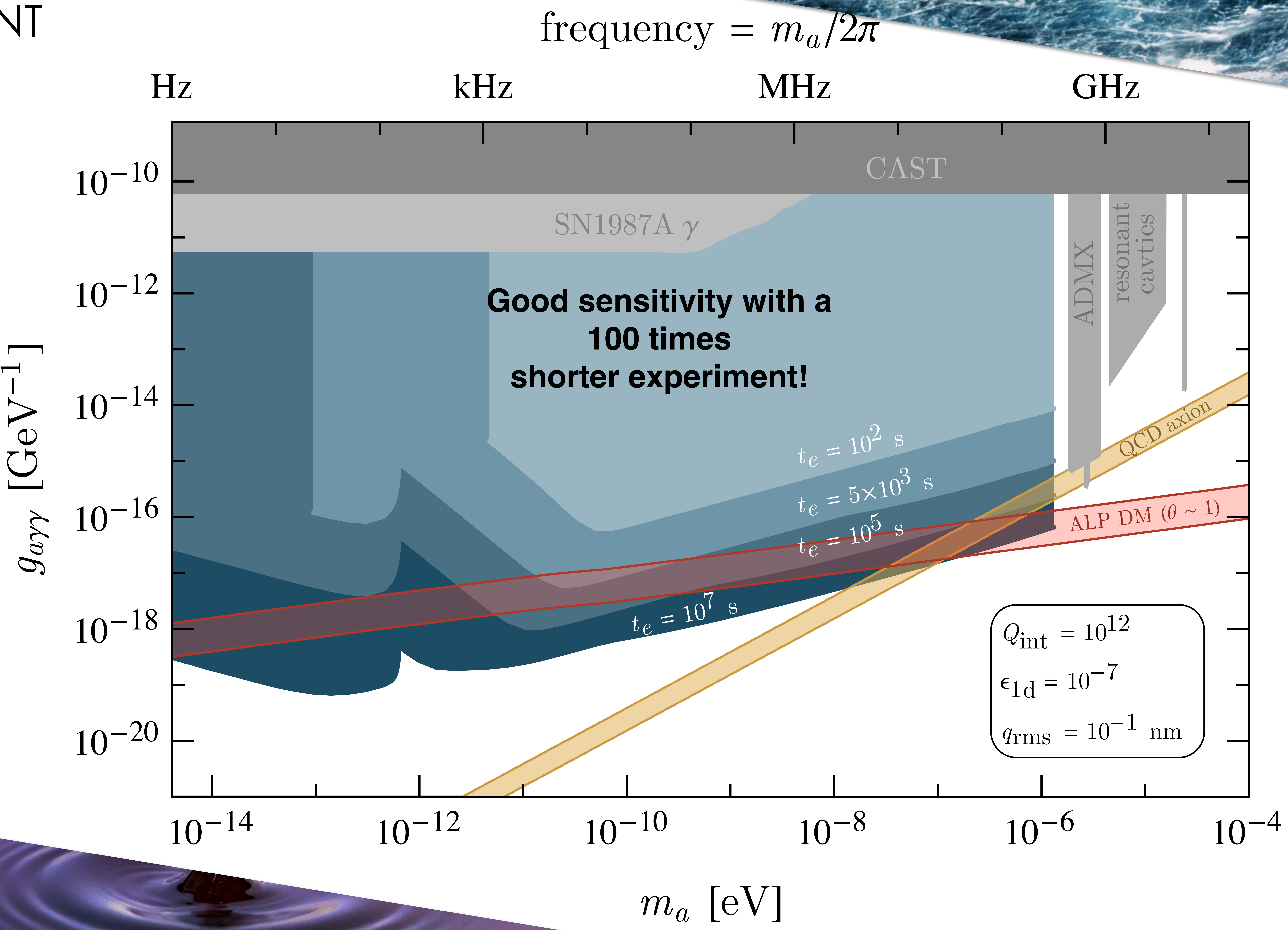
RESONANT



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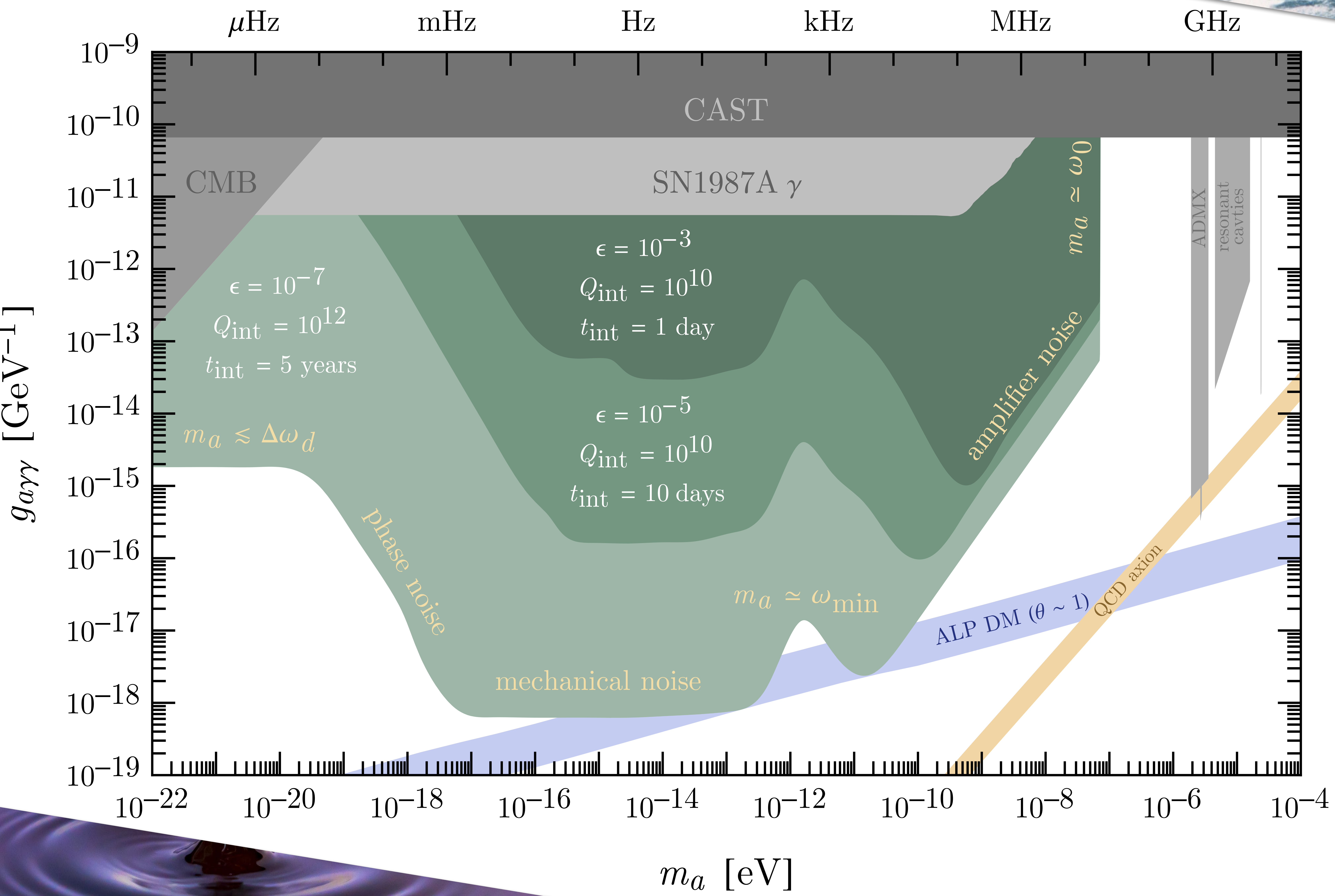
SN1987A γ
ADMX
resonant cavties
Good sensitivity with a 100 times shorter experiment!
$t_e = 10^2$ s
$t_e = 5\times10^3$ s
$t_e = 10^5$ s
$t_e = 10^7$ s
QCD axion
ALP DM ($\theta \sim 1$)
$Q_{int} = 10^{12}$
$\epsilon_{1d} = 10^{-7}$
$q_{rms} = 10^{-1}$ nm
$g_{a\gamma\gamma}$ [GeV^{-1}]
10^{-10}
10^{-12}
10^{-14}
10^{-16}
10^{-18}
10^{-20}
10^{-14}
10^{-12}
10^{-10}
10^{-8}
10^{-6}
10^{-4}
m_a [eV]

