



# Multiple Stellar Populations of Globular Clusters from Homogeneous Ca–CN–CH–NH Photometry. VI. M3 (NGC 5272) Is Not a Prototypical Normal Globular Cluster<sup>\*†</sup>

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## Abstract

We present Ca–CN–CH–NH photometry for the well-known globular cluster (GC) M3 (NGC 5272). We show new evidence for two M3 populations with distinctly different carbon and nitrogen abundances, seen in a sharp division between CN-weak and CN-strong red-giant branches (RGBs) in M3. The CN-strong population shows a C–N anticorrelation that is a natural consequence of the CN cycle, while the CN-weak population shows at most a very weak C–N anticorrelation. Additionally, the CN-weak population exhibits an elongated spatial distribution that is likely linked to its fast rotation. Our derived metallicities reveal bimodal distributions in both populations, with  $\langle \text{[Fe/H]} \rangle \approx -1.60$  and  $-1.45$ , which appear to be responsible for the discrete double RGB bumps in the CN-weak population and the large  $W_{\text{F275W} - \text{F814W}}^{1G}$  range. From this discovery, we propose that M3 consists of two GCs, namely C1 (23%,  $\langle \text{[Fe/H]} \rangle \approx -1.60$ ) and C2 (77%,  $\langle \text{[Fe/H]} \rangle \approx -1.45$ ), each of which has its own C–N anticorrelation and structural and kinematical properties, which are strong indications of independent systems in M3. The fractions of the CN-weak population for both C1 and C2 are high compared to Galactic GCs but they are in good agreement with GCs in the Magellanic Clouds. We suggest that M3 is a merger remnant of two GCs, most likely in a dwarf galaxy environment, and accreted to our Galaxy later in time. This is consistent with recent proposals of an ex situ origin for M3.

*Unified Astronomy Thesaurus concepts:* [Hertzsprung Russell diagram \(725\)](#); [Globular star clusters \(656\)](#); [Stellar abundances \(1577\)](#); [Stellar evolution \(1599\)](#)

## 1. Introduction

Galactic globular clusters (GCs) have been extensively studied to understand the formation and evolution of our Galaxy. For several decades, the existence of not only the in situ but also the accreted or ex situ GCs has been recognized (Searle & Zinn 1978; Zinn 1985, 1993; Lee & Carney 1999b) and interpreted in terms of cold dark matter cosmology that predicts a hierarchical structure formation in the universe. The Gaia satellite has radically improved the situation with astrometry and proper motion data for more than a billion stars (Gaia Collaboration et al. 2018). The previously challenging tasks of identifying merger events in our Galaxy can now be realized, and several remnants of ancient merger events have been proposed (e.g., see Helmi et al. 2018; Myeong et al. 2019).

It has been thought that younger halo GCs, including M3, were likely accreted by our Galaxy in the past to form components of the Galactic halo seen today (e.g., Zinn 1993; Lee & Carney 1999b). For example, Kruijssen et al. (2019) suggested that M3 is likely an ex situ GC and a member of a hypothetical galaxy, the so-called Kraken, that accreted into our Galaxy. Koppelman et al. (2019) argued that M3 belongs to

the Helmi Stream that was likely accreted about 5–8 Gyr ago, in which seven GCs exhibit a tight age–metallicity relation.

M3 has long been considered as a prototypical “normal” GC. But M3 lies relatively far from the Galactic center, and it has some interesting structural aspects. It is one of the most massive GCs in our Galaxy and, most interestingly, it is the most RR Lyrae (RRL)-rich GC in our Galaxy, possessing more than 240 RRLs, which makes M3 exceptional (Clement et al. 2001, the 2019 March version<sup>3</sup>). A second RRL-rich GC is  $\omega$  Cen, which is most likely the remnant of the core of the dwarf galaxy that accreted into our Galaxy and about 3.5 times more massive than M3 (Kruijssen et al. 2019).  $\omega$  Cen has about 200 RRLs, but it contains a significant fraction of metal-poor stars that do not pass through the instability strip during their core helium phase. A third RRL-rich GC is M5, which has a similar mass to M3 but has about 130 RRLs.

The first-generation (FG) stars of M3 has an abnormally large range of the Hubble Space Telescope (HST) pseudo-color index,  $\Delta_{\text{CF275W}, \text{F814W}}$ . This has led to investigations of the amount of helium spread in the M3 FG (Lardo et al. 2018; Tailo et al. 2019), which concluded that it cannot explain the HST  $\Delta_{\text{CF275W}, \text{F814W}}$  color index spread within our current understanding of the GC formation and chemical evolution scenarios. In our previous study of M3 CN band strengths (Lee 2019a), we found that the M3 CN-w population has an extended and tilted red-giant branch bump (RGBB), which later we will link to the large range of the  $\Delta_{\text{CF275W}, \text{F814W}}$  index.

Light-elemental abundance variations in GCs were first identified about 50 yr ago, and observational evidence for their

<sup>\*</sup> Based on observations made with the Kitt Peak National Observatory (KPNO) 0.9 m telescope, which is operated by WIYN Inc. on behalf of a consortium of partner universities and organizations.

<sup>†</sup> This work has made use of data from the European Space Agency (ESA) mission Gaia (<https://www.cosmos.esa.int/gaia>), processed by the Gaia Data Processing and Analysis Consortium (DPAC; <https://www.cosmos.esa.int/web/gaia/dpac/consortium>). Funding for the DPAC has been provided by national institutions, in particular the institutions participating in the Gaia Multilateral Agreement.

<sup>3</sup> See C. M. Clement, Catalog of GC variable stars, available at <http://www.astro.utoronto.ca/~cclement/cat/listngc.html>.

ubiquitous nature has steadily increased (e.g., Osborn 1971; Cohen 1978; Norris et al. 1981; Sneden et al. 1992). Explanations have centered on the idea of multiple stellar formation episodes, where the later generation of stars formed out of interstellar media polluted by the previous generation of stars (e.g., D’Ercole et al. 2008). During the past decade, a great deal of observational evidence of multiple populations (MPs) in GCs in the Milky Way has emerged (e.g., Carretta et al. 2009; Lee et al. 2009b; Milone et al. 2017; Gratton et al. 2019; Marino et al. 2019). In spite of many efforts, understanding the formation of GCs with MPs is not yet solved. So far, none of the proposed models can explain the observational constraints of GC MPs satisfactorily (e.g., Bastian & Lardo 2018). Furthermore, MPs exist in GCs of external low-mass Local Group galaxies, such as the Magellanic Clouds (MC; e.g., Mucciarelli et al. 2009; Milone et al. 2020) and the Fornax dwarf galaxy (e.g., Letarte et al. 2006), which poses another formidable problem of environmental effects on the formation of GCs with MPs. Revolutionary spectroscopic and photometric studies of MPs in extragalactic GCs will come with the advent of 30 m class telescope within the next decade. Until then, the best way to understand environmental effects on GC formation is to investigate the accreted or *ex situ* GCs that are located near us.

In this paper, we investigate M3 using our own photometric system optimized for measuring carbon and nitrogen abundances, which are key elements for GC MP studies. In spite of the importance of the carbon and nitrogen abundance in the evolution of metal-poor low-mass stars, their abundances in GCs are often poorly known due to the lack of measurable atomic absorption lines in the visual or infrared wavelength regime. Instead, their elemental abundances can be measured via diatomic molecules, such as NH, CN, CH, and CO.

Studies of C and N abundances from visual or infrared high-resolution spectroscopy have been restricted to very bright red-giant branch (RGB) or asymptotic giant branch (AGB) stars, where the surface C and N abundances can be significantly altered from their primordial values due to the onset of CN cycle accompanied by a noncanonical thermohaline mixing (Charbonnel & Zahn 2007). Low-resolution spectroscopy has been frequently applied to faint GC stars, but traditional spectroscopy cannot be used in very crowded regions, such as the central part of GCs, due to stellar blending. Even for the isolated stars, low-resolution spectroscopy is also vulnerable to the selection of continuum sidebands in the wavelength region of numerous strong absorption lines, which can compromise the measured spectrum indices (e.g., see Lee 2019a, 2019c). Our approach mitigates both of these problems through careful feature and passband choices for NH, CN, and CH photometric indices to provide reliable carbon and nitrogen abundances, and avoidance of stellar targets with potential contamination by other cluster members.

The work presented here will provide observational evidence that M3 is most likely a merger remnant of two GCs in a dwarf galaxy environment, where the relative velocity of the two GCs is smaller than their velocity dispersion (e.g., Gavagnin et al. 2016). The outline of this paper is as follows. We describe our new filter system, JWL34, which is designed to measure the absorption strength of the NH bands at  $\lambda 3360$  in Section 2. We present our observations and data reductions in Section 3. The definitions of our color indices, including the new  $nh_{JWL}$ , will be given in Section 4 along with a discussion of our new

strategy of populational tagging of M3 RGB stars. In Section 5, we show comparisons of our photometric indices with widely used color indices from HST photometry (Milone et al. 2015). We will discuss the metallicity dependency on the  $\Delta_{CF275W,F814W}$  index, which is the key to understand the large  $W_{F275W-F814W}^{1G}$  range of the M3 FG stars that has been pointed out in previous studies (Lardo et al. 2018; Tailo et al. 2019). In Section 6, we discuss synthetic model grids generated to interpret our observed indices into  $[Fe/H]$ ,<sup>4</sup>  $[C/Fe]$ , and  $[N/Fe]$  abundances. In particular, we will discuss the bimodal metallicity distributions for both CN-w and CN-s, which are important clues that require a new formation scenario for M3. We also discuss the separate C–N correlations between the two populations. Populational tagging for red horizontal-branch (RHB) stars will be presented in Section 7, where we will show that the populational number ratio of RHB stars is significantly different but can be naturally understood from the evolution of metal-poor low-mass stars. The structural and kinematical differences between MPs will be discussed in Section 8. The discussion on the discrete double RGBBs of the CN-w will be given in Section 9, where we will argue that the bimodal metallicity distribution is essential to explain the double RGBBs in the CN-w and the large extent of the  $W_{F275W-F814W}^{1G}$  of the M3 FG obtained with HST photometry. In Section 10, we argue that M3 is a merger remnant of two distinct GCs, having different elemental abundances, and structural and kinematical properties. Finally, a summary of our work will be provided in Section 11.

## 2. JWL34: A New NH Filter System

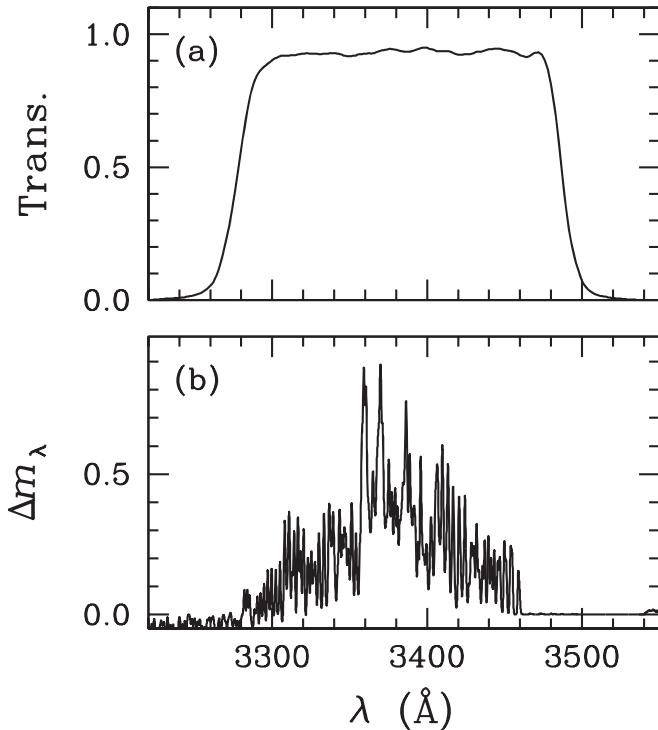
Our  $cn_{JWL}$  and  $ch_{JWL}$  indices (discussed in Section 4) can provide significant information on MPs in GCs (e.g., Lee 2017, 2018, 2019a, 2019c, 2020). Our  $cn_{JWL}$  index is satisfactory in most cases, but there are three issues that complicate attempts to transform  $cn_{JWL}$  indices into unambiguous nitrogen abundances.

First, the formation of CN molecules depends on both the carbon and nitrogen abundances. In low-luminosity, less evolved GC stars with abundance ratios  $n(N)/n(C) < 1$ , CN band strengths are mainly dependent on nitrogen abundances (e.g., Suntzeff 1981; Briley & Smith 1993). But this abundance ratio condition is not always true. For second-generation (SG) stars in GCs with abundance ratios  $n(N)/n(C) > 1$ , CN band strengths no longer scale monotonically with nitrogen abundance and can even decrease with depletion in the surface carbon abundances (e.g., Smith & Bell 1986).

Second, CN is a double-metal diatomic molecule, and its band absorption strengths rapidly decline with decreasing metallicity (e.g., Sneden 1974; Langer et al. 1992). On the other hand, absorption strengths of single-metal diatomic molecules (e.g., CH and NH) weaken only slowly with decreasing metallicity. They are still sensitive to variations in the carbon and nitrogen abundances in very metal-poor stars.

Third,  $[C/Fe]_{CN}$  and  $[N/Fe]_{CN}$  abundances derived from CN bands at  $\lambda 3883$  with prior  $[N/Fe]$  and  $[C/Fe]$  information supplied from  $nh_{JWL}$  and  $ch_{JWL}$  indices depend somewhat on  $[O/Fe]$  abundances and  $^{12}C/^{13}C$  ratios. In particular, the oxygen abundances of individual stars within a given

<sup>4</sup> We adopt the standard spectroscopic notation (Wallerstein & Helfer 1959) that for elements A and B,  $[A/B] \equiv \log_{10} (N_A/N_B)_* - \log_{10} (N_A/N_B)_\odot$ . We equate metallicity with the stellar  $[Fe/H]$  value.



**Figure 1.** (a) The filter transmission function of our new JWL34 filter. (b) Differences in monochromatic magnitude between the CN-w and CN-s using synthetic spectra for an arbitrary intermediate-metallicity ( $[\text{Fe}/\text{H}] = -1.5$ ) RGB stars at the level of the RGBB.

population have a substantial spread (e.g., see Carretta et al. 2009; Marino et al. 2019), which can result in non-negligible differences in  $[\text{C}/\text{Fe}]_{\text{CN}}$  and  $[\text{N}/\text{Fe}]_{\text{CN}}$  in the absence of  $[\text{O}/\text{Fe}]$  abundances of individual stars. As shown in Appendix A, at a metallicity of  $[\text{Fe}/\text{H}] = -1.5$  dex,  $\Delta[\text{O}/\text{Fe}] = +0.1$  dex can result in the overestimation of  $[\text{N}/\text{Fe}]_{\text{CN}}$  by about 0.08 dex (see Figure 28). At the bright RGB (bRGB),<sup>5</sup> the  $[\text{N}/\text{Fe}]_{\text{CN}}$  is also vulnerable to uncertain  $^{12}\text{C}/^{13}\text{C}$  ratios, which are often undetermined for GC RGB stars. The difference in  $[\text{N}/\text{Fe}]_{\text{CN}}$  with  $^{12}\text{C}/^{13}\text{C} = 5$  and 90 ( $\approx$  a solar carbon isotope ratio) can be as large as 0.5 dex (see Figure 29).

In an attempt to directly measure the NH band at  $\lambda 3360$ , we have developed a new filter, JWL34. In panel (a) of Figure 1, we show the transmission function of our JWL34 filter, which has a pivot wavelength of  $\approx \lambda 337$  nm and an FWHM of  $\Delta\lambda \approx 20.5$  nm. In panel (b) of this figure, we show the sensitivity of JWL34 to N abundance changes by plotting the difference in monochromatic magnitudes between two intermediate-metallicity ( $[\text{Fe}/\text{H}] - 1.50$ ) stars: a CN-w star (with assumed abundances  $[\text{C}/\text{Fe}] = -0.10$ ,  $[\text{N}/\text{Fe}] = 0.10$ ) and a CN-s star ( $[\text{C}/\text{Fe}] = -0.30$ ,  $[\text{N}/\text{Fe}] = 0.80$ ). The monochromatic magnitudes are those that were obtained with synthetic spectrum computations, as will be described in detail in Section 6.3.

### 3. Photometric Data

In 2017 and 2018, we obtained photometric data for M3 in 20 nights in 5 separate runs using the Half Degree Imager (HDI), which is equipped with an e2V  $4k \times 4k$  CCD chip,

<sup>5</sup> The bRGB stars and the faint RGB (fRGB) stars indicate stars brighter and fainter than the RGBB, respectively.

mounted on the KPNO 0.9 m telescope. The HDI provides a field of view (FOV) of  $30' \times 30'$ . The total integration times of Strömgren  $y$ ,  $b$ ,  $\text{Ca}_{\text{JWL}}$ , JWL39, and JWL43 for the M3 science field are given in Table 1 (see also Lee 2019a). Additional photometric data for M3 using our new JWL34 filter and Strömgren  $y$  and  $b$  filters were collected over six nights from 2019 June 27 to July 5. Integration times for JWL34, Strömgren  $y$ , and  $b$  filters were 19,400, 1220, and 2200 s. Due to the electronic problem with the HDI in 2019, we used the S2KB CCD cam, which is equipped with a  $2k \times 2k$  CCD chip and provides an FOV of  $21' \times 21'$ .

Our integration times for individual filters are long enough to perform accurate photometry of RGB stars in M3. Typical photometric measurement errors at the level of the RGBB in individual color indices are less than 0.01 mag, so the broad or bimodal RGB sequences in some color indices are real features in M3, not artifacts arising from measurement errors.

Limitations in the spatial angular resolving power of small-aperture ground-based telescopes, such as the KPNO 0.9 m telescope, lead to incomplete detection and large photometric measurement uncertainties in central parts of GCs. Deriving accurate populational number ratios requires detection of statistically robust samples with reliable measurements in the central part of GCs showing a strong radial populational gradient such as M3, but this is not an important point in our current study. Most of the analyses presented in this paper will be based on stars with very accurate measurements located at  $r \geq 1'$ , and we did not apply the method that we developed for using prior positional information from HST observations (e.g., see Appendix B of Lee 2017). As we will discuss below, we made use of the proper motion study from the second Gaia data release (Gaia Collaboration et al. 2018) in which the source detection incompleteness becomes large in the central part, which also hinders our populational tagging of the central part of the cluster.

The interstellar reddening of M3 is very small,  $E(B - V) = 0.01$  (Harris 1996, the 2010 edition).<sup>6</sup> Therefore, any differential reddening across the science field will be undetectably small and will not affect our results (see also Section 9.2).

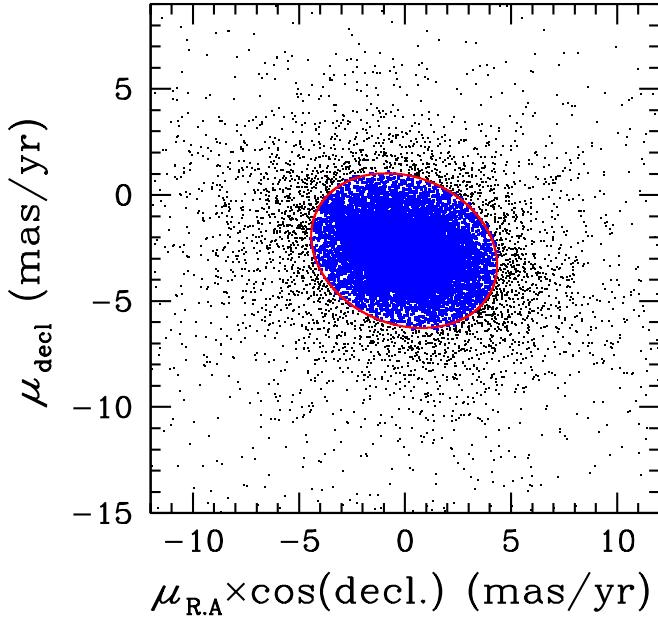
Reductions of our raw photometric data have been discussed in detail by Lee (2015) and will not be repeated here. The photometry of M3 and standard stars were analyzed using the DAOPHOTII, DAOGROW, ALLSTAR and ALLFRAME, and COLLECT-CCDAVE-NEWTRIAL packages (Stetson 1987, 1994; Lee & Carney 1999a). We derived astrometric solutions for individual stars using the data extracted from the Naval Observatory Merged Astrometric Dataset (NOMAD, Zacharias et al. 2004) and the IRAF IMCOORDS package.

We made use of the proper motions from the second Gaia data release to select the cluster's membership stars (Gaia Collaboration et al. 2018), following the method similar to those used in our previous studies (see, e.g., Lee 2020, and references therein). We derived the mean values of proper motions of M3 with iterative sigma-clipping calculations, finding that, in units of  $\text{mas yr}^{-1}$ ,  $(\mu_{\text{R.A.}} \times \cos \delta, \mu_{\text{decl.}}) = (-0.039, -2.622)$ , with standard deviations along the major axis of the ellipse of  $1.509 \text{ mas yr}^{-1}$  and along the minor axis of  $1.155 \text{ mas yr}^{-1}$ . We considered that stars within  $3\sigma$  from the mean values to be M3 member stars as shown in Figure 2. Then, we selected our target

<sup>6</sup> Available at <http://physwww.mcmaster.ca/~harris/mwgc.dat>.

**Table 1**  
Integration Times for M3 (in Seconds)

	<i>y</i>	<i>b</i>	$C_{\text{JWL}}$	JWL39	JWL43	JWL34	<i>V</i>	<i>B</i>	<i>I</i>	<i>U</i>
NGC 5272(M3)	8910	19,520	52,000	23,500	24,000	19,400	1380	3070	1220	11,350



**Figure 2.** Gaia DR2 proper motions of our FOV. The red ellipse indicates the boundary ( $3\sigma$ ) of the M3 member stars shown with blue dots.

RGB stars with  $-2 \text{ mag} \leq V - V_{\text{HB}} \leq 2 \text{ mag}$  from our multi-color photometry.

The Galactic latitude of the cluster is high,  $79^\circ$ . Therefore, the contamination by the off-cluster field stars will not be severe in our photometry. For example, we estimate that the total number of field stars with the magnitude range of our interest ( $-2 \text{ mag} \leq V - V_{\text{HB}} \leq 2 \text{ mag}$ ) toward our M3 science FOV is expected to be  $\approx 140$  (Ratnatunga & Bahcall 1985), which corresponds to only  $6.4\% \pm 0.6\%$  of our total number of proper motion membership stars in the magnitude range of our interest. In addition to proper motion study, as Lee (2015, 2020) demonstrated, our  $C_{\text{JWL}}$  photometry provides a powerful means to distinguish between GC RGB stars with low metallicities and Galactic disk stars with high metallicities. Therefore, we believe that the contribution from off-cluster field stars is too small to affect our results presented here.

#### 4. Populational Tagging: New Strategy

Throughout this work, we will use our own photometric indices, defined as

$$hk_{\text{JWL}} = (C_{\text{JWL}} - b) - (b - y), \quad (1)$$

$$nh_{\text{JWL}} = (JWL34 - b) - (b - y), \quad (2)$$

$$cn_{\text{JWL}} = JWL39 - C_{\text{JWL}}, \quad (3)$$

$$ch_{\text{JWL}} = (JWL43 - b) - (b - y). \quad (4)$$

These relationships have functional forms analogous to traditional color indices and will be treated as such throughout this paper.

Assuming constant [Ca/Fe] ratios (e.g., see Carney 1996; Marino et al. 2019), the  $hk_{\text{JWL}}$  index is a good photometric

measure of metallicity, with a weak CH band contamination (Anthony-Twarog et al. 1991; Lee et al. 2009a, 2009b; Lee 2015, 2019c). In Appendix C and Figure 30, we show the  $hk_{\text{JWL}}$  dependency on metallicity (see also Supplementary Information of Lee et al. 2009a). Also in Appendix D and Figure 31, we show the  $hk_{\text{JWL}}$  dependency on carbon abundance, which is almost nil for M3.

