

Wall crossing structure from quantum phenomena to Feynman Integrals

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Domodossola 14/07/2025

Based on

arXiv: 2506.03252

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UNIVERSITÀ DEGLI STUDI
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- Introduction and motivations
- Exponential integrals for holomorphic functions: Pearcey Integral
- Exponential integrals for closed forms: Legendre Family
- Conclusions and Outlooks

Exponential integrals

Consider the exponential integral:

$$I(\gamma) = \int_{\Gamma} e^{-\gamma f} \mu$$

Where:

- $f : X \rightarrow \mathbb{C}$
- Γ is a n -chain with $D_0 \supset \partial\Gamma \neq 0$
- $\gamma \in \mathbb{C}^*$
- μ is an algebraic volume n -form

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- Compute the intersection numbers
- Analyze the dependence on γ : wall crossing structure

Multiloop Feynman integrals in Baikov representation

$$I = \int_{\Delta} \mathcal{B}(x_i)^{-\gamma} \omega$$

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We lose geometric interpretation

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$$\int \exp[f] \mu : H_{\bullet}^{Betti, global}(X, D_0, f) \otimes H_{dR, global}^{\bullet}(X, D_0, f) \rightarrow \mathbb{C}$$

[Kontsevich and Soibelman: arXiv:2402.07343, 2024]

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Global Betti homology

$$H_{\text{Betti}, \text{global}}^{\bullet}((X, D_0), f, \mathbb{Z}) \equiv H^{\bullet}((X, D_0), f^{-1}(\infty), \mathbb{Z})$$

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Global Twisted de Rham cohomology

$$H_{dR}^{\bullet}(X, D_0, f) \cong \mathbb{H}^{\bullet}(X, \Omega_{X, D_0}^{\bullet}, \nabla_f)$$

With

$$(\Omega_{X, D_0}^{\bullet}, \nabla_f) : \Omega_{X, D_0}^0 \xrightarrow{\nabla_f} \Omega_{X, D_0}^1 \xrightarrow{\nabla_f} \dots \xrightarrow{\nabla_f} \Omega_{X, D_0}^n$$

$$\nabla_f = d - df \wedge$$

Holomorphic exponent: Morse (Picard-Lefschetz) Theory

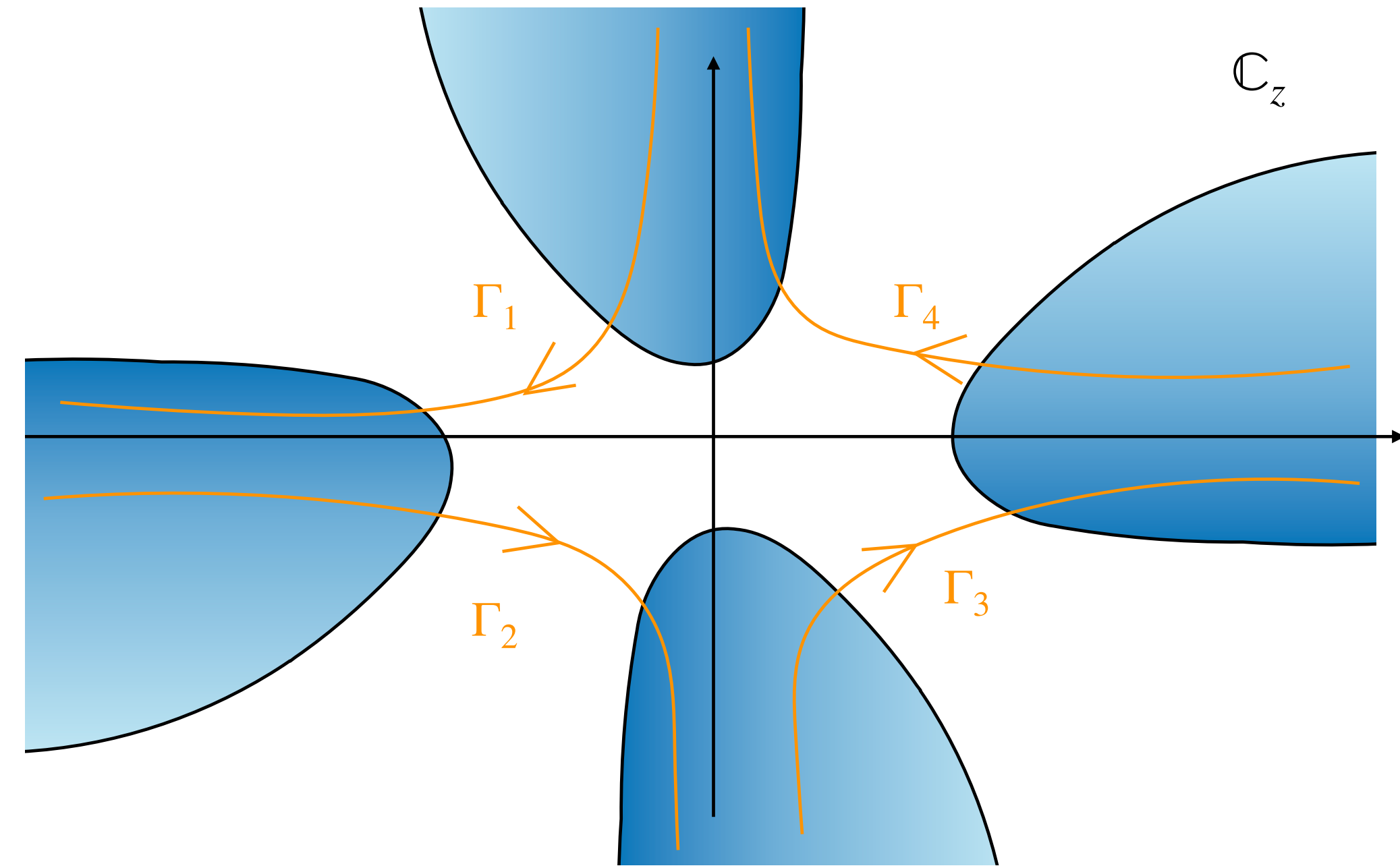
Let $X \cong \mathbb{C}^n$ and consider the set:

$$X \supset D_N = \{z \in \mathbb{C}^n \mid \operatorname{Re}(\gamma f(z)) \geq N\}$$

Any reasonable cycle Γ should connect two distinct regions in D_N :

$$\Gamma \in H_n(X, D_N, \mathbb{Z})$$

Moreover, in order to avoid oscillations $\operatorname{Im}(\gamma f)$ must remain constant along Γ