A.Berlin, RTD, S. Ellis, K. Zhou 2007.15656

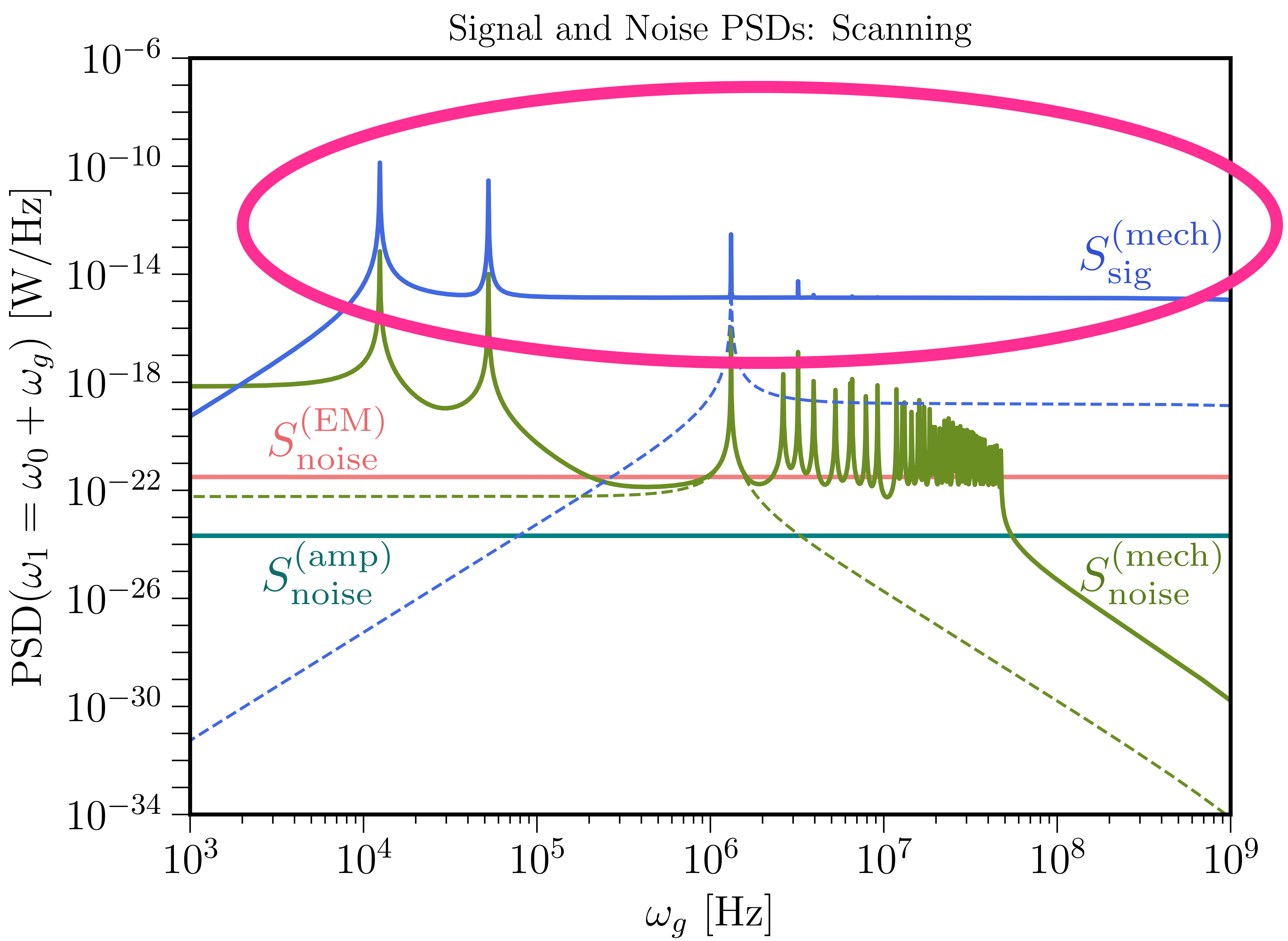


NOISE

Berlin, Blas, D'Agnolo, Ellis, Harnik, Kahn, Schutte-Engel,Wentzel '23

arXiv:2303.01518

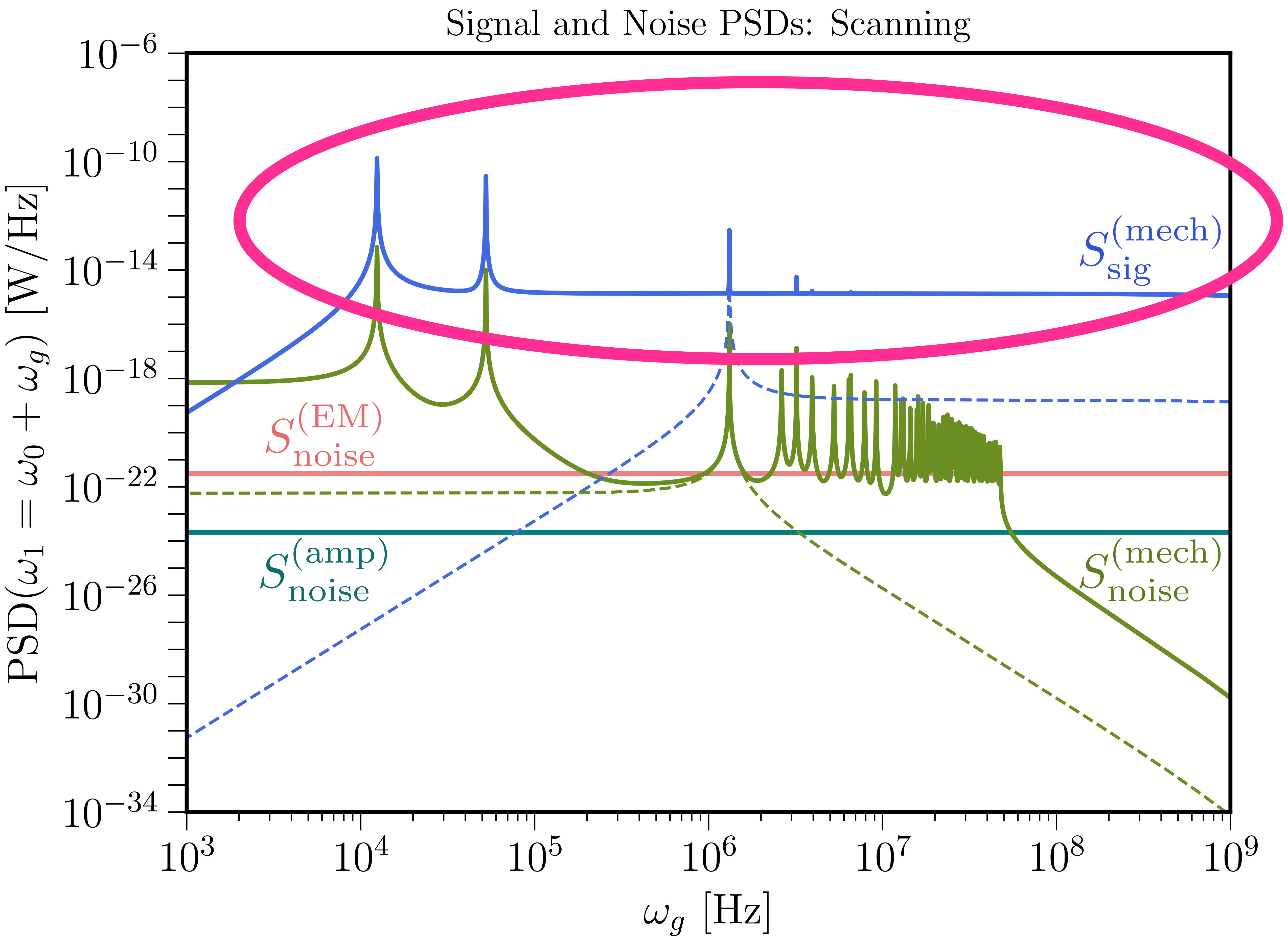
RESONANT



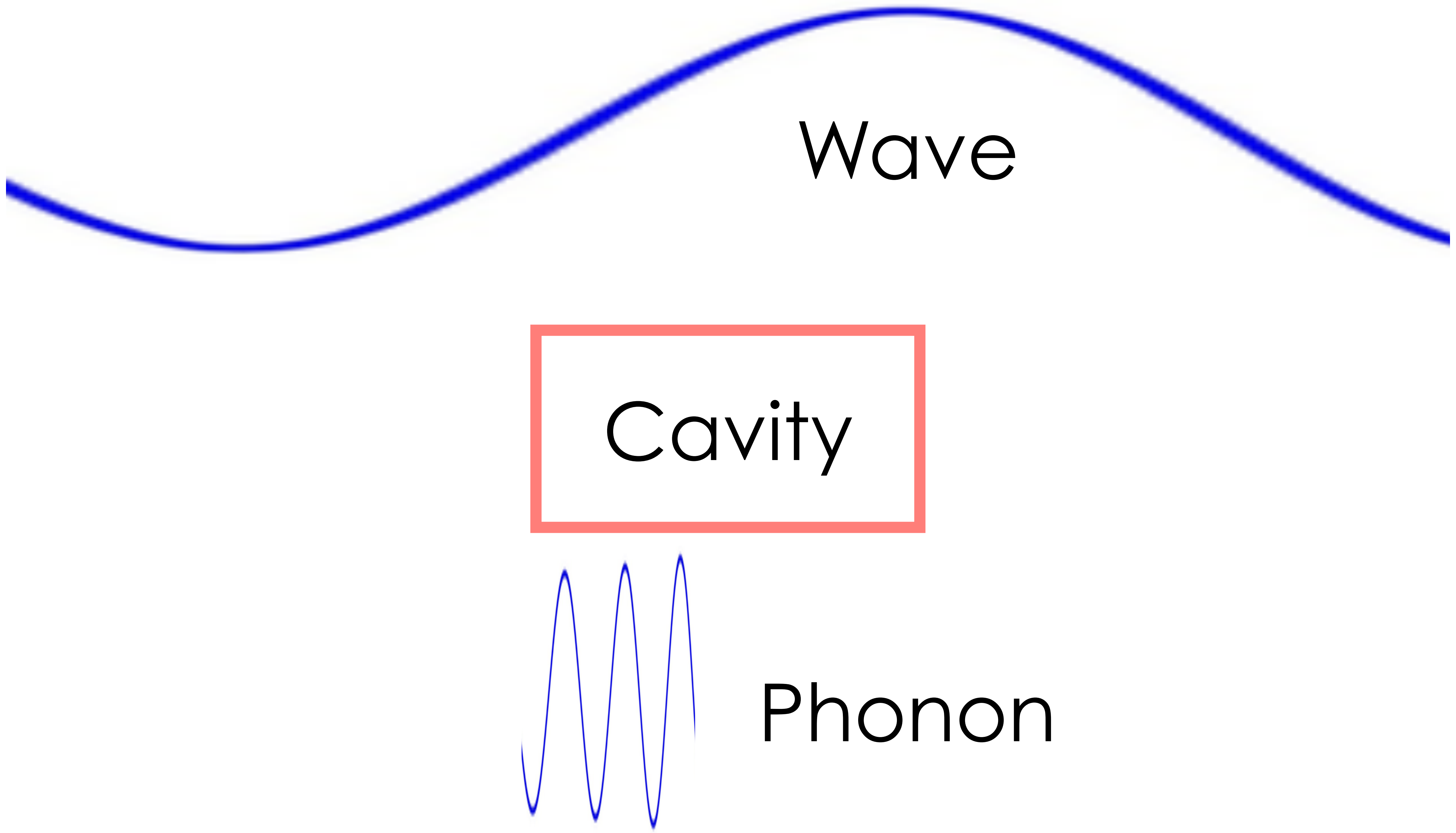
SIGNAL DOMINATED BY THE FIRST FEW MECHANICAL RESONANCES

Berlin, Blas, D'Agnolo, Ellis, Harnik, Kahn, Schutte-Engel,Wentzel '23
arXiv:2303.01518

RESONANT



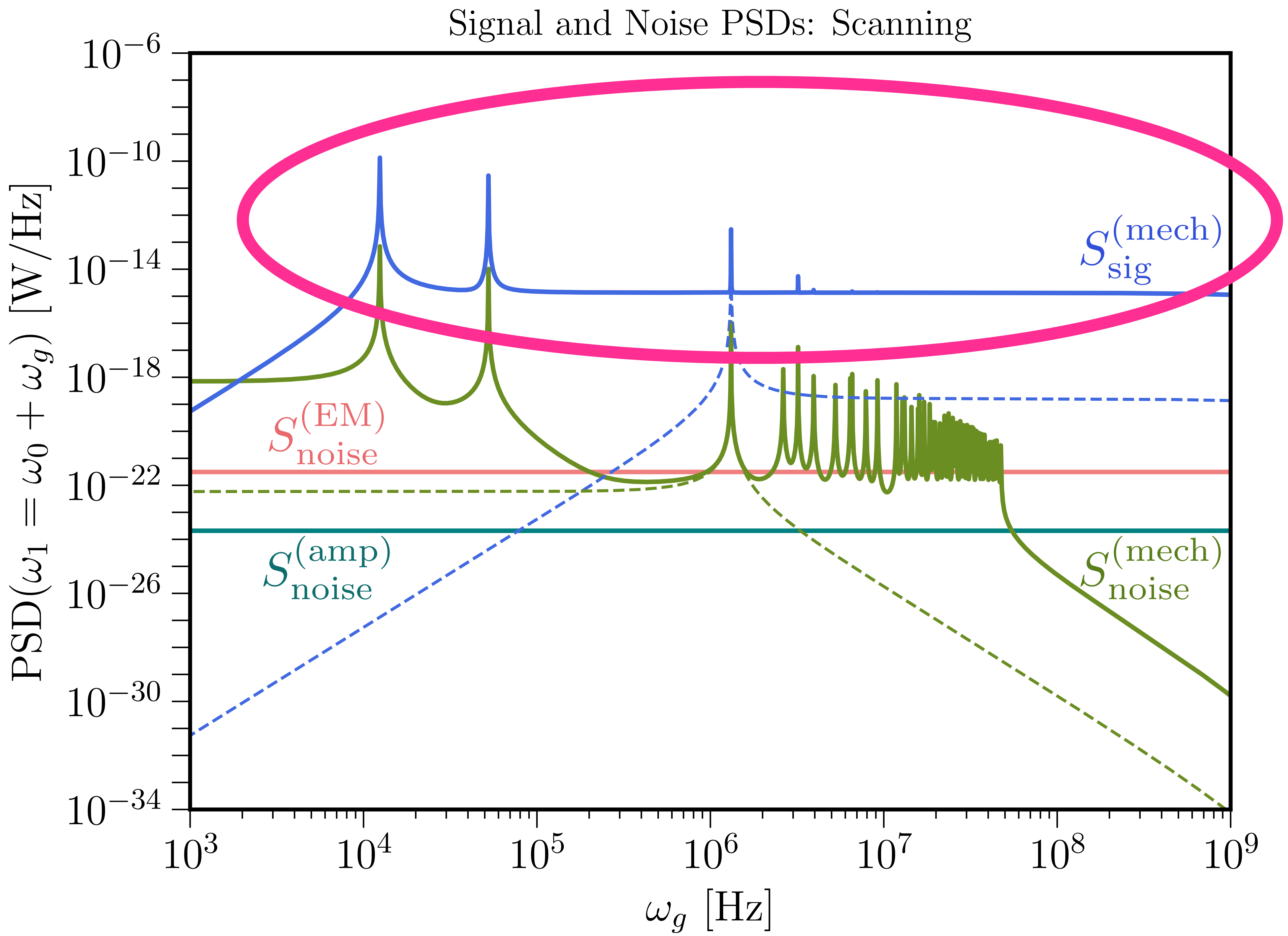
SIGNAL DOMINATED BY THE FIRST FEW MECHANICAL RESONANCES



Berlin, Blas, D'Agnolo, Ellis, Harnik, Kahn, Schutte-Engel,Wentzel '23

arXiv:2303.01518

RESONANT



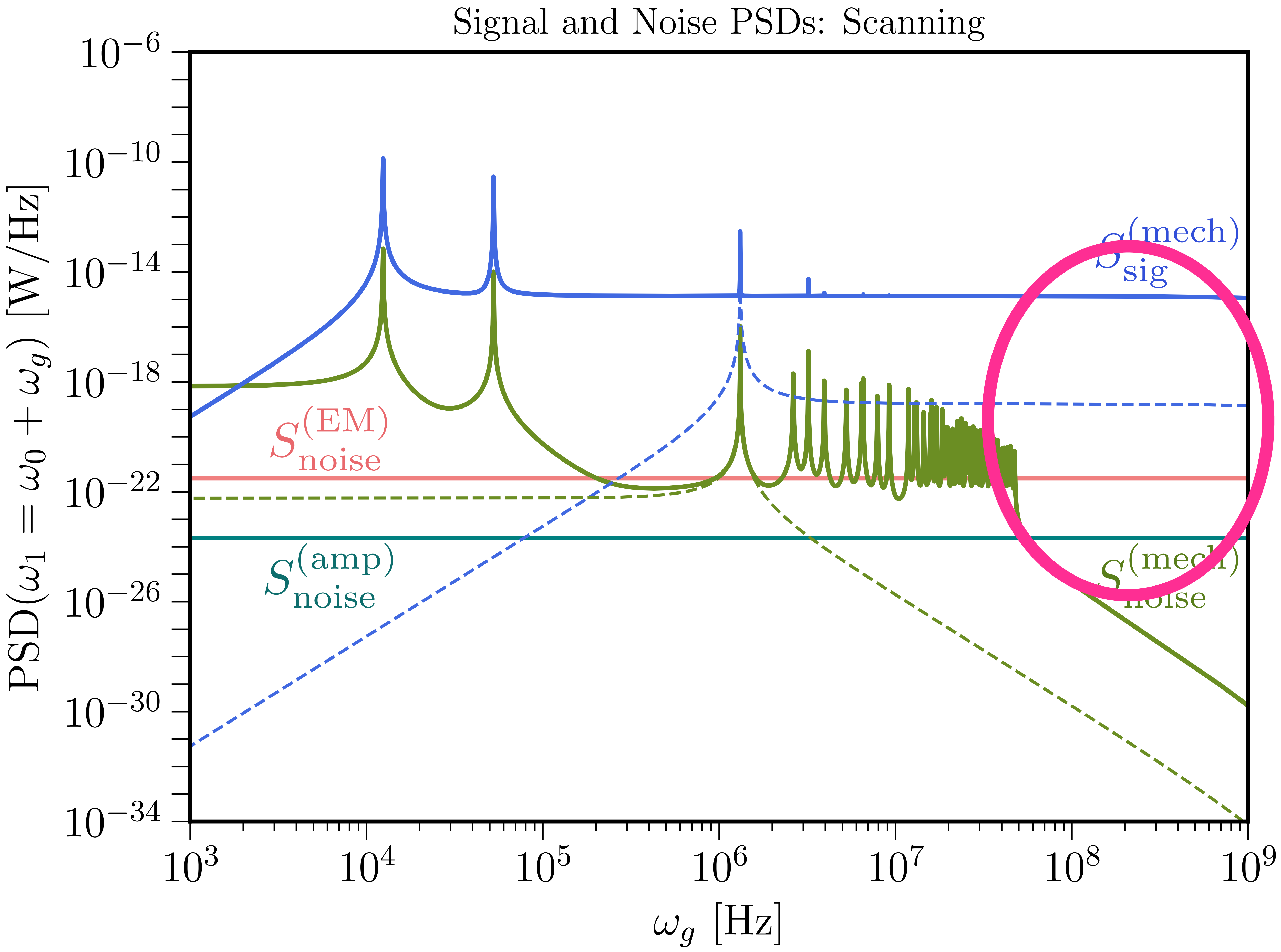
SIGNAL DOMINATED BY THE FIRST FEW MECHANICAL RESONANCES

