

# Gluings via Intersection Theory



Tobias Scherdin

Humboldt Universität zu Berlin  
Institut für Physik

Domoschool 2025  
July 18, 2025

based on 2411.07330 with G. Crisanti, B. Eden,  
M. Gottwald and P. Mastrolia and ongoing work

# Introduction

- $\mathcal{N} = 4$  SYM: correlators can be computed using the **hexagon ansatz**

# Introduction

- $\mathcal{N} = 4$  SYM: correlators can be computed using the **hexagon ansatz**
- method is based in integrability: operators are expressed as **spin chains** carrying excitations using **Bethe ansätze** [Minahan, Zarembo '02]

# Introduction

- $\mathcal{N} = 4$  SYM: correlators can be computed using the **hexagon ansatz**
- method is based in integrability: operators are expressed as **spin chains** carrying excitations using **Bethe ansätze** [Minahan, Zarembo '02]
- loop-corrections

# Introduction

- $\mathcal{N} = 4$  SYM: correlators can be computed using the **hexagon ansatz**
- method is based in integrability: operators are expressed as **spin chains** carrying excitations using **Bethe ansätze** [Minahan, Zarembo '02]
- loop-corrections:
  - **gluing corrections**

# Introduction

- $\mathcal{N} = 4$  SYM: correlators can be computed using the **hexagon ansatz**
- method is based in integrability: operators are expressed as **spin chains** carrying excitations using **Bethe ansätze** [Minahan, Zarembo '02]
- loop-corrections:
  - **gluing corrections**
  - alternatively: Lagrangian insertion [Eden, Gottwald, Le Plat, TS '23]

# Introduction

- $\mathcal{N} = 4$  SYM: correlators can be computed using the **hexagon ansatz**
- method is based in integrability: operators are expressed as **spin chains** carrying excitations using **Bethe ansätze** [Minahan, Zarembo '02]
- loop-corrections:
  - **gluing corrections**
  - alternatively: Lagrangian insertion [Eden, Gottwald, Le Plat, TS '23]
- gluing corrections can be reduced to rational integrals  
→ can be solved using **intersection theory**

# Introduction

- $\mathcal{N} = 4$  SYM: correlators can be computed using the **hexagon ansatz**
- method is based in integrability: operators are expressed as **spin chains** carrying excitations using **Bethe ansätze** [Minahan, Zarembo '02]
- loop-corrections:
  - **gluing corrections**
  - alternatively: Lagrangian insertion [Eden, Gottwald, Le Plat, TS '23]
- gluing corrections can be reduced to rational integrals  
→ can be solved using **intersection theory**



master decomposition



# Introduction

- $\mathcal{N} = 4$  SYM: correlators can be computed using the **hexagon ansatz**
- method is based in integrability: operators are expressed as **spin chains** carrying excitations using **Bethe ansätze** [Minahan, Zarembo '02]
- loop-corrections:
  - **gluing corrections**
  - alternatively: Lagrangian insertion [Eden, Gottwald, Le Plat, TS '23]
- gluing corrections can be reduced to rational integrals  
→ can be solved using **intersection theory**



master decomposition  
solving canonical differential equations for the masters

# Outline

## Hexagon Ansatz

Short Introduction

Gluing Corrections

## Intersection Theory

Twisted De Rham Co-Homology

Computing Intersection Numbers

## Results

Solving the Gluing Integrals

Conclusion

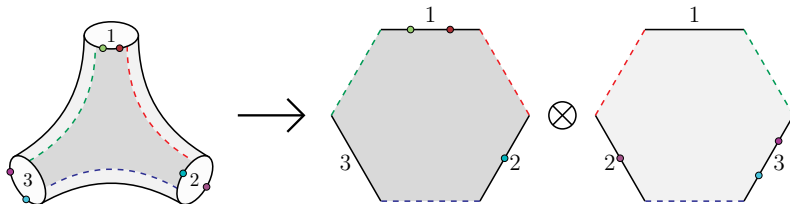
# The Hexagon Ansatz

- **goal:** compute correlators within integrability

# The Hexagon Ansatz

- **goal:** compute correlators within integrability
- **main idea:** cutting 3-point function into two hexagonal patches

[Basso, Komatsu, Vieira '15]



# The Hexagon Ansatz

- **goal:** compute correlators within integrability
- **main idea:** cutting 3-point function into two hexagonal patches  
[Basso, Komatsu, Vieira '15]
- sum over all ways to distribute the excitations, schematically:

$$\begin{aligned}
 \langle \mathcal{B}_{L_1} \mathcal{O}_{L_2} \mathcal{O}_{L_3} \rangle &\propto \text{Diagram 1} \otimes \text{Diagram 2} - e^{ip_2 h_{12}} \text{Diagram 3} \\
 &- S_{21} e^{ip_1 h_{12}} \text{Diagram 4} + \text{Diagram 5}
 \end{aligned}$$

The diagrams are hexagonal patches with vertices labeled 1, 2, 3. The patches are connected by tensor products ( $\otimes$ ) and summed with coefficients  $e^{ip_2 h_{12}}$  and  $S_{21} e^{ip_1 h_{12}}$ . The patches have colored edges: red dashed, green dashed, and blue dashed. The diagrams show different ways to distribute excitations (represented by dots on the edges) between the two hexagonal patches.