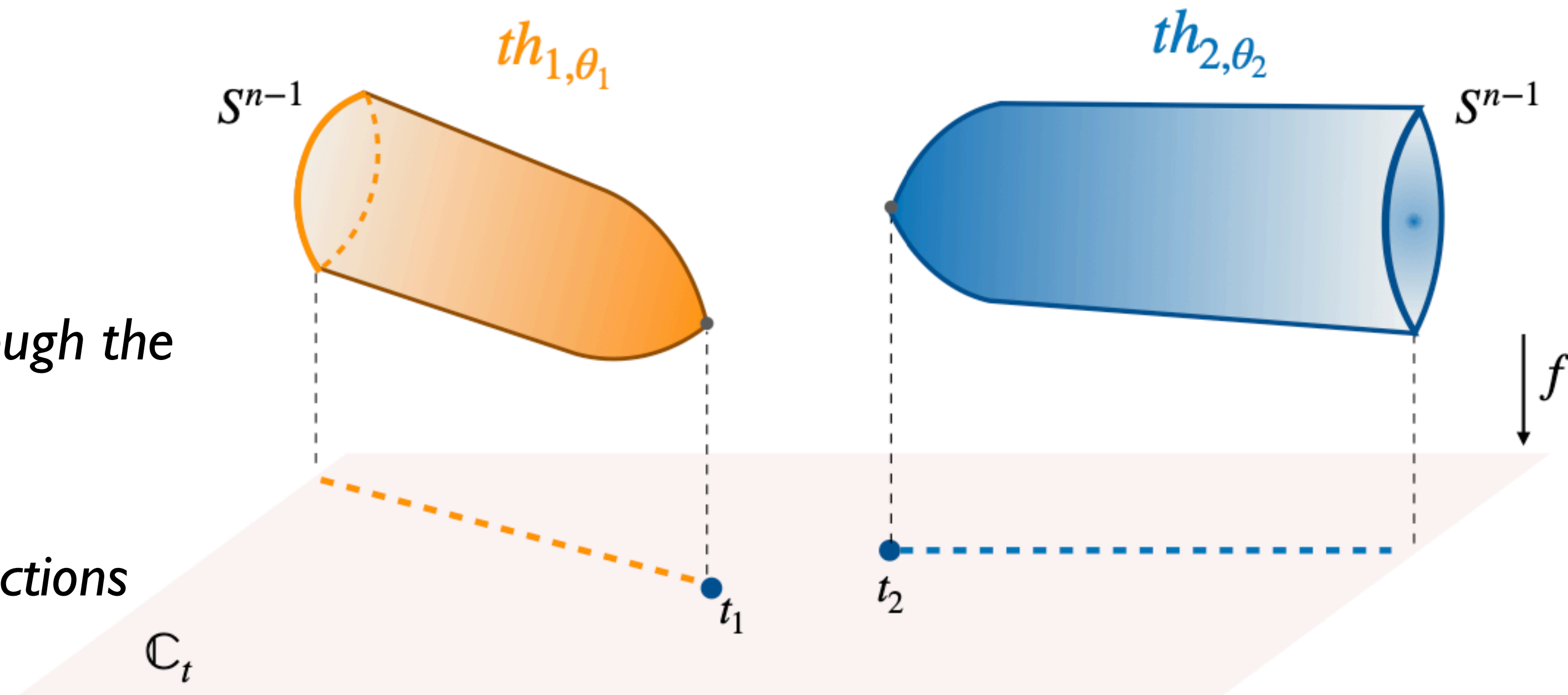


Thimbles decomposition

A good basis for $H_n(X, D_N, \mathbb{Z})$ is provided by the so called Lefschetz thimbles:



- Gradient flow lines with constant phase passing through the critical points of $h = \operatorname{Re}(\gamma f(z))$
- Ascendent paths Γ_i^+ and descendent paths Γ_i^-
- “Traces” of the vanishing cycles along vanishing directions



$$\operatorname{rank}[H_k(X, D_N, \mathbb{Z})] = \begin{cases} 0, & k < n, \\ \#\Sigma, & k = n. \end{cases}$$



Set of critical points

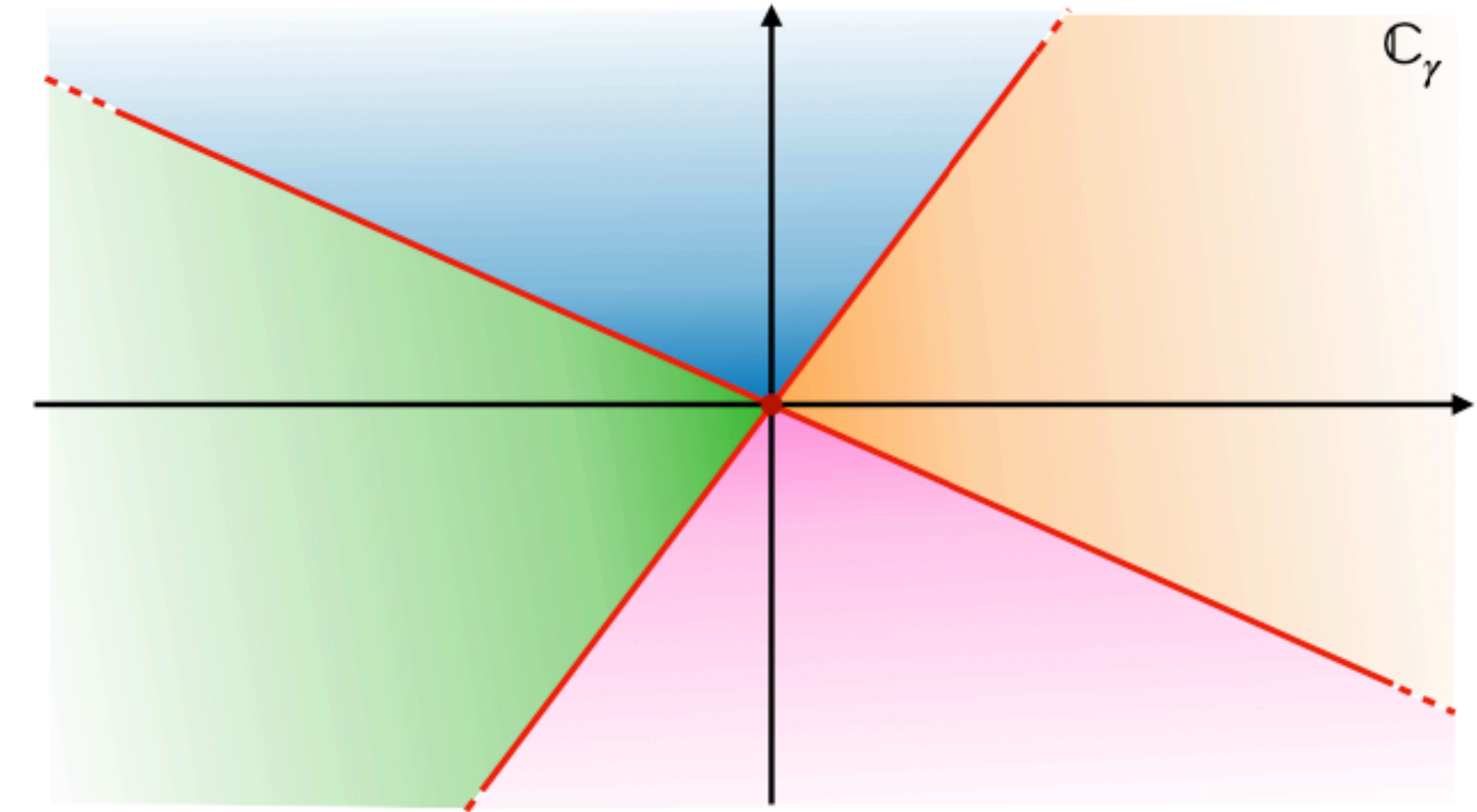
Stokes Phenomena

As γ varies it may happen that a thimble crosses more than one singular point:

$$l = \{\gamma \in \mathbb{C}^* \mid \operatorname{Im}(\gamma f(z))|_{\sigma_i} = \operatorname{Im}(\gamma f(z))|_{\sigma_j}\} \quad \text{Stokes' line}$$