$$\eta^g_{\mathrm{mech}} \sim \frac{1}{\omega_m^2}$$

$$Q_m \sim \frac{1}{\omega_m}$$

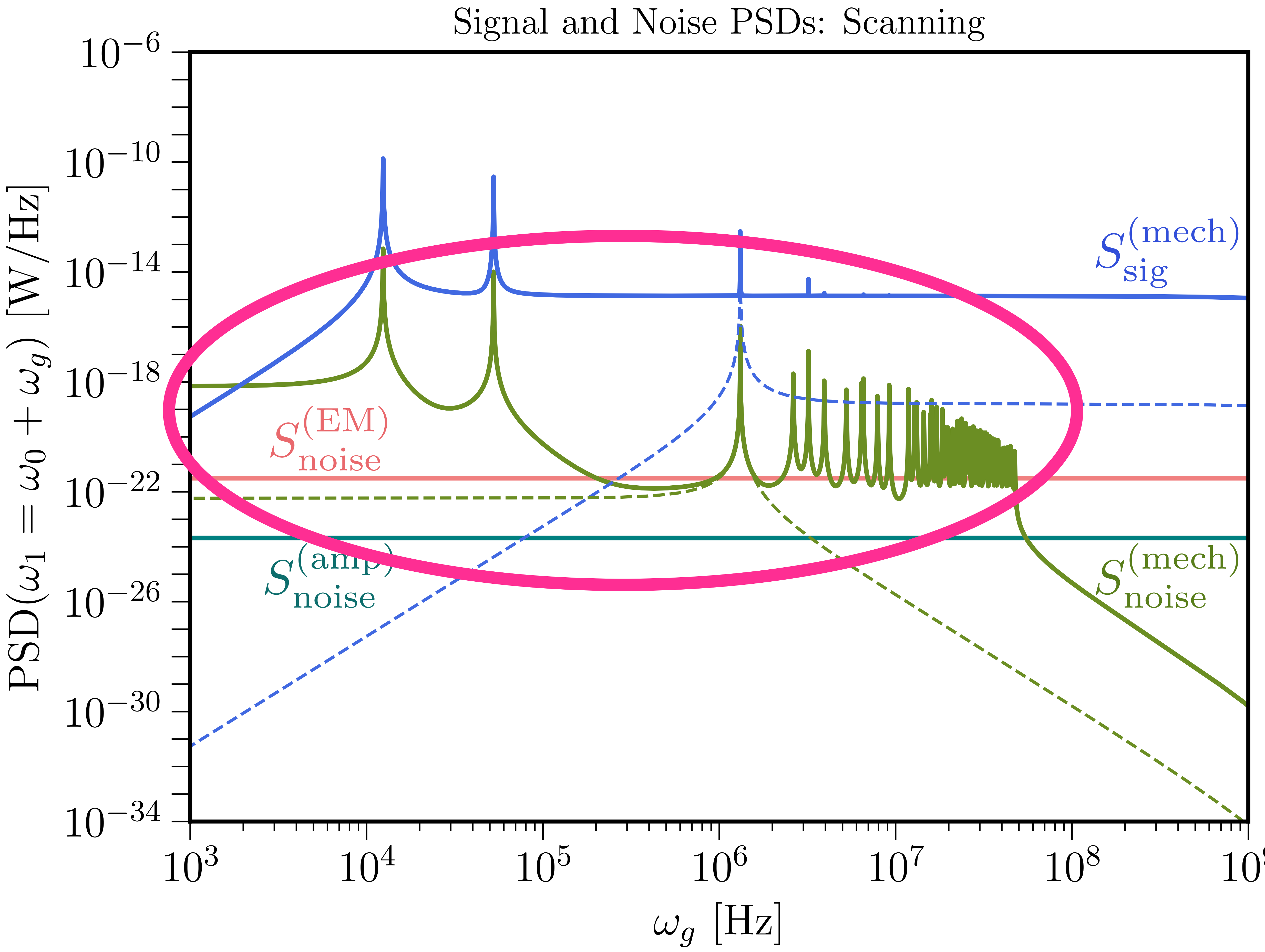
Berlin, Blas, D'Agnolo, Ellis, Harnik, Kahn, Schutte-Engel,Wentzel '23
arXiv:2303.01518

RESONANT



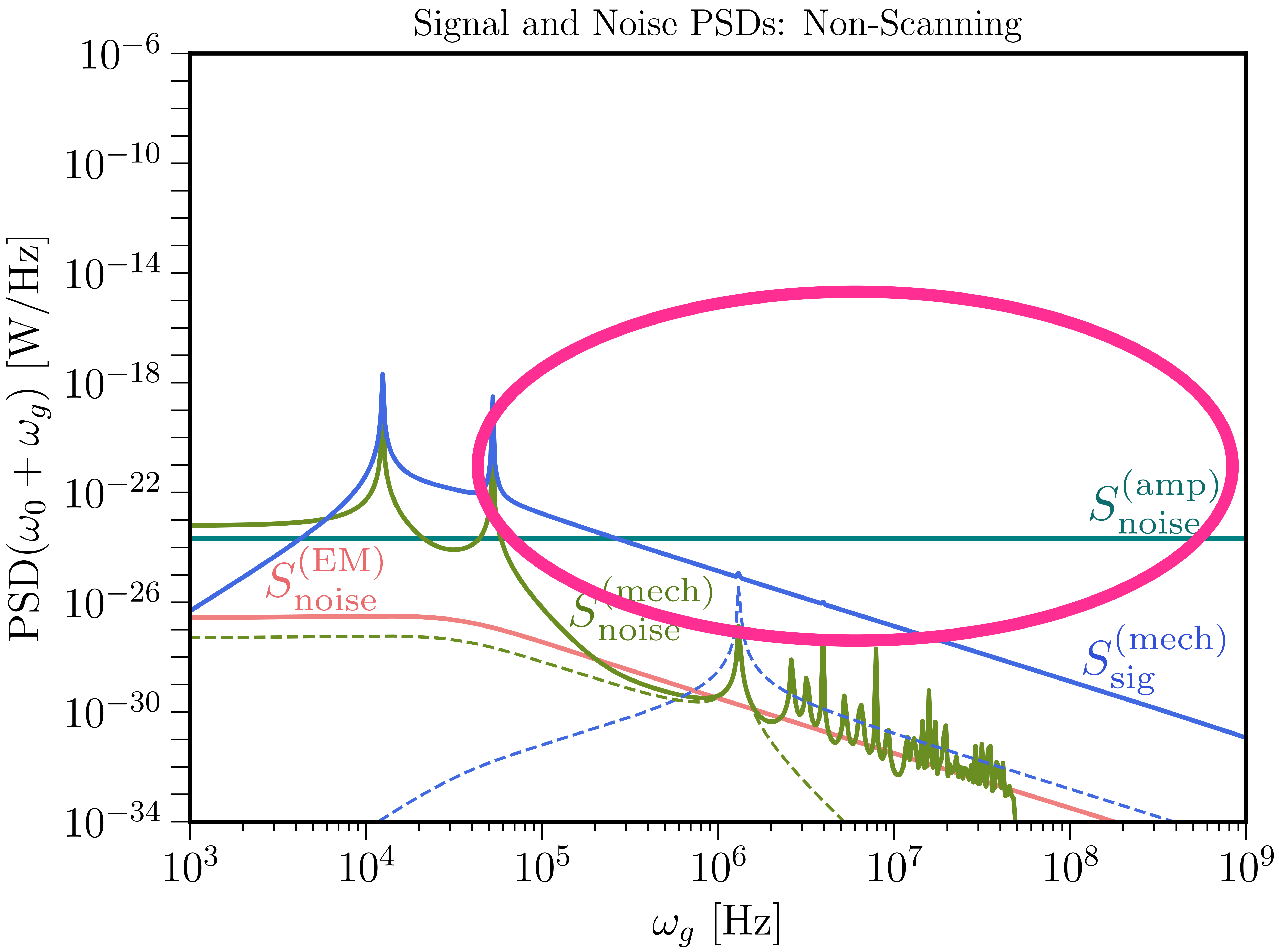
$$\mathrm{SNR}_{\mathrm{EM\ noise}} \simeq \frac{1}{64}\sqrt{\frac{\pi\, t_{\mathrm{int}}}{2\,\Delta\omega_{\mathrm{osc}}}}\; Q_{\mathrm{int}}^2\,|\eta_{\mathrm{mech}}^{g}|^2\,|\eta_{\mathrm{mech}}^{\mathrm{EM}}|^2\,\frac{P_{\mathrm{in}}}{T}\,h_0^2$$

RESONANT



$$\mathrm{SNR}_{\mathrm{mech\ noise}} \simeq \frac{1}{4}\sqrt{\frac{\pi^3\, t_{\mathrm{int}}}{2\,\Delta\omega_{\mathrm{osc}}}}\;|\eta^g_{\mathrm{mech}}(\omega_p^{\mathrm{sig}})|^2\,\frac{|\eta^{\mathrm{EM}}_{\mathrm{mech}}(\omega_p^{\mathrm{sig}})|^2}{|\eta^{\mathrm{EM}}_{\mathrm{mech}}(\omega_p^{\mathrm{noise}})|^2}\;\frac{M_{\mathrm{cav}}^2\, V_{\mathrm{cav}}^{2/3}}{S_{F_p}(\omega_g)}\;h_0^2\,\omega_g^4$$

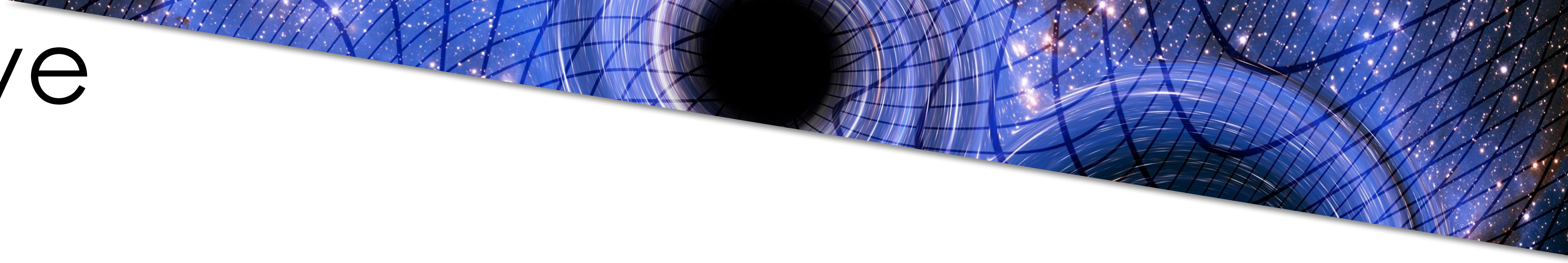
BROADBAND



$$\mathrm{SNR}^{(\mathrm{broadband})}_{\mathrm{amp\ noise}} \simeq \frac{1}{64}\sqrt{\frac{\pi\, t_{\mathrm{int}}}{2\,\Delta\omega_{\mathrm{osc}}}}\;\frac{Q_{\mathrm{int}}}{Q_{\mathrm{cpl}}}\,|\eta^{g}_{\mathrm{mech}}|^2\,|\eta^{\mathrm{EM}}_{\mathrm{mech}}|^2\;\frac{P_{\mathrm{in}}}{T}\,h_0^2\,\frac{\omega_0}{\omega_g^2}$$

Proper Detector Frame (PDF) Vs **Transverse Traceless (TT) Frame**

(Fermi Normal Coordinates)