# The Hexagon Ansatz

- **goal:** compute correlators within integrability
- **main idea:** cutting 3-point function into two hexagonal patches  
[Basso, Komatsu, Vieira '15]
- sum over all ways to distribute the excitations, schematically:

$$\begin{aligned}
 \langle \mathcal{B}_{L_1} \mathcal{O}_{L_2} \mathcal{O}_{L_3} \rangle \propto & \text{Diagram 1} \otimes \text{Diagram 2} - e^{ip_2 h_{12}} \text{Diagram 3} \\
 & - S_{21} e^{ip_1 h_{12}} \text{Diagram 4} + \text{Diagram 5}
 \end{aligned}$$

- can be extended towards:
  - general  $n$ -point function  $\rightarrow$  position dependence

# The Hexagon Ansatz

- **goal:** compute correlators within integrability
- **main idea:** cutting 3-point function into two hexagonal patches  
[Basso, Komatsu, Vieira '15]
- sum over all ways to distribute the excitations, schematically:

$$\begin{aligned}
 \langle \mathcal{B}_{L_1} \mathcal{O}_{L_2} \mathcal{O}_{L_3} \rangle \propto & \text{Diagram 1} \otimes \text{Diagram 2} - e^{ip_2 h_{12}} \text{Diagram 3} \\
 & - S_{21} e^{ip_1 h_{12}} \text{Diagram 4} + \text{Diagram 5}
 \end{aligned}$$

The diagrams consist of two hexagons joined at a vertex, with internal lines colored red, green, or blue. Diagram 1 and 2 are the base case. Diagram 3 has an additional arc connecting the top vertices. Diagram 4 has an additional arc connecting the top vertices and a different internal line configuration. Diagram 5 has an additional arc connecting the top vertices and a different internal line configuration.

- can be extended towards:
  - general  $n$ -point function  $\rightarrow$  position dependence
  - higher genus surfaces  $\rightarrow N$  expansion

# The Hexagon Ansatz

- **goal:** compute correlators within integrability
- **main idea:** cutting 3-point function into two hexagonal patches

[Basso, Komatsu, Vieira '15]

- sum over all ways to distribute the excitations, schematically:

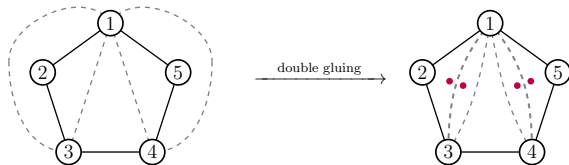
$$\begin{aligned}
 \langle \mathcal{B}_{L_1} \mathcal{O}_{L_2} \mathcal{O}_{L_3} \rangle \propto & \text{Diagram 1} \otimes \text{Diagram 2} - e^{ip_2 h_{12}} \text{Diagram 3} \\
 & - S_{21} e^{ip_1 h_{12}} \text{Diagram 4} + \text{Diagram 5}
 \end{aligned}$$

- can be extended towards:
  - general  $n$ -point function  $\rightarrow$  position dependence
  - higher genus surfaces  $\rightarrow N$  expansion
  - higher loop orders  $\rightarrow$  gluing corrections



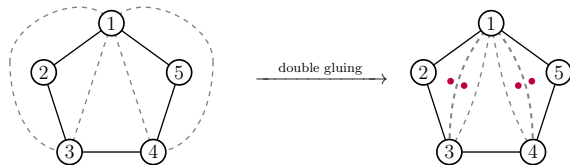
# Gluing Corrections

- “virtual particles” may propagate over edges



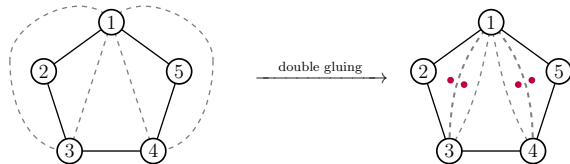
# Gluing Corrections

- “virtual particles” may propagate over edges
- included by summing over all possible states on each edge



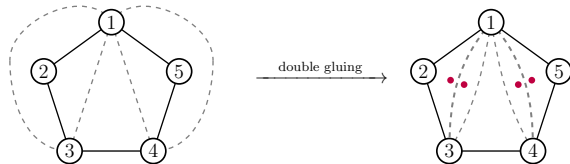
# Gluing Corrections

- “virtual particles” may propagate over edges
- included by summing over all possible states on each edge  
→ bound states with infinitely many fermions  $\psi^1, \psi^2$  and up to one of each of the bosons  $\phi^1, \phi^2$