The number of Thimbles for such values of γ is less than for generic γ

The space \mathbb{C}_γ splits into regions separated by Stokes lines

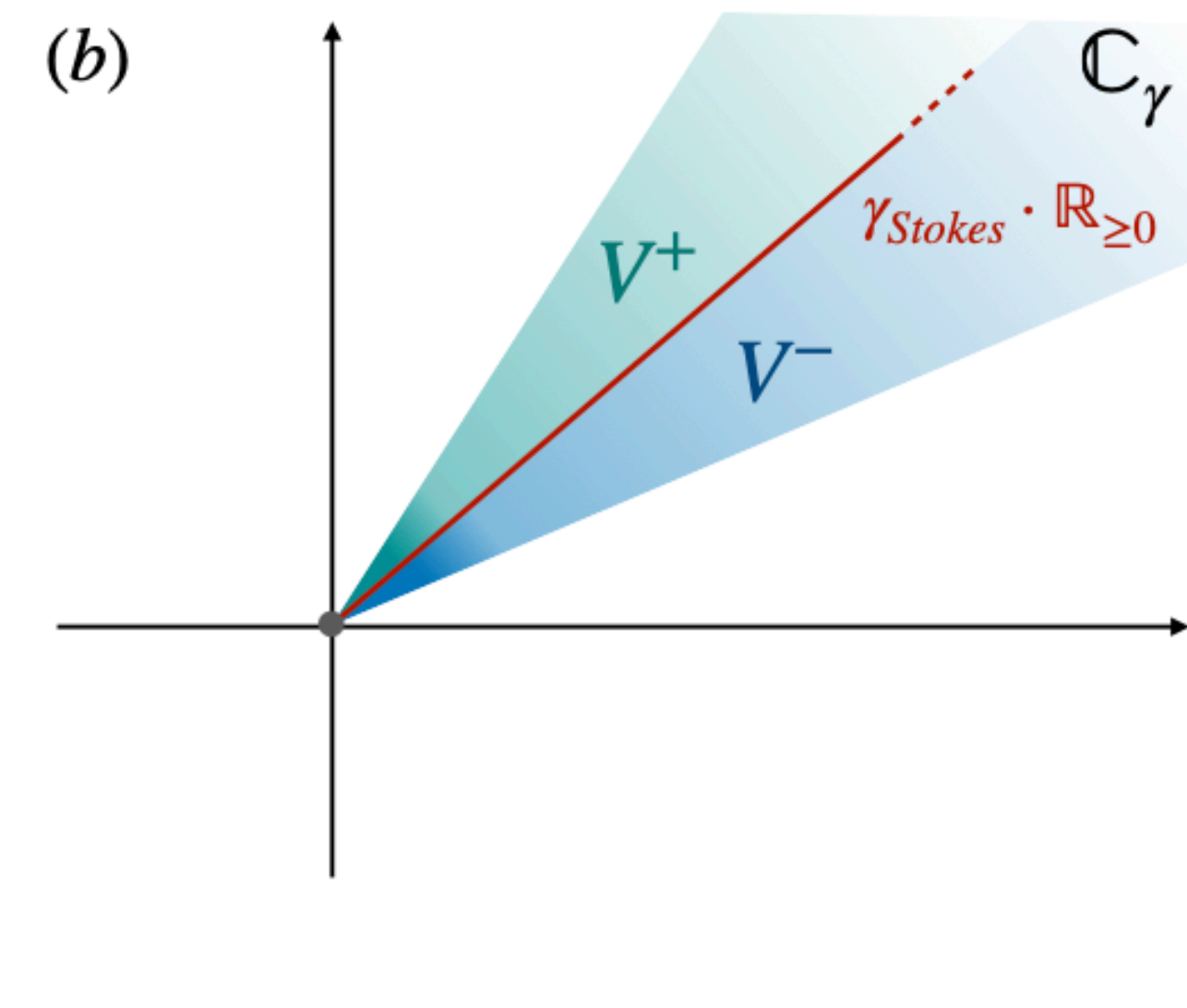
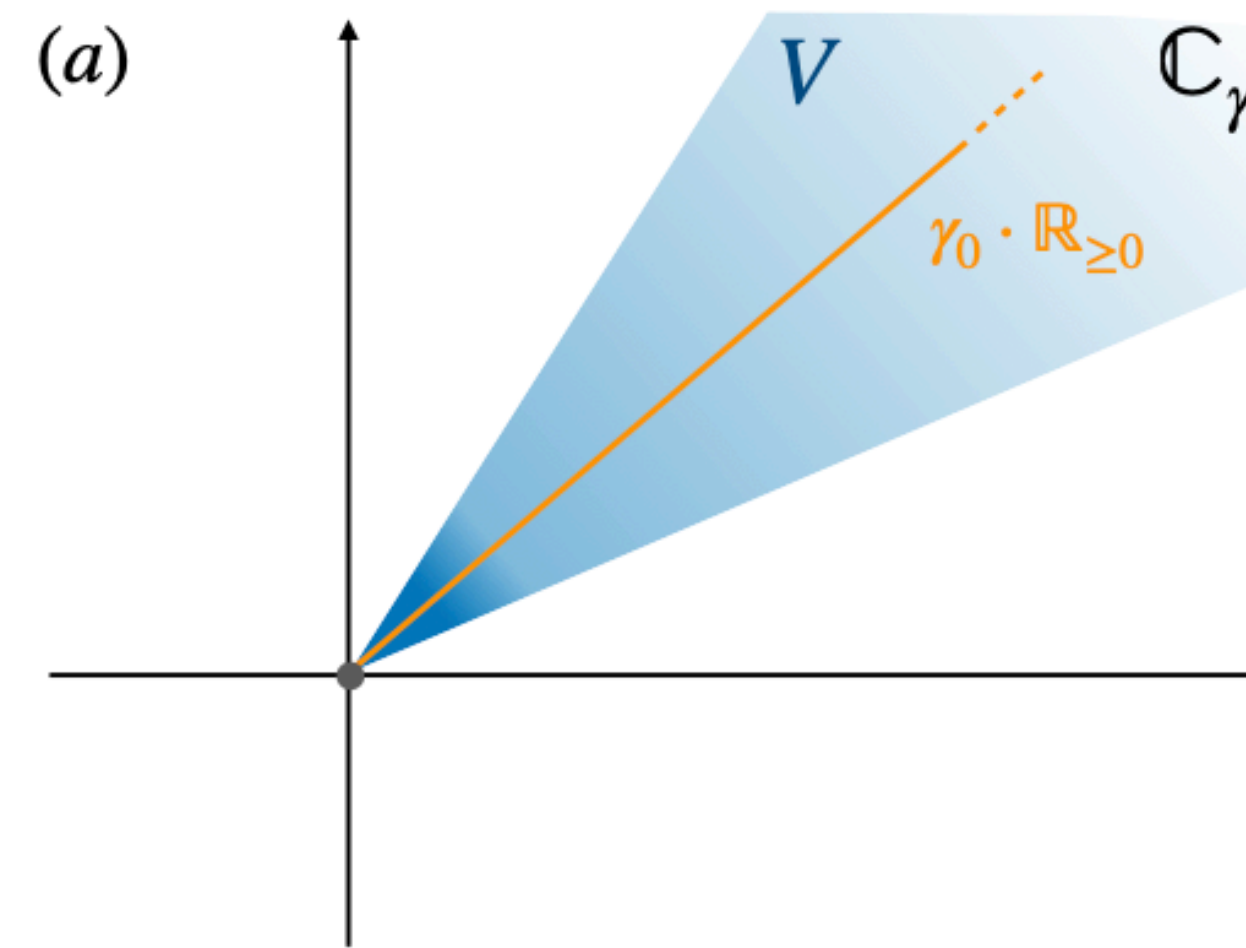


Wall crossing structure

As we cross a Stokes ray, associated with a Stokes line between the critical points σ_i and σ_j the corresponding thimbles Γ_j^+ and Γ_i^+ undergo a discontinuous jump to the adjacent region of the form:

$$\begin{pmatrix} \Gamma_i^{+(1)} \\ \Gamma_j^{+(1)} \end{pmatrix} = \begin{pmatrix} 1 & \Delta_{ij} \\ 0 & 1 \end{pmatrix} \begin{pmatrix} \Gamma_i^{+(0)} \\ \Gamma_j^{+(0)} \end{pmatrix}, \quad \text{for } h_{\sigma_i} < h_{\sigma_j}$$

Where $\Delta_{ij} = (\pm 1)\Delta_i \circ \Delta_j$.