TT gauge = comoving with the wave

$$(t_{\mathrm{TT}}, x_{\mathrm{TT}}, y_{\mathrm{TT}}, z_{\mathrm{TT}})$$

$$t_{\mathrm{TT}} \simeq t - \frac{i}{4}\omega_g\,(x^2 - y^2)\,h_+\,e^{i\omega_g t} \quad , \quad x_{\mathrm{TT}} \simeq x - \frac{1}{2}x\,(1 - i\omega_g z)\,h_+\,e^{i\omega_g t}$$

$$y_{\mathrm{TT}} \simeq y + \frac{1}{2}y\,(1 - i\omega_g z)\,h_+\,e^{i\omega_g t} \quad , \quad z_{\mathrm{TT}} \simeq z - \frac{i}{4}\omega_g\,(x^2 - y^2)\,h_+\,e^{i\omega_g t}$$

TT gauge = comoving with the wave

$$(t_{\mathrm{TT}}, x_{\mathrm{TT}}, y_{\mathrm{TT}}, z_{\mathrm{TT}})$$

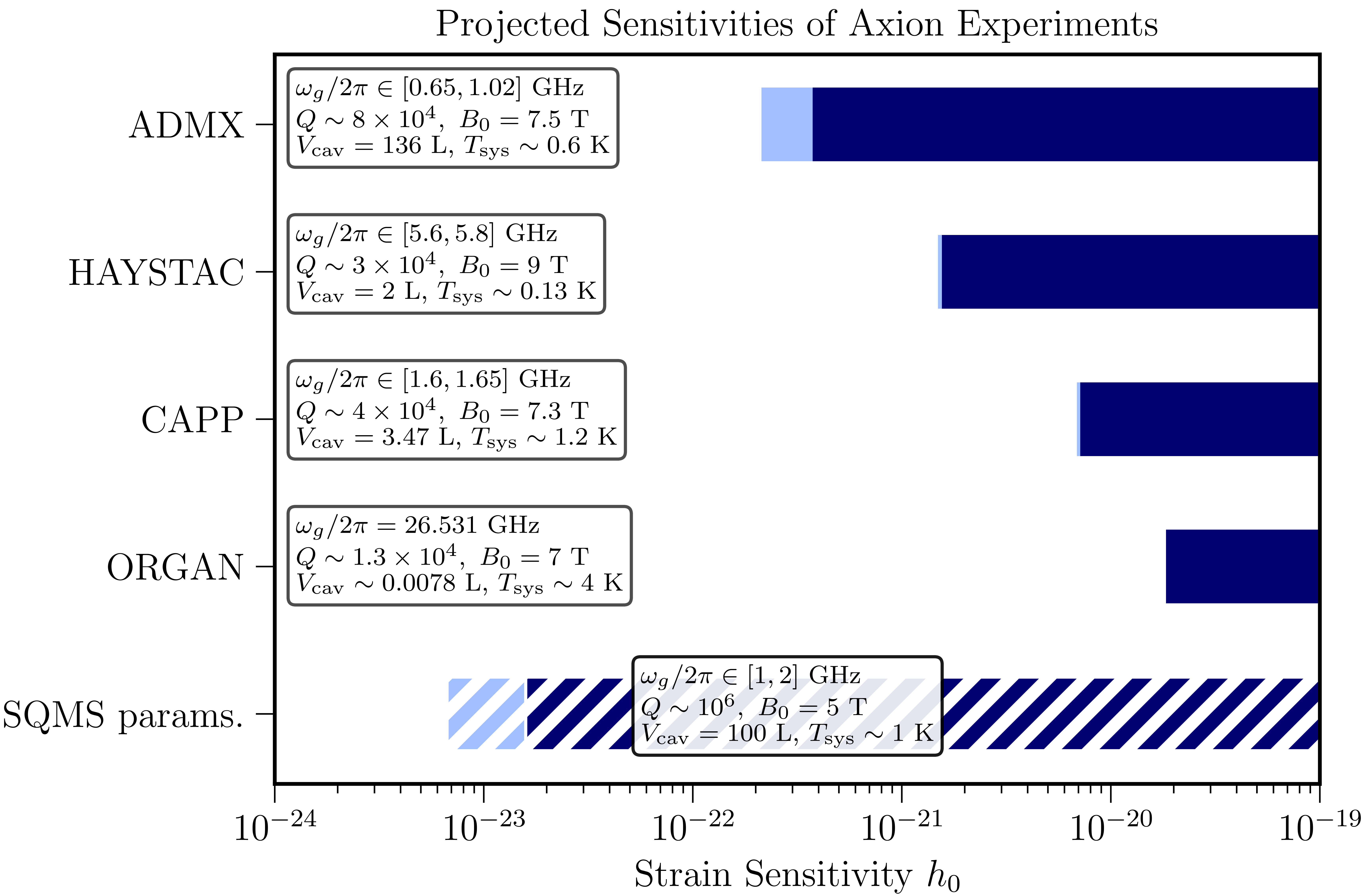
Wrong Conclusion

Theorem: $j^{\mu}_{\mathrm{eff},\mathrm{TT}} = 0$ → No signal

Doubly Wrong:

1. Impossible to prepare a uniform B-field in the TT frame
2. Even if you could do it, there would still be a signal (wire moving)

GRAVITATIONAL WAVE DETECTION



Berlin, Blas, D'Agnolo, Ellis, Harnik, Kahn, Schutte-Engel '21

SUPERRADIANCE

$$\omega_g \simeq 2m_a \simeq \mathrm{GHz}\,\left(\frac{10^{-4} M_\odot}{M_{\mathrm{BH}}}\right)$$

$$h = 10^{-24}\,\left(\frac{\Delta a_*}{0.1}\right)\left(\frac{1\,\mathrm{kpc}}{D}\right)\left(\frac{M_b}{10^{-1}\,M_\odot}\right)\left(\frac{\alpha}{0.2}\right)^7$$

Signal from:[Arvanitaki, Geraci ’12]

PRIMORDIAL BHs

$$\frac{N_{\rm cycles}}{Q} \simeq \left(\frac{\rm GHz}{f_g}\right)^{5/3} \left(\frac{10^{-9} M_\odot}{m_{\rm PBH}}\right)^{5/3} \qquad Q = 10^6$$

$$h \simeq 10^{-26} \left(\frac{1{\rm pc}}{D}\right) \left(\frac{m_{\rm PBH}}{10^{-9} M_\odot}\right)^{5/3} \left(\frac{\omega_g}{\rm GHz}\right)^{2/3}$$

If they are DM and are all in binaries (1 year) $\quad D \simeq 10^{-3}\ {\rm pc}$

If they are DM $\quad D \simeq 10\ {\rm pc} \left(\frac{m_{\rm PBH}}{10^{-9} M_\odot}\right)$

POWER FROM THE SUN

Power

$$\frac{dP}{dE} \sim \text{const.}$$

$$\frac{d\Gamma}{dE} \sim \frac{1}{E}$$

Number of gravitons

Signal from:[Weinberg '65]

POWER FROM THE SUN

$$\Gamma \sim \exp[B \log(E/\Lambda)], \quad B \ll 1$$

Signal from:[Weinberg ‘65]

POWER FROM THE SUN

Signal

$$P_{\text{tot}} \simeq 6 \times 10^{14}\ \frac{\text{erg}}{\text{s}} \simeq 10^{11}\ \text{eV}^2$$

Total emitted power (in gravitons)

$$P_{\text{exp}} \simeq 10^{-12}\ \text{eV}^2$$

Power reaching a m-sized experiment on Earth (in gravitons)

Signal from:[Weinberg '65]

POWER FROM THE SUN

Signal

$$P_{\rm tot} \simeq 6 \times 10^{14}\ \frac{\rm erg}{\rm s} \simeq 10^{11}\ {\rm eV}^2$$

Total emitted power (in gravitons)

$$P_{\rm exp} \simeq 10^{-12}\ {\rm eV}^2$$

Power reaching a m-sized experiment on Earth (in gravitons)