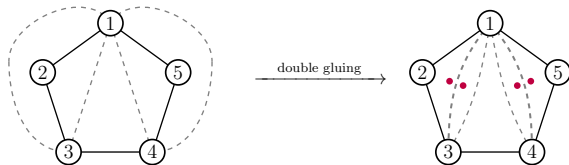
# Gluing Corrections

- “virtual particles” may propagate over edges
- included by summing over all possible states on each edge  
→ bound states with infinitely many fermions  $\psi^1, \psi^2$  and up to one of each of the bosons  $\phi^1, \phi^2$
- structure of  $S$ -matrix allows for power counting in  $g^2$



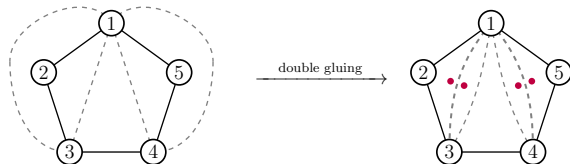
# Gluing Corrections

- “virtual particles” may propagate over edges
- included by summing over all possible states on each edge  
→ bound states with infinitely many fermions  $\psi^1, \psi^2$  and up to one of each of the bosons  $\phi^1, \phi^2$
- structure of  $S$ -matrix allows for power counting in  $g^2$ 
  - $g^{2(L+1)}$  for each excitation on an edge of size  $L$



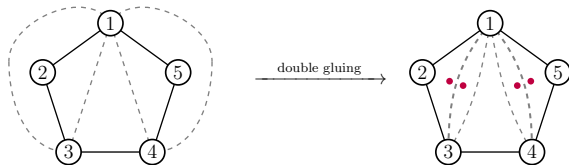
# Gluing Corrections

- “virtual particles” may propagate over edges
- included by summing over all possible states on each edge  
→ bound states with infinitely many fermions  $\psi^1, \psi^2$  and up to one of each of the bosons  $\phi^1, \phi^2$
- structure of  $S$ -matrix allows for power counting in  $g^2$ 
  - $g^{2(L+1)}$  for each excitation on an edge of size  $L$
  - $g^{-2}$  for each pair of excitations on different edges of the same hexagon



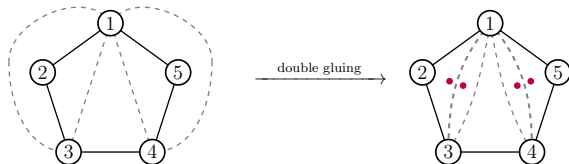
# Gluing Corrections

- “virtual particles” may propagate over edges
- included by summing over all possible states on each edge  
→ bound states with infinitely many fermions  $\psi^1, \psi^2$  and up to one of each of the bosons  $\phi^1, \phi^2$
- structure of  $S$ -matrix allows for power counting in  $g^2$ 
  - $g^{2(L+1)}$  for each excitation on an edge of size  $L$
  - $g^{-2}$  for each pair of excitations on different edges of the same hexagon
- we try to evaluate the first correction to  $\langle \mathcal{O}_2 \mathcal{O}_2 \mathcal{O}_2 \mathcal{O}_2 \mathcal{O}_2 \rangle$



# Gluing Corrections

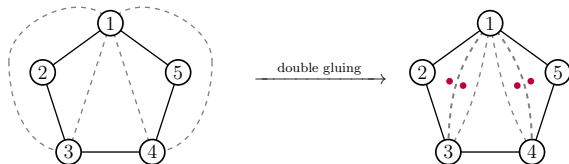
- “virtual particles” may propagate over edges
- included by summing over all possible states on each edge  
→ bound states with infinitely many fermions  $\psi^1, \psi^2$  and up to one of each of the bosons  $\phi^1, \phi^2$
- structure of  $S$ -matrix allows for power counting in  $g^2$ 
  - $g^{2(L+1)}$  for each excitation on an edge of size  $L$
  - $g^{-2}$  for each pair of excitations on different edges of the same hexagon
- we try to evaluate the first correction to  $\langle \mathcal{O}_2 \mathcal{O}_2 \mathcal{O}_2 \mathcal{O}_2 \mathcal{O}_2 \rangle$ 
  - single-gluing channels can be evaluated rather straight forward





# Gluing Corrections

- “virtual particles” may propagate over edges
- included by summing over all possible states on each edge  
→ bound states with infinitely many fermions  $\psi^1, \psi^2$  and up to one of each of the bosons  $\phi^1, \phi^2$
- structure of  $S$ -matrix allows for power counting in  $g^2$ 
  - $g^{2(L+1)}$  for each excitation on an edge of size  $L$
  - $g^{-2}$  for each pair of excitations on different edges of the same hexagon
- we try to evaluate the first correction to  $\langle \mathcal{O}_2 \mathcal{O}_2 \mathcal{O}_2 \mathcal{O}_2 \mathcal{O}_2 \rangle$ 
  - single-gluing channels can be evaluated rather straight forward
  - double-gluing is more complicated