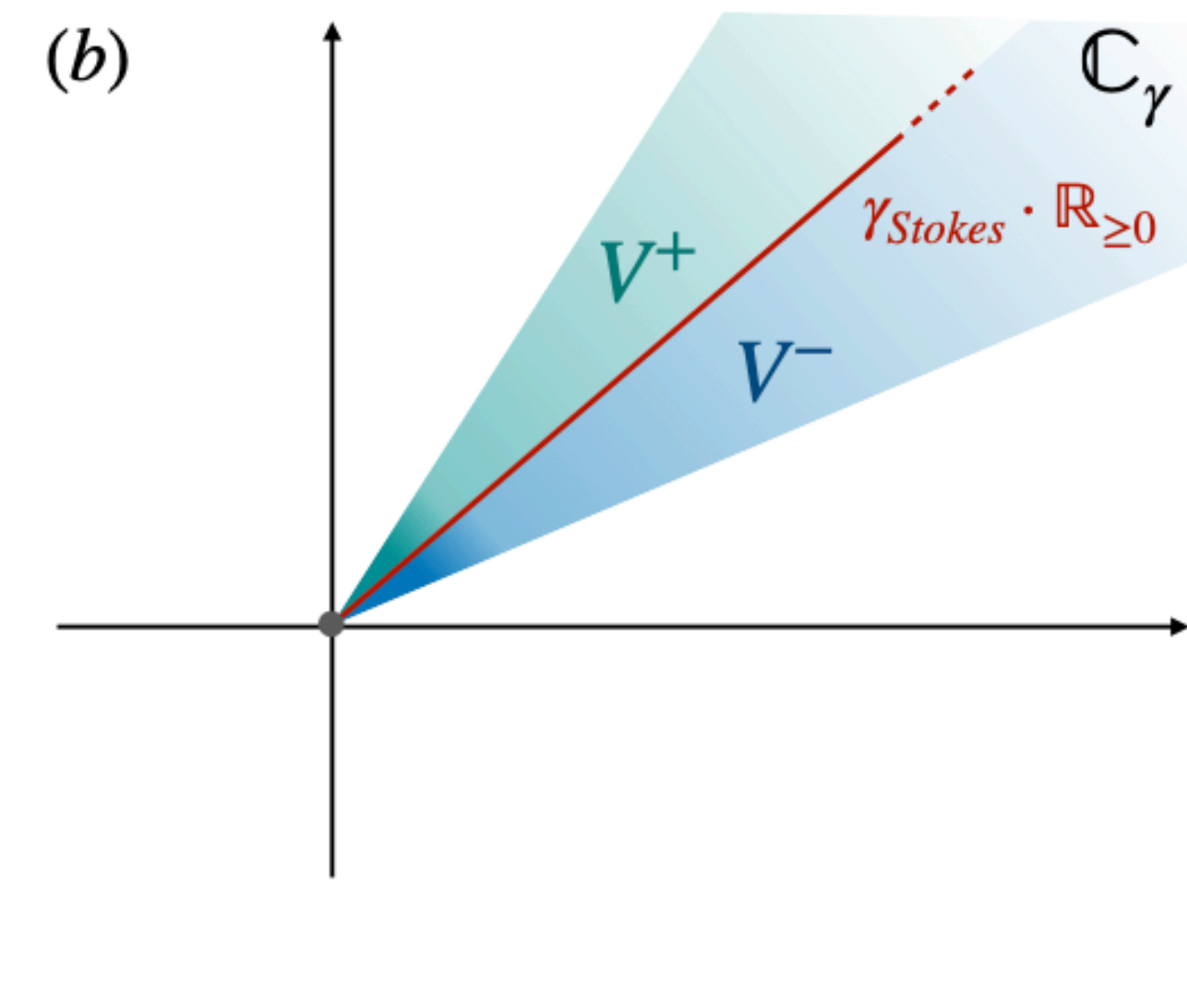
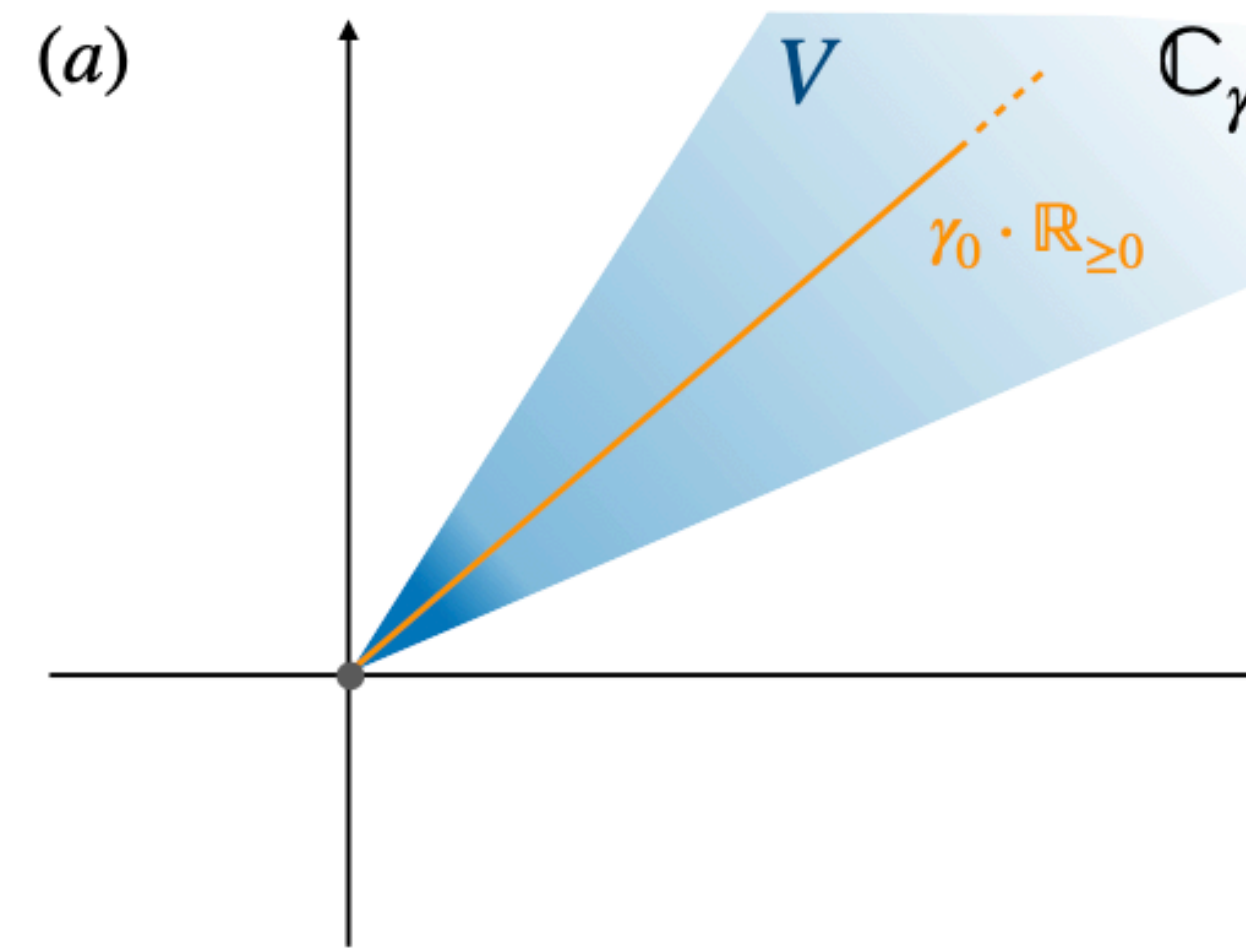


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Where $\Delta_{ij} = (\pm 1) \Delta_i \circ \Delta_j$  Intersection among the corresponding vanishing cycles