As we discussed in our previous works,  $cn_{\text{JWL}}$  and  $ch_{\text{JWL}}$  are excellent photometric measures of the CN band at  $\lambda 3883$  and the CH *G* band at  $\lambda 4250$ , respectively, for cool stars (see Lee 2020 and references therein).

In our study of M3, the  $cn_{\text{JWL}}$  index is still powerful in classifying MPs, as we will show below. For cool GC RGB stars, the emergent stellar surface flux near  $\lambda 3883$  is much greater than that near  $\lambda 3360$ . At the same time, the degree of interstellar and atmospheric extinctions for JWL39 is much smaller than those for JWL34 (e.g., see Figure 2 of Lee 2017). Therefore, more accurate photometry can be attained with  $cn_{\text{JWL}}$  than with  $nh_{\text{JWL}}$ . In addition, populational tagging from  $cn_{\text{JWL}}$  can be more readily accomplished than by using  $ch_{\text{JWL}}$  or  $nh_{\text{JWL}}$  indices, as long as the C–N anticorrelation holds. For example, the CH distributions of GC RGB stars may not show a clear bimodality (e.g., see Figure 6 of Norris & Smith 1984), but the CN bimodality in GCs has been known for decades and frequently studied for GC stars (e.g., Osborn 1971).

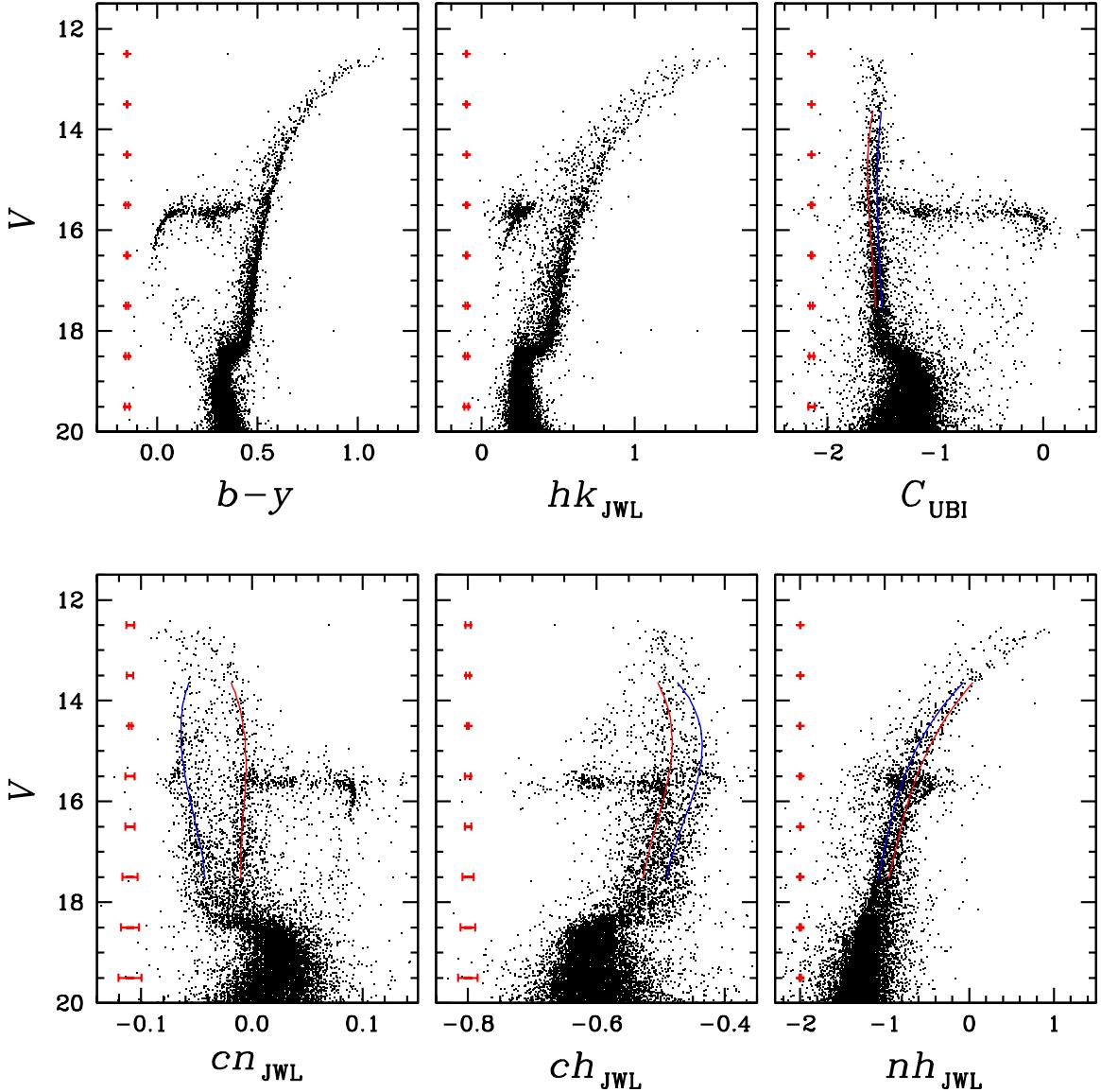
In Figure 3, we show color–magnitude diagrams (CMDs) of M3 member stars based on the proper motion study of Gaia DR2 (Gaia Collaboration et al. 2018). We also show measurement uncertainties for each color index, suggesting that the discrete double RGB sequences or the spread in color indices in the magnitude range of our interest ( $-2 \text{ mag} \leq V - V_{\text{HB}} \leq 2 \text{ mag}$ ) are real features; they are much larger than photometric errors. Our  $cn_{\text{JWL}}$  CMD (bottom-left panel) shows conspicuous double RGB sequences in M3; these have been known for decades (e.g., Suntzeff 1981; Briley & Smith 1993). Our  $ch_{\text{JWL}}$ ,  $nh_{\text{JWL}}$ , and  $C_{\text{UBI}}$  CMDs also show weak bimodalities or a large spread in their color indices, suggestive of heterogeneous CNO abundances among RGB stars.

In order to remove the luminosity effect on individual color indices, we use parallelized color indices (also see Lee 2019a, 2019c). The RGB sequences in the individual color indices were parallelized using the following relation,

$$\|CI(x) \equiv \frac{CI(x) - CI_{\text{red}}}{CI_{\text{red}} - CI_{\text{blue}}}, \quad (5)$$

where  $CI(x)$  is the color index of the individual stars and  $CI_{\text{red}}$  and  $CI_{\text{blue}}$  are the color indices for the fiducials of the red and blue sequences of the individual color indices (see also Milone et al. 2017).

To perform a populational tagging of RGB stars in the magnitude range of  $-2 \text{ mag} \leq V - V_{\text{HB}} \leq 2 \text{ mag}$ , we employed an expectation maximization (EM) algorithm for a multiple-component Gaussian mixture distribution model followed by principal component analysis using the programming language R (R Core Team 2017). To make the best use of our measurements, we used  $\|cn_{\text{JWL}}$ ,  $\|ch_{\text{JWL}}$ , and  $\|nh_{\text{JWL}}$  to calculate principal



**Figure 3.** CMDs with various color axes of M3 member stars based on the proper motion study of Gaia DR2. The red error bars represent the mean photometric measurement uncertainties at given magnitude bins. The fiducial sequences to calculate the parallelized color indices are also shown.

components as shown in Figure 4. We obtained the populational number ratio of  $n(\text{CN-w}):n(\text{CN-s}) = 48:52 (\pm 3)$  by employing the EM algorithm along the PC1, which is exactly the same value that we obtained from the EM algorithm for the multiple-component Gaussian mixture distribution model along the  $\|cn_{\text{JWL}}$  distribution (see Table 5 of Lee 2019a). Note that the eigenvector of the  $\|cn_{\text{JWL}}$  is parallel to the PC1-axis that maximizes the variance of the projected data, confirming our previous argument that  $cn_{\text{JWL}}$  is still important in populational tagging.

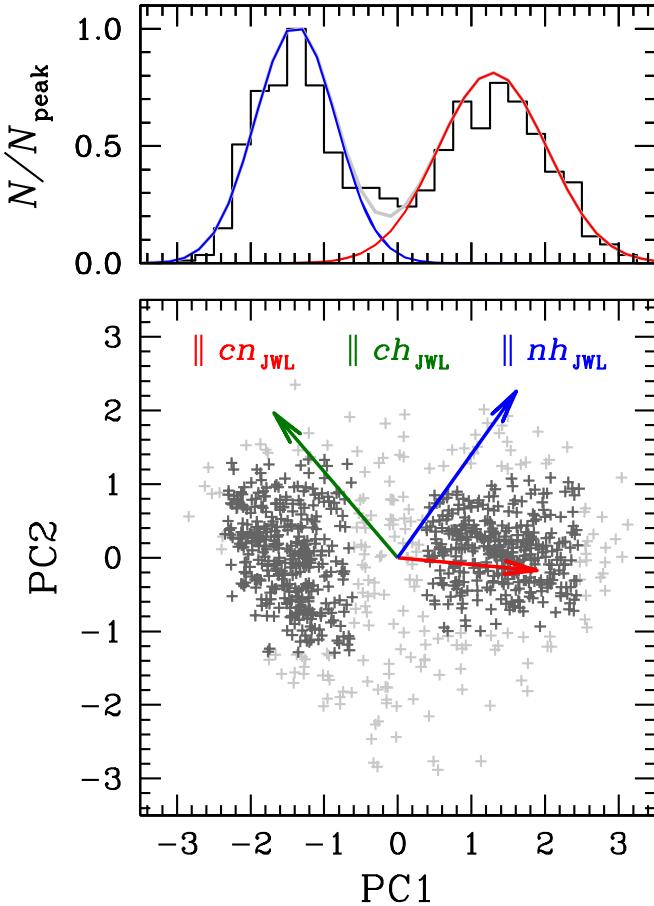
In order to circumvent potential confusion from population intruders and outliers in our analyses performed in Sections 4–6, we used the well-behaved RGB stars located within  $2\sigma$  from the centers of each population on the PC1–PC2 plane as shown by the dark gray color in Figure 4.

As we noted in Lee (2019a), our populational number ratio of M3 RGB stars is significantly different from that by Milone et al. (2017), who obtained an FG fraction of  $0.305 \pm 0.014$  based on the HST observations of the central part of the cluster. The discrepancy between our work and that by Milone et al. (2017) is solely due to a very strong radial gradient in the

populational number ratio in M3, in the sense that the CN-s population is more centrally concentrated and Milone et al. (2017) relied on the central part of the cluster, where the contribution of the CN-s population is greater, while our study tends to be weighted more in the outer part of the cluster (e.g., see Figure 8 of Lee 2019a).

Because our photometry is incomplete in the central part of the cluster due to limitations in the spatial angular resolving power of the small-aperture ground-based telescope, while that by Milone et al. (2017) is missing in the outer part of the cluster due to the small FOV of the HST, we attempted to merge the data sets. Assuming our CN-w population corresponds to the FG group by Milone et al. (2017), we obtained the populational number ratio of  $n(\text{CN-w}):n(\text{CN-s}) \approx 38:62 (\pm 2)$  from the merged data.

In panels (a)–(c) of Figure 5, we show the  $\|cn_{\text{JWL}}$  versus  $\|ch_{\text{JWL}}$ ,  $\|nh_{\text{JWL}}$  versus  $\|ch_{\text{JWL}}$ , and  $\|nh_{\text{JWL}}$  versus  $\|cn_{\text{JWL}}$  relations for M3 RGB stars in the magnitude range of  $-2 \text{ mag} \leq V - V_{\text{HB}} \leq 2 \text{ mag}$ . All of these relations exhibit discontinuities in the populational transition between the



**Figure 4.** (Bottom panel) A plot of PC1 vs. PC2, calculated using  $\|cn_{JWL}\|$ ,  $\|ch_{JWL}\|$ , and  $\|nh_{JWL}\|$ . Individual eigenvectors are also shown. All stars are shown by light gray plus signs, while dark gray plus signs indicate stars within  $2\sigma$  from the center of each population in each axis. (Top panel) Populational tagging using the expectation maximization method along the PC1.

CN-w and CN-s domains (see also Smith et al. 2013; Lee 2017, 2018, 2019a, 2019c, 2020). In no case does the photometric index comparison appear to connect the CN-w and CN-s populations, supported by the Pearson’s correlation coefficients and  $p$  values of the fit between our color indices as shown in Table 2. Our results show that the chemical evolution from the CN-w population to the CN-s population is not continuous in M3, as already discovered in other GCs in our previous studies (Lee 2017, 2018, 2019c, 2020). As we will discuss later, different index versus index trends between the CN-w and CN-s populations shown in Figure 5 are related to different [C/Fe] versus [N/Fe] relations, most likely due to different physical environments during the formation of each population.

Panels (d)–(f) of Figure 5 show the variation of the  $\|cn_{JWL}\|$ ,  $\|ch_{JWL}\|$ , and  $\|nh_{JWL}\|$  indices with  $V$  magnitude. The sharp distinction between CN-w and CN-s populations is easily seen in  $\|cn_{JWL}\|$  at all giant branch luminosities. However,  $\|ch_{JWL}\|$  and  $\|nh_{JWL}\|$  do not show a clear separation between the two groups.

## 5. Comparisons with HST Photometry

We compared our color indices with those of the pseudo-color indices devised by Milone et al. (2015),  $\Delta_{CF275W,F336W,F438W}$  and  $\Delta_{CF275W,F814W}$ . In Figure 6, we show the results for stars with the radial distance of  $r \geq 1'$  to avoid the blending effect in our results.

The figure indicates that the CN-w and CN-s populations have different correlations with  $\Delta_{CF275W,F336W,F438W}$  and  $\Delta_{CF275W,F814W}$ , which is not surprising because individual stellar populations on the so-called chromosome map of individual GCs of Milone et al. (2017) are not in linear relations. These different behaviors of individual populations have been discussed in previous papers in this series (Lee 2017, 2018, 2019c). We calculated the Pearson’s correlation coefficients and  $p$  values of the fit, and we show our results in Table 3. Our results are as follows:

1. Our  $\|nh_{JWL}\|$  index, i.e., the nitrogen abundance, does not appear to correlate with either  $\Delta_{CF275W,F336W,F438W}$  or  $\Delta_{CF275W,F814W}$ .
2. Our  $\|ch_{JWL}\|$  index, i.e., the carbon abundance, is anticorrelated with  $\Delta_{CF275W,F336W,F438W}$ , while it is positively correlated with  $\Delta_{CF275W,F814W}$ , with a different degree of correlation between the CN-w and CN-s populations.
3. Our  $\|cn_{JWL}\|$  index is positively correlated with  $\Delta_{CF275W,F336W,F438W}$ , which is consistent with the results from our previous works (Lee 2017, 2018).
4. Finally, our  $\|hk_{JWL}\|$  (i.e., the metallicity index) of the CN-w population is correlated with the  $\Delta_{CF275W,F814W}$  in M3. This suggests that the large extent of the  $\Delta_{CF275W,F814W}$  of the FG of stars in M3 pointed out by Lardo et al. (2018) and Tailo et al. (2019) is likely due to the spread in metallicity, as we will discuss below.

## 6. Photometric Elemental Abundances

### 6.1. Synthetic Grids

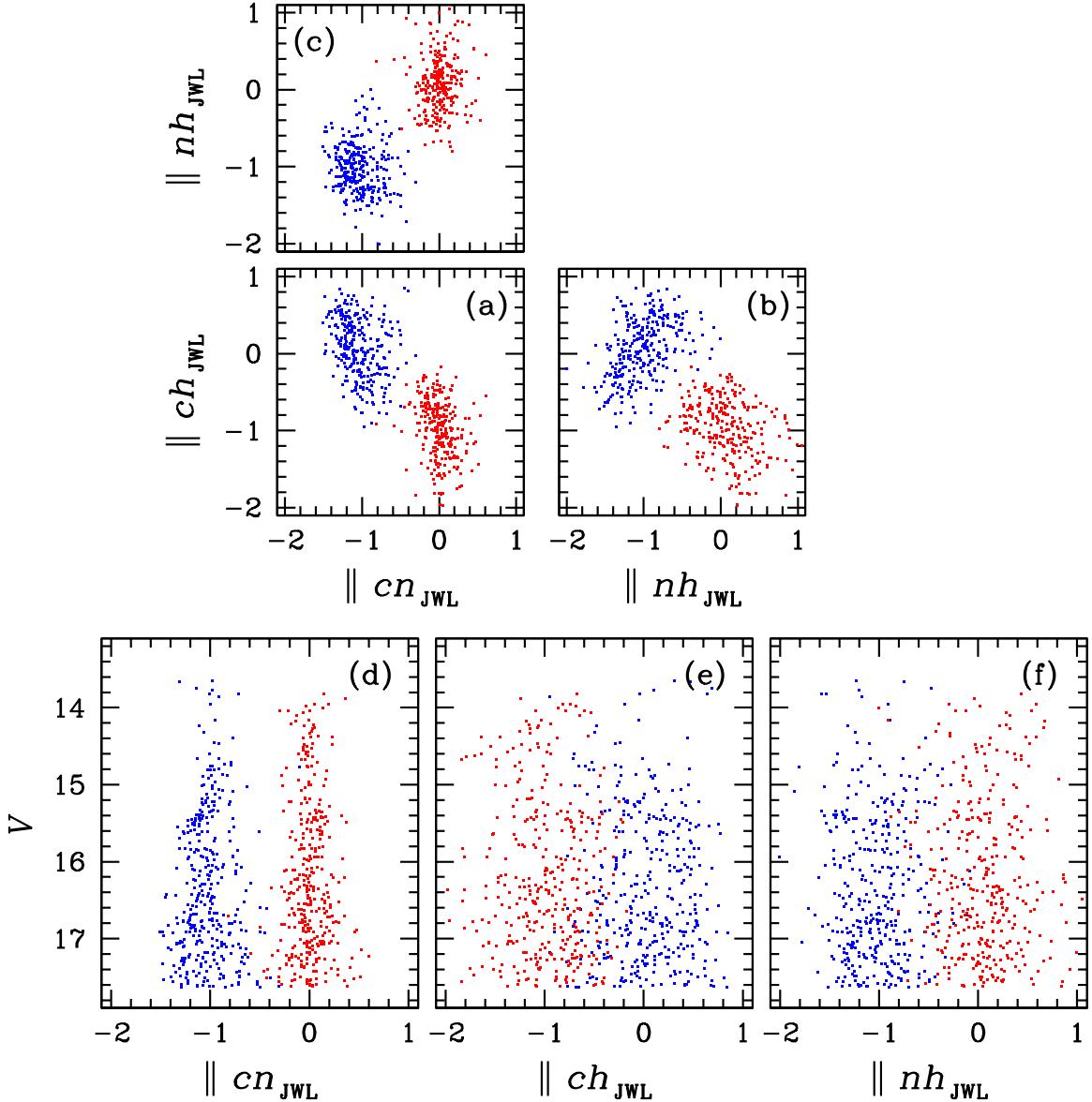
We derive photometric [Fe/H], [C/Fe], and [N/Fe] from our photometric indices by comparing synthetic grids constructed using various input parameters.

First, we obtained the Dartmouth model isochrones for [Fe/H] =  $-1.4$ ,  $-1.5$ , and  $-1.6$  with  $[\alpha/\text{Fe}] = +0.4$  and the age of 12.5 Gyr (Dotter et al. 2008). We interpolated the effective temperatures and surface gravities from  $M_V = 3.5$  to  $-2.5$  mag with a magnitude step size of  $\Delta M_V = 0.5$  mag. Using these stellar parameters, we constructed series of synthetic spectra with varying carbon and nitrogen abundances with an abundance step size of 0.1 dex for carbon and 0.15 dex for nitrogen.

For our synthetic spectrum calculations, we generated atomic/molecular line lists with the linemake facility.<sup>7</sup> We computed the spectra with the 2011 version of the local thermodynamic equilibrium line analysis code MOOG<sup>8</sup> (Sneden 1973; Sobeck et al. 2011) that includes continuum-scattering calculations. For metal-poor stars, the negative hydrogen ion’s bound-free ( $H_{bf}^-$ ) absorption dominates continuum opacities in the optical and infrared wavelength domains. However, in the blue and UV regions, Rayleigh scattering from neutral hydrogen atoms (RSNH) becomes the dominant source of continuum opacity, especially in cool RGB atmospheres (e.g., see Suntzeff 1981; Sobeck et al. 2011). Among many versions of MOOG, only that by Sobeck et al. (2011) takes proper care of RSNH, which is important to

<sup>7</sup> Available at <https://github.com/vmplacco/linemake>; line lists built by linemake begin with those in the Kurucz (2011) atomic line compendium and merge them with laboratory atomic line data from the University of Wisconsin and molecular line data from Old Dominion University. References to specific data sources are given in the linemake website.

<sup>8</sup> Available at <http://www.as.utexas.edu/~chris/moog.html>.



**Figure 5.** (a) The photometric  $\|cn_{JWL}\|$  vs.  $\|ch_{JWL}\|$  relation for RGB stars in the magnitude range  $-2 \text{ mag} \leq V - V_{\text{HB}} \leq 2 \text{ mag}$  in M3. (b) The  $\|nh_{JWL}\|$  vs.  $\|ch_{JWL}\|$  relation. (c) The  $\|cn_{JWL}\|$  vs.  $\|nh_{JWL}\|$  relation. (d) The  $\|cn_{JWL}\|$  vs.  $V$  CMD. (e) The  $\|ch_{JWL}\|$  vs.  $V$  CMD. (f) The  $\|nh_{JWL}\|$  vs.  $V$  CMD. In all panels, the blue dots denote the CN-w population while the red dots denote the CN-s population.

**Table 2**  
Pearson's Correlation Coefficients and  $p$  Values of the Fit between Our Color Indices

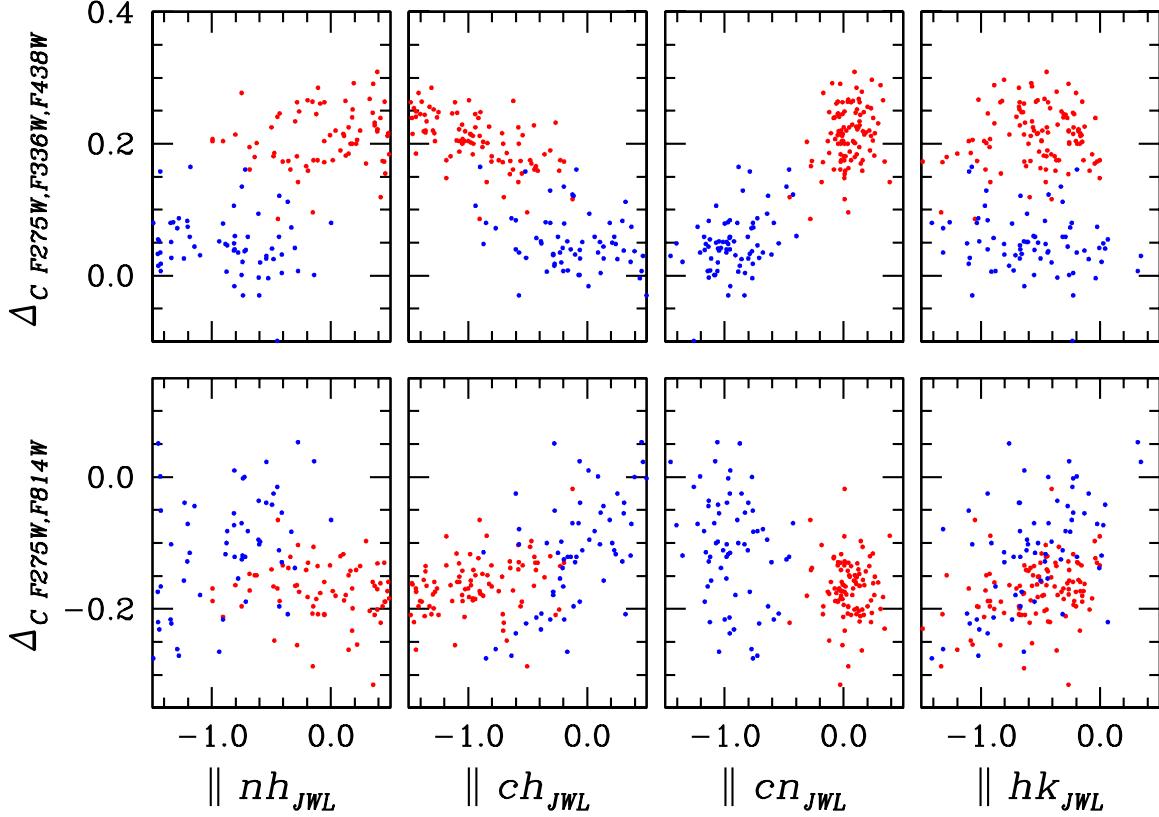
CN-w				CN-s				
$\ cn_{JWL}\ $		$\ ch_{JWL}\ $		$\ cn_{JWL}\ $		$\ ch_{JWL}\ $		
$\rho$	$p$ Value	$\rho$	$p$ Value	$\rho$	$p$ Value	$\rho$	$p$ Value	
$\ nh_{JWL}\ $	-0.060	0.206	0.525	$2.2 \times 10^{-16}$	0.153	$7.9 \times 10^{-4}$	-0.008	0.849
$\ ch_{JWL}\ $	-0.285	$9.5 \times 10^{-10}$	...	...	-0.329	$1.5 \times 10^{-13}$	...	...

calculate continuum opacities for our short-wavelength indices such as our JWL34 (the RSNH cross-section has a  $\lambda^{-4}$  dependency).