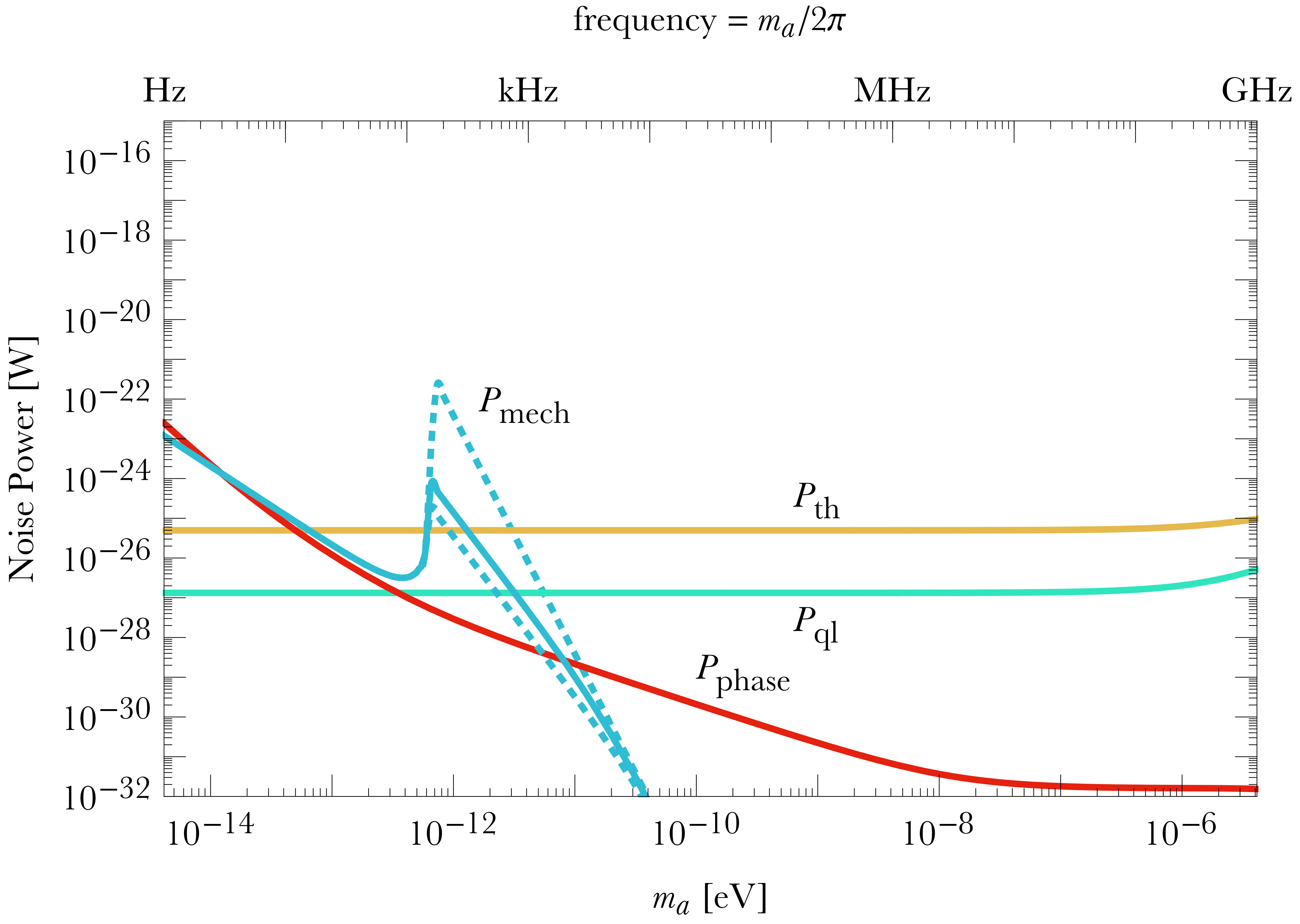
Noise

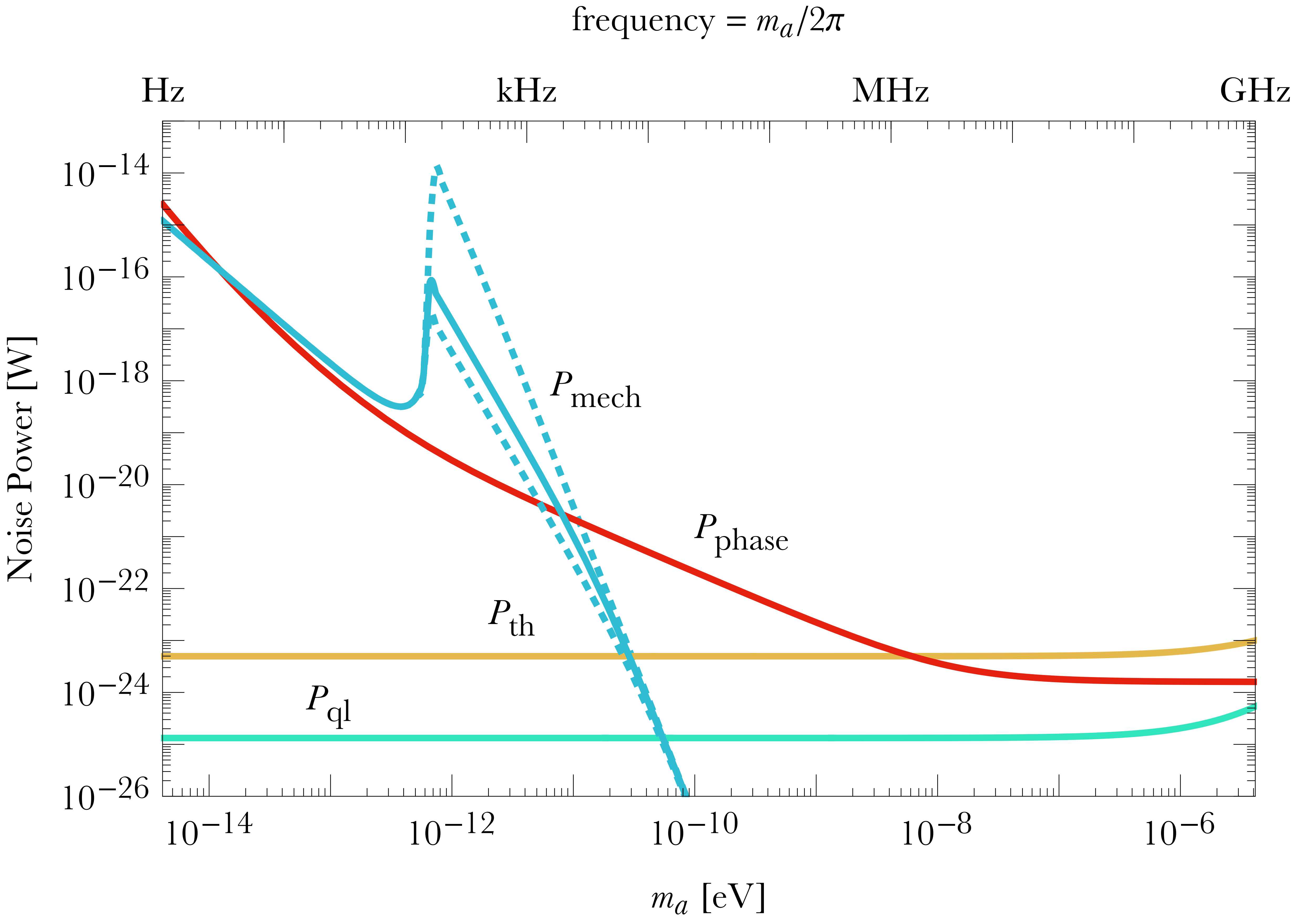
$$P_{\rm th} \simeq T\Delta\omega \simeq {\rm eV}^2 \left(\frac{T}{\rm K}\right)$$

Thermal Noise in the bandwidth of the signal (in photons)

Signal from:[Weinberg '65]

NOISE PSDs

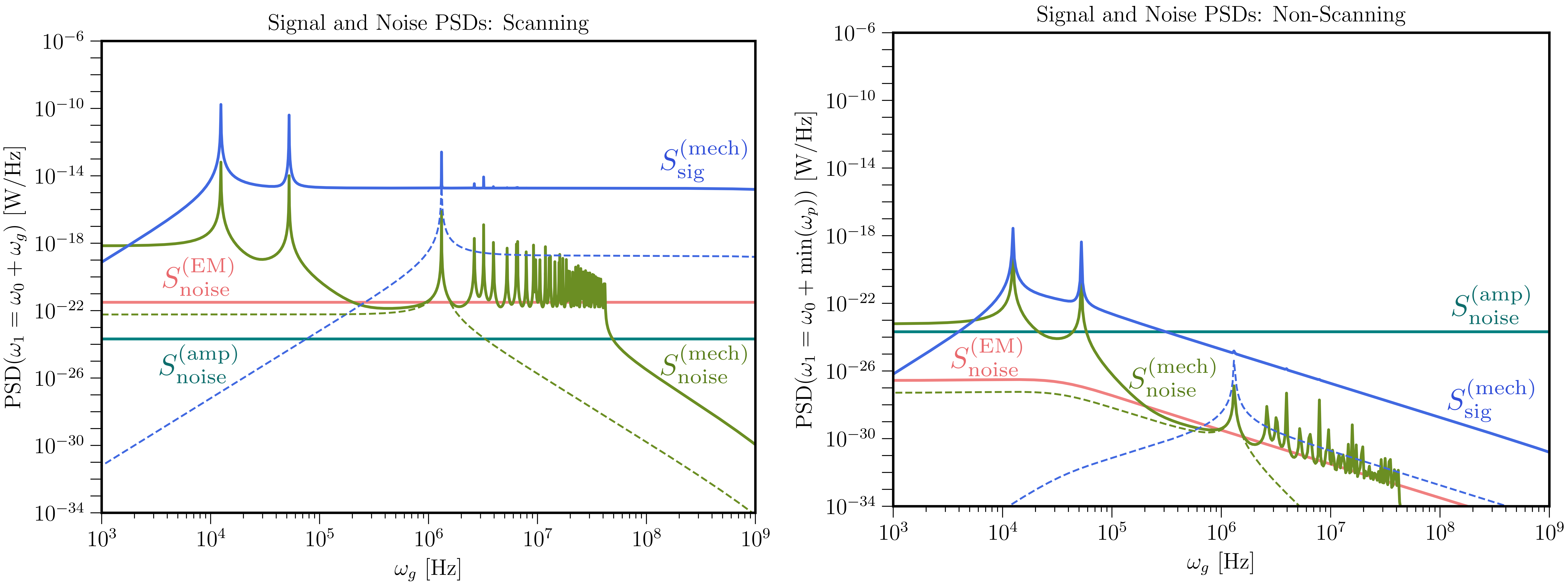




$$\epsilon_{1d} = 10^{-7}, \quad Q = 10^{12}$$

$$\epsilon_{1d} = 10^{-5}, \quad Q = 10^{10}$$

NOISE PSDs GWs



MADMAX and LAMPOST

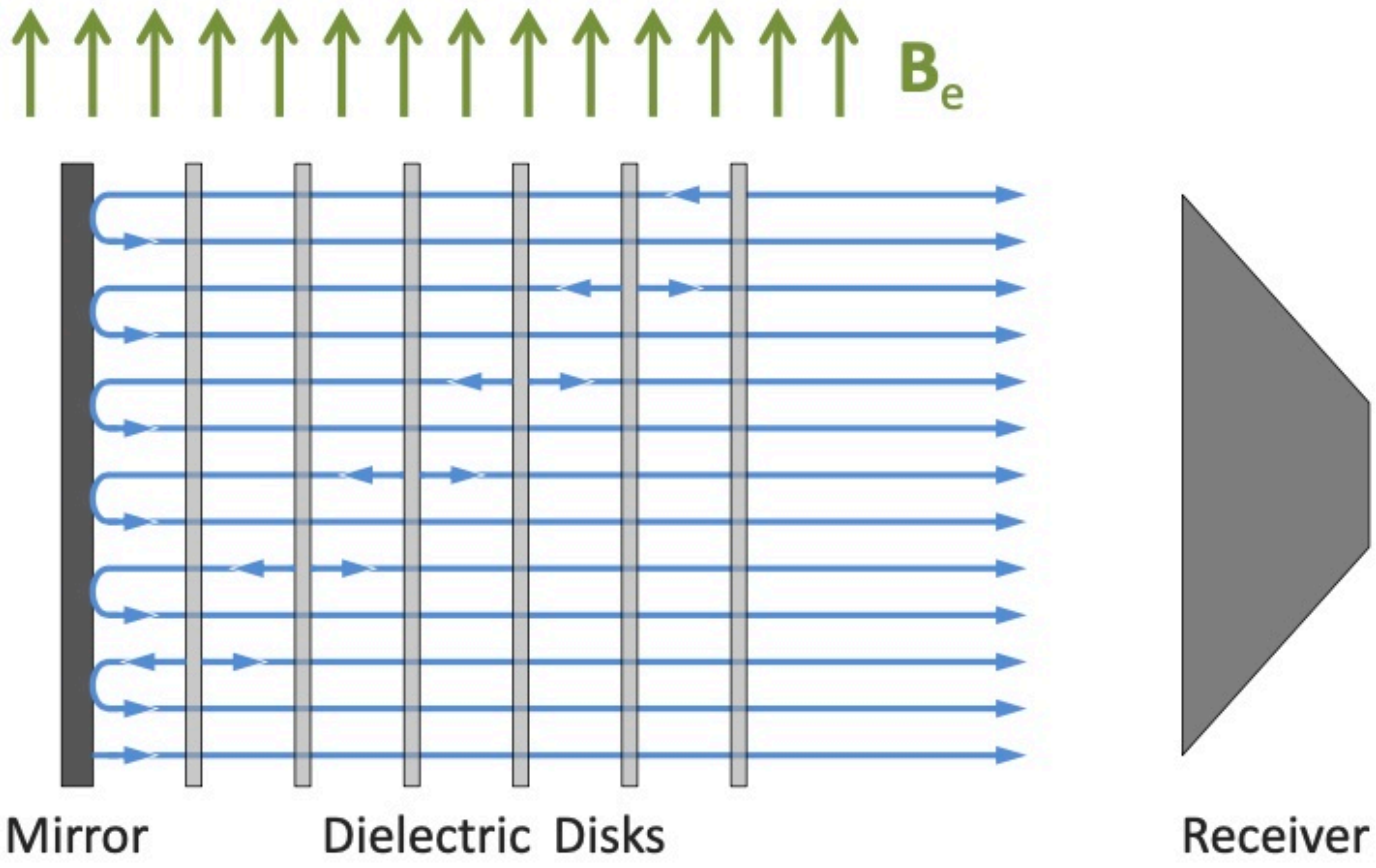
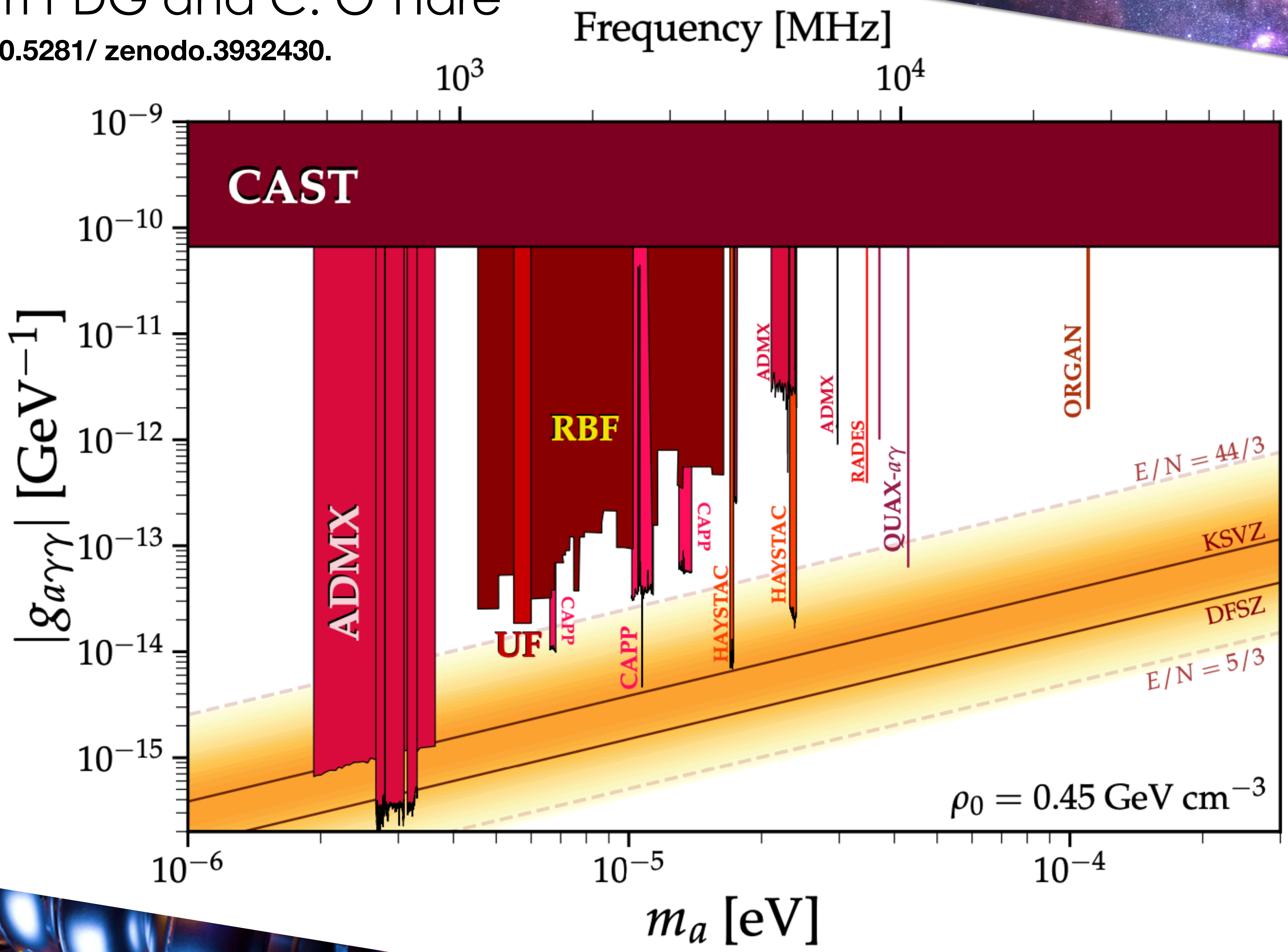