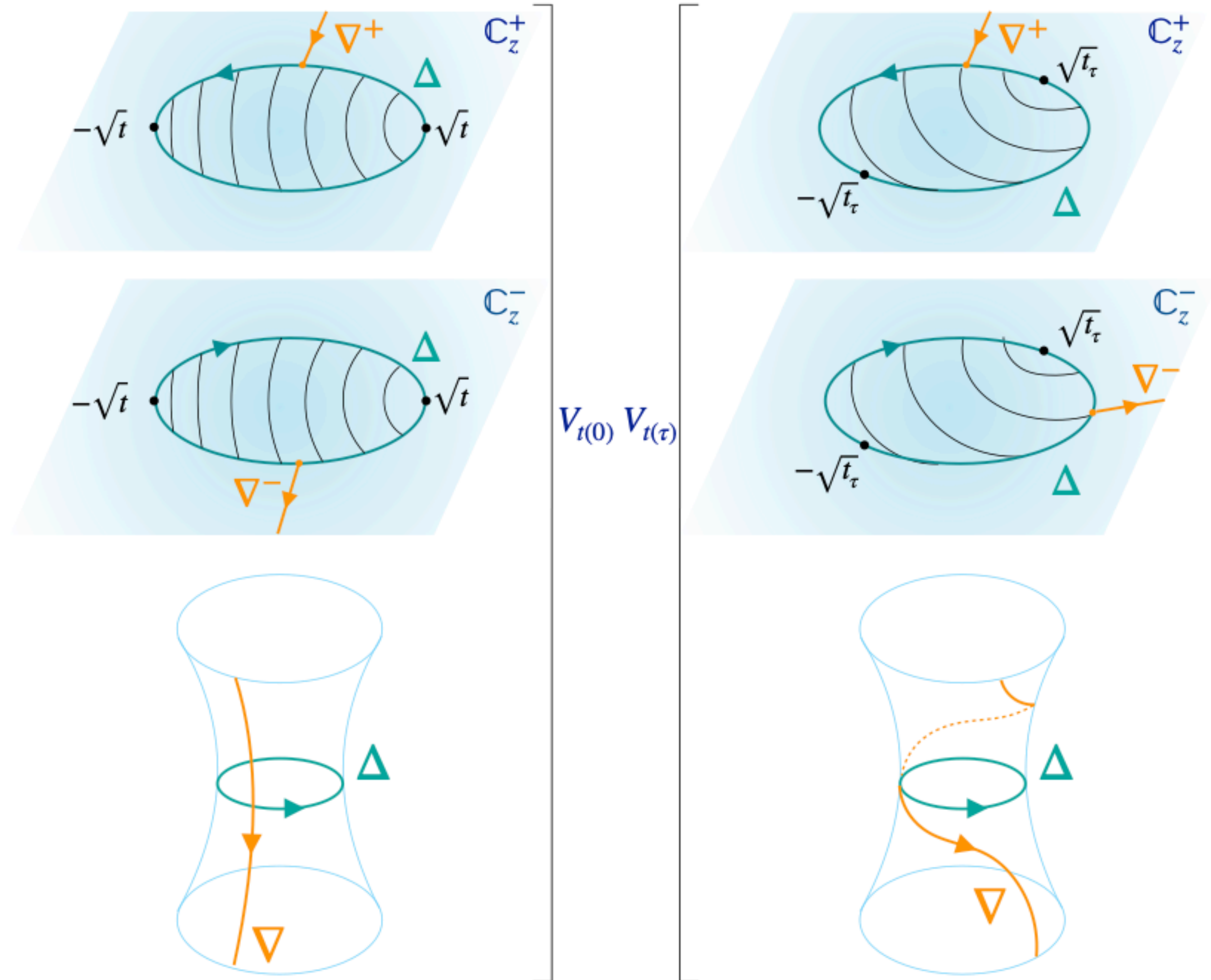


Thimbles intersection

Picard-Lefschetz theorem:

$$h_\tau(a) = a + (-1)^{n(n+1)/2} (a \circ \Delta) \Delta$$

Representation of the monodromy acting on vanishing cycles



Let consider $f = \mathcal{P}_\ell \in \mathbb{C}[z]$ and extend it to $\overline{\mathcal{P}}_\ell : \mathbb{P}^1 \rightarrow \mathbb{P}^1$.

We want to compute the Hypercohomology of the complex:

$$(\Omega_{\mathbb{P}^1, p}^\bullet, \nabla) : 0 \rightarrow \mathcal{O}_{\mathbb{P}^1}(*p) \xrightarrow{\nabla} \Omega_{\mathbb{P}^1}^1(*p) \rightarrow 0$$

With respect to the differential:

$$\nabla \equiv (\gamma^{-1}d + d\mathcal{P}_\ell \wedge)$$

Holomorphic exponent: Twisted de Rham cohomology

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Skipping the details, the final result is

$$H_{dR}^0(\mathbb{C}, d\mathcal{P}_\ell) \cong 0$$

$$H_{dR}^1(\mathbb{C}, d\mathcal{P}_\ell) \cong \mathbb{C}^{\ell-1}$$

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This result depends on the global Jacobian ring

$$J_{\mathcal{P}_\ell} \cong \mathbb{C}^{\ell-1}$$

that is NOT sensitive to the coalescence of critical points! If one point as multiplicity m , by considering the middle extension of the connexion we get

$$H_{dR}^1(\mathbb{C}, d\mathcal{P}_\ell)_{(m)} \cong \mathbb{C}^{\ell-m}$$

$$P(\gamma) = \int_{\Gamma} e^{-\gamma(z^4 + bz^2 + cz + d)} dz$$

Describing the Grand-canonical partition function of gauge Skyrme models for nuclear matter

[Cacciatori, Canfora, Lagos, Muscolino, Vera: JHEP 12:150, 2021]