Synthetic spectra can be calculated with the code SYNTHE (Castelli 2005; Kurucz 2005). It should be noted that the advantage of using MOOG is in its fast computations and its flexible treatment of input elemental isotopic ratios. For example, MOOG is several times faster than SYNTHE, which greatly reduced our total computation time for about 13,500

synthetic spectra with various input elemental abundances and stellar parameters. During our calculations, we compared continua returned from MOOG and SYNTHE for selective models and we found that both results are consistent. Therefore, the results from our MOOG should be adequate for our purpose.

Finally, individual synthetic spectra were convolved with our filter transmission functions, and then they were converted to our photometric system.



**Figure 6.** Comparisons of our color indices with those of Milone et al. (2017). The blue and red dots denote the CN-w and CN-s RGB stars in M3.

**Table 3**  
Pearson's Correlation Coefficients and  $p$  Values of the Fit between Our Color Indices and Those of Milone et al. (2017)

	Δ <sub>CF275W,F336W,F438W</sub>				Δ <sub>CF275W,F814W</sub>			
	CN-w		CN-s		CN-w		CN-s	
	$\rho$	$p$ Value	$\rho$	$p$ Value	$\rho$	$p$ Value	$\rho$	$p$ Value
nh <sub>JWL</sub>	-0.166	0.159	0.010	0.308	0.333	0.004	-0.008	0.935
ch <sub>JWL</sub>	-0.346	0.002	-0.614	$7.7 \times 10^{-13}$	0.507	$4.1 \times 10^{-6}$	0.422	$3.9 \times 10^{-6}$
cn <sub>JWL</sub>	0.296	0.011	0.107	0.265	-0.169	0.149	-0.172	0.071
hk <sub>JWL</sub>	-0.191	0.102	0.049	0.611	0.574	$9.1 \times 10^{-8}$	0.299	0.001

We emphasize that using our  $hk_{JWL}$ ,  $ch_{JWL}$ , and  $nh_{JWL}$  makes [Fe/H], [N/Fe], and [C/Fe] abundance estimates mathematically simple. At the atmospheric line formation depth, the nitrogen enhancement does not affect the CH molecule formation rate, while the carbon enhancement very weakly influences the NH molecule formation rate through the enhanced formation of the CN molecule. Therefore, as we will discuss below, we can estimate [Fe/H], [C/Fe], and [N/Fe] as follows:

$$[\text{Fe}/\text{H}] \approx f_1(hk_{JWL}, M_V), \quad (6)$$

$$[\text{C}/\text{Fe}] \approx f_2(ch_{JWL}, [\text{Fe}/\text{H}], M_V), \quad (7)$$

$$[\text{N}/\text{Fe}] \approx f_3(nh_{JWL}, [\text{Fe}/\text{H}], M_V), \quad (8)$$

where [Fe/H] on the right-hand side of Equations (7) and (8) is fed from Equation (6). The distance modulus and interstellar reddening value for M3 are well known and, therefore, we can obtain accurate  $M_V$  of individual RGB stars from our photometry. At a given  $M_V$ , our synthetic model grids can be

interpolated to match the observed  $hk_{JWL}$ ,  $ch_{JWL}$ , and  $nh_{JWL}$  to derive [Fe/H], [C/Fe], and [N/Fe].

Similarly, the [C/Fe] and [N/Fe] from the  $cn_{JWL}$ , [C/Fe]<sub>CN</sub>, and [N/Fe]<sub>CN</sub>, can be given as

$$[\text{C}/\text{Fe}]_{\text{CN}} \approx f_4(cn_{JWL}, [\text{Fe}/\text{H}], [\text{N}/\text{Fe}], M_V), \quad (9)$$

$$[\text{N}/\text{Fe}]_{\text{CN}} \approx f_5(cn_{JWL}, [\text{Fe}/\text{H}], [\text{C}/\text{Fe}], M_V), \quad (10)$$

where [N/Fe] and [C/Fe] on the right-hand side of Equations (9) and (10) are fed from Equations (8) and (7), respectively. At given  $M_V$ , [Fe/H], [N/Fe], and [C/Fe], our synthetic model grids were interpolated to match observed  $cn_{JWL}$  to derive [C/Fe]<sub>CN</sub> and [N/Fe]<sub>CN</sub>.

## 6.2. Metallicity

We explore the metallicity spread of M3 using our Ca II H & K photometry. Calcium is synthesized in high-mass stars that die explosively, not in the kind of medium-mass stars whose AGB ejecta probably seeded the CN-s stars of a GC's later population. This makes our  $hk_{JWL}$  an excellent measure of

metallicity with a fixed [Ca/Fe] abundance (Anthony-Twarog et al. 1991; Lee et al. 2009a; Lee 2015). One of the caveats for using the Ca II H & K lines is that these stellar lines arise from the Ca II ground state, but so do absorption lines from gas in the interstellar medium. Thus, interstellar contamination could affect the stellar absorption strengths. In our previous study (see Supplementary Information of Lee et al. 2009a), we demonstrated that the absorption of Ca II H & K lines by interstellar media having an interstellar reddening of  $E(B-V) = 0.32$  mag contributes only  $\Delta h_k = 0.01$  mag to our results. Because M3 does not show any perceptible differential reddening effect, the contribution of interstellar Ca II H & K absorption can be completely ignored.

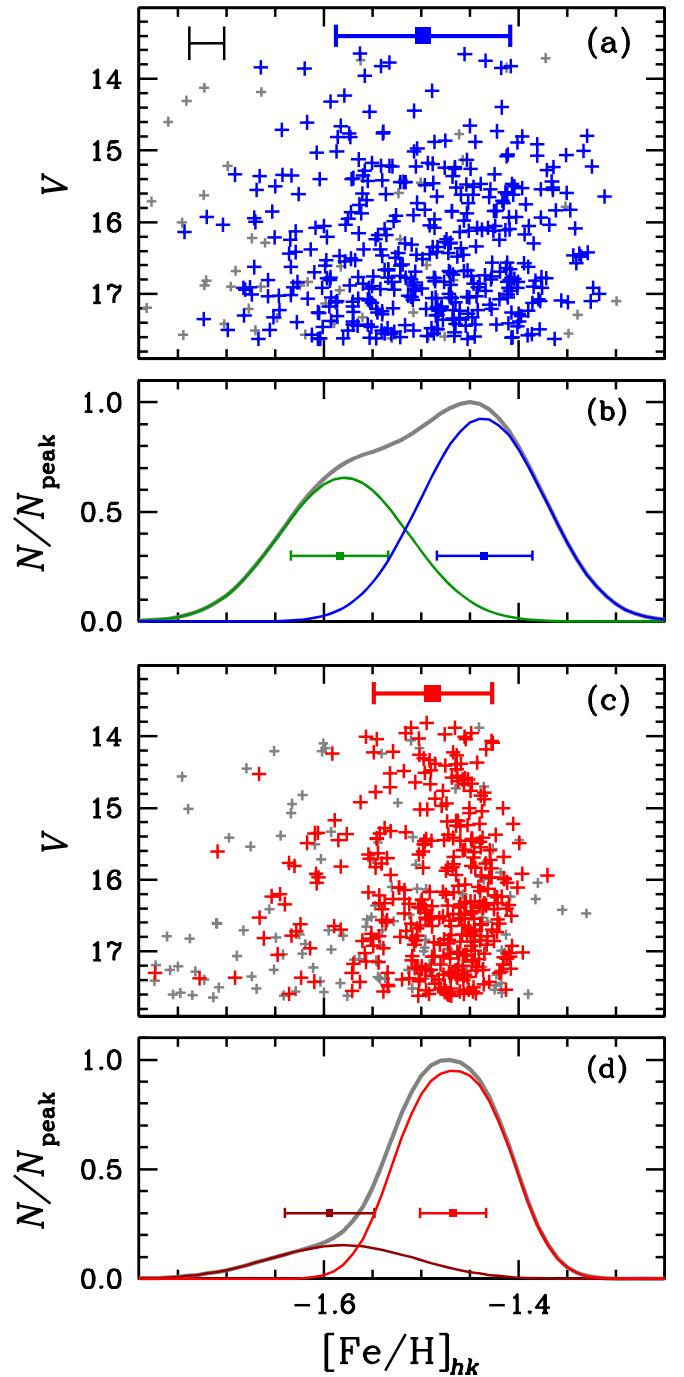
In Figure 7, we show plots of the [Fe/H] derived from our  $h_{k_{JWL}}$  index,  $[\text{Fe}/\text{H}]_{hk}$ , versus the  $V$  magnitude for the CN-w and CN-s populations.  $h_{k_{JWL}}$  is an almost reddening-independent index,  $E(h_{k_{JWL}}) = -0.12 \times E(B-V)$ . Therefore, the  $[\text{Fe}/\text{H}]_{hk}$  dispersions shown in Figure 7 cannot be due to differential reddening (Anthony-Twarog et al. 1991; Lee et al. 2009a). For this figure, we selected only those RGB stars with very small measurement errors,  $\sigma(h_k) \leq 0.01$  mag, in each population in order to ensure that their  $h_{k_{JWL}}$  and subsequently  $[\text{Fe}/\text{H}]_{hk}$  values are not affected by observational errors. Using Equation (6), we obtained  $[\text{Fe}/\text{H}]_{hk}$  from our  $h_{k_{JWL}}$  measurements. Metallicity gradients with  $V$  are absent in both populations in the figure, which is natural because calcium is not synthesized or destroyed along the RGB phase of low-mass stars. We obtained the mean photometric metallicities of  $\langle [\text{Fe}/\text{H}]_{hk} \rangle = -1.498 \pm 0.089 \pm 0.004$  for the CN-w population and  $-1.488 \pm 0.061 \pm 0.003$  for the CN-s population, a nearly identical value (the errors denote standard deviations and standard errors).

The most interesting and unexpected aspect of Figure 7 is that our metallicity distributions for each population are well described by bimodal distributions. In Table 4, we show the peak [Fe/H] values for individual subpopulations and the subpopulational number ratios returned from our EM estimators. The peak metallicities of subpopulations are in excellent agreement between the CN-w and CN-s populations. However, the subpopulational number ratios are very different. The CN-w population has  $n(\text{SP1}):n(\text{SP2}) = 36:64$  ( $\pm 4$ ), while the CN-s population has 17:83 ( $\pm 3$ ), where SP1 refers to the metal-poor component.

As we will argue later, the CN-w population shows a distinctive double RGBB distribution, which is related to the double metallicity distribution of the CN-w population. On the other hand, we think that the conspicuous double RGBBs cannot be seen in the CN-s population due to the low fraction of the metal-poor component, as we will discuss below.

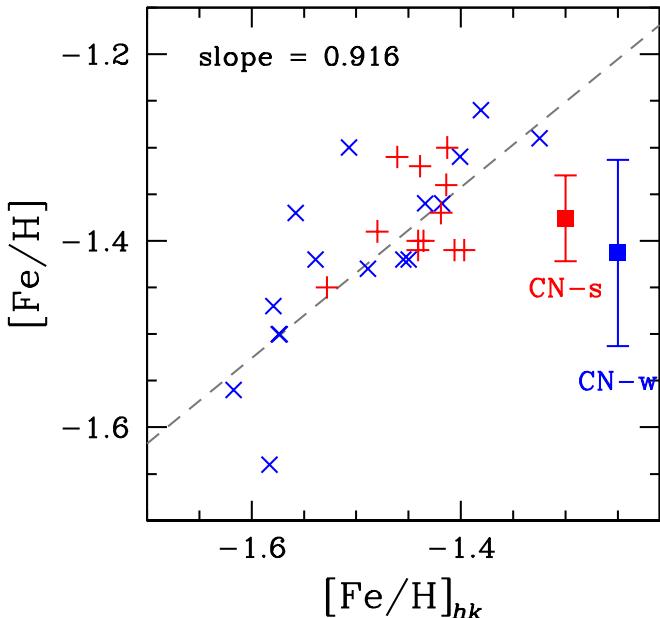
Our results clearly show that the  $[\text{Fe}/\text{H}]_{hk}$  (i.e., the  $h_{k_{JWL}}$ ) dispersion of the CN-w RGB stars is larger than that of the CN-s population due to the low fraction of the metal-poor component of the CN-s population, in accordance with the results from high-resolution spectroscopy that the metallicity spread of the CN-w population is larger than that of the CN-s population (e.g., Sneden et al. 2004; Mészáros et al. 2015).

Sneden et al. (2004) studied 28 RGB stars in M3 using high-resolution visual spectroscopy. Unfortunately, because their sample stars are either too bright ( $V < V - V_{\text{HB}} - 2$  mag) or too close to the cluster's center ( $r < 1'$ ), none of the stars studied by Sneden et al. (2004) fall within our sample RGB selection criteria ( $-2 \text{ mag} \leq V - V_{\text{HB}} \leq 2 \text{ mag}$  and  $r \geq 1'$ ). However, Marino et al. (2019) suggested that the metallicity spread for



**Figure 7.** (a) Photometric metallicity from our  $h_{k_{JWL}}$ ,  $[\text{Fe}/\text{H}]_{hk}$ , vs.  $V$  magnitude of the CN-w population. The blue plus signs denote the CN-w RGB stars with low measurement errors,  $\sigma(h_{k_{JWL}}) \leq 0.01$  mag, while the gray plus signs indicate RGB stars with large measurement errors,  $\sigma(h_{k_{JWL}}) > 0.01$  mag, those are not used in our analysis. The blue filled square and the error bar indicate the mean  $[\text{Fe}/\text{H}]_{hk}$  values and the standard deviation for the CN-w population shown with blue plus signs. The black error bar in the upper-left corner denotes the  $\pm 1\sigma$  uncertainty in our  $[\text{Fe}/\text{H}]_{hk}$  measurements. (b) Generalized histograms returned from our EM estimator for each population. We also show the mean  $[\text{Fe}/\text{H}]_{hk}$  values and standard deviations of each subpopulation. (c) Same as (a), but for the CN-s population. (d) Same as (b), but for the CN-s population.

the CN-w population is larger. They obtained  $\langle [\text{Fe}/\text{H}] \rangle = -1.63 \pm 0.07$  for FG, which is equivalent to the CN-w, and  $-1.57 \pm 0.04$  for SG, i.e., CN-s, for M3 using the results by Sneden et al. (2004).



**Figure 8.** Comparison of our  $[\text{Fe}/\text{H}]_{hk}$  with  $[\text{Fe}/\text{H}]$  of Mészáros et al. (2015). The blue and red dots denote the CN-w and CN-s RGB stars. The error bars indicate the mean  $[\text{Fe}/\text{H}]$  values and standard deviations of high-resolution spectroscopic measurements by Mészáros et al. (2015).

**Table 4**  
[Fe/H] of Subpopulations and Number Ratios

	[Fe/H]		$n(\text{SP1})$ : $n(\text{SP2})$
	SP1	SP2	
CN-w	$-1.584 \pm 0.050 \pm 0.004$	$-1.435 \pm 0.049 \pm 0.003$	36:64 ( $\pm 4$ )
CN-s	$-1.594 \pm 0.046 \pm 0.006$	$-1.467 \pm 0.034 \pm 0.002$	17:83 ( $\pm 2$ )

More recently, Mészáros et al. (2015) performed a high-resolution infrared spectroscopic study of M3. In Figure 8, we show comparisons of our  $[\text{Fe}/\text{H}]_{hk}$  measurements with those of Mészáros et al. (2015). As shown in the figure, our  $[\text{Fe}/\text{H}]_{hk}$  measurements correlate well with the  $[\text{Fe}/\text{H}]$  measurements by Mészáros et al. (2015), confirming again that our  $h_{\text{JWL}}$  is a good measure of metallicity. We calculate Pearson's correlation coefficients and  $p$  values of the fit, obtaining  $\rho = 0.749$  and  $p$  value of  $4.5 \times 10^{-6}$ . Inspection of data by Mészáros et al. (2015) also suggests that the metallicity spreads of CN-w stars are larger than those of the CN-s, and statistics of the two samples support this notion. For the CN-w group  $\langle [\text{Fe}/\text{H}] \rangle = -1.41 \pm 0.10 \pm 0.03$ , while for the CN-s group  $\langle [\text{Fe}/\text{H}] \rangle = -1.38 \pm 0.05 \pm 0.01$ . We depict the  $[\text{Fe}/\text{H}]$  sample statistics with vertical lines drawn in Figure 8 at arbitrary  $[\text{Fe}/\text{H}]_{hk}$  positions for clarity. This result does not appear to conform to simple ideas of the sequential chemical evolution from the CN-w to the CN-s populations, in which the later generation of stars, i.e., the CN-s population in our study, would be naturally expected to have larger scatter in elemental abundances.

Finally, note in Figure 8 that the number of the CN-s stars studied by Mészáros et al. (2015) is smaller than that of the CN-w stars,  $n(\text{CN-w}):n(\text{CN-s}) = 16:12$ . This may support our result that the CN-s population is more centrally concentrated than the CN-w population because the sample RGB stars

studied by Mészáros et al. (2015) are located in the outer part of the cluster.

### 6.3. Carbon and Nitrogen

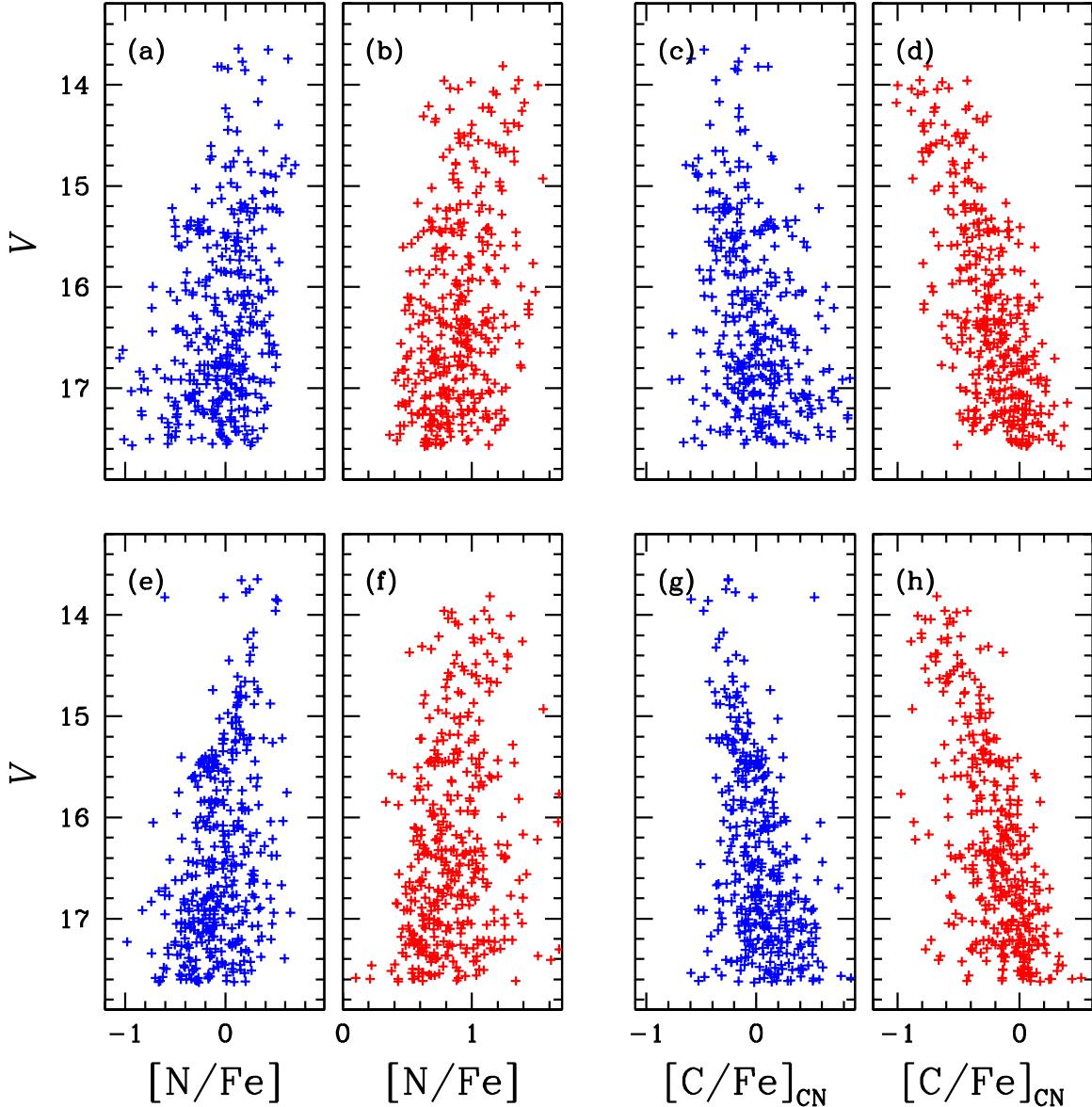
We derive photometric  $[\text{C}/\text{Fe}]$  and  $[\text{N}/\text{Fe}]$  abundances using our photometric measurements in two ways: (1) direct photometric  $[\text{C}/\text{Fe}]$  and  $[\text{N}/\text{Fe}]$  abundances from the  $ch_{\text{JWL}}$  and  $nh_{\text{JWL}}$  indices using Equations (7) and (8), and (2) indirect photometric  $[\text{C}/\text{Fe}]_{\text{CN}}$  and  $[\text{N}/\text{Fe}]_{\text{CN}}$  abundances from  $cn_{\text{JWL}}$  with prior  $[\text{C}/\text{Fe}]$  and  $[\text{N}/\text{Fe}]$  measurements from the  $ch_{\text{JWL}}$  and  $nh_{\text{JWL}}$  using Equations (9) and (10), respectively.

Carbon and nitrogen abundances derived from the CN band absorption strengths (as in our approach or as in low- to intermediate-resolution spectroscopy) depend on oxygen abundances due to the competition of CO molecules for carbon. At fixed carbon and nitrogen abundances, as oxygen abundance increases, CO formation is slightly enhanced. Then, the CN and CH formations are slightly suppressed due to a decrease in available atomic carbon. On the other hand, the NH band strength (i.e.,  $[\text{N}/\text{Fe}]$  from our  $nh_{\text{JWL}}$ ) remains almost intact. Our calculations show that the NH populations at the optical depth of the line formation do not change with the variation of oxygen abundances. The study of Yong et al. (2008) nicely illustrates the situation. They obtained nitrogen abundances of a dozen RGB stars in NGC 6752 based on high-resolution spectroscopy of the NH band at  $\lambda 3360$ , and they argued that their nitrogen abundances are significantly different from those derived from the CN band by others. Yong et al. (2008) attributed the discrepancy in nitrogen abundances to the inappropriate assignment of oxygen abundances during  $[\text{N}/\text{Fe}]$  derivation from the CN band by others.