[Fleury, Komatsu '18]

# The Double-Gluing Process

- for simplicity, we only consider the so-called  $\mathcal{X}$ -element of the  $S$ -matrix by restricting to  $\phi_1^1 \phi_2^0$ -states: [De Leeuw, Eden, Le Plat, Meier, Sfondrini '19]

$$I_{\mathcal{X}} = \sum_{K, L=1}^{\infty} \sum_{k, l=0}^{K-1, L-1} \int \frac{du dv}{4\pi} \frac{K L g^2(v^+ - u^-) \Sigma^{KL}(u, v) \hat{\mathcal{X}}_k^{k, l}(u, v) W(u, v)}{(u^-)^2 (v^+)^2 u^+ v^- (u^+ - v^-)}$$

# The Double-Gluing Process

- for simplicity, we only consider the so-called  $\mathcal{X}$ -element of the  $S$ -matrix by restricting to  $\phi_1^1 \phi_2^0$ -states: [De Leeuw, Eden, Le Plat, Meier, Sfondrini '19]

$$I_{\mathcal{X}} = \sum_{K, L=1}^{\infty} \sum_{k, l=0}^{K-1, L-1} \int \frac{du dv}{4\pi} \frac{K L g^2(v^+ - u^-) \Sigma^{KL}(u, v) \hat{\mathcal{X}}_k^{k, l}(u, v) W(u, v)}{(u^-)^2 (v^+)^2 u^+ v^- (u^+ - v^-)}$$

- picking residua, we can decompose this as follows

$$I_{\mathcal{X}} = (S_1 + S_2) + S_W + S_{mes} + S_{mat}$$

$$S_1 = \sum a^{\sigma_{lm}} b^{\sigma_{jkm}} y^{\sigma_{Lm}} z^{\sigma_{jK}} \frac{\sigma_{kKm}}{j\sigma_{jKk}} * \frac{\Gamma_{1km} \Gamma_{1lm} \Gamma_{Km} \Gamma_{Lm} \Gamma_{jKLm} \Gamma_{1jklLM}}{\Gamma_{1k} \Gamma_{1l} \Gamma_K \Gamma_L \Gamma_{1m}^2 \Gamma_{1ILM} \Gamma_{1jKm} \Gamma_{1jklLM}},$$

$$S_2 = \sum a^{\sigma_{jlm}} b^{\sigma_{km}} y^{\sigma_{jLm}} z^K \frac{\sigma_{ILM}}{j\sigma_{jILM}} * \frac{\Gamma_{1km} \Gamma_{1lm} \Gamma_{Km} \Gamma_{Lm} \Gamma_{jKLm} \Gamma_{1jkKm} \Gamma_{1jklLM}}{\Gamma_{1k} \Gamma_{1l} \Gamma_K \Gamma_L \Gamma_{1m}^2 \Gamma_{1kKm} \Gamma_{1jKm} \Gamma_{1jklLM} \Gamma_{1jILM}}$$

# The Double-Gluing Process

- for simplicity, we only consider the so-called  $\mathcal{X}$ -element of the  $S$ -matrix by restricting to  $\phi_1^1 \phi_2^0$ -states: [De Leeuw, Eden, Le Plat, Meier, Sfondrini '19]

$$I_{\mathcal{X}} = \sum_{K, L=1}^{\infty} \sum_{k, l=0}^{K-1, L-1} \int \frac{du dv}{4\pi} \frac{K L g^2 (v^+ - u^-) \Sigma^{KL}(u, v) \hat{\mathcal{X}}_k^{k, l}(u, v) W(u, v)}{(u^-)^2 (v^+)^2 u^+ v^- (u^+ - v^-)}$$

- picking residua, we can decompose this as follows

$$I_{\mathcal{X}} = (S_1 + S_2) + S_W + S_{mes} + S_{mat}$$

$$S_1 = \sum a^{\sigma_{lm}} b^{\sigma_{jkm}} y^{\sigma_{Lm}} z^{\sigma_{jK}} \frac{\sigma_{kKm}}{j\sigma_{jKkM}} * \frac{\Gamma_{1km} \Gamma_{1lm} \Gamma_{Km} \Gamma_{Lm} \Gamma_{jKLm} \Gamma_{1jklLM}}{\Gamma_{1k} \Gamma_{1l} \Gamma_K \Gamma_L \Gamma_{1m}^2 \Gamma_{1ILM} \Gamma_{1jKm} \Gamma_{1jkLM}},$$

$$S_2 = \sum a^{\sigma_{jlm}} b^{\sigma_{km}} y^{\sigma_{jLm}} z^K \frac{\sigma_{ILM}}{j\sigma_{jILM}} * \frac{\Gamma_{1km} \Gamma_{1lm} \Gamma_{Km} \Gamma_{Lm} \Gamma_{jKLm} \Gamma_{1jkKm} \Gamma_{1jklLM}}{\Gamma_{1k} \Gamma_{1l} \Gamma_K \Gamma_L \Gamma_{1m}^2 \Gamma_{1kKm} \Gamma_{1jKm} \Gamma_{1jkLM} \Gamma_{1jILM}}$$