Figure from PDG and C. O'Hare

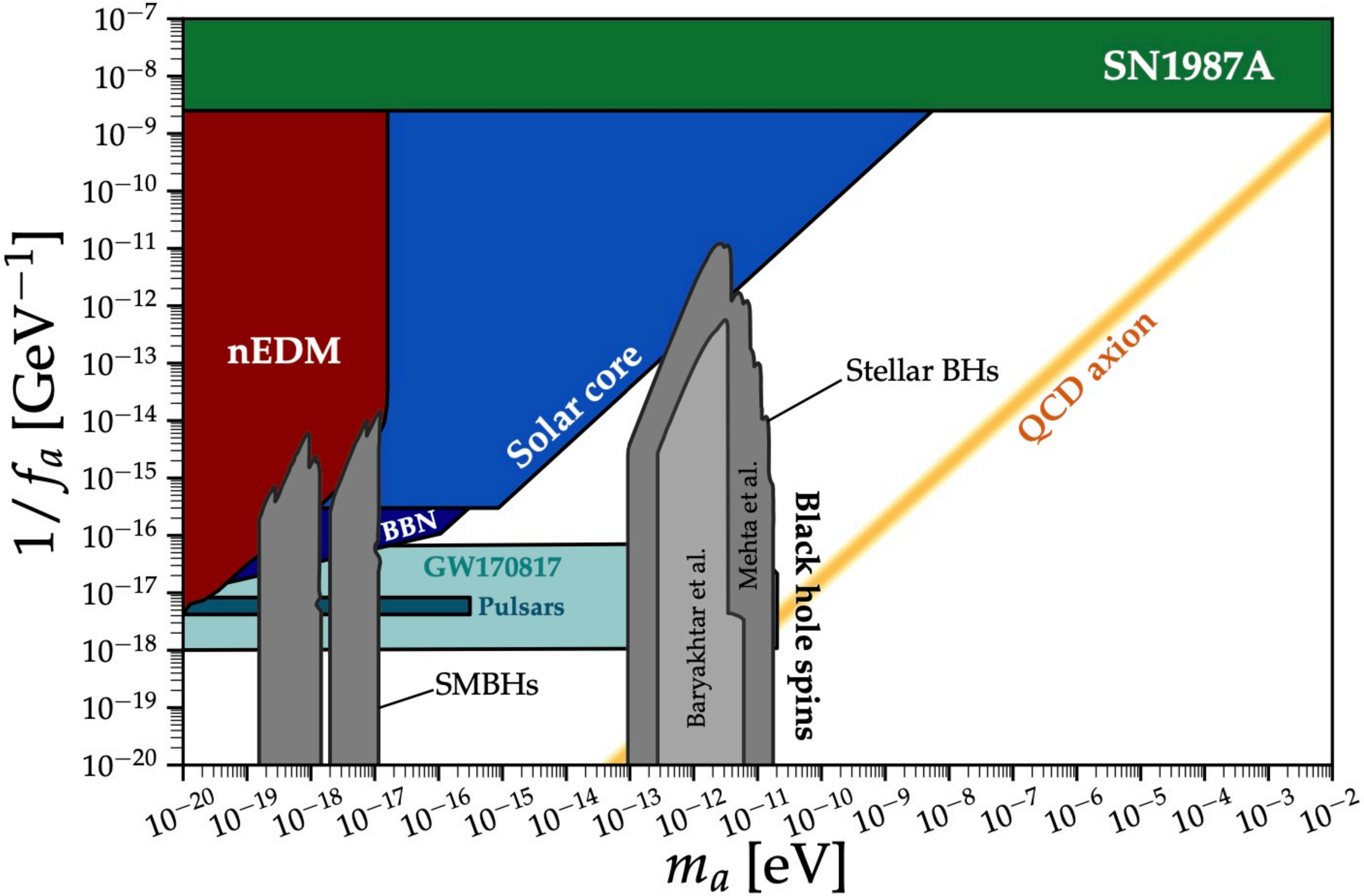
https://doi.org/10.5281/ zenodo.3932430.



Astro & Lab Constraints
DM Radio
HETERODYNE
Axion DM
ALP DM
CASPER
QCD
HP+QCD+DM
10^{-10}
10^{-15}
10^{-20}
10^{-25}
10^{-30}
$g_{a\gamma\gamma}(\mathrm{GeV}^{-1})$
10^{7}
10^{12}
10^{17}
10^{22}
10^{27}
$f_a(\mathrm{GeV})$
10^{-20}
10^{-15}
10^{-10}
10^{-5}
1
$m_a(\mathrm{eV})$

Figure from PDG and C. O'Hare

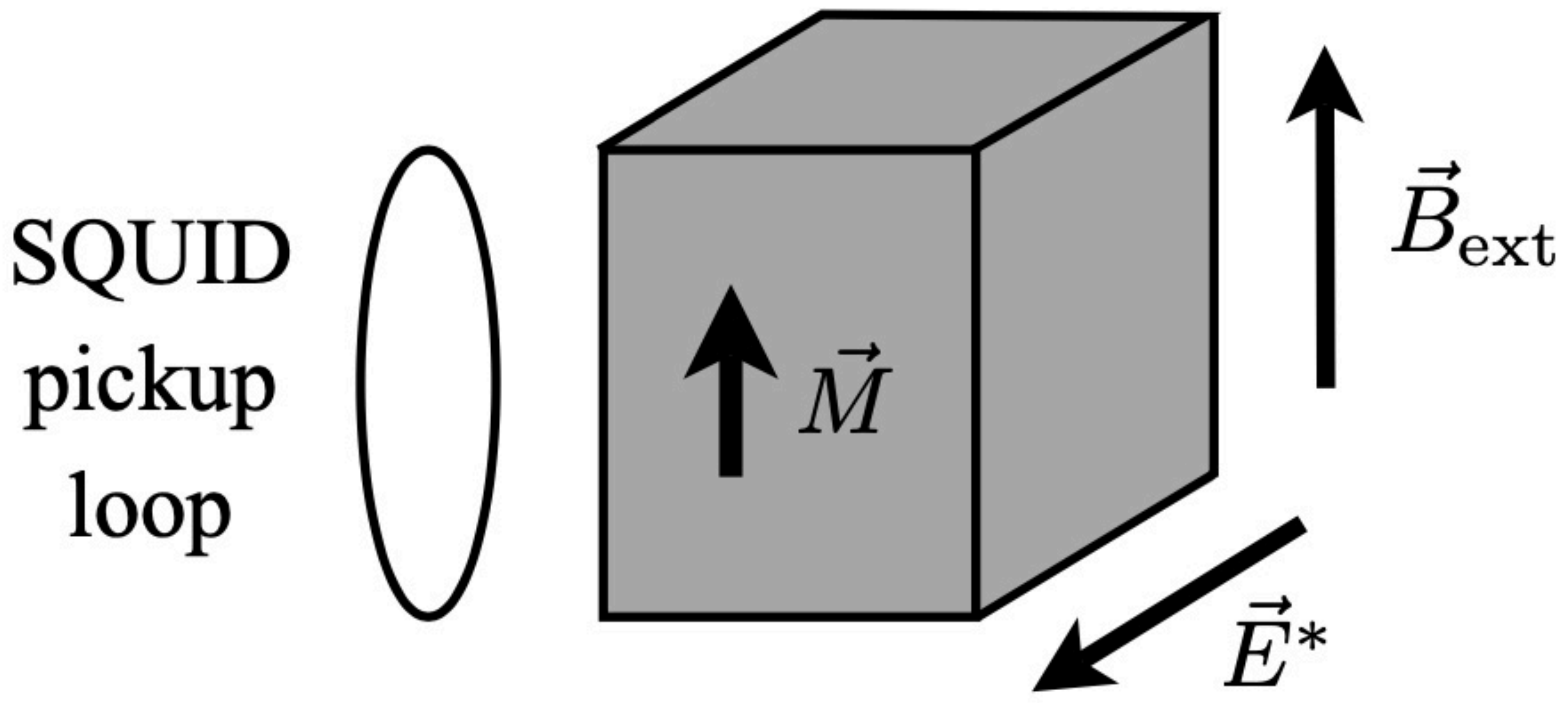
https://doi.org/10.5281/ zenodo.3932430.



Valid for QCD Axion!

CASPER

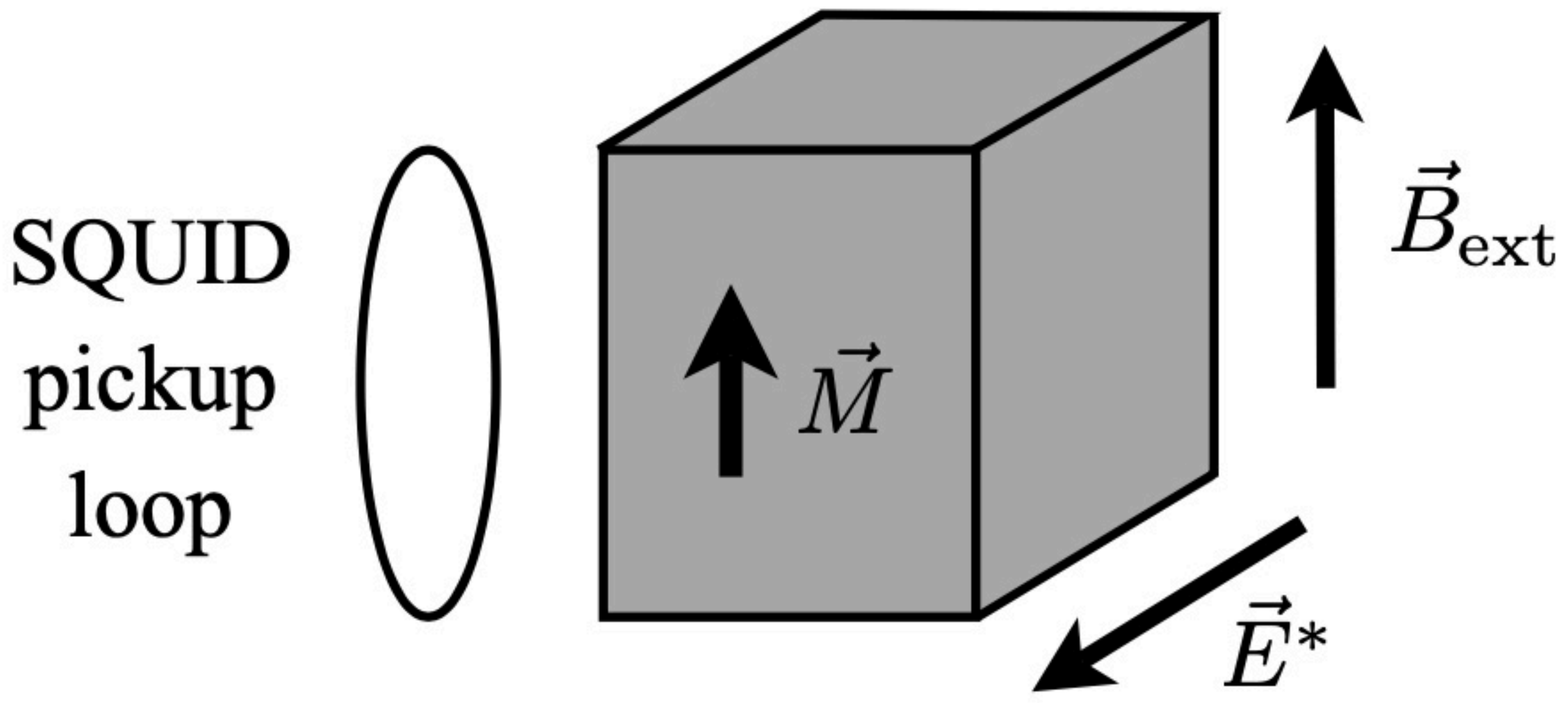
$$H \simeq \frac{a}{m_N f_a} \vec{\sigma} \cdot \vec{E}$$



Without the axion you have a magnetisation component precessing around B (Larmor frequency)