To compare our photometric carbon and nitrogen abundances with those of Mészáros et al. (2015), we computed synthetic spectra assuming  $[\text{O}/\text{Fe}] = 0.50$  for CN-w and 0.15 for CN-s populations. These mean  $[\text{O}/\text{Fe}]$  values are consistent with those of three stars for the CN-w and two stars for the CN-s populations by Mészáros et al. (2015). Note that their oxygen abundances are about 0.2–0.3 dex larger than those of Sneden et al. (2004), who did not measure carbon and nitrogen abundances in M3. Using the results by Sneden et al. (2004), we calculated  $\langle [\text{O}/\text{Fe}] \rangle = 0.16 \pm 0.06$  (12 stars) for stars with  $[\text{Na}/\text{Fe}] < 0$ , which roughly correspond to our CN-w population, and  $-0.07 \pm 0.12$  for stars with  $[\text{Na}/\text{Fe}] \geq 0$  (i.e.,  $\approx$ CN-s in our study). Our mean oxygen abundances are comparable to those of Marino et al. (2019), who obtained  $\langle [\text{O}/\text{Fe}] \rangle = 0.14 \pm 0.04$  (four stars) for the FG group and  $0.00 \pm 0.10$  (eight stars) for the SG group using the results of Sneden et al. (2004).

We are puzzled by this discrepancy because the oxygen abundance from the IR OH lines is supposed to be consistent with that from the  $[\text{O}/\text{I}]$  line (e.g., see Balachadran & Carney 1996). The origin of this oxygen abundance discrepancy is beyond the scope of our current study. However, we caution that carbon and nitrogen abundances by Mészáros et al. (2015) could be incorrect because the oxygen abundances may affect the subsequent carbon and nitrogen abundance determinations through CO and CN bands.

Our  $[\text{N}/\text{Fe}]_{\text{CN}}$  values are based on the same method as that employed in typical  $[\text{N}/\text{Fe}]$  derivations from CN bands used for low- to intermediate-resolution spectroscopy. But computations of  $[\text{C}/\text{Fe}]_{\text{CN}}$  are less frequent in high-resolution spectroscopic studies. Therefore, we derived  $[\text{C}/\text{Fe}]_{\text{CN}}$  values in order

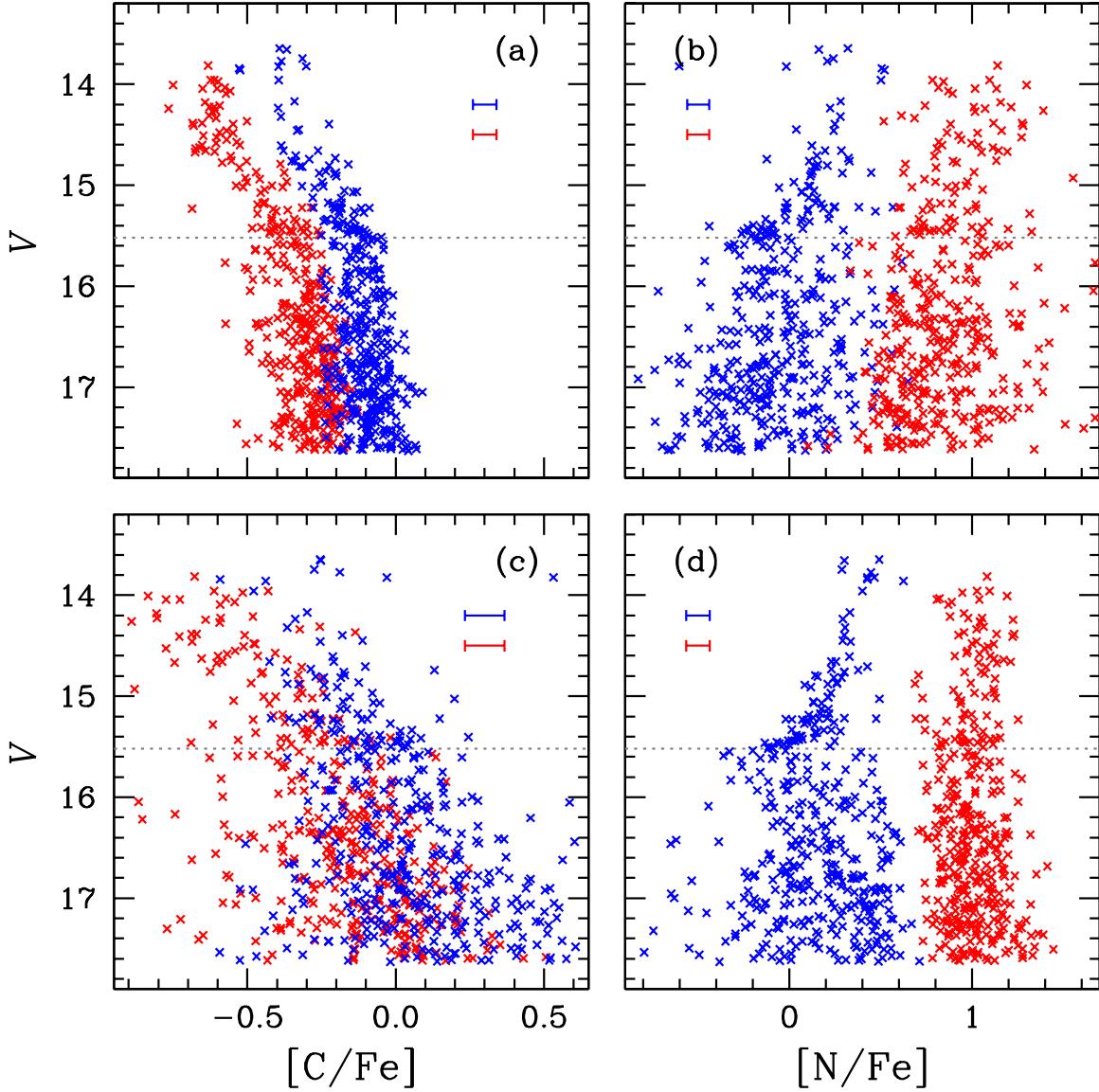


**Figure 9.** (a)–(b)  $[\text{N}/\text{Fe}]$  from  $nh_{\text{JWL}}$ . The blue and red colors indicate the CN-w and CN-s stars. (c)–(d)  $[\text{C}/\text{Fe}]$  from  $cn_{\text{JWL}}$  with  $[\text{N}/\text{Fe}]$  from  $nh_{\text{JWL}}$ . (e)–(f) Same as (a)–(b) but the metallicity effect is corrected. (g)–(h) Same as (c)–(d) but the metallicity effect is corrected.

to check the internal consistency between carbon abundances from the  $ch_{\text{JWL}}$  and  $cn_{\text{JWL}}$  indices.

Our photometric  $[\text{N}/\text{Fe}]$  measurements from the  $nh_{\text{JWL}}$  index depend on metallicity as enhanced metallicity affects not only the line opacity of metallic lines, in particular in the shorter-wavelength regime, but also the continuum opacity. As we will discuss in Appendix C,  $\Delta[\text{Fe}/\text{H}] = 0.1$  dex can result in  $\Delta[\text{N}/\text{Fe}] \approx 0.1\text{--}0.2$  dex and the CN-w group is more vulnerable to such effect because it shows a substantial metallicity spread. We attempted to correct the metallicity effect on our photometric nitrogen abundance using our metallicity estimated from the  $hk_{\text{JWL}}$  already shown in Figure 7. In Figure 9, we show the  $[\text{N}/\text{Fe}]$  distributions with and without the metallicity correction, where a narrower RGB sequence in the  $[\text{N}/\text{Fe}]$  distribution of the CN-w can be seen when a metallicity correction is applied. It is also true for the  $[\text{C}/\text{Fe}]_{\text{CN}}$  of the CN-w, because  $[\text{N}/\text{Fe}]$  is used to derive  $[\text{C}/\text{Fe}]_{\text{CN}}$ . For the CN-s population, the improvement is not as great as those that can be seen in the CN-w population.

In Figure 10 and Table 5, we show carbon and nitrogen distributions of M3 RGB stars. In the figure, the measurement uncertainties ( $\pm 1\sigma$ ) were estimated by the bootstrap method for 200,000 resampled data for each case. For fRGB stars, we obtained  $\langle [\text{C}/\text{Fe}] \rangle = -0.10 \pm 0.07 \pm 0.00$  and  $\langle [\text{N}/\text{Fe}] \rangle = 0.10 \pm 0.28 \pm 0.01$  for the CN-w population, while  $\langle [\text{C}/\text{Fe}] \rangle = -0.30 \pm 0.08 \pm 0.00$  and  $\langle [\text{N}/\text{Fe}] \rangle = 0.82 \pm 0.27 \pm 0.02$  for the CN-s population. Both populations have nearly identical abundance dispersions for  $[\text{C}/\text{Fe}]$  and  $[\text{N}/\text{Fe}]$ . Also note that the carbon and nitrogen abundances from the CN band at  $\lambda 3883$  are higher than those from the CH and NH bands. If we attribute this discrepancy to uncertain oxygen abundances, our high  $[\text{C}/\text{Fe}]_{\text{CN}}$  and  $[\text{N}/\text{Fe}]_{\text{CN}}$  abundances may indicate that our assumed oxygen abundances from Mészáros et al. (2015) for each population ( $[\text{O}/\text{Fe}] = 0.5$  dex for the CN-w and 0.15 dex for the CN-s) are about 0.2 dex larger than they should be (see Figure 28). If so, the resultant oxygen abundances ( $[\text{O}/\text{Fe}] = 0.30$  dex for the CN-w and  $-0.05$  dex for the CN-s) are in good agreement with those of Sneden et al. (2004).



**Figure 10.** (a)  $[\text{C}/\text{Fe}]$  from  $ch_{\text{JWL}}$ . The blue crosses are for the CN-w stars, while the red crosses are for the CN-s stars. The blue and red error bars denote measurement uncertainties ( $\pm 1\sigma$ ) for the CN-w and CN-s stars, respectively. (b)  $[\text{N}/\text{Fe}]$  from  $nh_{\text{JWL}}$ . (c)  $[\text{C}/\text{Fe}]$  from  $cn_{\text{JWL}}$  with  $[\text{N}/\text{Fe}]$  from panel (b) (i.e.,  $[\text{N}/\text{Fe}]$  from  $nh_{\text{JWL}}$ ). (d)  $[\text{N}/\text{Fe}]$  from  $cn_{\text{JWL}}$  with  $[\text{C}/\text{Fe}]$  from panel (a) (i.e.,  $[\text{C}/\text{Fe}]$  from  $ch_{\text{JWL}}$ ).

**Table 5**  
Mean Photometric  $[\text{C}/\text{Fe}]$  and  $[\text{N}/\text{Fe}]$  of M3 RGBs

	All		fRGB( $V \geq V_{\text{bump}}$ )	
	$[\text{C}/\text{Fe}]$	$[\text{N}/\text{Fe}]$	$[\text{C}/\text{Fe}]$	$[\text{N}/\text{Fe}]$
CN-w	$-0.12 \pm 0.09 \pm 0.00$	$-0.06 \pm 0.28 \pm 0.01$	$-0.10 \pm 0.07 \pm 0.00$	$0.10 \pm 0.28 \pm 0.01$
CN-s	$-0.35 \pm 0.13 \pm 0.01$	$0.85 \pm 0.26 \pm 0.02$	$-0.30 \pm 0.08 \pm 0.00$	$0.82 \pm 0.27 \pm 0.02$
CN-w <sup>a</sup>	$0.02 \pm 0.32 \pm 0.02$	$0.12 \pm 0.39 \pm 0.02$	$0.06 \pm 0.33 \pm 0.02$	$0.10 \pm 0.42 \pm 0.02$
CN-s <sup>a</sup>	$-0.21 \pm 0.26 \pm 0.01$	$1.01 \pm 0.14 \pm 0.01$	$-0.13 \pm 0.23 \pm 0.01$	$1.01 \pm 0.14 \pm 0.01$

**Note.**

<sup>a</sup>  $[\text{C}/\text{Fe}]_{\text{CN}}$  and  $[\text{N}/\text{Fe}]_{\text{CN}}$ .

We compare our photometric carbon and nitrogen abundances with those of Mészáros et al. (2015) and Smith (2002), showing our results in Figure 11 and Table 6. In the table, we show standard deviations and standard errors for each case. For the CN-w population, our carbon and nitrogen abundances are in good agreement with those of Mészáros et al. (2015), while

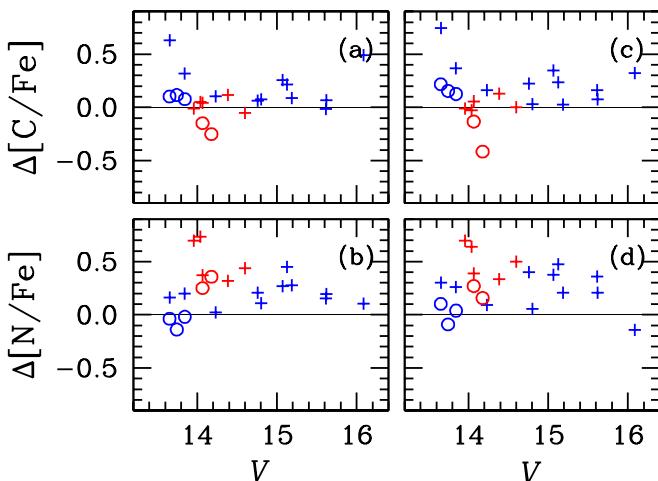
our  $[\text{C}/\text{Fe}]$  is slightly smaller and  $[\text{N}/\text{Fe}]$  is slightly larger than those of Mészáros et al. (2015) in the CN-s. Our  $[\text{C}/\text{Fe}]_{\text{CN}}$  and  $[\text{N}/\text{Fe}]_{\text{CN}}$  are larger than  $[\text{C}/\text{Fe}]$  and  $[\text{N}/\text{Fe}]$  from the  $ch_{\text{JWL}}$  and  $nh_{\text{JWL}}$  indices and, as consequences, they are larger than those of Mészáros et al. (2015), except for the  $[\text{C}/\text{Fe}]_{\text{CN}}$  of the CN-s population. The differences between our results and those

**Table 6**  
Comparisons of Photometric [C/Fe] and [N/Fe] to Mészáros et al. (2015) and Smith (2002)

	Mészáros et al. (2015)		Smith (2002)	
	[C/Fe]	[N/Fe]	[C/Fe]	[N/Fe]
CN-w	$0.10 \pm 0.02 \pm 0.01$	$-0.07 \pm 0.05 \pm 0.03$	$0.21 \pm 0.19 \pm 0.06$	$0.20 \pm 0.11 \pm 0.03$
CN-s	$-0.24 \pm 0.06 \pm 0.04$	$0.13 \pm 0.03 \pm 0.02$	$0.03 \pm 0.06 \pm 0.03$	$0.51 \pm 0.17 \pm 0.07$
CN-w <sup>a</sup>	$+0.20 \pm 0.05 \pm 0.04$	$+0.30 \pm 0.05 \pm 0.04$	$0.25 \pm 0.20 \pm 0.06$	$0.24 \pm 0.17 \pm 0.05$
CN-s <sup>a</sup>	$-0.27 \pm 0.14 \pm 0.10$	$+0.21 \pm 0.06 \pm 0.04$	$0.03 \pm 0.06 \pm 0.03$	$0.51 \pm 0.14 \pm 0.06$

**Note.**

<sup>a</sup> [C/Fe] and [N/Fe] from  $cn_{JWL}$ .



**Figure 11.** (a) Comparisons of [C/Fe] from  $ch_{JWL}$  vs. [C/Fe] by Mészáros et al. (2015, open circles) and by Smith (2002, plus signs). The blue and red colors indicate the CN-w and CN-s populations. (b) Comparisons of [N/Fe] from  $nh_{JWL}$  vs. [C/Fe] by Mészáros et al. (2015, open circles) and by Smith (2002, plus signs). (c) Same as (a) but for  $[C/Fe]_{CN}$ . (d) Same as (b) but for  $[N/Fe]_{CN}$ .

of Smith (2002) are larger, but this discrepancy is not a major issue. As Smith (2002) noted, their results were based on merged data in the literature, and all of their [N/Fe] and most of their [C/Fe] were taken from Suntzeff (1981), whose input model atmospheres and spectral line information must have been very different than those used in our work. We simply note that the absence of the gradient in the abundance difference against the  $V$  magnitude is encouraging.

Finally, we show a plot of [N/Fe] versus [C/Fe] of the M3 fRGB stars in Figure 12. We also calculated Pearson's correlation coefficients and  $p$  values of the fit between the [C/Fe] versus [N/Fe], and we obtained  $\rho = -0.117$  and a  $p$  value of 0.041 for the CN-w and  $\rho = -0.360$  and a  $p$  value of  $1.8 \times 10^{-10}$  for the CN-s populations. The figure clearly shows the two separate relations between the carbon and nitrogen abundances. For the CN-w population, the carbon abundance is weakly anticorrelated with (or, perhaps, remains flat against) the nitrogen abundance, which may indicate that the CN-w fRGB stars formed out of gas experienced no or little CN-cycle hydrogen burning, while the carbon and nitrogen abundances are more strongly anticorrelated in the CN-s population, which is a natural consequence of the CN-cycle hydrogen burning.

## 7. Populational Tagging for RHB

GC RHB stars are not warm enough to completely suppress the formation of diatomic molecules, such as NH, CN, and CH,

and they can be used in populational tagging. For example, Norris & Freeman (1982) and Smith & Penny (1989) studied CN bimodal distributions of RHB stars in metal-rich GCs, 47 Tuc and M71, respectively.

In Figure 13, we show a plot of  $(b - y)_0$  versus  $cn_{JWL}$  for RHB stars in M3 and M71. Note that the photometry for M71 is from our unpublished work. As shown in the figure, the RHB stars in the metal-rich GC, M71,  $[Fe/H] \approx -0.8$  dex, have a discrete bimodal  $cn_{JWL}$  distribution, which is consistent with the results by Smith & Penny (1989; see their Figures 2 and 3). Using our photometry, we obtained the RHB populational ratio of  $n(\text{CN-w}):n(\text{CN-s}) = 46:54 (\pm 11)$  for M71.

It is evident that the distribution of M3 RHB stars is different from that of M71. In particular, the  $cn_{JWL}$  indices of M3 RHB stars increase with decreasing  $(b - y)_0$  color. This is due to increasing contributions with effective temperature from  $H_{\zeta}$  and  $H_{\eta}$  at 3889.05 Å and 3835.38 Å, respectively, which reside in the JWL38 passband. The six RHB stars that deviate far from the main body of the M3 RHB stars are thought to be members of the CN-s population. If so, the M3 RHB stars are mainly composed of the CN-w population,  $94\% \pm 3\%$ , which is consistent with the results from cumulative radial distributions as will be shown below. We also note that our results are in good agreement with the results of synthetic HB model simulations by Tailo et al. (2019) that the CN-w population is the major component for M3 RHB stars.

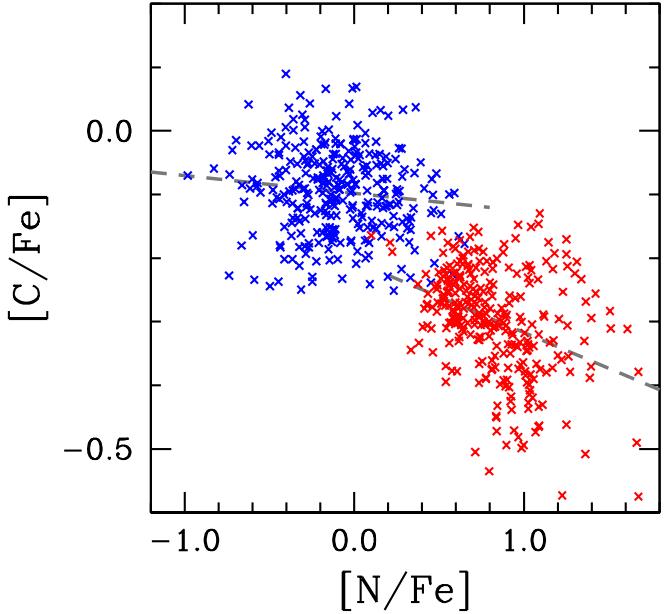
## 8. Structural Differences between Multiple Stellar Populations

### 8.1. Cumulative Radial Distributions

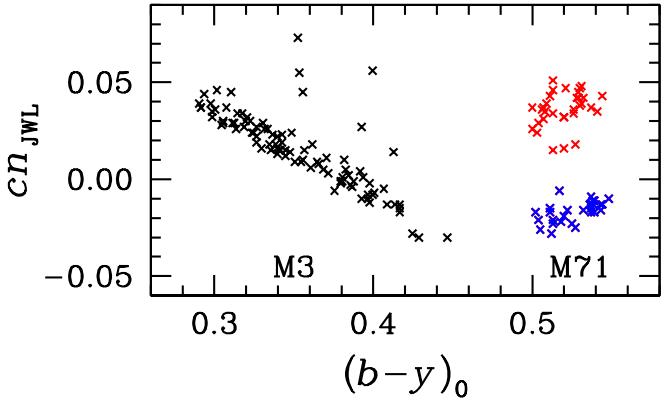
The cumulative radial distributions (CRDs) of individual populations in GCs are frequently used to characterize populational properties and can also provide important information on their kinematics (e.g., Vesperini et al. 2013). In our previous study, we showed that the M3 CN-s population is more centrally concentrated than the CN-w population and M3 has a strong radial gradient in the populational number ratio (Lee 2019a).

In Figure 14, we show the CRDs of the CN-w and CN-s RGB populations, RHB, RRL, and blue horizontal-branch (BHB) stars in M3. We have not used stars in the central part of the cluster ( $r \leq 1'$ ), due to detection-completeness issues. We calculated the horizontal-branch ratio (HBR)<sup>9</sup> from our photometry of HB stars with  $r \geq 1'$ , and we obtained  $HBR = 0.08 \pm 0.02$ , which is in excellent agreement with that of Harris (1996, the 2003 edition), 0.08. We performed

<sup>9</sup>  $HBR = (B - R)/(B + V + R)$ , where  $B$ ,  $R$ , and  $V$  denote the numbers of the BHB, RHB, and RRL, respectively.



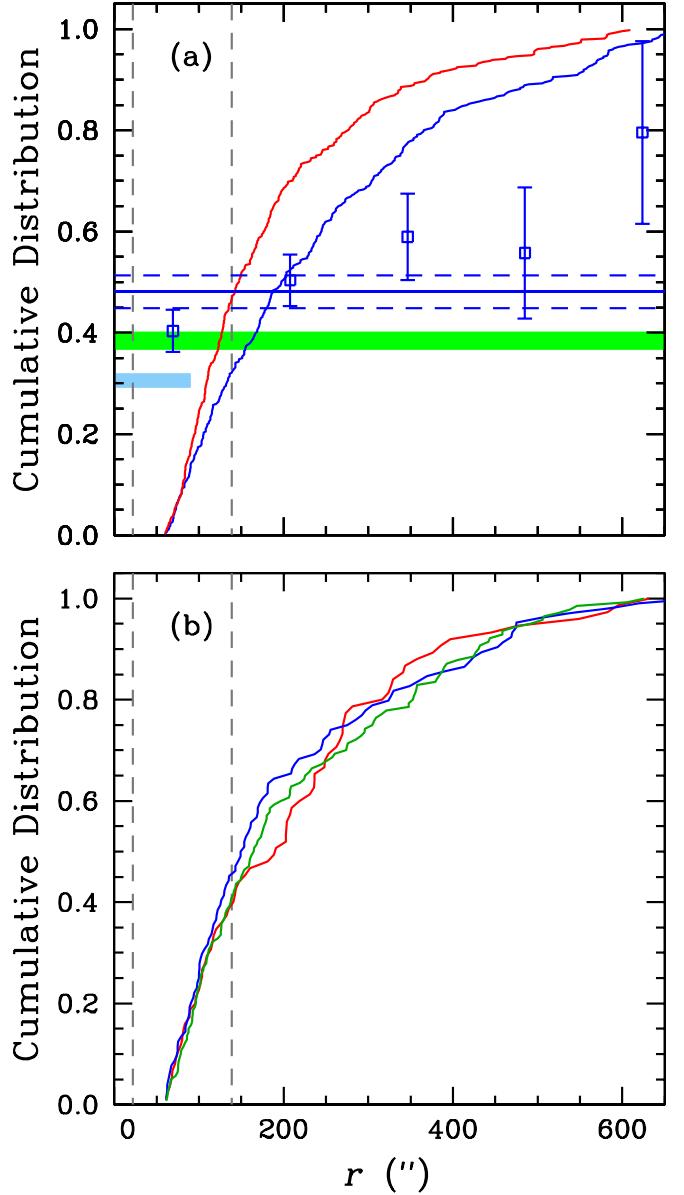
**Figure 12.** Plot of  $[C/Fe]$  vs.  $[N/Fe]$  for M3 fRGB stars in the CN-w (blue) and the CN-s (red) populations. The dashed lines show least-squares regression. The carbon abundance is weakly anticorrelated with (or remains flat against) the nitrogen abundance in the CN-w population, while the carbon and nitrogen abundances are more strongly anticorrelated in the CN-s population.