- the other parts are easily resummed

# The Double-Gluing Process

- for simplicity, we only consider the so-called  $\mathcal{X}$ -element of the  $S$ -matrix by restricting to  $\phi_1^1 \phi_2^0$ -states: [De Leeuw, Eden, Le Plat, Meier, Sfondrini '19]

$$I_{\mathcal{X}} = \sum_{K, L=1}^{\infty} \sum_{k, l=0}^{K-1, L-1} \int \frac{du dv}{4\pi} \frac{K L g^2(v^+ - u^-) \Sigma^{KL}(u, v) \hat{\mathcal{X}}_k^{k, l}(u, v) W(u, v)}{(u^-)^2 (v^+)^2 u^+ v^- (u^+ - v^-)}$$

- picking residua, we can decompose this as follows

$$I_{\mathcal{X}} = (S_1 + S_2) + S_W + S_{mes} + S_{mat}$$

$$S_1 = \sum a^{\sigma_{lm}} b^{\sigma_{jkm}} y^{\sigma_{Lm}} z^{\sigma_{jK}} \frac{\sigma_{kKm}}{j \sigma_{jKkM}} * \frac{\Gamma_{1km} \Gamma_{1lm} \Gamma_{Km} \Gamma_{Lm} \Gamma_{jKLm} \Gamma_{1jklLM}}{\Gamma_{1k} \Gamma_{1l} \Gamma_K \Gamma_L \Gamma_{1m}^2 \Gamma_{1ILM} \Gamma_{1jKm} \Gamma_{1jklM}},$$

$$S_2 = \sum a^{\sigma_{jlm}} b^{\sigma_{km}} y^{\sigma_{jLm}} z^K \frac{\sigma_{ILM}}{j \sigma_{jILM}} * \frac{\Gamma_{1km} \Gamma_{1lm} \Gamma_{Km} \Gamma_{Lm} \Gamma_{jKLm} \Gamma_{1jkKm} \Gamma_{1jklLM}}{\Gamma_{1k} \Gamma_{1l} \Gamma_K \Gamma_L \Gamma_{1m}^2 \Gamma_{1kKm} \Gamma_{1jKm} \Gamma_{1jklM} \Gamma_{1jILM}}$$

- the other parts are easily resummed
- $a, b, y, z$  are defined by cross-ratios of the planar kinematics

# Resumming in Terms of Rational Integrals

- these sums are of  ${}_{p+1}F_p$ -type  $\rightarrow$  can be rewritten as integrals:

# Resumming in Terms of Rational Integrals

- these sums are of  $_{p+1}F_p$ -type  $\rightarrow$  can be rewritten as integrals:

$$S_1 = \int_0^1 ds \int_0^1 dt \frac{byz^2(1-t)}{q_1(s,0)q_1(s,1)q_1(s,t)},$$

$$S_2 = l_1 + l_2,$$

$$l_1 = \int_0^1 ds \int_0^1 dt \frac{n_1(s,t)}{q_2(s,t)q_3(s,t)},$$

$$l_2 = \int_0^1 dr \int_0^1 ds \int_0^1 dt \frac{n_2(r,s,t)}{(q_4(r,s,t))^2}$$

# Resumming in Terms of Rational Integrals

- these sums are of  $_{p+1}F_p$ -type  $\rightarrow$  can be rewritten as integrals:

$$S_1 = \int_0^1 ds \int_0^1 dt \frac{byz^2(1-t)}{q_1(s,0)q_1(s,1)q_1(s,t)},$$

$$S_2 = l_1 + l_2,$$

$$l_1 = \int_0^1 ds \int_0^1 dt \frac{n_1(s,t)}{q_2(s,t)q_3(s,t)},$$

$$l_2 = \int_0^1 dr \int_0^1 ds \int_0^1 dt \frac{n_2(r,s,t)}{(q_4(r,s,t))^2}$$

- the polynomials  $q_i$  are in general quadratic in all variables



# Basics of Intersection Theory

- twisted period integrals:  $\int_{C_R} u \varphi_L =: \langle \varphi_L | C_R \rangle$