The singular locus is the set

$$\Sigma = \{z \in \mathbb{C} \mid f'(z) = 0\} = \{\sigma_i(\Delta)\}$$

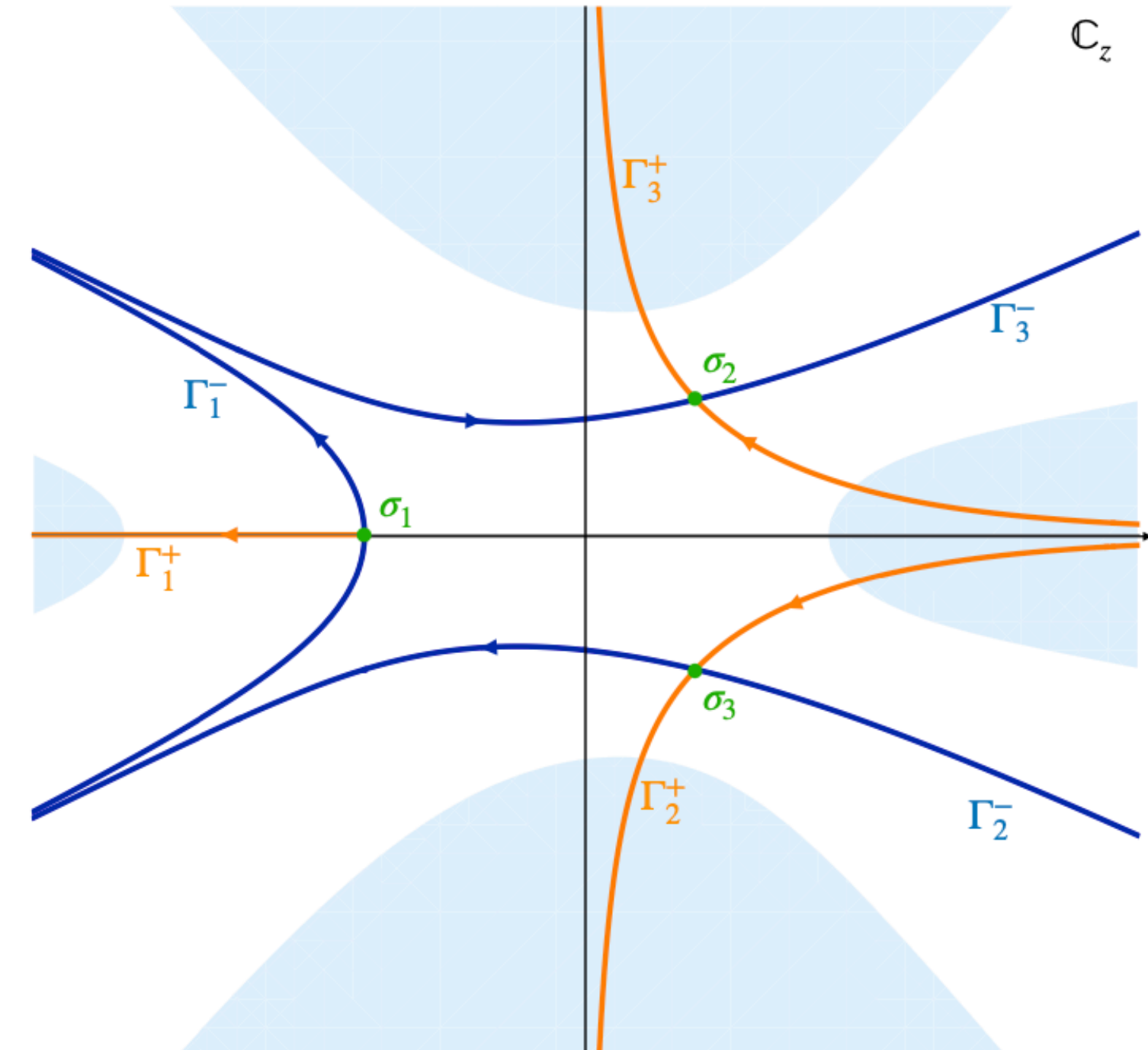
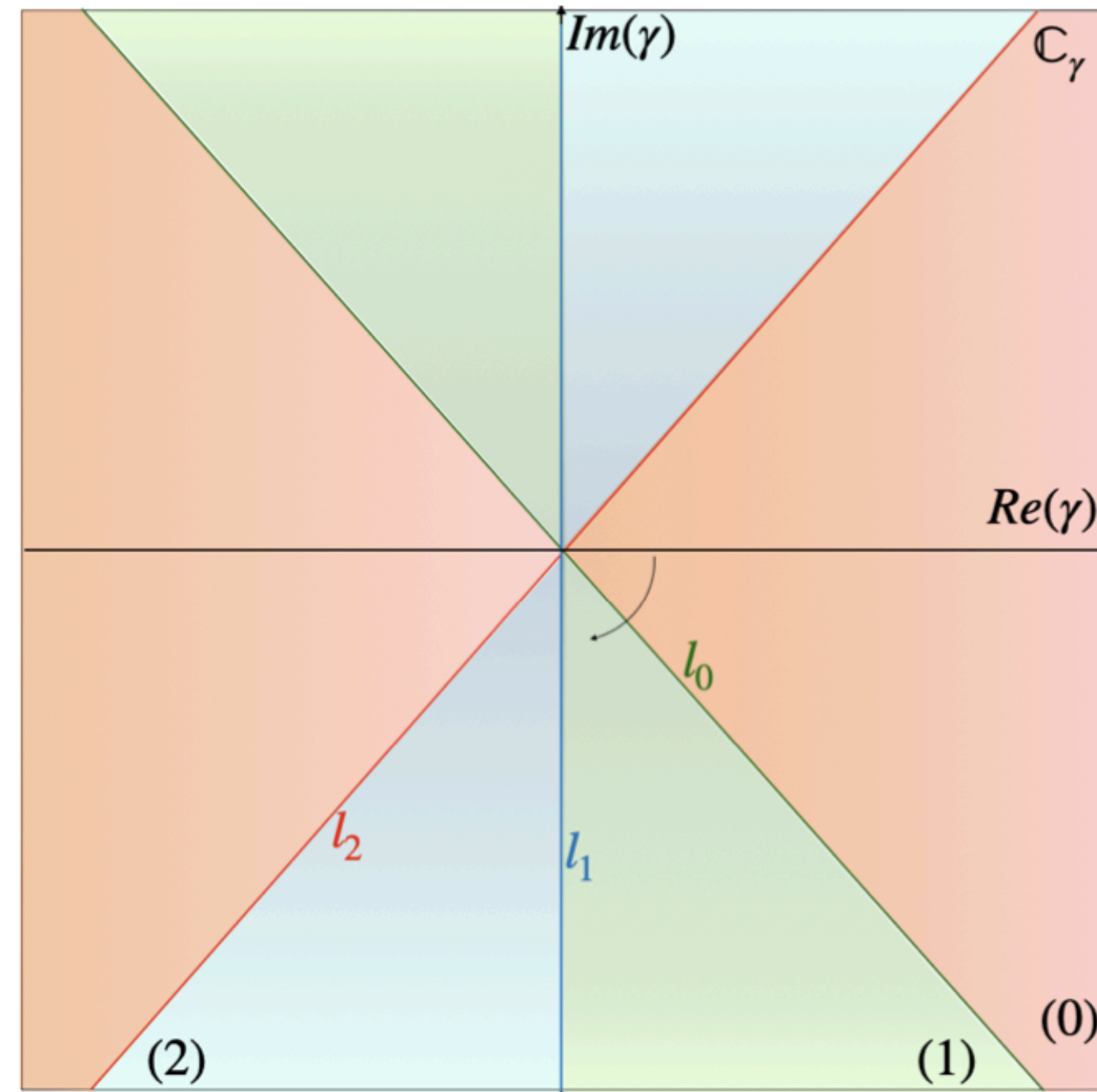
$$\text{Where } \Delta = 8b^2 + 27c^2$$

$$\Delta \equiv \begin{cases} > 0 & 1 \text{ real and 2 complex conjugate solutions,} \\ < 0 & 3 \text{ real different solutions,} \\ = 0 & 3 \text{ real solutions with at least a multiple root.} \end{cases}$$

Pearcey Integral $\Delta > 0$

$$\begin{aligned}
 l_0 : Re(\gamma) &= -\frac{11}{16}\sqrt{\frac{3}{2}}Im(\gamma), & \text{where } Im(\gamma f(z))|_{\sigma_1} &= Im(\gamma f(z))|_{\sigma_2}, \\
 l_1 : Re(\gamma) &= 0, & \text{where } Im(\gamma f(z))|_{\sigma_2} &= Im(\gamma f(z))|_{\sigma_3}, \\
 l_2 : Re(\gamma) &= \frac{11}{16}\sqrt{\frac{3}{2}}Im(\gamma), & \text{where } Im(\gamma f(z))|_{\sigma_1} &= Im(\gamma f(z))|_{\sigma_3}.
 \end{aligned}$$

Fix $\gamma \notin l_i$



Consider the preimage of the level set $f(z) = t$

$$f^{-1}(t) = \begin{pmatrix} z_1(t) \\ z_2(t) \\ z_3(t) \\ z_4(t) \end{pmatrix}$$

Consider the preimage of the level set $f(z) = t$

Vanishing cycles

$$\Delta_1 = \{z_3\} - \{z_4\} \quad \text{and} \quad \Delta_2 = \Delta_3 = \{z_1\} - \{z_4\}$$

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Monodromy matrices

$$M_1 = \begin{pmatrix} -1 & 0 & 0 \\ -1 & 1 & 0 \\ 1 & 0 & 1 \end{pmatrix} \quad \text{and} \quad M_2 = M_3 = \begin{pmatrix} 1 & -1 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & -1 \end{pmatrix}$$