**Figure 13.** Plot of  $(b-y)_0$  vs.  $cn_{JWL}$  for RHB stars in M3 and the metal-rich GC M71. The CN-w and CN-s M71 RHB stars are denoted with blue and red crosses, respectively. The six RHB stars separated far from the main body of the M3 RHB stars are thought to be the CN-s RHB stars.

Kolmogorov–Smirnov (K-S) tests between individual RGB and HB populations to see if they have the same CRDs and the  $p$  values from our K-S tests are given in Table 7. Our results are as follows:

1. The CN-w and CN-s RGB stars have completely different radial distributions.
2. The CN-w RGB and RHB are most likely drawn from the same parent distribution, suggesting that the bulk of the RHB stars are the progeny of the CN-w RGB population. This is consistent with our result presented in Section 7 that the major component of the M3 RHB is the CN-w population. Similarly, CN-s RGB and BHB are most likely drawn from the same parent distribution, and most of the BHB stars are most likely the progeny of the CN-s RGB population.



**Figure 14.** (a) CRDs of CN-w (the blue solid line) and CN-s (the red solid line) in the radial zone of  $1' \leq r \leq 10'$ . The blue open squares and error bars indicate the mean CN-w fractions at given radial zones. The horizontal blue solid and dashed lines denote the mean fraction for the CN-w population with  $\pm 1\sigma$ , and the vertical gray dashed lines denote the core and half-light radii of M3. The cyan shaded area indicates the fraction of the FG population with  $\pm 1\sigma$  and the FOV of Milone et al. (2017), while the green shaded area indicates the fraction of the CN-w population estimated from merged data. (b) CRDs of the RHB (the red solid line), RRL (the green solid line), and BHB (the blue solid line) stars in M3 with  $r \geq 1'$ .

3. The CN-w RGB and BHB are not most likely drawn from the same parent distribution. The same is true for the CN-s RGB and RHB stars.
4. The RRL stars are likely a mixed population, i.e., the progeny of both the CN-w and CN-s RGB populations.
5. Unlike RGB populations, the RHB and BHB stars appear to share a similar CRD. It is possible that small numbers of RHB and BHB stars may result in a slightly ambiguous  $p$  value in our K-S test.

The similarity in the structural properties between the CN-w RGB and RHB stars can also be found in Figure 15. In the

**Table 7**  
 $p$  Values Returned from the K-S Tests for CRDs with  $r \geq 1'$

	$p$ Value
CN-w vs. CN-s	$4.3 \times 10^{-10}$
CN-w vs. RHB	0.367
CN-w vs. RRL	0.080
CN-w vs. BHB	0.008
CN-s vs. RHB	0.029
CN-s vs. RRL	0.181
CN-s vs. BHB	0.724
RHB vs. RRL	0.646
RHB vs. BHB	0.249
RRL vs. BHB	0.893

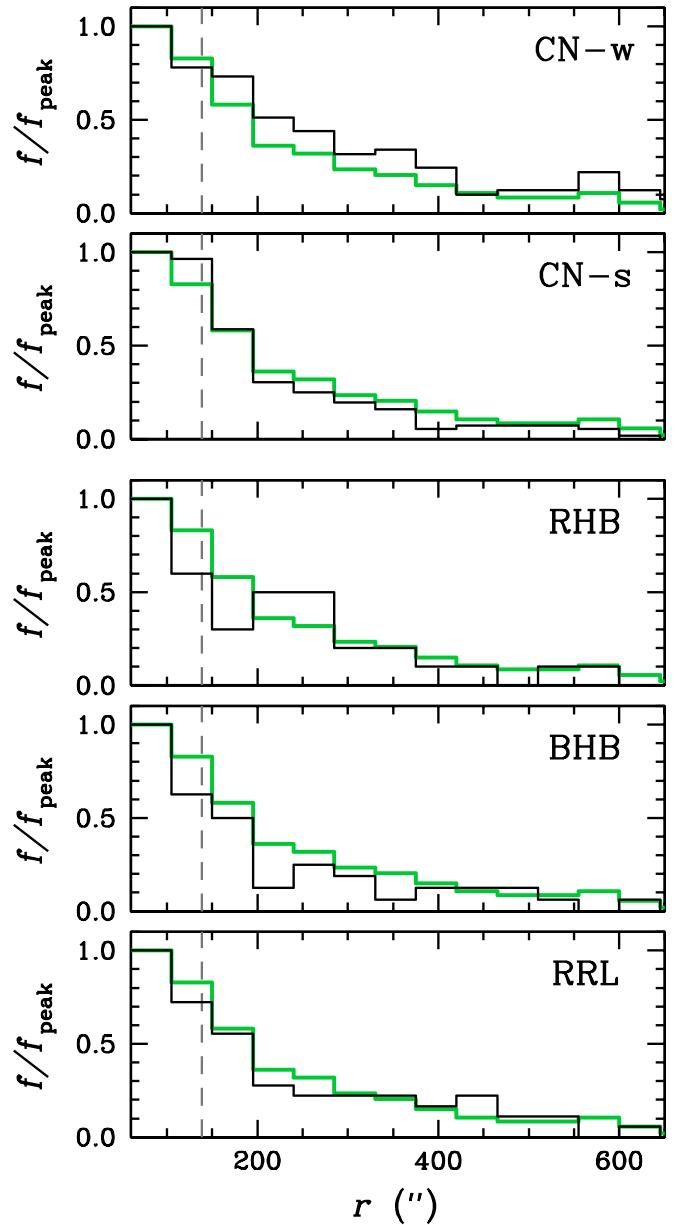
figure, we show histograms of individual populations against the radial distance. In a radial zone from  $r \approx 200''$  and  $300''$ , both the CN-w RGB and RHB stars are slightly overpopulated compared to the average distribution of stars, a strong indication that complete homogenization had not been achieved in the CN-w population.

In Figure 16, we show the moving averages of the RGB stars for the adjacent 50 points of the  $\|nh_{JWL}\|$ ,  $\|ch_{JWL}\|$ , and  $\|cn_{JWL}\|$  indices and  $[Fe/H]_{hk}$ ,  $[N/Fe]$ , and  $[C/Fe]$  against radial distance. The figure shows that M3 exhibits a metallicity gradient against the radial distance, which is an unexpected result. In Section 10, we will discuss that M3 appears to consist of two GCs with slightly different metallicities, namely, C1 and C2. As we will show later, this metallicity gradient in M3 is consistent with the idea that the metal-poor C1 is more centrally concentrated than the metal-rich C2. On the other hand, due to the domination of the CN-s population in the central part of M3, as already shown in Figure 14, the large values of the  $\|nh_{JWL}\|$  and  $\|cn_{JWL}\|$  indices and the small value of the  $\|ch_{JWL}\|$  index and, subsequently, the large value of  $[N/Fe]$  and the small value of  $[C/Fe]$  can be found in the inner part of the clusters, reminiscent of the strong radial CN variation in 47 Tuc (Chun & Freeman 1979). From the radial distance of about  $200''$ – $300''$ , bumps in all color indices and  $[Fe/H]_{hk}$ ,  $[N/Fe]$ , and  $[C/Fe]$  are noticeable due to the existence of the overpopulation of the CN-w RGB stars in this radial zone as we already showed in Figure 15.

### 8.2. Spatial Distributions

The GC morphology can be affected by both internal processes, such as relaxation and cluster rotation, and external processes, such as tidal fields. For example, the ellipticity of GCs can be affected by the dynamical relaxation that leads the stars with excess angular momentum to evaporate (e.g., Shapiro & Marchant 1976). On the other hand, Bianchini et al. (2013) argued that the well-relaxed GC 47 Tuc can be explained very well with its internal rotation.

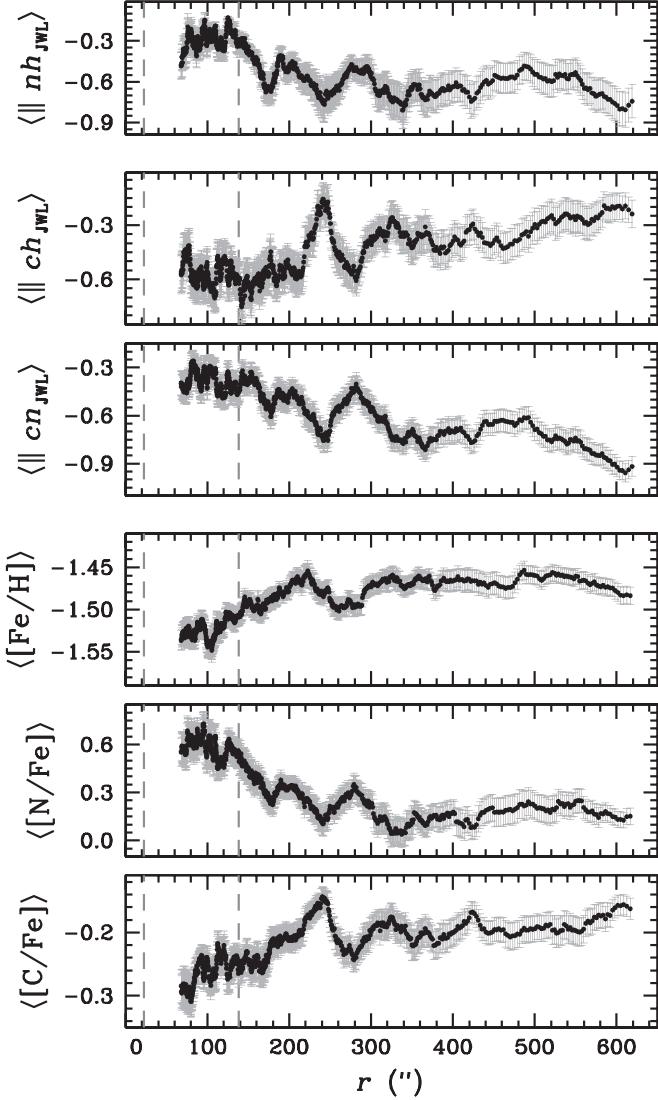
Here, we investigate the morphology of each population. If both populations were dynamically evolved together, the dynamical relaxation and the tidal effect should be similar for both populations. Therefore, it can be thought that the morphological difference between the two populations is due to other effects, such as internal rotation, as we will discuss below. The  $N$ -body simulations by Vesperini et al. (2013), for example, showed that the complete mixing of GCs can be achieved in about at least 20 half-mass relaxation time. As the half-mass relation time for M3 is about 2 Gyr (Harris 1996,



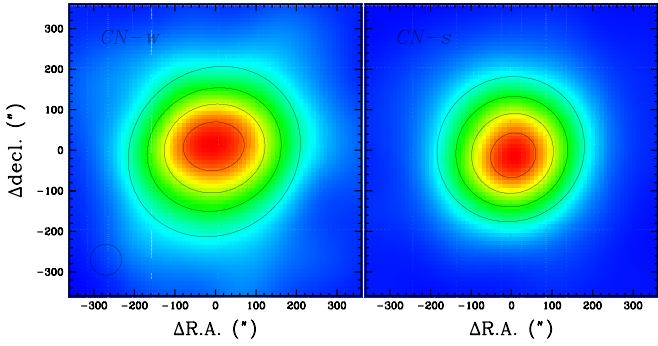
**Figure 15.** Radial distributions of RGB stars in M3. The green solid lines represent the histogram for the whole sample of stars. The vertical gray dashed lines denote the half-light radius of M3. Note that CN-w RGB and RHB populations appear to be slightly overdense in the region between  $r \approx 200''$  and  $300''$ .

2003 version), complete mixing has not been in M3 within the Hubble time, and the two stellar populations in M3 may keep their initial kinematical properties.

In order to investigate the spatial distributions of the CN-w and CN-s RGB populations, we constructed smoothed density maps following the same method described in our previous work (Lee 2015 and see references therein). We present our results in Figure 17, showing that the CN-w population has a more elongated spatial distribution. Evidence for the spatially more elongated structure of the CN-w population can be clearly seen in Figure 18, where we show radial distributions of the axial ratio,  $b/a$ , and the ellipticity,  $e = (1 - b/a)$ , of the CN-w and CN-s populations. As shown in the figure, the ellipticity of the CN-w population is significantly larger than the CN-s

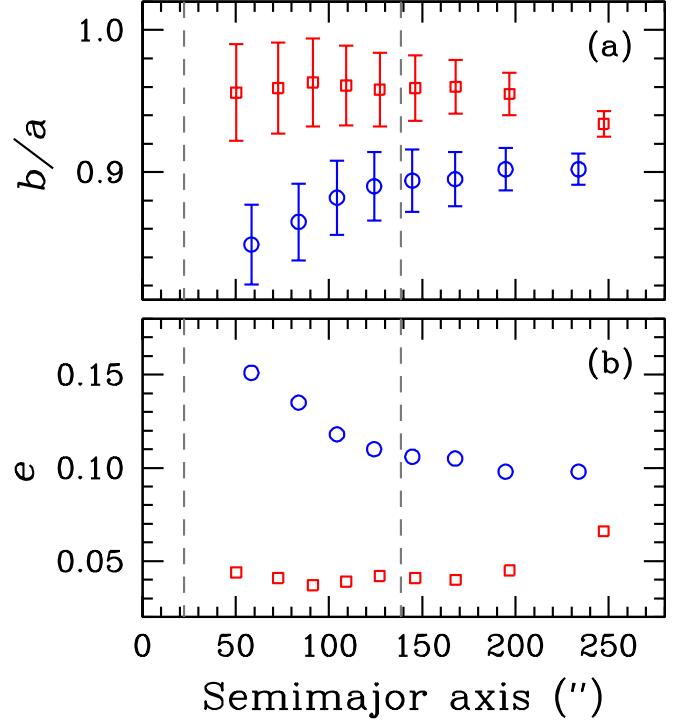


**Figure 16.** Moving averages of the adjacent 50 points for each color index of the RGB stars plotted as functions of radial distance from the M3 center. The vertical thin gray error bars denote the standard error of the mean, and the vertical gray dashed lines denote the core and the half-light radii of M3.



**Figure 17.** Smoothed density maps of the CN-w (left) and CN-s (right) RGB populations. The FWHM of our Gaussian kernel is shown with a circle, and isodensity contour lines for 90%, 70%, 50%, and 30% of the peak values of each population are also shown. Note that the absolute peak values are different for each population.

population up to  $r \approx 1.5r_h$ . At a large radial distance, the axial ratios of the CN-w and CN-s appear to converge on  $b/a \approx 0.90\text{--}0.94$ , consistent with that of Chen & Chen (2010), who



**Figure 18.** The run of the axial ratio  $b/a$  and ellipticity  $e$  of the CN-w (blue) and CN-s (red) populations in M3 against the semimajor axis  $a$ . The vertical gray dashed lines denote the core and half-light radii of M3. The CN-w population is more elongated.

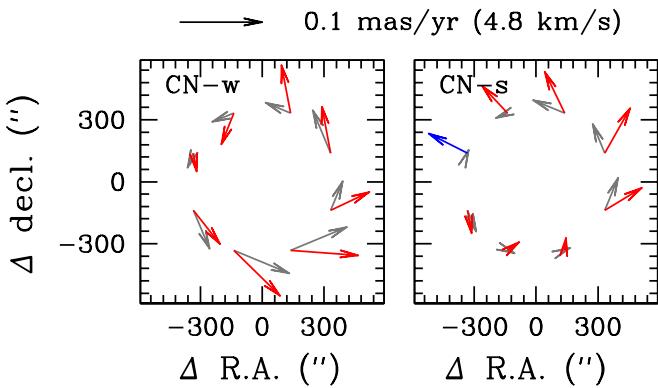
obtained the axial ratio of  $b/a = 0.94 \pm 0.01$  at  $a \approx 11'$  for M3 using the 2MASS point sources.

### 8.3. Internal Rotations

Some GCs, such as 47 Tuc, M15, and  $\omega$  Cen, show substantial internal rotations, reaching maximum line-of-sight velocities of  $\approx 3\text{--}6 \text{ km s}^{-1}$ . Previous studies of M3's rotation using radial velocities showed that M3 has a weak rotational component, on the order of  $1.0\text{--}1.5 \text{ km s}^{-1}$  (Fabricius et al. 2014; Ferraro et al. 2018). Note that the internal rotation of GCs is not an invariant parameter, but it can change with time. For example, numerical simulations by Wang et al. (2016) showed that the internal rotation of GCs can gradually decrease with time due to the loss of angular momentum via two-body relaxation and mass loss, although the timescale for such processes appears to be somewhat uncertain.

In our previous studies, we reported that the structural and kinematical properties between MPs in GCs are decoupled (e.g., see Lee 2015, 2017, 2018, 2020). In the same spirit, we have now investigated the tangential rotation of the CN-w and CN-s populations using the proper motion study of the Gaia DR2 (Gaia Collaboration et al. 2018).

We divided the sphere into eight different slices in a single radial zone of  $1' \leq r \leq 10'$ , and we calculated the mean proper motion vectors in each slice. In Figure 19, we show the tangential rotation of each population. We note that both populations share the same direction in their rotation. On the celestial sphere, they appear to rotate in a counterclockwise sense ( $E \rightarrow N \rightarrow W \rightarrow S \rightarrow E$ ), with the CN-w population having a more well-defined tangential rotation as shown by the gray arrows. Assuming a circular projected rotation for both populations, we calculated mean rotational components of



**Figure 19.** Mean proper motion vectors of eight slices in the radial zone of  $1' \leq r \leq 10'$ . The red arrows denote a counterclockwise rotation ( $E \rightarrow N \rightarrow W \rightarrow S \rightarrow E$ ), while the blue arrow denotes a clockwise rotation at a given position vector. The gray arrows represent the rotating component assuming a circular rotation. The arrow shown above both panels indicates a proper motion of  $0.1 \text{ mas yr}^{-1}$  or  $4.8 \text{ km s}^{-1}$  assuming  $10.2 \text{ kpc}$  for the distance to M3.

$0.068 \pm 0.031 \text{ mas}$  and  $0.033 \pm 0.029 \text{ mas}$  for the CN-w and CN-s populations, respectively, which correspond to  $3.28 \pm 1.51 \text{ km s}^{-1}$  and  $1.60 \pm 1.41 \text{ km s}^{-1}$ , respectively, assuming the distance to M3 of  $10.2 \text{ kpc}$  (Harris 1996). Our results show that the CN-w population exhibits a larger, statistically significant, net tangential rotation than the CN-s population, suggesting that its large ellipticity in the M3 CN-w population is likely induced by its fast internal rotation, as can be seen in M5, M22, and NGC 6752 in our previous studies (Lee 2015, 2017, 2018, and references therein).

## 9. The Double Red-giant Branch Bumps of the M3 CN-w

In the course of low-mass star evolution, RGB stars experience slower evolution and a temporary drop in luminosity when the very thin H-burning shell crosses the discontinuity in the chemical composition and the lowered mean molecular weight left by the deepest penetration of the convective envelope during the ascent of the RGB, the so-called RGBB (e.g., see Cassisi & Salaris 2013). It is well known that the RGBB luminosity increases with helium abundance and decreases with metallicity at a given age (e.g., see Bjork & Chaboyer 2006; Valcarce et al. 2012).

The helium abundance of GC stars cannot be measured directly due to the lack of helium absorption lines. Instead, the RGBB magnitude can be a powerful tool to probe the helium content of a given stellar group at a fixed age and metallicity. For example,  $\Delta Y \approx 0.02$  can be translated into  $\Delta[\text{He/H}] \approx 0.03 \text{ dex}$ , which is almost impossible to detect in high-resolution spectroscopy. On the other hand, this small helium difference can affect the RGBB magnitude by  $\Delta V \approx 0.05 \text{ mag}$ , which can be easily detected, unless there exists no degeneracy with metallicity (e.g., see Lee 2015, 2017, 2018, 2019c; Lagioia et al. 2018; Milone et al. 2018).

As mentioned above, in M3, the extent of the  $\Delta_{\text{CF275W}, \text{F814W}}$  distribution of the CN-w population (i.e., the FG of stars in the frame of the sequential formation of the MPs in GCs) is unusually large (Milone et al. 2017; Lardo et al. 2018). Based on the dependence of synthetic color indices with varying He abundances but a fixed metallicity, Lardo et al. (2018) claimed that the large extent of the  $\Delta_{\text{CF275W}, \text{F814W}}$  distribution of the M3 CN-w RGB stars is due to helium enhancement as large as

$\Delta Y \approx 0.024$  and a small range of nitrogen abundance dispersion within the CN-w population. However, our results show that CN-w and CN-s have identical nitrogen abundance dispersions as shown in Table 5,  $\sigma[\text{N/Fe}] \approx 0.28 \text{ dex}$ .

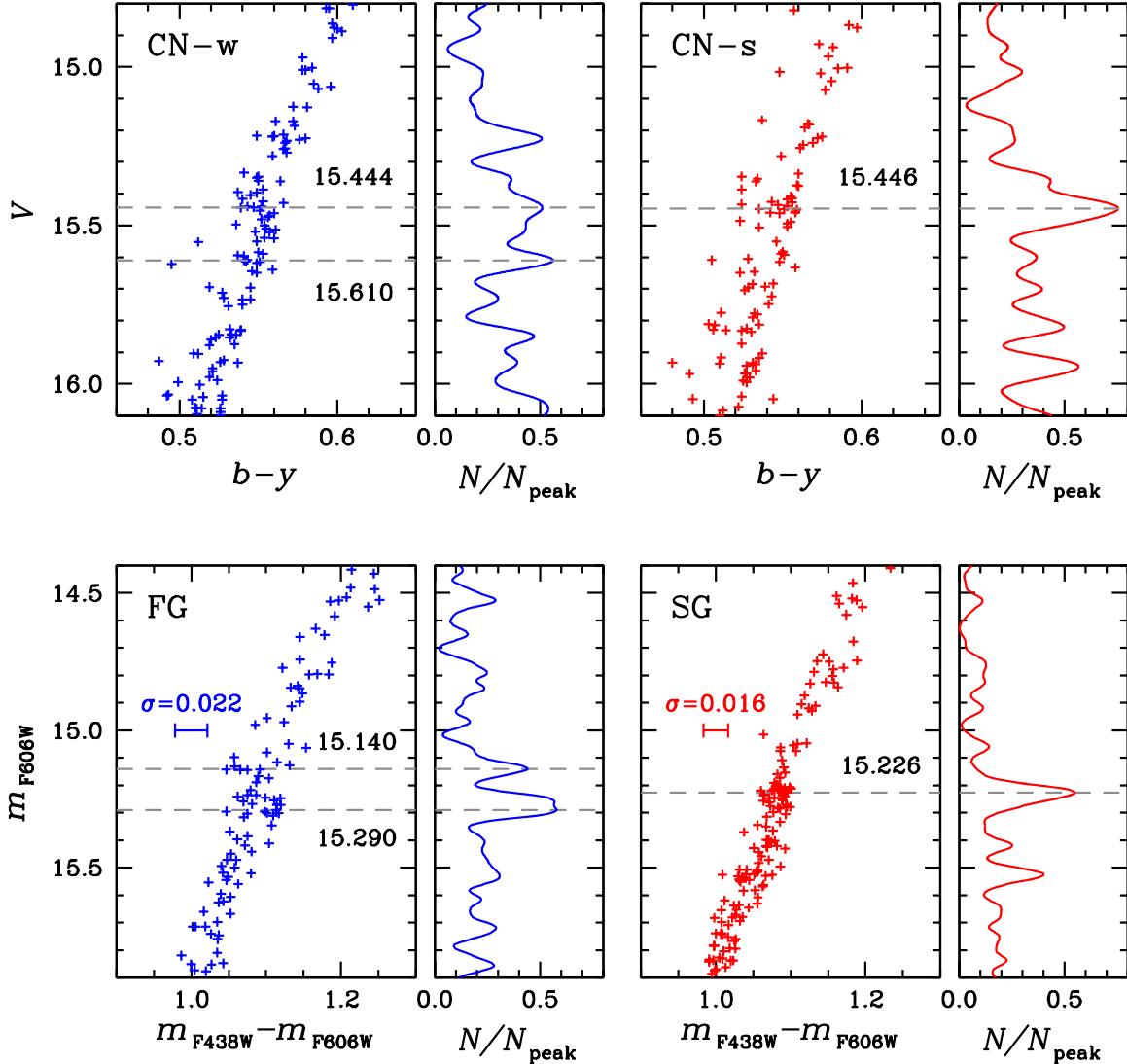
More recently, Tailo et al. (2019) performed synthetic HB model simulations, finding that a model without helium enhancement in the FG of stars can best match with the observed M3 HB type, RRL period distribution, and the MS distribution. However, they noted that they were not able to address the color spread of the M3 FG star satisfactorily. Tailo et al. (2019) also did not consider the metallicity spread among the M3 FG stars.