# Basics of Intersection Theory

- twisted period integrals:  $\int_{\mathcal{C}_R} u \varphi_L =: \langle \varphi_L | \mathcal{C}_R \rangle$   
 $\rightarrow \varphi_L = \widehat{\varphi}_L(z) dz := \widehat{\varphi}_L(z) dz_1 \wedge \dots \wedge dz_n$ : rational  $n$ -form

# Basics of Intersection Theory

- twisted period integrals:  $\int_{\mathcal{C}_R} u \varphi_L =: \langle \varphi_L | \mathcal{C}_R \rangle$ 
  - $\varphi_L = \widehat{\varphi}_L(z) dz := \widehat{\varphi}_L(z) dz_1 \wedge \dots \wedge dz_n$ : rational  $n$ -form
  - $u$ : multivalued function, vanishing at the boundaries of  $\mathcal{C}_R$  and poles of  $\varphi_L$

# Basics of Intersection Theory

- twisted period integrals:  $\int_{\mathcal{C}_R} u \varphi_L =: \langle \varphi_L | \mathcal{C}_R \rangle$ 
  - $\varphi_L = \widehat{\varphi}_L(z) dz := \widehat{\varphi}_L(z) dz_1 \wedge \dots \wedge dz_n$ : rational  $n$ -form
  - $u$ : multivalued function, vanishing at the boundaries of  $\mathcal{C}_R$  and poles of  $\varphi_L$
  - define twist  $d \log(u) =: \omega =: \sum \widehat{\omega}_i dz_i$

# Basics of Intersection Theory

- twisted period integrals:  $\int_{\mathcal{C}_R} u \varphi_L =: \langle \varphi_L | \mathcal{C}_R \rangle$ 
  - $\varphi_L = \widehat{\varphi}_L(z) dz := \widehat{\varphi}_L(z) dz_1 \wedge \dots \wedge dz_n$ : rational  $n$ -form
  - $u$ : multivalued function, vanishing at the boundaries of  $\mathcal{C}_R$  and poles of  $\varphi_L$
  - define twist  $d \log(u) =: \omega =: \sum \widehat{\omega}_i dz_i$
- integrals are invariant under shifts  $\varphi \rightarrow \varphi + \nabla_\omega \phi := \varphi + d\phi + \omega \wedge \phi$

# Basics of Intersection Theory

- twisted period integrals:  $\int_{\mathcal{C}_R} u \varphi_L =: \langle \varphi_L | \mathcal{C}_R \rangle$ 
  - $\varphi_L = \widehat{\varphi}_L(z) dz := \widehat{\varphi}_L(z) dz_1 \wedge \dots \wedge dz_n$ : rational  $n$ -form
  - $u$ : multivalued function, vanishing at the boundaries of  $\mathcal{C}_R$  and poles of  $\varphi_L$
  - define twist  $d \log(u) =: \omega =: \sum \widehat{\omega}_i dz_i$
- integrals are invariant under shifts  $\varphi \rightarrow \varphi + \nabla_\omega \phi := \varphi + d\phi + \omega \wedge \phi$ 
  - **twisted De Rham cohomology** with finite basis of **masters** of size  $\nu = \#$  of solutions to  $(\omega_i = 0)$

# Basics of Intersection Theory

- twisted period integrals:  $\int_{\mathcal{C}_R} u \varphi_L =: \langle \varphi_L | \mathcal{C}_R \rangle$ 
  - $\varphi_L = \widehat{\varphi}_L(z) dz := \widehat{\varphi}_L(z) dz_1 \wedge \dots \wedge dz_n$ : rational  $n$ -form
  - $u$ : multivalued function, vanishing at the boundaries of  $\mathcal{C}_R$  and poles of  $\varphi_L$
  - define twist  $d \log(u) =: \omega =: \sum \hat{\omega}_i dz_i$
- integrals are invariant under shifts  $\varphi \rightarrow \varphi + \nabla_\omega \phi := \varphi + d\phi + \omega \wedge \phi$ 
  - **twisted De Rham cohomology** with finite basis of **masters** of size  $\nu = \#$  of solutions to  $(\omega_i = 0)$
- master decomposition formula: [Mastrolia, Mizera '19]  
[Frellesvig, Gasparotto, Laporta, Mandal, Mastrolia, Mattiazzi, Mizera '19]

$$\langle \varphi_L | = \sum_{i=1}^{\nu} c_i \langle e_i | \quad \text{with} \quad c_i = \langle \varphi_L | h_j \rangle (\mathbf{C}^{-1})_{ji} \quad \text{and} \quad \mathbf{C}_{ij} := \langle e_i | h_j \rangle$$