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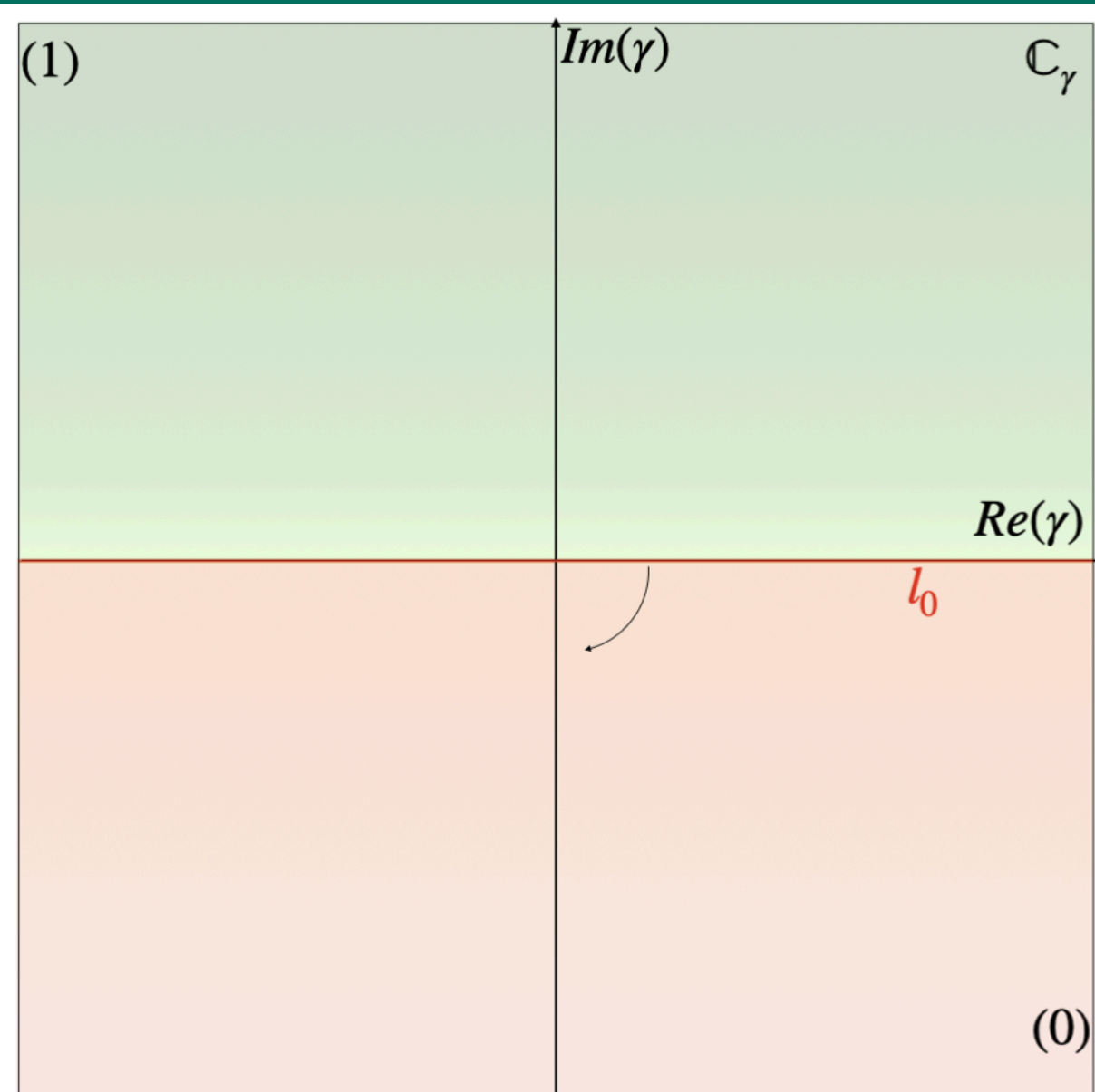
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Jump matrices

$$T^{(0)} = \begin{pmatrix} 1 & 1 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix} \quad T^{(1)} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & -2 & 1 \end{pmatrix} \quad T^{(2)} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 1 & 0 & 1 \end{pmatrix}$$