The large helium spread of  $\Delta Y \approx 0.024$  in the FG of stars without any perceptible CNO abundance spreads by Lardo et al. (2018) is difficult to reconcile with the current understanding of the GC formation scenario. As Lardo et al. (2018) argued, their results imply that  $p-p$  chain hydrogen burning induced helium enhancement within the FG of stars. If so, the FG of such a GC requires a very long timescale, and it does not conform to the currently believed sequential two-step formation scenario with a relatively short timescale,  $\lesssim 100 \text{ Myr}$ , or the small age dispersion in M3 seen from the HST observations,  $\pm 0.5 \text{ Gyr}$  (Dotter et al. 2010).

Here, we investigate the helium content and metallicity of the M3 CN-w population using the RGBB. From a stellar evolution perspective, the RGBB is less complicated than the HB and AGB. For example, the HB morphology of GCs depends not only on the helium abundance but also on the mass loss during the RGB phase. At the same time, typical GC HB stars are too faint to perform populational tagging from high-resolution spectroscopy. Tagging HB stars from low-resolution spectroscopy or from narrowband photometry, such as ours, would also be difficult for hot BHB stars, where the most informative diatomic molecular bands are absent. In addition to statistical fluctuations due to the small number statistics of AGB stars, the problem unsolvable with a populational number ratio study of AGB stars is the missing populations that eventually evolve into the AGB-manqué (e.g., see Figures 20 and 24 of Lee 2017). On the other hand, populational tagging for entire RGBB populations can be readily done with our new color indices.

### 9.1. Double RGBBs of the M3 CN-w

As we mentioned above, the extended and tilted RGBB of the CN-w population is one of the interesting features of M3. In our current work, we investigate the differential luminosity functions (LFs) of each population to understand the underlying stellar populations in M3. In Figure 20, we show the CMDs and LFs of both populations, finding  $V_{\text{bump}} = 15.610 \pm 0.025 \text{ mag}$  and  $15.444 \pm 0.025 \text{ mag}$  for the CN-w population and  $15.446 \pm 0.025 \text{ mag}$  for the CN-s population. Using all RGB stars from both populations, we obtained  $15.445 \pm 0.02 \text{ mag}$ , and our RGBB  $V$  magnitude is in excellent agreement with that of Ferraro et al. (1997), who reported  $V_{\text{bump}} = 15.45 \pm 0.05 \text{ mag}$  for M3. The double RGBBs of the M3 CN-w population from our ground-based observations can also be supported by the HST photometry by Nardiello et al. (2018). In Figure 20, we also show CMDs and differential LFs of the FG and SG groups. We note again that the FG and SG groups correspond to our CN-w and CN-s populations, respectively. The figure clearly shows that the FG group has discrete double RGBBs, consistent with our results for the



**Figure 20.** (Top panels) CMDs and differential LFs of M3 from our observations. The blue and red colors denote the CN-w and CN-s. The gray dashed lines represent the  $V$  magnitude of the RGBBs. (Bottom panels) CMDs and differential LFs of M3 using the HST photometry by Nardiello et al. (2018). The  $m_{F606W}$  values for RGBB are shown with gray dashed lines. The RGB widths around the RGBB region are also shown with error bars.

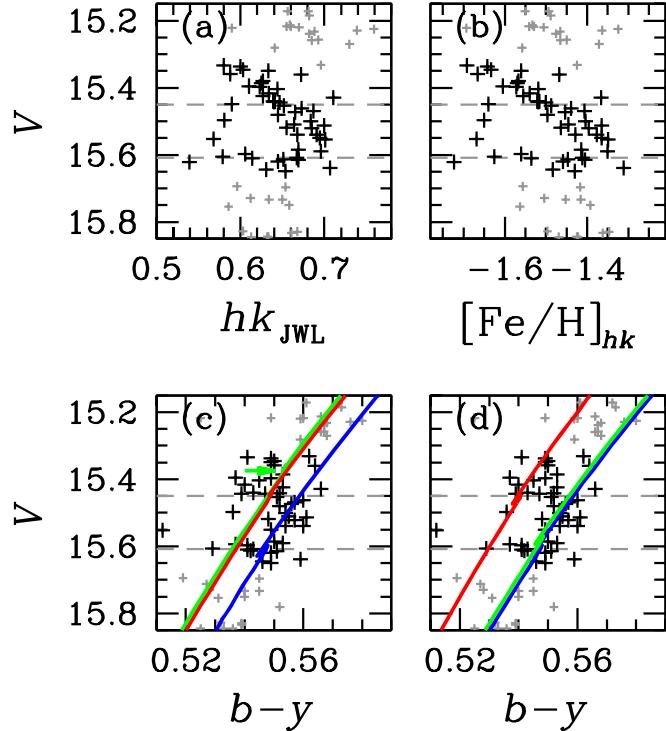
CN-w population. Interestingly, the magnitude difference between the two RGBBs of the HST photometry is  $0.150 \pm 0.020$  mag, in excellent agreement with our ground-based observation,  $0.166 \pm 0.035$  mag. We derived fourth-order polynomial fits to measure the RGB width of the FG and SG groups. We calculated scatters around the fitted lines, and we obtained  $\sigma(m_{F438W} - m_{F606W}) = 0.022$  mag for the FG and  $0.016$  mag for the SG. The broader RGB width of the FG group is consistent with our previous result shown in Figure 7.

We note that the RGBB magnitude of the SG group from the HST photometry does not agree with our results. The CN-s RGBB magnitude is almost identical to that of the bright CN-w RGBB. However, the RGBB magnitude of the SG group from the HST observations is  $0.086$  mag fainter than the bright RGBB of the FG group. The cause of this discrepancy is unclear. (1) It could originate from the different behavior in different photometric systems with significantly different bandwidths (i.e., the Strömgren  $y$  versus the HST  $F606W$ ) for certain chemical abundances. But our calculations show that the difference in the filter passbands between the Strömgren  $y$

and the HST  $F606W$  can only explain the magnitude difference of  $0.005$  mag. (2) Perhaps there might exist a radial gradient of the CN-s RGBB in the sense that the CN-s RGBB gets fainter with decreasing radial distance, probably due to the increasing metallicity of the main body of the CN-s. We examined our photometry within the radial zone of  $1' \leq r \leq 10'$ , but we did not find any hint of a radial gradient of the CN-s RGBB magnitude. It is interesting to note that the double RGBBs cannot be seen in the CN-s population, which we will discuss later. Future studies of this issue would be welcome.

At face value, if we assume that only helium abundance affects the RGBB magnitude, the difference in the magnitude between the faint CN-w RGBB and CN-s RGBB,  $\Delta V = 0.164 \pm 0.035$  mag, in our observations can be translated into a helium spread of  $\Delta Y = 0.066 \pm 0.014$ , in the sense that the CN-s is significantly enhanced in helium abundance, which does not seem very plausible for the M3 HB morphology.

This is an apparently straightforward conclusion, but the interpretation of the detailed substructures of the RGBB may not be so simple. The double RGBBs in the CN-w, with the magnitude



**Figure 21.** (a)  $hk_{JWL}$  vs.  $V$  CMD around the RGBB region of the CN-w population. The gray dashed lines indicate the  $V$  magnitudes of the double RGBBs, and the black plus signs denote the stars near the double RGBBs. (b) Same as (a), but for  $[Fe/H]_{hk}$ . (c)  $(b-y)$  vs.  $V$  CMD. The blue solid line represents the model isochrone for  $Y = 0.25$  and  $[Fe/H] = -1.45$  dex, while the red solid line represents that of  $[Fe/H] = -1.50$  dex and  $Y = 0.29$ . The green solid line and the arrow indicate the isochrone and the RGBB  $V$  magnitude for  $[Fe/H] = -1.45$  dex and  $Y = 0.33$ . (d) The blue and red solid lines represent the model isochrones for  $Y = 0.25$  and  $[Fe/H] = -1.45$  and  $-1.60$  dex, respectively, while the green solid line represents that for  $Y = 0.26$  and  $[Fe/H] = -1.45$  dex.

difference between the two peaks of  $\Delta V_{\text{bump}} \approx 0.16$  mag, can arise from at least three different physical sources: the differential reddening effect, and the helium and metallicity spreads.

### 9.2. Differential Reddening

Differential reddening across the M3 field could lead to apparent double RGBBs, although no such effect has been reported so far. We can argue in two ways against double RGBBs due solely to differential reddening. First, if the magnitude difference between the two RGBBs in the CN-w is due to differential reddening, the spread in the interstellar reddening becomes  $\Delta E(B-V) \approx \Delta V_{\text{bump}}/3.1 \approx 0.053$  mag, which becomes  $\Delta E(b-y) \approx 0.04$  mag, assuming  $E(b-y) = 0.74 \times E(B-V)$ . In other words, the intrinsic  $(b-y)$  color (i.e., the reddening-corrected color) of the lower part of RGBB at  $(b-y) \approx 0.55$  and  $V \approx 15.6$  mag in Figure 21 should be 0.04 mag bluer at  $V = 15.44$  mag, i.e.,  $(b-y)_0 = 0.51$  mag. If so, the width of the dereddened CN-w RGB population becomes larger,  $\Delta(b-y) \gtrsim 0.03$  mag, than the reddened RGB width,  $\Delta(b-y) \approx 0.02$  mag.

In panel (a) of Figure 21, we show the  $hk_{JWL}$  versus  $V$  CMD. The interstellar reddening of the  $hk_{JWL}$  index is given by  $E(hk_{JWL}) = -0.16 \times E(b-y) = -0.12 \times E(B-V)$  (see Supplementary Information of Lee et al. 2009a, and references therein). If differential reddening is responsible for the double RGBBs, then the dereddened  $hk_{JWL}$  color of the faint CN-w RGBB, 0.70, becomes about 0.05 mag redder, making the

width of  $hk_{JWL}$  for the CN-w stars too broad for being a simple stellar population. Therefore, it is most likely that other effects are responsible for the double RGBBs in the CN-w population, as discussed below.

### 9.3. Monte Carlo Simulations for the CN-w Double RGBBs

In order to understand the origin of the double RGBBs with a slightly broader RGB width in the CN-w population, we performed Monte Carlo simulations by constructing evolutionary population synthesis (EPS) models similar to those of our previous studies (e.g., Lee & Carney 1999b; Lee 2015, 2019b).

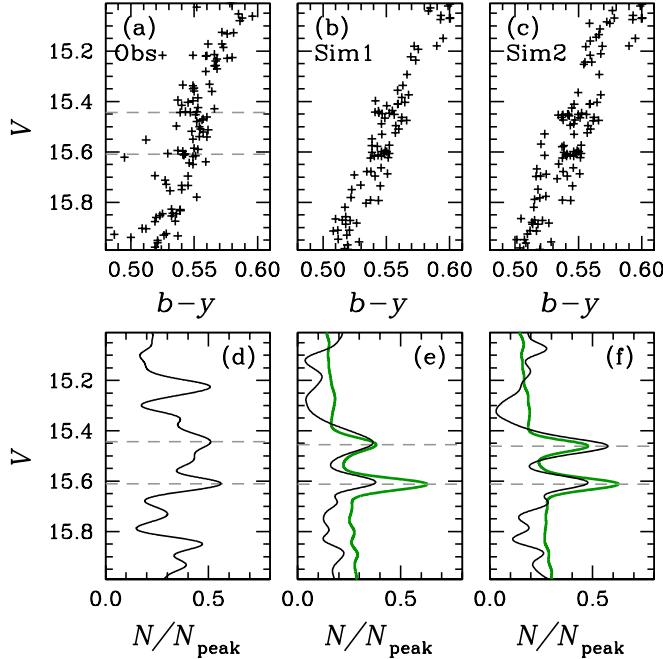
We have shown that at a fixed age, the RGBB magnitude is sensitively dependent both on the helium abundance and on the metallicity. In Figure 21, panels (c)–(d), we present comparisons of model isochrones of Valcarce et al. (2012) to our observations. Note that the  $V$ -magnitude levels of the RGBB of the model isochrones by Valcarce et al. (2012) are about 0.40 mag brighter than our observations. Recall that our RGBB  $V$  magnitude for M3 is in excellent agreement with that of Ferraro et al. (1997). In Figure 21, we add an additional 0.40 mag to the model isochrones to match the location of the RGBB, and the figure suggests that both helium enhancement and a metallicity spread could explain the double RGBBs in the CN-w population. Using model isochrones by Valcarce et al. (2012), we constructed two sets of EPS models: (1) an EPS model with  $\Delta Y = 0.04$  and  $\Delta[Fe/H] = -0.05$  dex to reproduce the bright RGBB (Sim1, hereafter), and (2) an EPS model with  $\Delta[Fe/H] = -0.15$  dex without the helium enhancement to reproduce the bright RGBB (Sim2, hereafter). Finally, we convolved our observational uncertainties as shown in Figure 3 to generate synthetic CMDs. We performed 1000 simulations for both sets to avoid artifacts that arise from small sample sizes.

In our simulations, we adopt a fixed age for both Sim1 and Sim2, 12.5 Gyr. At a fixed metallicity and helium abundance, the RGBB luminosity decreases with age, and, therefore, the magnitude difference between the two RGBBs can be slightly tweaked with ages between the two presumed stellar populations. However, previous HST photometry suggested that M3 does not appear to show any perceptible age spread,  $12.5 \pm 0.50$  Gyr (Dotter et al. 2010).

### 9.4. Helium Abundance: Sim1

In Figure 21(c), we show model isochrones for Sim1:  $[Fe/H] = -1.45$  dex,  $Y = 0.25$ , and  $[Fe/H] = -1.50$  dex,  $Y = 0.29$  (Valcarce et al. 2012). Note that a small metallicity spread is required to reproduce the observed RGB width. Our input models can explain the  $V$  magnitudes of the double RGBBs satisfactorily, but the RGB width between the two isochrones is slightly narrower than our observations. To reproduce the RGB width of the isochrone for  $[Fe/H] = -1.50$  dex,  $Y = 0.29$  without metallicity and age spreads, a model with  $Y \approx 0.33$  is required as shown in the figure. However, its RGBB  $V$  magnitude is about 0.08 mag brighter than that of the bright RGBB in our observations. An increase in age of 1.0–1.5 Gyr can reproduce the observed bright RGBB magnitude, but it is against the very narrow or no age spread of M3 from the HST photometry (Dotter et al. 2010).

In Figure 22(b) and Table 8, we show the results from our Monte Carlo simulations. In the figure, we show a synthetic CMD returned from one particular simulation and its



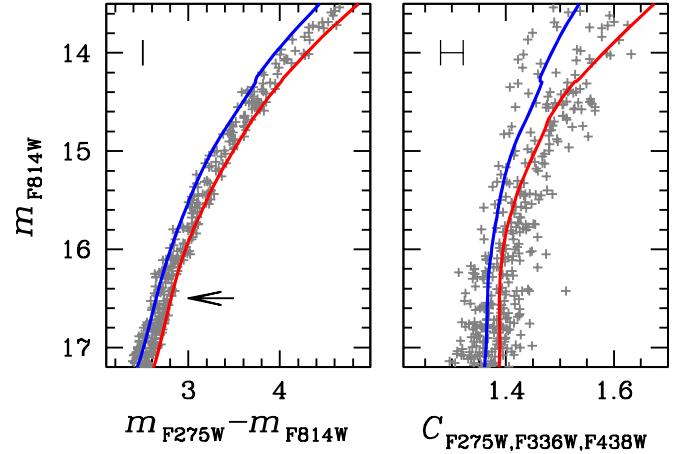
**Figure 22.** (a) CMD of the M3 CN-w RGB stars around the RGBB. The gray dashed lines indicate the double RGB peaks. (b) Synthetic CMD for Sim1,  $\Delta Y = 0.04$  and  $\Delta[\text{Fe}/\text{H}] = 0.05$  dex. (c) Synthetic CMD for Sim2,  $\Delta[\text{Fe}/\text{H}] = 0.15$  dex. (d) Differential LF of (a). (e) The black solid line shows the differential LF of (b), while the green solid line, and the gray dashed lines represent the mean differential LF of 1000 trials and the  $V$  magnitude of the double RGBBs. (f) Same as (e), but for (c).

differential LF. As mentioned earlier, Sim1 shows a slightly narrower RGB width than our observations. We also show the average LF from 1000 simulations, showing two RGBBs whose  $V$  magnitudes are consistent with our observations.

In terms of the  $V$  magnitudes of the double RGBBs and the CN-w RGB width, the helium-enhanced model can reproduce our observations. However, this large amount of helium enhancement in the CN-w population is hard to reconcile with the M3 HB morphology (see also Tailo et al. 2019). Also, assuming the metallicity is a chronometer, it is difficult to understand the chemical evolution from stars having low metallicity with high helium abundance ( $[\text{Fe}/\text{H}] = -1.50$ ,  $Y = 0.29$ ) to stars having slightly high metallicity but significantly low helium abundance ( $[\text{Fe}/\text{H}] = -1.45$ ,  $Y = 0.25$ ). The helium abundance can be slightly decreased by dilution with pristine gas but the metallicity should be decreased, too. Without invoking a metallicity spread, it requires a large amount of helium enhancement,  $\Delta Y = 0.08$ , and a large age spread, more than 1 Gyr, which are against the HB morphology (Tailo et al. 2019) and the undetectably small age spread of M3 (Dotter et al. 2010). Therefore, variations in the helium abundance do not appear to be responsible for the double RGBBs of the CN-w population.

### 9.5. Metallicity: Sim2

The M3 CN-w population has a metallicity spread of  $\approx 0.07$ – $0.10$  dex from high-resolution spectroscopy (Sneden et al. 2004; Mészáros et al. 2015), and we have derived a bimodal metallicity distribution with  $\Delta[\text{Fe}/\text{H}] \approx 0.15$  dex from  $hk_{\text{JWL}}$  photometry. In Figure 21(c), we show the model isochrones for Sim2:  $[\text{Fe}/\text{H}] = -1.60$  and  $-1.45$  dex with a



**Figure 23.** HST pseudo-color CMDs for the M3 FG RGB stars using photometry by Nardiello et al. (2018). The blue and red solid lines are for the Dartmouth isochrone with  $[\text{Fe}/\text{H}] = -1.60$  and  $-1.45$  dex, respectively. The arrow in the left panel indicates the  $m_{\text{F}814}$  magnitude where  $W_{\text{F}275\text{W}-\text{F}814\text{W}}^{\text{IG}}$  is measured, 2 mag brighter than the main-sequence turn-off. The thin error bars indicate the contribution of the spread in nitrogen abundance by  $\Delta[\text{N}/\text{Fe}] = 0.1$  dex.

**Table 8**  
The  $V$  Magnitudes of RGBBs

	Bright RGBB	Faint RGBB
Observations	15.444	15.610
Sim1	15.455	15.612
Sim2	15.461	15.612

fixed helium abundance,  $Y = 0.25$ . The figure shows that our input models can explain the double RGBBs and the RGB width reasonably well. In the figure, we also show the model isochrone with a slight helium enhancement by  $\Delta Y = 0.01$  at the metal-rich regime. It does not disagree with our observations but suggests that helium enhancement may not be the essential factor for understanding the double RGBBs in M3.

In Figure 22(c) and Table 8, we show the results of our Monte Carlo simulations, indicating that both the  $V$  magnitudes of the double RGBBs and RGB width are consistent with our observations. Therefore, we suggest that the discrete bimodal metallicity distribution must be the main reason for the extended and tilted RGBB and the width of the RGB of the CN-w population.

Finally, we also compare the HST photometry by Nardiello et al. (2018) to model isochrones by Dotter et al. (2008), because Valcarce et al. (2012) does not provide magnitudes for the passbands required to construct the chromosome map. In Figure 23, we show the Dartmouth model isochrones with  $[\text{Fe}/\text{H}] = -1.60$  and  $-1.45$  along with the M3 FG RGB stars using the results of Nardiello et al. (2018). Both model isochrones have  $Y = 0.25$  and  $[\alpha/\text{Fe}] = 0.25$  dex and the distance modulus of 15.07 mag and the offset value of  $-0.10$  mag in the  $(m_{\text{F}275\text{W}} - m_{\text{F}814\text{W}})$  color were assumed. Milone et al. (2017) measured the RGB width of the  $(m_{\text{F}275\text{W}} - m_{\text{F}814\text{W}})$  color at 2 mag brighter than the main-sequence turn-off, obtaining  $W_{\text{F}275\text{W}-\text{F}814\text{W}}^{\text{IG}} = 0.244 \pm 0.014$  for the M3 FG population. As shown in the figure, model isochrones with different metallicities can nicely explain the M3 FG’s  $(m_{\text{F}275\text{W}} - m_{\text{F}814\text{W}})$  and  $[(m_{\text{F}275\text{W}} - m_{\text{F}336\text{W}}) - (m_{\text{F}336\text{W}} - m_{\text{F}438\text{W}})]$  widths without invoking helium enhancement. We also calculate the contribution

of the nitrogen abundance enhancement to  $W_{\text{F275W}-\text{F814W}}^{1G}$ , and our result confirms that nitrogen abundance does not affect  $W_{\text{F275W}-\text{F814W}}^{1G}$ , as Milone et al. (2017, and references therein) suggested. On the other hand, the  $[(m_{\text{F275W}} - m_{\text{F336W}}) - (m_{\text{F336W}} - m_{\text{F438W}})]$  color is sensitively dependent on nitrogen abundance as the F336W passband contains the NH band at  $\lambda 3360$ .

We conclude that the M3 CN-w population has its double RGBBs because of the bimodal metallicity distribution.

## 10. M3 as the Merger Remnant of Two Globular Clusters

### 10.1. Discrete Bimodal $[\text{Fe}/\text{H}]$ Distributions and a Sequential Formation Scenario

As we presented in Figure 7 and Table 4, both the CN-w and CN-s populations in M3 show discrete bimodal  $[\text{Fe}/\text{H}]_{hk}$  distributions,  $-1.58 \pm 0.05$  (CN-w:SP1) and  $-1.44 \pm 0.05$  (CN-w:SP2) with the subpopulational number ratio of 36:64 ( $\pm 4$ ) for the CN-w population and  $-1.59 \pm 0.05$  (CN-s:SP1) and  $-1.47 \pm 0.03$  (CN-s:SP2) with the subpopulational number ratio of 17:83 ( $\pm 2$ ) for the CN-s population. In each population, the mean metallicities of the two subpopulations do not overlap within a  $1\sigma$  level.