# Basics of Intersection Theory

- twisted period integrals:  $\int_{\mathcal{C}_R} u \varphi_L =: \langle \varphi_L | \mathcal{C}_R \rangle$ 
  - $\varphi_L = \widehat{\varphi}_L(z) dz := \widehat{\varphi}_L(z) dz_1 \wedge \dots \wedge dz_n$ : rational  $n$ -form
  - $u$ : multivalued function, vanishing at the boundaries of  $\mathcal{C}_R$  and poles of  $\varphi_L$
  - define twist  $d \log(u) =: \omega =: \sum \widehat{\omega}_i dz_i$
- integrals are invariant under shifts  $\varphi \rightarrow \varphi + \nabla_\omega \phi := \varphi + d\phi + \omega \wedge \phi$ 
  - **twisted De Rham cohomology** with finite basis of **masters** of size  $\nu = \#$  of solutions to  $(\omega_i = 0)$

- master decomposition formula: [Mastrolia, Mizera '19]  
[Frellesvig, Gasparotto, Laporta, Mandal, Mastrolia, Mattiazzi, Mizera '19]

$$\langle \varphi_L | = \sum_{i=1}^{\nu} c_i \langle e_i | \quad \text{with} \quad c_i = \langle \varphi_L | h_j \rangle (\mathbf{C}^{-1})_{ji} \quad \text{and} \quad \mathbf{C}_{ij} := \langle e_i | h_j \rangle$$

- we can also decompose derivatives of masters:  $\partial_x \langle e_i | = \langle \nabla_{\sigma_x} e_i | = \Omega_{ij} \langle e_j |$



# Computing Intersection Numbers

- 1-forms:

# Computing Intersection Numbers

- 1-forms:

$$\langle \varphi_L | \varphi_R \rangle = \sum_{p \in \mathcal{P}_\omega} \text{Res}_{z=p}(\psi_p \varphi_R)$$

$$\nabla_\omega \psi_p = \varphi_L \qquad \psi_p = \sum_{n=\min}^{\max} c_i (z-p)^n + \mathcal{O}(z-p)^{\max+1}$$

# Computing Intersection Numbers

## ■ 1-forms:

$$\langle \varphi_L | \varphi_R \rangle = \sum_{p \in \mathcal{P}_\omega} \text{Res}_{z=p}(\psi_p \varphi_R)$$

$$\nabla_\omega \psi_p = \varphi_L \qquad \psi_p = \sum_{n=\min}^{\max} c_i (z-p)^n + \mathcal{O}(z-p)^{\max+1}$$

## ■ $n$ -forms:

# Computing Intersection Numbers

- 1-forms:

$$\langle \varphi_L | \varphi_R \rangle = \sum_{p \in \mathcal{P}_\omega} \text{Res}_{z=p}(\psi_p \varphi_R)$$

$$\nabla_\omega \psi_p = \varphi_L \quad \psi_p = \sum_{n=\min}^{\max} c_i (z-p)^n + \mathcal{O}(z-p)^{\max+1}$$

- $n$ -forms:

- can iteratively project out 1-form bases

# Computing Intersection Numbers

- 1-forms:

$$\langle \varphi_L | \varphi_R \rangle = \sum_{p \in \mathcal{P}_\omega} \text{Res}_{z=p}(\psi_p \varphi_R)$$

$$\nabla_\omega \psi_p = \varphi_L \quad \psi_p = \sum_{n=\min}^{\max} c_i (z-p)^n + \mathcal{O}(z-p)^{\max+1}$$

- $n$ -forms:

- can iteratively project out 1-form bases  
→ we can reduce everything to 1-form intersection numbers

# Computing Intersection Numbers

- 1-forms:

$$\langle \varphi_L | \varphi_R \rangle = \sum_{p \in \mathcal{P}_\omega} \text{Res}_{z=p}(\psi_p \varphi_R)$$

$$\nabla_\omega \psi_p = \varphi_L \quad \psi_p = \sum_{n=\min}^{\max} c_i (z-p)^n + \mathcal{O}(z-p)^{\max+1}$$

- $n$ -forms:

- can iteratively project out 1-form bases
  - we can reduce everything to 1-form intersection numbers
- doing so introduces multiplicatively growing total basis

# Computing the Gluing Integrals

- our integrals have no twist  $\rightarrow$  can be introduced by defining
$$u = (s(1-s)t(1-t) \prod p_i(s, t))^\gamma$$

## Computing the Gluing Integrals

- our integrals have no twist  $\rightarrow$  can be introduced by defining
$$u = (s(1-s)t(1-t) \prod p_i(s, t))^\gamma$$
$$\rightarrow \text{the limit } \gamma \rightarrow 0 \text{ yields the correct integral}$$



# Computing the Gluing Integrals

- our integrals have no twist  $\rightarrow$  can be introduced by defining
$$u = (s(1-s)t(1-t) \prod p_i(s, t))^\gamma$$
$$\rightarrow \text{the limit } \gamma \rightarrow 0 \text{ yields the correct integral}$$
- dealing with the quadratic polynomials requires us to write them in a factorized form  $p(s, t) = (s - s^+(t))(s - s^-(t))$