CASPER

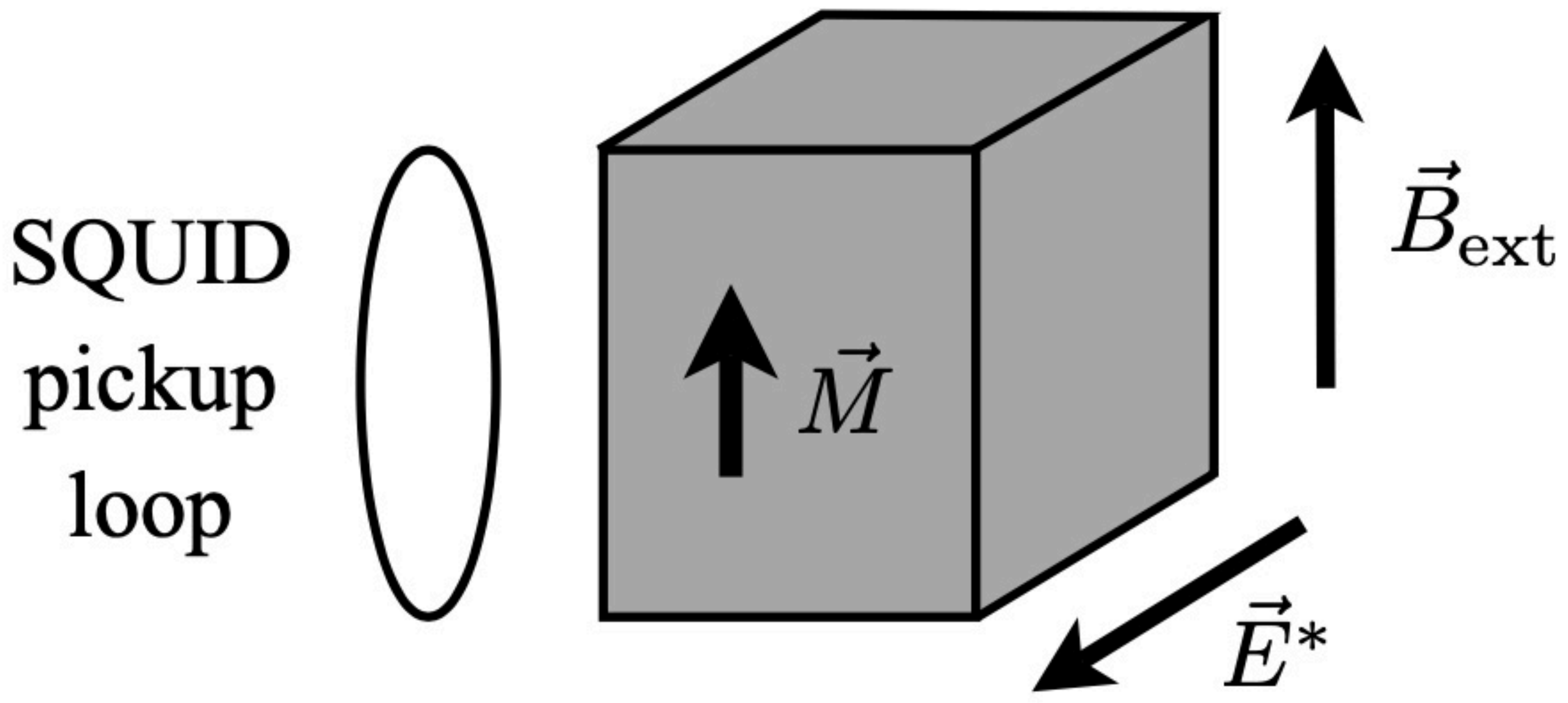
$$H \simeq \frac{a}{m_N f_a} \vec{\sigma} \cdot \vec{E}$$



$$M(t) \approx np\mu E^* \epsilon_S d_n \frac{\sin\left[\left(\frac{2\mu B_{\rm ext} - m_a c^2}{\hbar}\right) t\right]}{\frac{2\mu B_{\rm ext} - m_a c^2}{\hbar}} \sin\left(2\mu B_{\rm ext} t\right)$$

CASPER

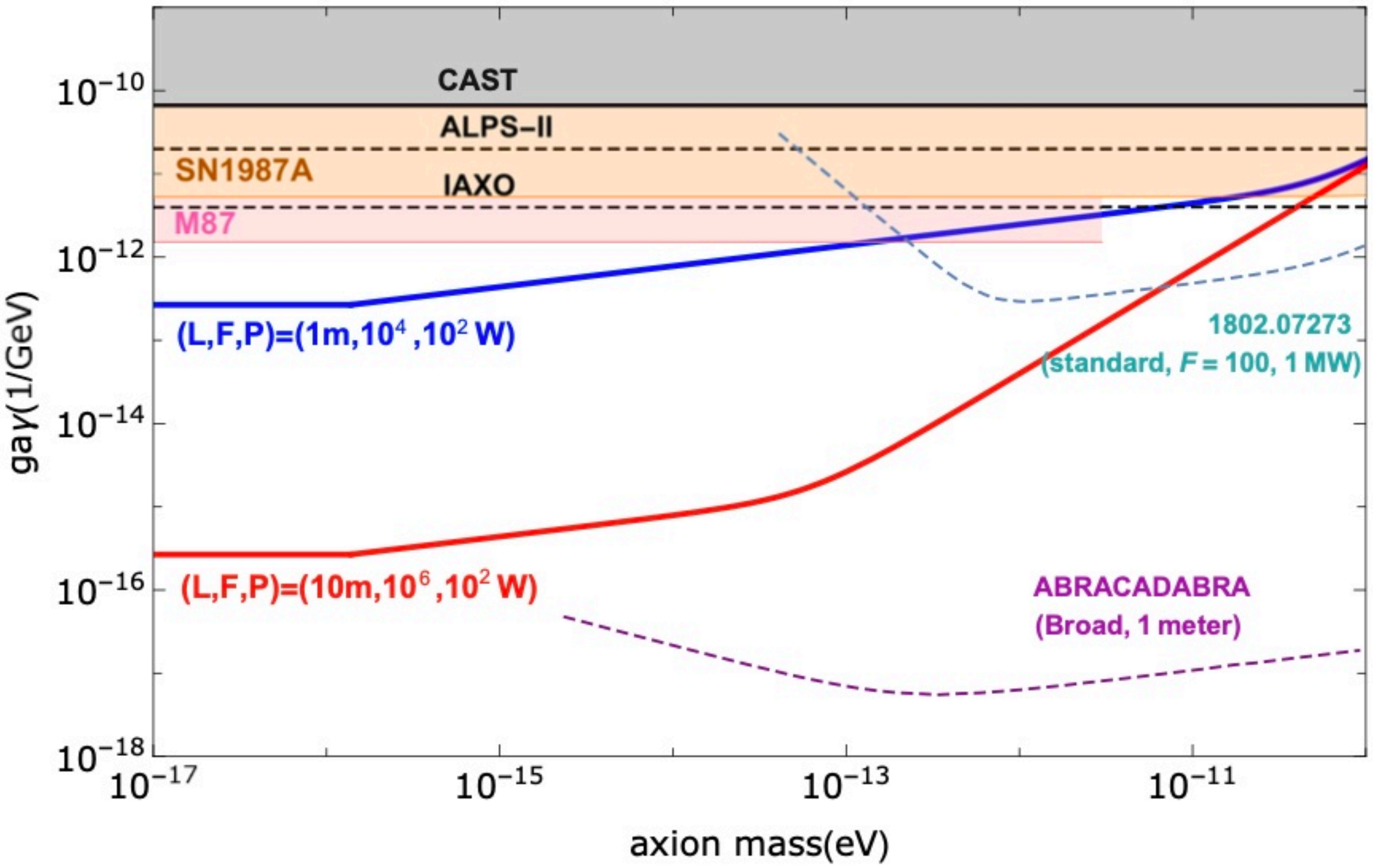
$$H \simeq \frac{a}{m_N f_a} \vec{\sigma} \cdot \vec{E}$$



	n	E^*	p	T_2	Max. B_{ext}
Phase 1	$10^{22}\ \frac{1}{\text{cm}^3}$	$3 \times 10^8\ \frac{\text{V}}{\text{cm}}$	10^{-3}	1 ms	10 T
Phase 2	$10^{22}\ \frac{1}{\text{cm}^3}$	$3 \times 10^8\ \frac{\text{V}}{\text{cm}}$	1	1 s	20 T

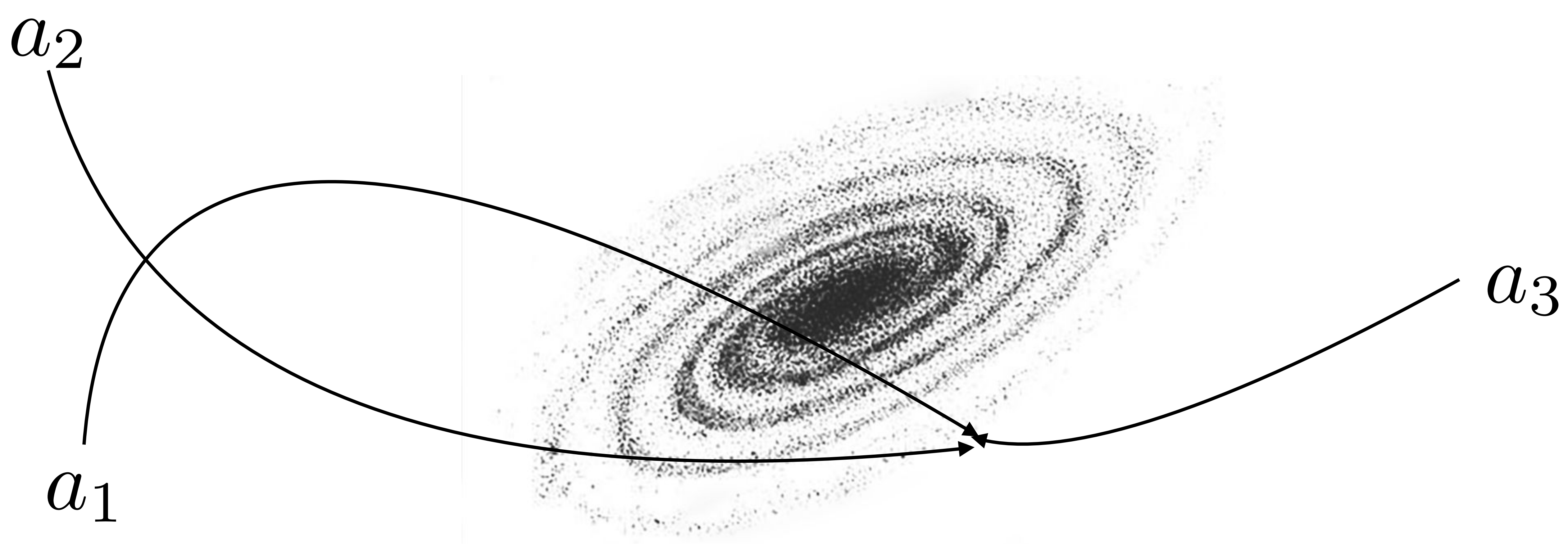
DANCE

Obata, Fujita, Michimura, '18



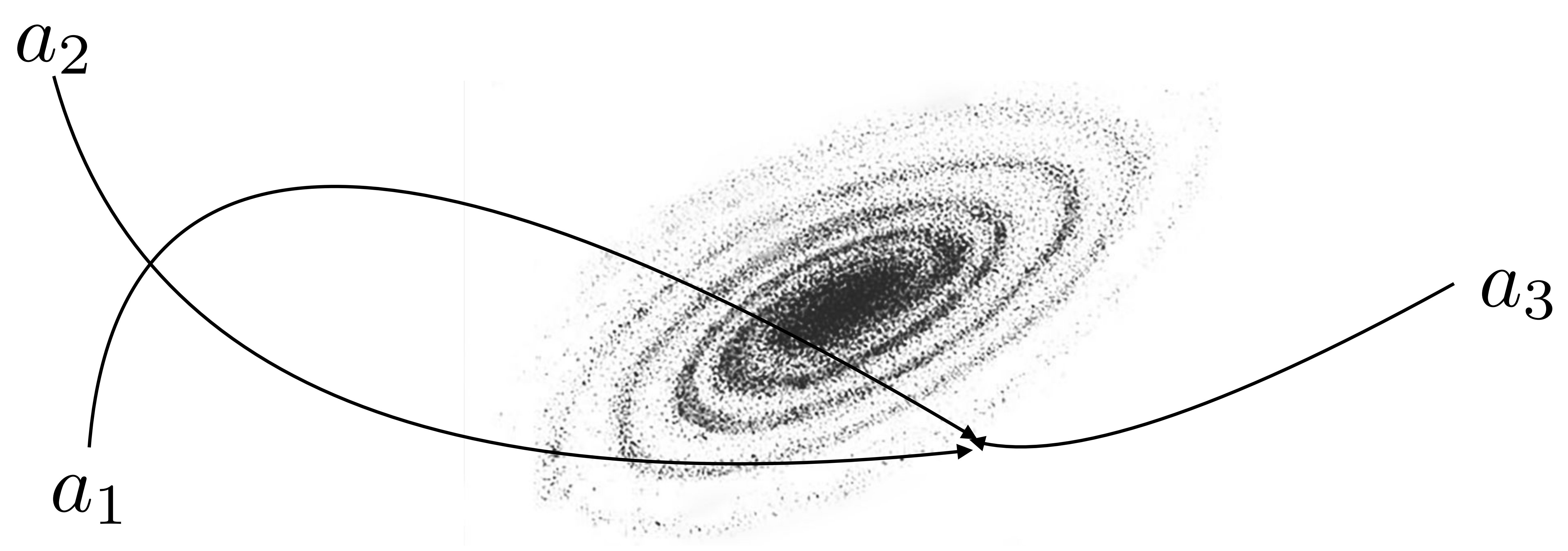
ALP DARK MATTER IN THE LAB

In each experimental bin we are summing over a multitude of plane waves with different phases



ALP DARK MATTER IN THE LAB

In each experimental bin we are summing over a multitude of plane waves with different phases