Discrete bimodal metallicity distributions in both populations cannot be produced by incomplete mixing in their protocluster clouds, and they require episodic star formations. Furthermore, we find it impossible to explain our results of the elemental abundances in M3 with a sequential star formation scenario (see also Lee 2015). Assuming our CN-w population is the FG and the CN-s is the SG, the sequential formation scenario for M3 may be described in the following steps: (1) the formation of the CN-w:SP1 with high  $[\text{C}/\text{Fe}]$  and low  $[\text{N}/\text{Fe}]$  abundances; (2) Type II supernova (SNe II) explosion that enriched the metallicity of CN-w:SP2 without changing the  $[\text{C}/\text{Fe}]$  and  $[\text{N}/\text{Fe}]$  abundances; (3) the formation of the CN-s:SP1 out of the gas experienced CN-cycle hydrogen burning and an unknown process that decreased metallicity from  $[\text{Fe}/\text{H}] = -1.45$  to  $-1.60$ ; and (4) SN II explosion that enriched the metallicity of the CN-s:SP2 without changing the  $[\text{C}/\text{Fe}]$  and  $[\text{N}/\text{Fe}]$  abundances.

Alternative scenarios with some variations could also be possible, but we expect at least four difficulties in any sequential formation scenarios: (1) the entire formation time-scale should be no larger than 0.5 Gyr, as the HST photometry of Dotter et al. (2010) suggested. (2) It requires some unknown processes that reduce metallicity before the formation of the CN-s population. (3) To retain ejecta from energetic SN II explosions requires a very massive system for M3 in the past, most probably the relic of a more massive primeval dwarf galaxy (e.g., Lee et al. 2009a; Lee 2015). (4) If the metal-rich components were the SGs of M3, they tend to be more centrally concentrated (e.g., see Bellini et al. 2009, for  $\omega$  Cen), which is against our results as discussed below.

### 10.2. A Merger Scenario

A more simple and plausible explanation can be found in a merger of two GCs, à la M22. We have shown that the peak metallicity values and metallicity dispersions coincide both in the CN-w and CN-s,  $[\text{Fe}/\text{H}]_{hk} \approx -1.60$  and  $-1.45$ , with the metal-rich subpopulations being the major components in both populations. If we divide individual subpopulations with metallicity, the formation of M3 can be depicted in a simple way: the

merger of two GCs, most probably, in a dwarf galaxy environment where the GC merger rates are significantly higher than those of our Galaxy, because the relative velocity of the two GCs in a dwarf galaxy is smaller than the velocity dispersions (e.g., Thurl & Johnston 2002; Lee et al. 2009a; Bekki & Yong 2012; Lee 2015; Gavagnin et al. 2016). Let us suppose that individual subpopulations can be rearranged as follows: (1) the C1 system with  $[\text{Fe}/\text{H}] = -1.60$ , which is composed of the CN-w:SP1 and the CN-s:SP1. In C1, the CN-w:SP1 corresponds to the C1:CN-w and the CN-s:SP1 to the C1:CN-s, with the number ratio of  $n(\text{CN-w}):n(\text{CN-s}) = 81:19 (\pm 5)$ . (2) The C2 system with  $[\text{Fe}/\text{H}] = -1.45$ , which is composed of the CN-w:SP1 and the CN-s:SP1. Similar to C1, in C2, the CN-w:SP2 corresponds to the C2:CN-w and the CN-s:SP2 to the C2:CN-s, with the number ratio of  $43:57 (\pm 4)$ . We caution that our subpopulational number ratios for C1 and C2 could be different from their true values, because we relied on the stars located at  $r \geq 1''$  and  $\sigma(hk) \leq 0.01$  mag. The latter constraint also prefers stars in the outer region, where the degree of crowdedness is less severe, resulting in small measurement errors.

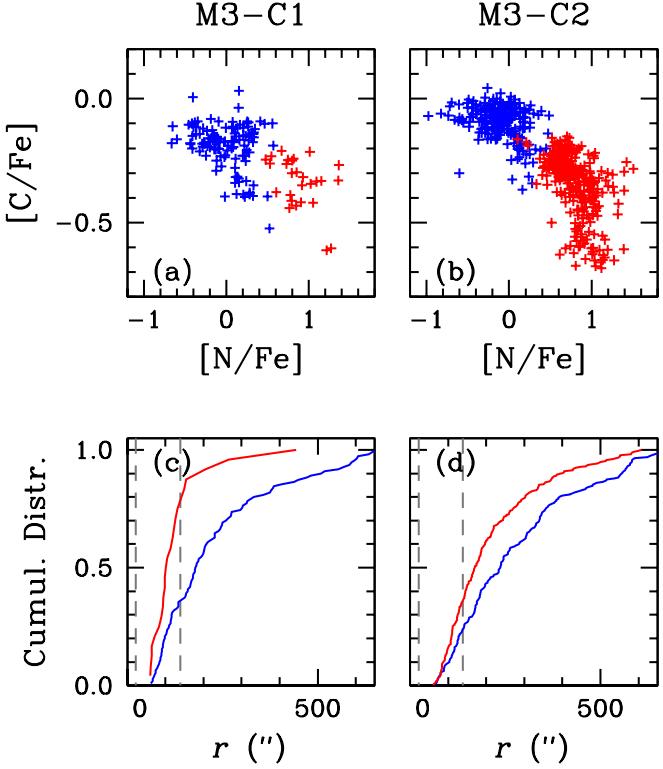
Recently, Milone et al. (2020) studied GC MPs both in our Galaxy and the MC, finding that the FG fractions of MC clusters are significantly higher for their masses, from about 50% up to 80%, than Galactic GCs, most likely reflecting environmental effect. The current mass of M3 from Baumgardt & Hilker (2018) is  $\approx 3.9 \times 10^5 M_{\odot}$ . If we adopt the number ratio between the two systems,  $n(\text{C1}):n(\text{C2}) = 23:77 (\pm 2)$ , then the total masses for C1 and C2 become  $\approx 0.9 \times 10^5 M_{\odot}$  and  $3.0 \times 10^5 M_{\odot}$ , respectively. If our CN-w is the same as the FG of Milone et al. (2020), the fractions of the FG of C1 and C2, 81% and 43%, respectively, appear to follow the trend of the MC GCs (see their Figures 7 and 12).

In Figures 24(a)–(b), we show the  $[\text{C}/\text{Fe}]$  versus  $[\text{N}/\text{Fe}]$  anticorrelations for C1 and C2. C1 shows a less extended C–N anticorrelation than C2, which may be consistent with the fact that less massive GC systems show less extended Na–O anticorrelations (e.g., Carretta et al. 2009).

Being a less massive system, the internal helium enhancement of C1 must have been smaller than that of C2. At the current mass of C1, the results of Milone et al. (2018) suggest that the helium enhancement of the C1:CN-s (=CN-s:SP1) subpopulation is almost nil, which may explain the absence of double RGBBs in our CN-s population as shown in Figure 20. Without helium enhancement, the RGB  $V$  magnitudes of the CN-w:SP1 and CN-s:SP1 would be the same, consistent with our observations,  $V = 15.444$  and  $15.446$  mag. In addition, the fraction of the metal-poor component of the CN-s population, the C1:CN-s (=CN-s:SP1), appears to be too low to produce conspicuous double RGBBs in the CN-s population.

Figures 24(c)–(d) show CRDs of individual subpopulations in C1 and C2. In both systems, the CN-s subpopulations are more centrally concentrated than the CN-w, consistent with those of other Galactic GCs. In Table 9, we show the  $p$  values returned from the K-S tests for the CRDs of C1 and C2. The CN-w:SP1 (i.e., the C1:CN-w) and CN-s:SP2 (the C2:CN-s) show a high  $p$  value. It is not clear whether they are physically linked together or their similar CRDs are pure coincidence. The  $p$  values for other combinations are very low, and their CRDs are probably not related to each other.

We also examine the proper motion of C1 and C2, presenting our results in Figure 25. C1 does not appear to show any aligned projected rotation, and it appears to be a random-motion-dominated system, while C2 shows a well-aligned projected



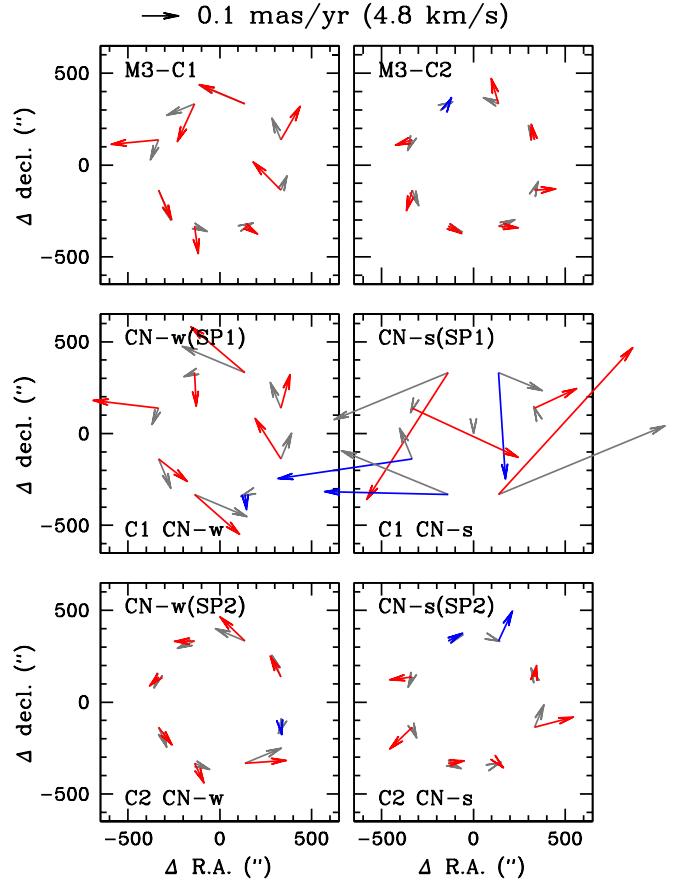
**Figure 24.** (a)  $[C/Fe]$  vs.  $[N/Fe]$  distribution for C1. The blue and red plus signs denote the C1:CN-w (i.e., CN-w:SP1) and C1:CN-s (i.e., CN-s:SP1), respectively. (b)  $[C/Fe]$  vs.  $[N/Fe]$  distribution for C2. The blue and red plus signs denote the C2:CN-w (i.e., CN-w:SP2) and C2:CN-s (i.e., CN-s:SP2), respectively. (c) The CRDs of the C1:CN-w (blue) and C1:CN-s (red). The vertical gray dashed lines denote the core and half-light radii of M3. (d) The CRDs of the C2:CN-w (blue) and C2:CN-s (red).

**Table 9**  
 $p$  Values Returned from the K-S Tests for CRDs of M3 C1 and C2

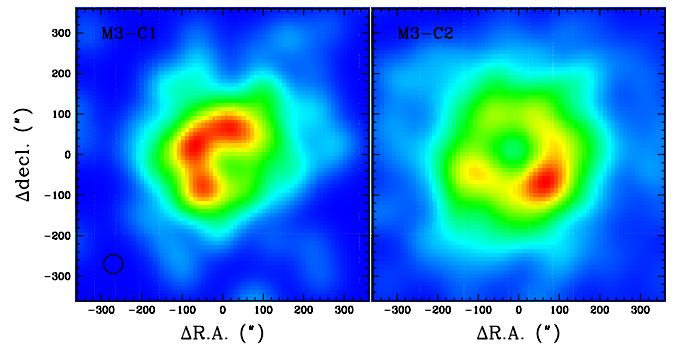
	$p$ Value
CN-w(SP1) vs. CN-w(SP2)	0.005
CN-s(SP1) vs. CN-s(SP2)	$2.3 \times 10^{-7}$
CN-w(SP1) vs. CN-s(SP1)	$6.8 \times 10^{-5}$
CN-w(SP1) vs. CN-s(SP2)	0.852
CN-w(SP2) vs. CN-s(SP1)	$2.4 \times 10^{-8}$
CN-w(SP2) vs. CN-s(SP2)	$2.4 \times 10^{-5}$
CN-w vs. CN-w(SP1)	0.548
CN-w vs. CN-w(SP2)	0.009
CN-w vs. CN-s(SP1)	$7.8 \times 10^{-6}$
CN-w vs. CN-s(SP2)	0.010
CN-s vs. CN-w(SP1)	0.004
CN-s vs. CN-w(SP2)	$1.2 \times 10^{-14}$
CN-s vs. CN-s(SP1)	0.015
CN-s vs. CN-s(SP2)	0.001

circular motion. Our results strongly suggest that C1 and C2 have different internal motions and they are not kinematically homogenized.

Finally, we examined the spatial distributions, and we show the smoothed density map of C1 and C2 in Figure 26. The figure clearly shows that the centers of C1 and C2 are most likely different. We tried to derive their centers but, unfortunately, the coordinates of each center without stars in the central part of M3 are significantly different depending on the methods that we adopted in our previous study (Lee 2015)



**Figure 25.** Mean proper motion vectors of eight slices in the radial zone of  $1' \leq r \leq 10'$ . The red arrows denote a counterclockwise rotation ( $E \rightarrow N \rightarrow W \rightarrow S \rightarrow E$ ), while the blue arrow denotes a clockwise rotation at a given position vector. The gray arrows represent the rotating component assuming a circular rotation. The arrow shown above the panels indicates a proper motion of  $0.1 \text{ mas yr}^{-1}$  or  $4.8 \text{ km s}^{-1}$  assuming  $10.2 \text{ kpc}$  for the distance to M3. (Top panels) M3-C1 and M3-C2. (Middle panels) C1:CN-w (i.e., CN-w:SP1) and C1:CN-s (i.e., CN-s:SP1). (Bottom panels) C2:CN-w (i.e., CN-w:SP2) and C2:CN-s (i.e., CN-s:SP2).



**Figure 26.** Smoothed density maps of the C1 (left) and C2 (right) RGB populations. The FWHM of our Gaussian kernel are shown with a circle. Note that the absolute peak values are different.

and failed to yield reliable coordinates. Note that Figure 26 can be somewhat misleading. Due to the lack of stars in  $r < 1'$ , asymmetries both in C1 and C2 are exaggerated during our smoothing process. In the future, it would be very desirable to investigate the central part of M3 using our photometric system, which will reveal the structural properties of C1 and C2.

## 11. Summary and Conclusion

As part of our ongoing effort to investigate MPs in GCs based on homogeneous photometry, we developed a new filter, JWL34, that can measure the NH bands at  $\lambda 3360$  of cool stars and a new photometric index,  $nh_{JWL}$ , a photometric measure of nitrogen abundance. Our new photometric system combined with the robust and self-consistent theoretical fine model grids for various parameter sets using synthetic spectra allows us to measure key elements in stellar populations, [Fe/H], [C/Fe], and [N/Fe], even in the extremely crowded fields, such as the central part of GCs, where traditional spectroscopic observations cannot be performed. Unlike other broadband photometric systems, such as those used in the HST photometry, our photometric indices,  $hk_{JWL}$ ,  $nh_{JWL}$ , and  $ch_{JWL}$ , suffer a lesser degree of degeneracy of various elemental abundances, and they can provide straightforward solutions. Our approach appears to be a good pathway to study GC MPs even with small-aperture ground-based telescopes. It is a powerful tool to secure large sample sizes that permit us to reveal substructures in M3 that cannot be seen with the more limited samples available in the past.

In this paper, we investigated the MPs of the prototypical normal GC M3. We derived the populational number ratio based on the merged data (i.e., our photometry in the outer part and the HST photometry in the central part of M3), obtaining  $n(\text{CN-w}):n(\text{CN-s}) = 38:62 (\pm 2)$ , which is different from our previous estimate,  $48:52 (\pm 3)$ , due to the incomplete detection of stars in the central part of M3 from our ground-based observations with a small-aperture telescope. Our revised populational number is still different from that of Milone et al. (2017), who obtained the FG fraction of  $0.305 \pm 0.014$ . It is an apparently simple number-counting task, but our results for the populational number ratio vividly show the importance of securing an unbiased sample from a large sampling area and complete detection of sources, which always has been a core difficulty in astronomy.

Comparisons of our color indices with those of the HST (Milone et al. 2017) showed that the CN-w and CN-s have different correlations with the HST color indices, mainly due to different elemental abundances between the two populations and the different degrees of sensitivity of individual passbands with significantly different bandwidths on elemental abundances.

The  $[\text{Fe}/\text{H}]_{hk}$  derived from our  $hk_{JWL}$  showed bimodal metallicity distributions for both the CN-w and CN-s populations, with the peak metallicities of  $[\text{Fe}/\text{H}]_{hk} = -1.60$  and  $-1.45$ , which we believed to be the mean metallicities of the two GC populations in M3.

Assuming the  $[\text{O}/\text{Fe}]$  from Mészáros et al. (2015), we obtained  $[\text{C}/\text{Fe}] = -0.10$  dex,  $[\text{N}/\text{Fe}] = 0.10$  dex for the less evolved CN-w fRGB, while  $[\text{C}/\text{Fe}] = -0.30$  dex,  $[\text{N}/\text{Fe}] = 0.80$  dex for the less evolved CN-s fRGB. Our carbon and nitrogen abundances from the CH and NH bands are in reasonable agreement with those of other low-resolution spectroscopy studies. However, our carbon and nitrogen abundances from the CN band at  $\lambda 3883$  are about 0.2 dex higher than those from the CH and NH bands. Our calculations showed that the high input oxygen abundances that we adopted from Mészáros et al. (2015) are most likely responsible for our high carbon and nitrogen abundances from the CN band.

Evidence from many studies, including our previous papers, have shown that the chemical evolution from the CN-w population to the CN-s population is not continuous. For

example, the plot of  $[\text{C}/\text{Fe}]$  versus  $[\text{N}/\text{Fe}]$  clearly shows two separate relations. Carbon abundance does not appear to be strongly correlated with nitrogen abundance in the CN-w population, indicating that the CN-w population formed out of interstellar media that experienced little or no CN-cycle hydrogen burning. On the other hand, the CN-s population shows a C–N anticorrelation, likely due to a natural consequence of the CN-cycle that occurred in the previous generation.

We showed that the CN-w and CN-s populations are also different in their structural and kinematical properties. We confirmed our previous result that the M3 CN-s population is more centrally concentrated than the CN-w. On the other hand, the CN-w shows a more elongated morphology with a well-aligned circular motion.

As an extension of the metal-rich RHB studies by others (e.g., Norris & Freeman 1982; Smith & Penny 1989), we showed that the  $cn_{JWL}$  index can be used to study the RHB stellar population of intermediate-metallicity GCs. We derived the populational number ratio for the M3 RHB, finding  $n(\text{CN-w}):n(\text{CN-s}) = 94:6 (\pm 3)$ , which is significantly different from that of the RGB but can be naturally explained in the evolution of metal-poor low-mass stars. The helium-normal CN-w M3 RGB stars evolve into the RHB or RRLs, while the helium-enhanced CN-s RGB stars evolve into the BHB or RRLs, which is also supported by the synthetic HB model simulations of Tailo et al. (2019). Our K-S tests for the M3 RGB and HB showed that the RHB is a CN-w population, while the BHB is CN-s. The RRL is a mixed population.

We discussed the existence of discrete double RGBBs in the CN-w population. Lee (2019b) suggested that the CN-w population exhibits an extended and tilted RGBB. However, our current study clearly shows that the CN-w RGBB can be well described by discrete double RGBBs, which are consistent with the HST photometry by Nardiello et al. (2018).

To understand the origin of discrete double RGBBs, we performed Monte Carlo simulations using our EPS models, finding that M3's bimodal metallicity distribution,  $[\text{Fe}/\text{H}] = -1.60$  and  $-1.45$ , can nicely reproduce our observations. The enhanced helium model cannot reproduce our observations without some metallicity span ( $\Delta[\text{Fe}/\text{H}] \geq 0.05$  dex) or a large age spread ( $\Delta\text{Age} \geq 1$  Gyr), which is also supported by the results of Tailo et al. (2019) that helium enhancement cannot fully explain the M3 HB morphology. At the same time, the bimodal metallicity distribution can explain the large M3 FG  $W_{\text{F}275\text{W}-\text{F}814\text{W}}^{1G}$  range that has puzzled others (Lardo et al. 2018; Tailo et al. 2019).

Finally, we discussed the new formation history of M3. From our new discovery of the bimodal metallicity distributions for both the CN-w and CN-s populations, which is supported by the discrete double RGBBs and the large  $W_{\text{F}275\text{W}-\text{F}814\text{W}}^{1G}$  range, we propose that M3 can be well described by two GCs, namely, C1 (23% of our sample with  $\langle [\text{Fe}/\text{H}] \rangle \approx -1.60$ ) and C2 (77% of our sample with  $\langle [\text{Fe}/\text{H}] \rangle \approx -1.45$ ). In both systems, the CN-s components are more centrally concentrated, which is a common property of Galactic GCs with MPs. At the same time, the CN-w and CN-s components in each system share common kinematical properties: C1 appears to be a random-motion-dominated system, while C2 shows a more well-aligned rotation. We derived the subpopulational number ratios, finding that  $n(\text{CN-w}):n(\text{CN-s}) = 81:19 (\pm 5)$  for C1 and  $43:57 (\pm 4)$  for C2. Given the current mass of M3, the subpopulational number ratios of C1 and C2

compared to Galactic GCs are large for their individual masses, but they are in good agreement with those of GCs in the MC (e.g., see Milone et al. 2020). Therefore, it is most likely that M3 is a merger remnant of two GCs, likely in a less massive environment, such as dwarf galaxies where the GC merger rates are higher than our Galaxy (Thurl & Johnston 2002; Bekki & Yong 2012; Lee 2015; Bekki 2019), and accreted to our Galaxy later in time, which is consistent with the origin of the Helimi Stream as proposed by Koppelman et al. (2019).

Although it is somewhat speculative, a merger scenario could help us understand the neglected mystery of the RRL richness of M3. For example, a merger of two GCs, such as M5 ( $\approx 130$  RRLs and  $[\text{Fe}/\text{H}] = -1.47$ ) and IC 4499 ( $\approx 100$  RRLs and  $[\text{Fe}/\text{H}] = -1.53$ ), can explain the RRL richness of M3 ( $\approx 240$  RRLs and  $[\text{Fe}/\text{H}] = -1.50$ ).

Due to limitations in the spatial angular resolving power of the small-aperture ground-based telescope that we employed, we were not able to perform a reliable study in the central part of M3 ( $r \leq 1'$ ), where the main body of M3 resides. A follow-up study using our new photometric system and large-aperture telescopes to obtain high-precision photometry of the central part of M3 will shed more light on our new discovery presented in this paper.