# Computing the Gluing Integrals

- our integrals have no twist  $\rightarrow$  can be introduced by defining
$$u = (s(1-s)t(1-t) \prod p_i(s, t))^\gamma$$
$$\rightarrow \text{the limit } \gamma \rightarrow 0 \text{ yields the correct integral}$$
- dealing with the quadratic polynomials requires us to write them in a factorized form  $p(s, t) = (s - s^+(t))(s - s^-(t))$
- most intersection numbers only need the leading order in this pole, however for differential matrices we need second order

# Computing the Gluing Integrals

- our integrals have no twist  $\rightarrow$  can be introduced by defining
$$u = (s(1-s)t(1-t) \prod p_i(s, t))^\gamma$$
$$\rightarrow \text{the limit } \gamma \rightarrow 0 \text{ yields the correct integral}$$
- dealing with the quadratic polynomials requires us to write them in a factorized form  $p(s, t) = (s - s^+(t))(s - s^-(t))$
- most intersection numbers only need the leading order in this pole, however for differential matrices we need second order  
 $\rightarrow$  very tedious expansion

# Computing the Gluing Integrals

- our integrals have no twist  $\rightarrow$  can be introduced by defining
$$u = (s(1-s)t(1-t) \prod p_i(s, t))^\gamma$$
$$\rightarrow \text{the limit } \gamma \rightarrow 0 \text{ yields the correct integral}$$
- dealing with the quadratic polynomials requires us to write them in a factorized form  $p(s, t) = (s - s^+(t))(s - s^-(t))$
- most intersection numbers only need the leading order in this pole, however for differential matrices we need second order
  - $\rightarrow$  very tedious expansion
  - $\rightarrow$  actually, everything beyond leading order turns out to be just a consistency condition

# Computing the Gluing Integrals

- our integrals have no twist  $\rightarrow$  can be introduced by defining
$$u = (s(1-s)t(1-t) \prod p_i(s, t))^\gamma$$
$$\rightarrow \text{the limit } \gamma \rightarrow 0 \text{ yields the correct integral}$$
- dealing with the quadratic polynomials requires us to write them in a factorized form  $p(s, t) = (s - s^+(t))(s - s^-(t))$
- most intersection numbers only need the leading order in this pole, however for differential matrices we need second order
  - $\rightarrow$  very tedious expansion
  - $\rightarrow$  actually, everything beyond leading order turns out to be just a consistency condition
- we find canonical differential equations for the masters that can be integrated

# Computing the Gluing Integrals

- our integrals have no twist  $\rightarrow$  can be introduced by defining
$$u = (s(1-s)t(1-t) \prod p_i(s, t))^\gamma$$
$$\rightarrow \text{the limit } \gamma \rightarrow 0 \text{ yields the correct integral}$$
- dealing with the quadratic polynomials requires us to write them in a factorized form  $p(s, t) = (s - s^+(t))(s - s^-(t))$
- most intersection numbers only need the leading order in this pole, however for differential matrices we need second order
  - $\rightarrow$  very tedious expansion
  - $\rightarrow$  actually, everything beyond leading order turns out to be just a consistency condition
- we find canonical differential equations for the masters that can be integrated
- putting everything together in the end, we can confirm the expected identity
$$S_2 = S_1(a \leftrightarrow b, y \leftrightarrow z)$$

# Summary

- we solved first gluing integrals (without inherent twist) using intersection theory

# Summary

- we solved first gluing integrals (without inherent twist) using intersection theory
- proof of principle



# Summary

- we solved first gluing integrals (without inherent twist) using intersection theory
- proof of principle
  - allows generalisation to more general correlators

# Summary

- we solved first gluing integrals (without inherent twist) using intersection theory
- proof of principle  
→ allows generalisation to more general correlators
- other matrix elements (i.e. bound state compositions) can be computed similarly

# Summary

- we solved first gluing integrals (without inherent twist) using intersection theory
- proof of principle
  - allows generalisation to more general correlators
- other matrix elements (i.e. bound state compositions) can be computed similarly
- doable in reasonable time-frame on a laptop

# Summary

- we solved first gluing integrals (without inherent twist) using intersection theory
- proof of principle
  - allows generalisation to more general correlators
- other matrix elements (i.e. bound state compositions) can be computed similarly
- doable in reasonable time-frame on a laptop
  - seconds to few minutes for most intersection numbers

# Summary

- we solved first gluing integrals (without inherent twist) using intersection theory
- proof of principle
  - allows generalisation to more general correlators
- other matrix elements (i.e. bound state compositions) can be computed similarly
- doable in reasonable time-frame on a laptop
  - seconds to few minutes for most intersection numbers
  - 10s of minutes to few hours for differential-equation matrices

# Thank You!