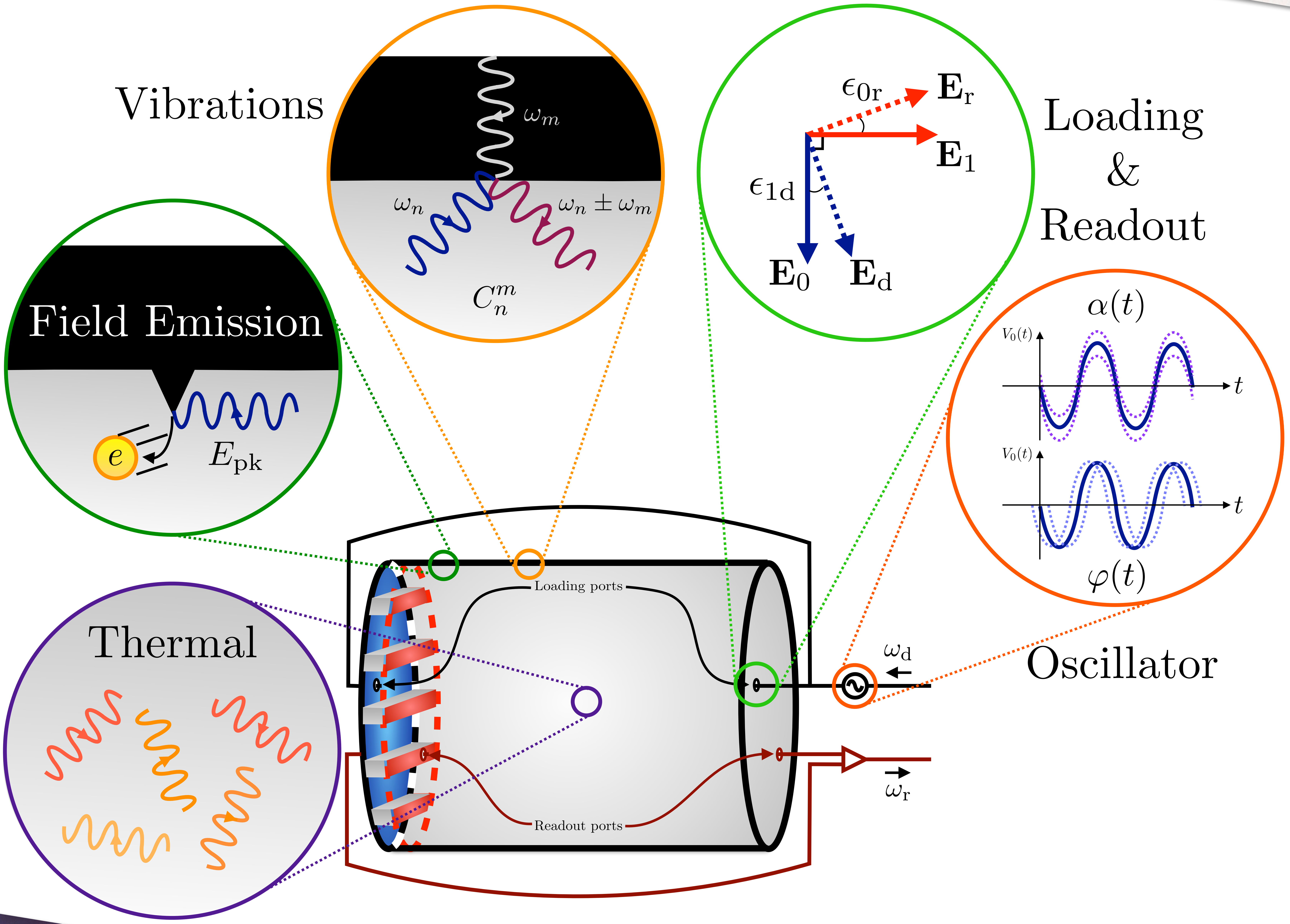
$$a(t) = a_0 \left[\cos\left(m_a\left(1+\frac{v_1^2}{2}\right)t+\phi_1\right)+\cos\left(m_a\left(1+\frac{v_2^2}{2}\right)t+\phi_2\right)+\ldots\right]$$

$$\simeq a_0\cos(m_a t+\phi)\left[\cos(\delta\omega_a t+\phi')+\ldots\right]$$

$$\delta\omega_a \simeq \frac{1}{m_a\langle v_{\rm DM}^2\rangle}\simeq\frac{10^6}{m_a}$$

Effectively: very **slow modulation of** an approximately **monochromatic field**

ULTRALIGHT AXION-LIKE DARK MATTER



Vibrations
ω_m
ω_n
$\omega_n \pm \omega_m$
C_n^m
Field Emission
e
E_{pk}
ϵ_{0r}
$\mathbf{E}_r$
$\mathbf{E}_1$
ϵ_{1d}
$\mathbf{E}_0$
$\mathbf{E}_d$
Loading & Readout
$\alpha(t)$
$V_0(t)$
t
$\varphi(t)$
Thermal
Loading ports
Readout ports
ω_d
ω_r
Oscillator

LEAKAGE NOISE

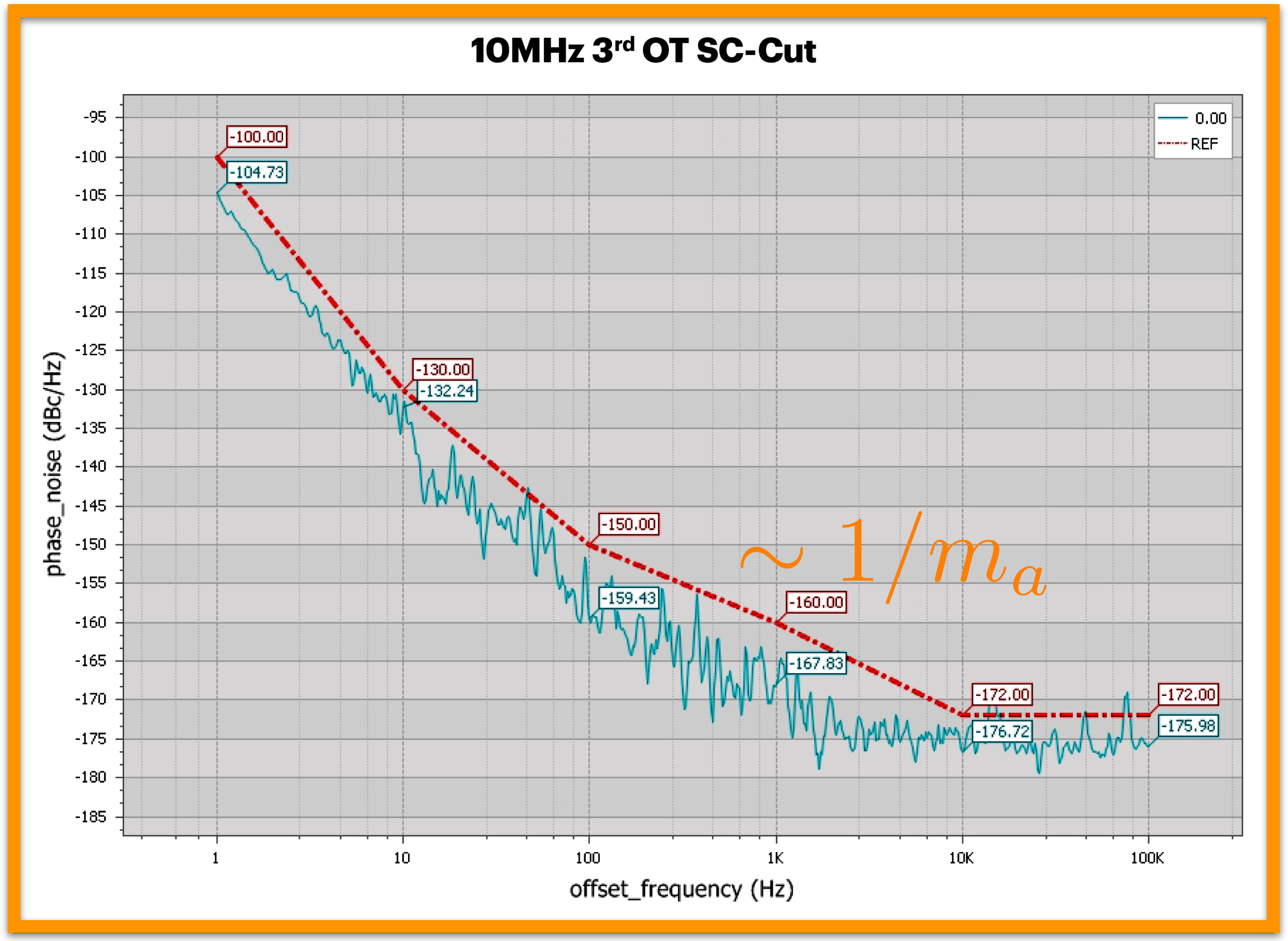
$$S_{\rm phase}(\omega) \simeq \frac{1}{2}\,\epsilon_{\rm 1d}^2\, S_\phi(\omega-\omega_0)\, \boxed{\frac{(\omega\,\omega_1/Q_1)^2}{(\omega^2-\omega_1^2)^2+(\omega\,\omega_1/Q_1)^2}}\, \frac{\omega_0 Q_1}{\omega_0 Q_0}\, P_{\rm in}$$

Cavity Response

LEAKAGE NOISE

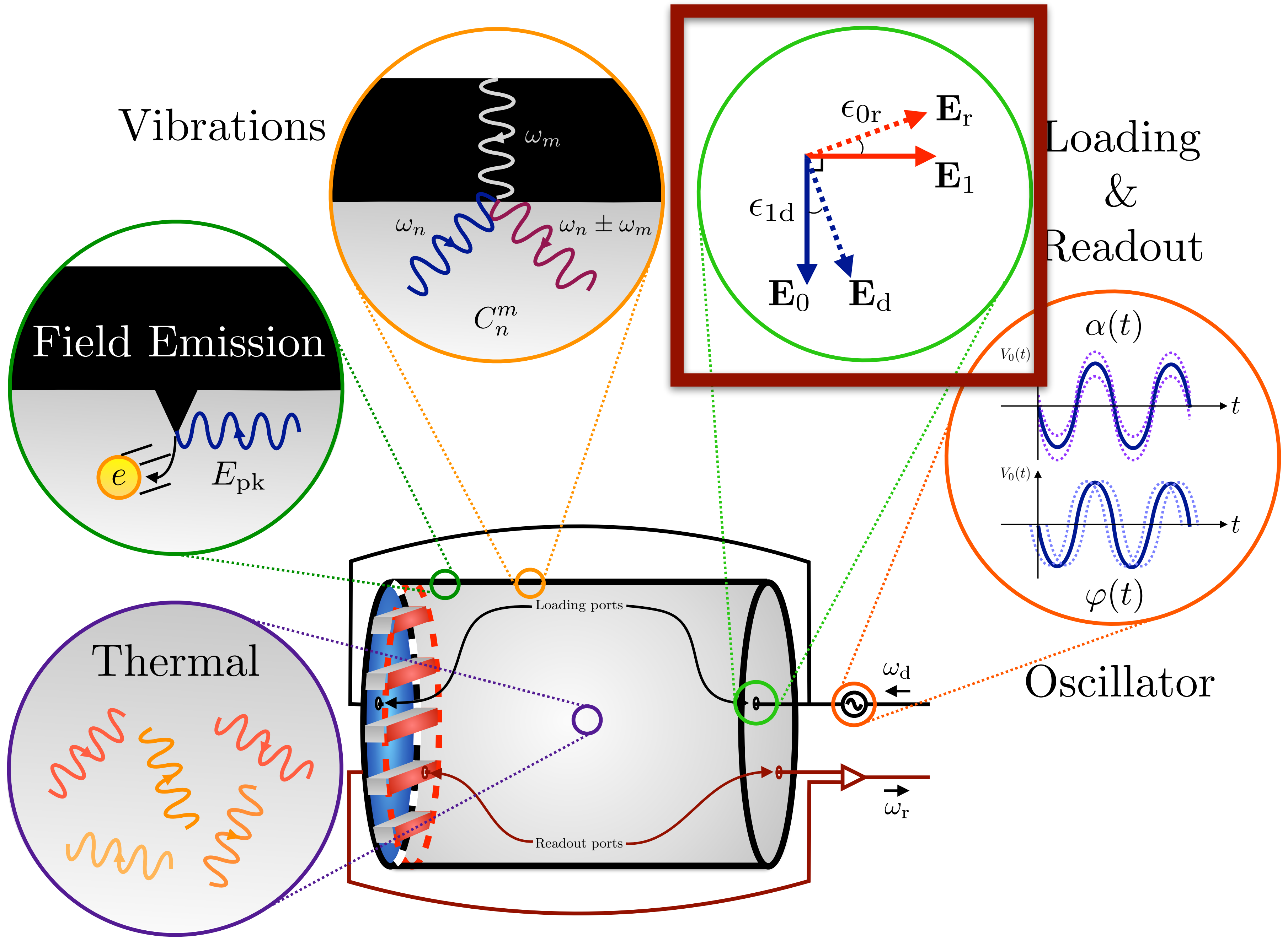
$$S_{\text{phase}}(\omega) \simeq \frac{1}{2}\,\epsilon_{1\text{d}}^2\, \boxed{S_\phi(\omega-\omega_0)}\, \frac{(\omega\,\omega_1/Q_1)^2}{(\omega^2-\omega_1^2)^2+(\omega\,\omega_1/Q_1)^2}\, \frac{\omega_0 Q_1}{\omega_0 Q_0}\, P_{\text{in}}$$

$\sim 1/m_a$



LEAKAGE NOISE

$$S_{\text{phase}}(\omega) \simeq \frac{1}{2}\epsilon_{1\text{d}}^2 S_\phi(\omega-\omega_0)\frac{(\omega\,\omega_1/Q_1)^2}{(\omega^2-\omega_1^2)^2+(\omega\,\omega_1/Q_1)^2}\frac{\omega_0 Q_1}{\omega_0 Q_0}P_{\text{in}}$$



From MAGO
and other similar cavities