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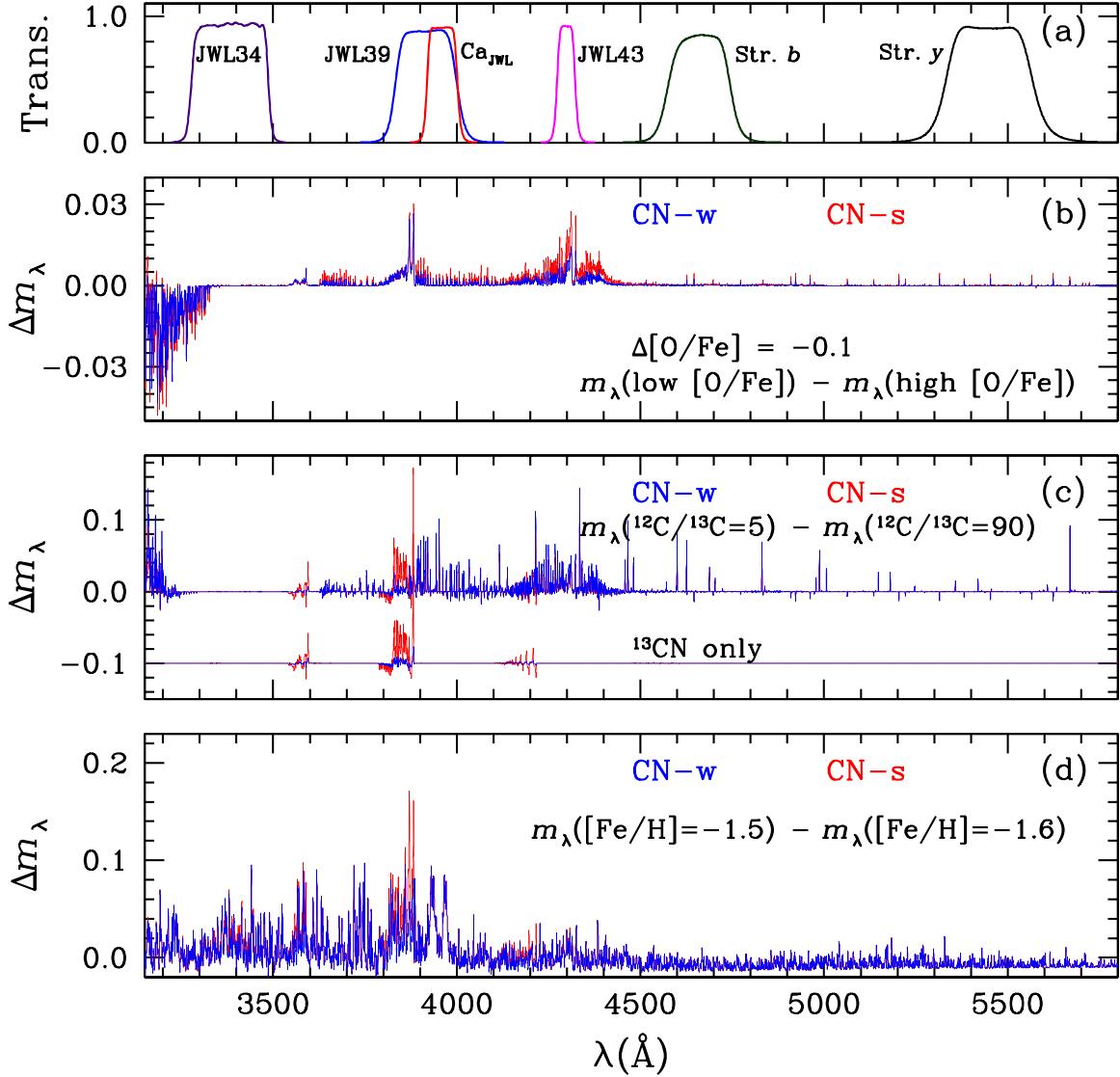
## Appendix A Oxygen Abundance

The synthetic model grids used in our study depend on several parameters. Here, we investigate three input parameters ( $[\text{O}/\text{Fe}]$ ,  $^{12}\text{C}/^{13}\text{C}$ , and  $[\text{Fe}/\text{H}]$ ) for synthetic model grid calculations that can affect our photometric  $[\text{C}/\text{Fe}]$ ,  $[\text{N}/\text{Fe}]$ , and  $[\text{Fe}/\text{H}]_{hk}$ . In addition, we also investigate the effect of carbon abundance on  $[\text{Fe}/\text{H}]_{hk}$ . Note that elemental abundance measurements from any low- to intermediate-resolution spectroscopy are also vulnerable to some uncertainties discussed here.

Molecular CN and CH band strengths in RGB stars can be affected by not only the carbon and nitrogen abundances but also the oxygen abundance through the formation of the CO molecule in their atmospheres. In order to examine the influence of oxygen, we calculated individual color indices using the method described in Section 6 with different input oxygen abundances. For our calculations, we assumed  $[\text{Fe}/\text{H}] = -1.50$  for both populations with different CNO abundances:  $[\text{C}/\text{Fe}] = -0.10$ ,  $[\text{N}/\text{Fe}] = 0.10$ , and  $[\text{O}/\text{Fe}] = 0.50$  and 0.40 for the CN-w population and  $[\text{C}/\text{Fe}] = -0.30$ ,  $[\text{N}/\text{Fe}] = 0.80$ , and  $[\text{O}/\text{Fe}] = 0.15$  and 0.05 for the CN-s population (see Table 5).

In Figures 27(a)–(b), we show filter transmission functions and differences in monochromatic magnitudes, in the sense of those of low  $[\text{O}/\text{Fe}]$  minus those of high  $[\text{O}/\text{Fe}]$ . The different oxygen abundance affects our color indices in two ways: (1) in the wavelength region shorter than  $\approx 3300 \text{ \AA}$ , the enhanced oxygen abundance strengthens the OH band features, part of which lie within the JWL34 bandwidth, as shown in Figure 27(b). (2) It weakens the CN and CH band features because a considerable amount of carbon atoms are tied up in CO molecules. On the other hand, the variation of the CN abundance due to varying oxygen abundance does not affect the NH populations at optical depths of the absorption line formation. Therefore, our  $nh_{\text{JWL}}$  is affected only by the OH band features, as we mentioned above.

Figure 28 shows how oxygen abundance affects our color indices and photometric elemental abundances. As shown, oxygen abundance does not affect the  $(b - y)$  and  $hk_{\text{JWL}}$ . On the other hand, our  $nh_{\text{JWL}}$  index is slightly affected by oxygen abundances, not by different NH populations in the atmosphere but by the OH contribution in the JWL34 bandwidth, as we mentioned above.  $cn_{\text{JWL}}$  and  $ch_{\text{JWL}}$  can be more seriously affected by oxygen abundances. At a fixed carbon abundance, as oxygen abundance increases, the  $cn_{\text{JWL}}$  and  $ch_{\text{JWL}}$  index values decrease due to a decrease in the CN and CH populations. As a consequence, an underestimated oxygen abundance would result in a higher photometric  $[\text{C}/\text{Fe}]_{\text{CN}}$  from the  $ch_{\text{JWL}}$  and  $cn_{\text{JWL}}$  indices and a higher photometric  $[\text{N}/\text{Fe}]_{\text{CN}}$  from the  $cn_{\text{JWL}}$  index. For bright RGB stars,  $\Delta[\text{O}/\text{Fe}] = 0.1 \text{ dex}$  can result in  $\Delta[\text{C}/\text{Fe}] \approx 0.06\text{--}0.07 \text{ dex}$  from  $ch_{\text{JWL}}$  and  $cn_{\text{JWL}}$ , and  $\Delta[\text{N}/\text{Fe}]_{\text{CN}} \approx 0.10 \text{ dex}$  from  $cn_{\text{JWL}}$  at the bright RGB regime for both populations. On the other hand, the oxygen abundance does not affect the  $[\text{N}/\text{Fe}]$  from  $nh_{\text{JWL}}$ , affecting it by no more than 0.02 dex.



**Figure 27.** (a) Filter transmission functions. (b) The difference in the monochromatic magnitude between the model atmosphere with  $\Delta[\text{O}/\text{Fe}] = -0.1$ . We show the results for  $M_V = 0.0$  mag. The blue and red solid lines denote the CN-w and CN-s, respectively. (c) Same as (b) but for  $^{12}\text{C}/^{13}\text{C} = 5$  and 90. The lower lines show differences due only to the CN contribution, with an offset of  $-0.1$  mag. (d) Same as (b) but for  $\Delta[\text{Fe}/\text{H}] = 0.1$  dex.

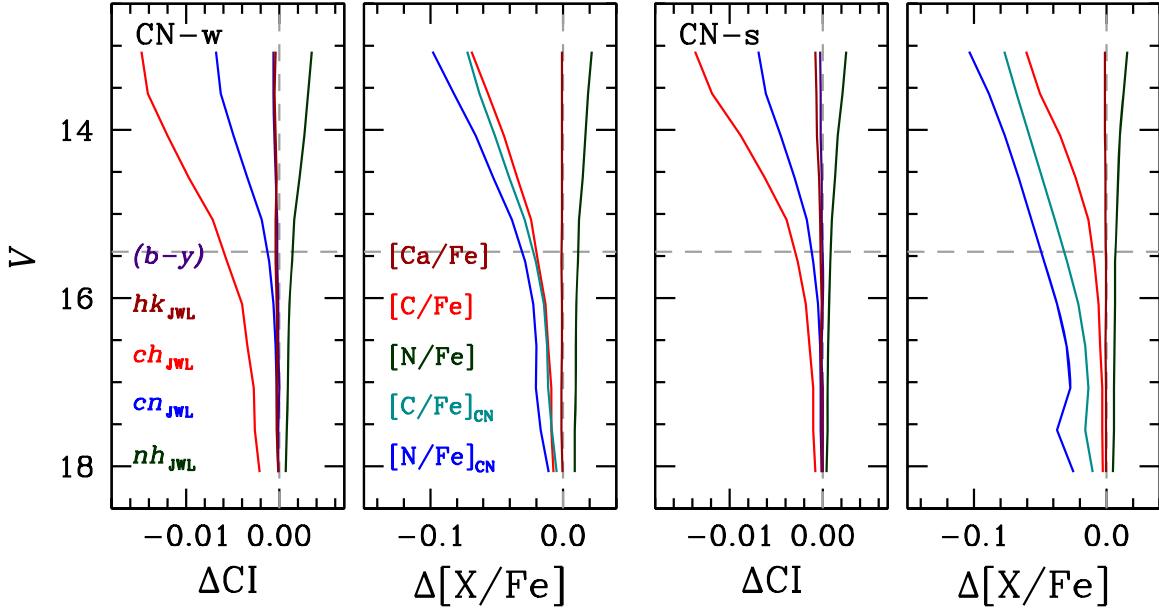


Figure 28. Differences in color indices and inferred elemental abundances due to different [O/Fe].

## Appendix B $^{12}\text{C}/^{13}\text{C}$ Ratios

The carbon isotope ratio,  $^{12}\text{C}/^{13}\text{C}$ , can also affect some of our color indices (e.g., see Briley et al. 1989; Lee 2019c). In GC RGB stars, as a result of the CN process at high temperature accompanied by a noncanonical thermohaline mixing (e.g., see Charbonnel & Zahn 2007), the surface carbon  $^{13}\text{C}$  isotope abundance increases, and the  $^{12}\text{C}/^{13}\text{C}$  ratio decreases. Inhomogeneous  $^{12}\text{C}/^{13}\text{C}$  ratios among different populations in a given GC also can be produced depending on the contribution of the CN processes in the previous generation of stars, supported by the existence of inhomogeneous nitrogen abundances in GC stars. Unfortunately,  $^{12}\text{C}/^{13}\text{C}$  ratios in large-enough samples of GC stars are largely unknown.

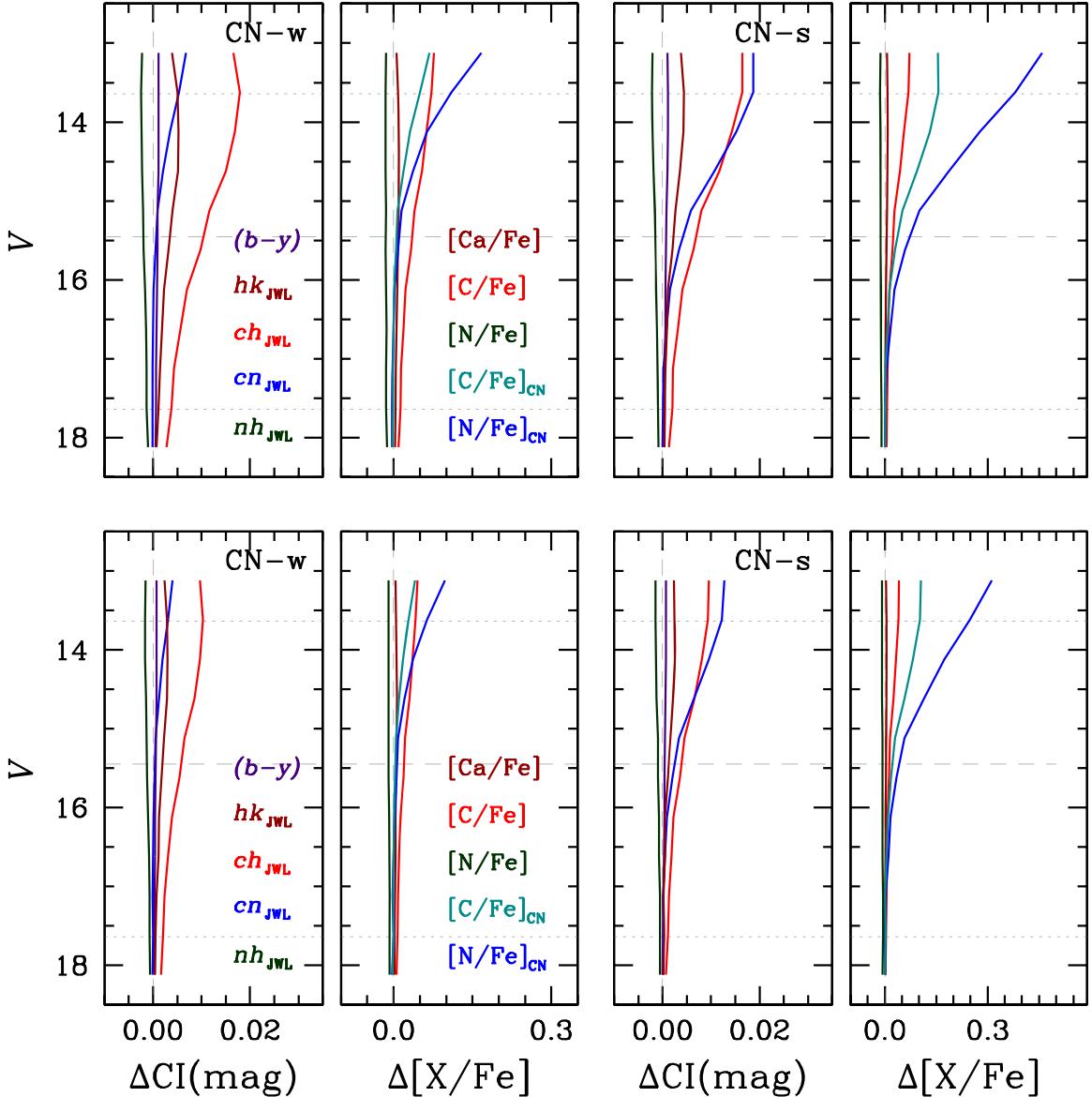
In Figure 27(c), we show the differences in monochromatic magnitude between synthetic spectra with an enhanced  $^{13}\text{C}$  abundance ( $^{12}\text{C}/^{13}\text{C} = 5$ ) and a solar  $^{13}\text{C}$  abundance ( $^{12}\text{C}/^{13}\text{C} = 90$ ). In that panel, we show two cases: (1) those including contributions from both CN and CH molecules, and (2) just those from the CN molecule ( $\text{C}_2$  transitions are undetectably weak). The contribution from  $^{13}\text{CN}$  clearly is much smaller than that from  $^{13}\text{CH}$ . Only in the CN  $\lambda 3883$  band

does the enhanced  $^{13}\text{C}$  abundance make any spectroscopic difference (as shown by Briley et al. 1989).

Our JWL34 filter is almost free from  $^{13}\text{CH}$  contamination and, as a consequence, our  $[\text{N}/\text{Fe}]$  measurements from the  $nh_{\text{JWL}}$  index are hardly affected by the varying  $^{13}\text{C}$  abundance. The  $nh_{\text{JWL}}$  index is very weakly dependent on the  $^{12}\text{C}/^{13}\text{C}$  ratios due to some weak CH band features in the Strömgren  $b$  and  $h$  passbands.

There exist some  $^{13}\text{CH}$  features within the  $\text{Ca}_{\text{JWL}}$  band. However, the  $hk_{\text{JWL}}$  index is much more sensitively dependent on changes in metallicity with a fixed  $[\text{Ca}/\text{Fe}]$  than it is to  $^{13}\text{CH}$ .  $[\text{Fe}/\text{H}]_{hk}$  is essentially independent of  $^{13}\text{CH}$  contamination.

Figure 29 shows how  $^{12}\text{C}/^{13}\text{C}$  ratios affect the color indices and elemental abundances inferred from color indices. In the figure, we show two cases: (1)  $^{12}\text{C}/^{13}\text{C} = 5$  and 90, and (2)  $^{12}\text{C}/^{13}\text{C} = 5$  and 10. The enhanced  $^{13}\text{CN}$  and  $^{13}\text{CH}$  abundances can make substantial differences in both the  $[\text{N}/\text{Fe}]_{\text{CN}}$  and  $[\text{C}/\text{Fe}]_{\text{CN}}$  values measured from  $cn_{\text{JWL}}$ , in particular for bright RGB stars. It also affects the  $[\text{C}/\text{Fe}]$  from  $ch_{\text{JWL}}$ . As shown in the figure, the difference in the inferred  $[\text{N}/\text{Fe}]_{\text{CN}}$  can be as large as 0.5 dex, in the sense that the enhanced  $^{13}\text{CH}$  abundance can make stars be seen more  $[\text{N}/\text{Fe}]$  rich.



**Figure 29.** Differences in color indices and inferred elemental abundances due to different  $^{12}\text{C}/^{13}\text{C}$  ratios. The upper panels are for  $^{12}\text{C}/^{13}\text{C} = 5$  and 90, and the lower panels are for  $^{12}\text{C}/^{13}\text{C} = 5$  and 10.

### Appendix C Metallicity

Metallicity also affects our color indices. In Figure 27(d), we show monochromatic magnitude differences for  $\Delta[\text{Fe}/\text{H}] = 0.1$  dex. Metallicity affects our color indices in two ways: (1) the continuum opacity from the  $\text{H}^-$  ion and Rayleigh scattering from neutral hydrogen depends on metallicity, as the formation of the  $\text{H}^-$  ion diminishes as metallicity decreases. (2) The overall absorption strengths increase with metallicity. The former effect can be minimized by using local continuum sidebands as practiced in low- to intermediate-resolution spectroscopy. However, inappropriate local continuum sideband assignment could bring a serious problem by setting incorrect continuum levels in the forest of heavily absorption

lines (e.g., see Lee 2019c for the case of the CH  $G$  band). Our  $nh_{\text{JWL}}$  index, for example, can be affected not only by individual absorption strengths but also by the shape of the continuum because the baseline of our  $nh_{\text{JWL}}$  (the difference in the pivot wavelengths between the JWL34 and Strömgren  $y$ ) is as large as  $\Delta\lambda \approx 2000 \text{ \AA}$ .

In Figure 27(c), we show the differences in monochromatic magnitude between synthetic spectra with different metallicities, where the slight difference in continuum opacity can be seen. Figure 30 shows the influence of metallicity on our color indices. As shown,  $[\text{N}/\text{Fe}]$  estimated from the  $nh_{\text{JWL}}$  is affected most, in particular in the faint RGB stars. In our simulations,  $\Delta[\text{Fe}/\text{H}] = 0.1$  dex can result in  $\Delta[\text{N}/\text{Fe}] = 0.2$  dex. Happily, the change in  $[\text{Ca}/\text{Fe}]$  causes the change in  $[\text{Fe}/\text{H}]$ , confirming  $hk_{\text{JWL}}$  as a good metallicity index.

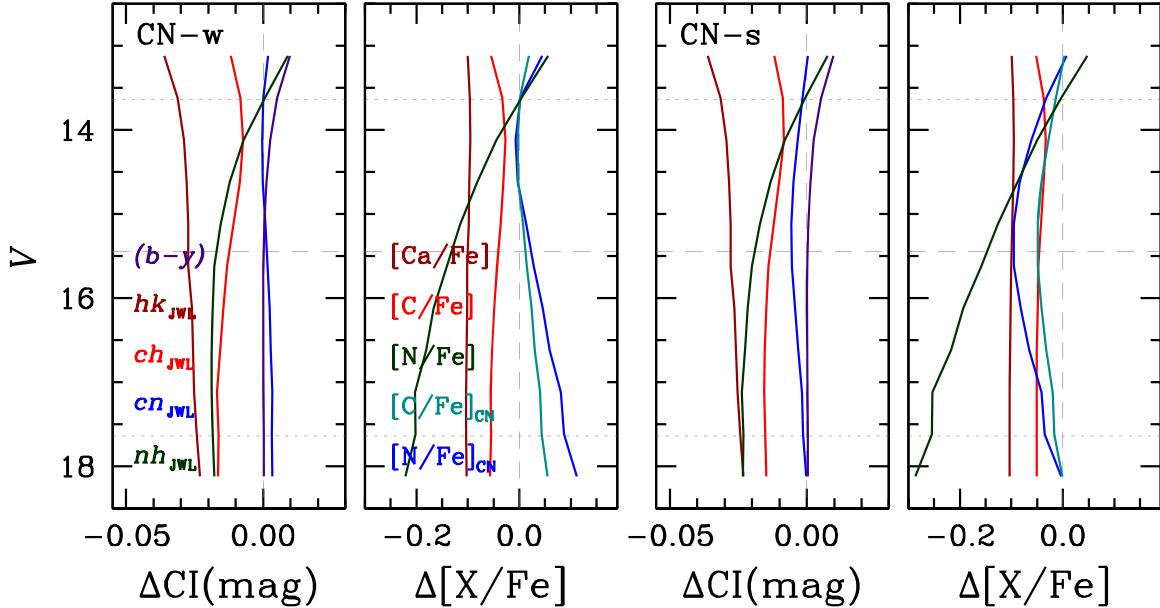


Figure 30. Differences in color indices and inferred elemental abundances due to different metallicities.

#### Appendix D

##### Effect of Carbon Abundance on $hk_{JWL}$ and $[Fe/H]_{hk}$

In our previous work (Lee 2019c), we discussed that the passband of our  $C_{JWL}$  filter contains weak CH lines, which is also shown in Figures 27(b)–(c). In Figure 31, we show the effect of the

carbon abundance enhancement by  $\Delta[C/Fe] = +0.2$  dex on our  $hk_{JWL}$  and  $[Fe/H]_{hk}$ . Our results strongly suggest that the spreads in carbon abundance in both populations,  $\sigma[C/Fe] \leq 0.13$  dex, are not responsible for the bimodal  $[Fe/H]_{hk}$  distribution with the difference between the two peaks of  $\Delta[Fe/H]_{hk} = 0.15$  dex.

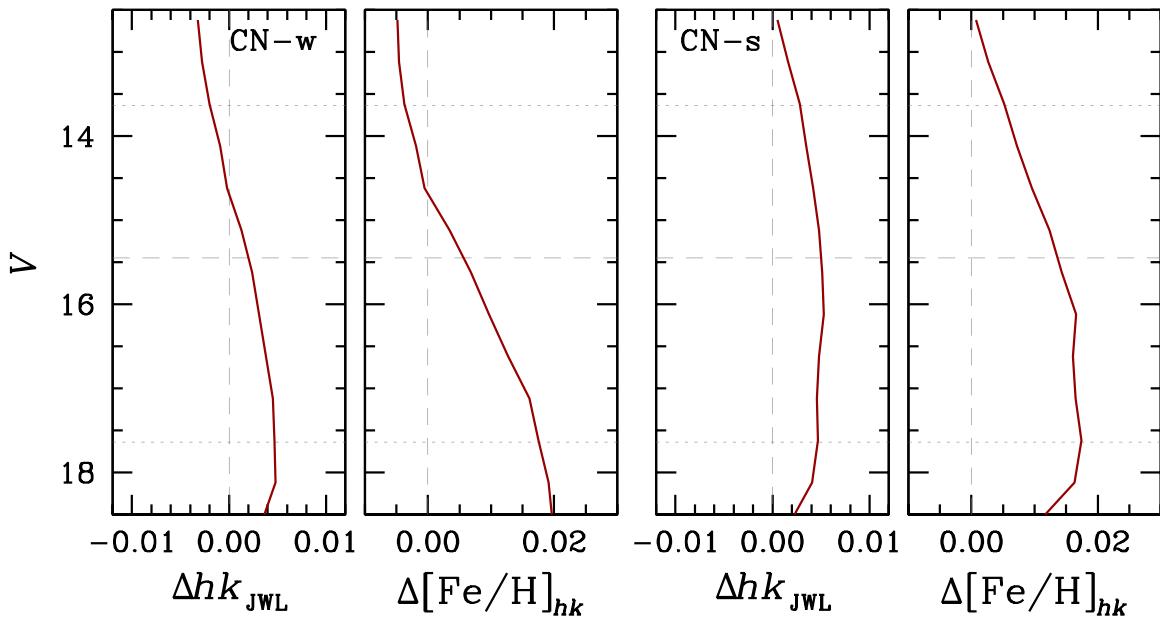


Figure 31. Differences in the  $hk_{JWL}$  and  $[Fe/H]_{hk}$  due to the enhancement of carbon abundance by  $\Delta[C/Fe] = 0.2$  dex.

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