Pearcey Integral: $\Delta < 0$



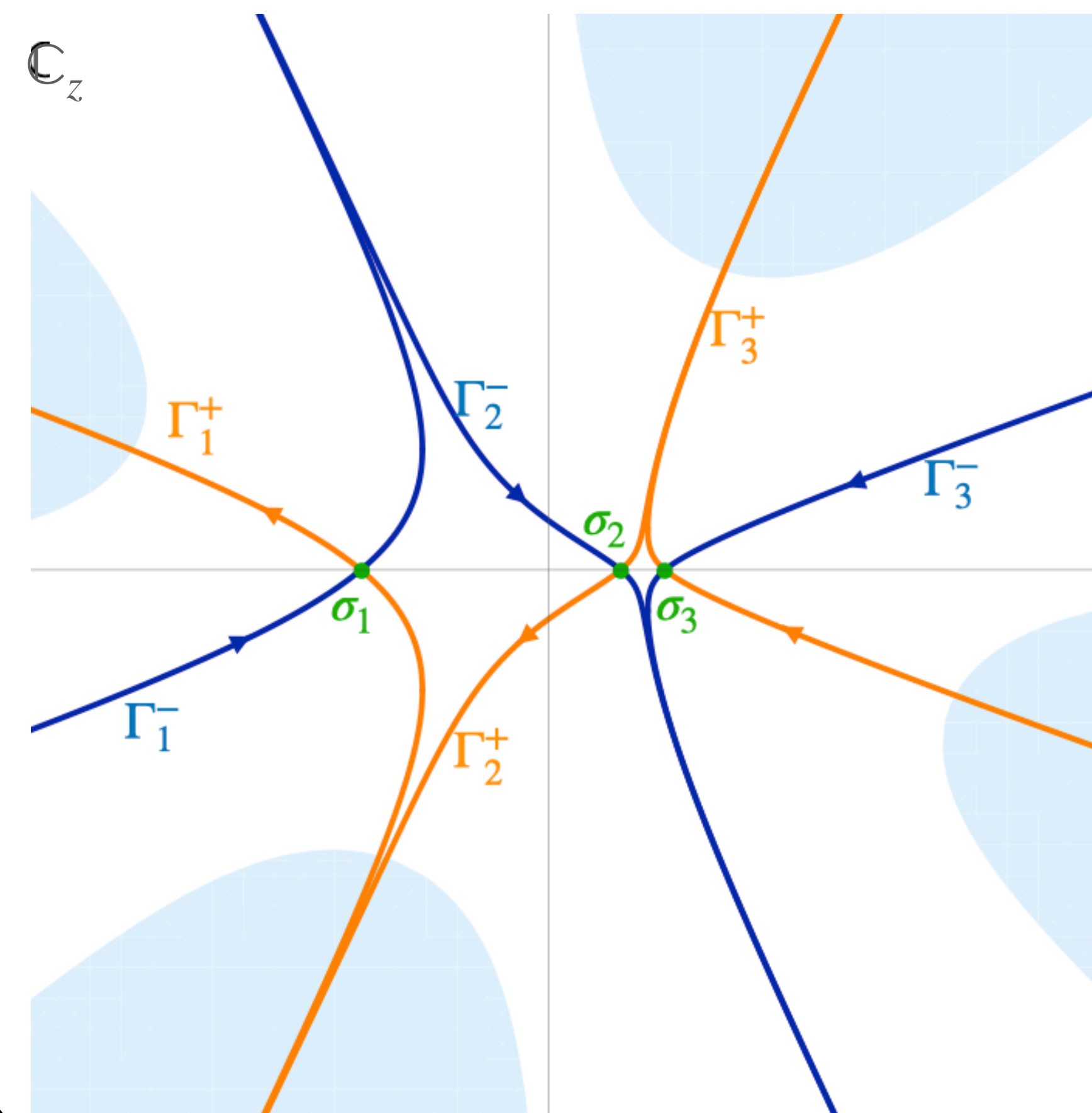
Stoke's line: $\text{Im}(\gamma) = 0$

Vanishing cycles

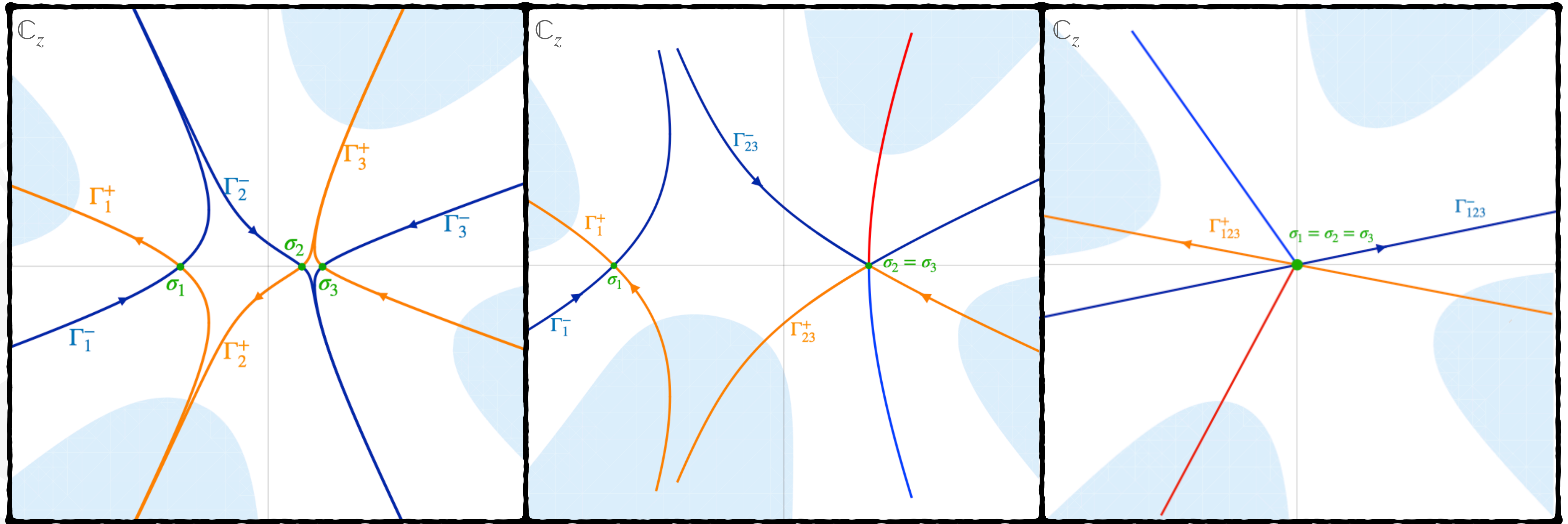
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Pearcey Integral



$$H_1(X, D_N, \mathbb{Z}) = \text{span}\{\Gamma_1^+, \Gamma_2^+, \Gamma_3^+\} \cong \mathbb{Z}^3$$

$$H_1(X, D_N, \mathbb{Z})^\vee = \text{span}\{\Gamma_1^-, \Gamma_2^-, \Gamma_3^-\} \cong \mathbb{Z}^3$$

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Logarithmic exponent

Let us consider now

$$f(z_1, \dots, z_n) = \log \mathcal{B}(z_1, \dots, z_n)$$

$$f \text{ is defined on } X = \mathbb{C}^n \setminus \{\mathcal{B} = 0\} \quad \longrightarrow \quad \bar{X} = \mathbb{P}^n = \mathbb{C}^n \cup \mathbb{P}^{n-1}$$

We classify the added divisor according to the behavior of $\overline{\mathcal{B}}$, the extension of \mathcal{B} over \bar{X} :

$$\bar{X} - X = D_h \cup D_v \cup D_{\log}$$

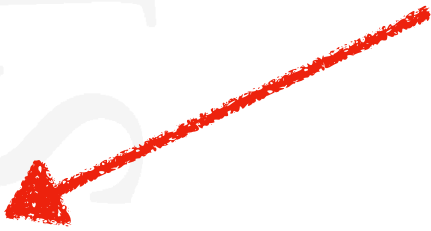
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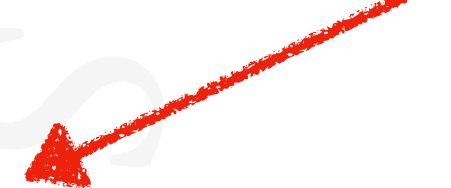
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Horizontal divisor: $\overline{\mathcal{B}}$ is finite



Vertical divisor(at infinity): $\overline{\mathcal{B}}$
has poles

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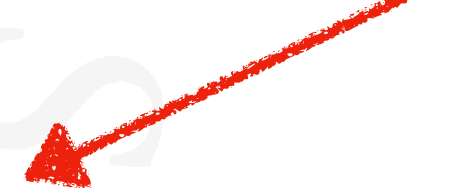
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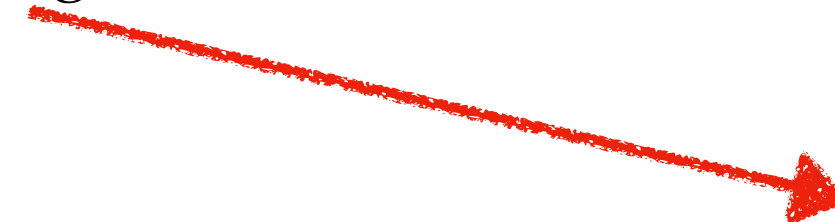
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Horizontal divisor: $\overline{\mathcal{B}}$ is finite



Vertical divisor (at infinity): $\overline{\mathcal{B}}$ has poles



Logarithmic divisor: $d\overline{\mathcal{B}}$ has logarithmic poles

Consider the local system.

$$\mathcal{L}_{\alpha,\gamma}(U) = \ker \nabla|_U, \quad U \subset X$$

with $\nabla = d + \gamma\alpha$

The global Betti cohomology is defined as

$$H^\bullet(X, \alpha) \cong H^\bullet\left(\widetilde{X}, D_v^{\mathbb{R}+} \cup D_{\log}^{\mathbb{R}-}, \Pi_*(\mathcal{L}_{\alpha,\gamma})\right)$$

We get:

$$H^\bullet(X, \alpha) \cong \begin{cases} H^\bullet\left(\tilde{X}, D_\infty^{\mathbb{R}}, \Pi_*(\mathcal{L}_{\alpha,\gamma})\right), & \operatorname{Re}(\gamma) > 0, \\ H^\bullet\left(\tilde{X}, D_{\mathcal{B}}^{\mathbb{R}}, \Pi_*(\mathcal{L}_{\alpha,\gamma})\right), & \operatorname{Re}(\gamma) < 0. \end{cases}$$

Consider the integral

$$\mathcal{J} = \int_{\Gamma} \frac{dx \wedge dy}{[y^2 + x(x-1)(x-\lambda)]^{\gamma}} = \int_{\Gamma} e^{-\gamma \log[y^2 + x(x-1)(x-\lambda)]} dx \wedge dy = \int_{\Gamma} e^{-\gamma \log \mathcal{B}(x,y;\lambda)} dx \wedge dy$$

We extend \mathcal{B} to \mathbb{P}^2 :

$$\overline{\mathcal{B}}(x, y, \eta; \lambda) = y^2 \eta - x(x - \eta)(x - \eta \lambda)$$

And define the close form

$$d \log \overline{\mathcal{B}} = \frac{2\eta y dy + [y^2 + x^2 + x\lambda(x - 2\eta)]d\eta + [-3x^2 - \eta^2\lambda + 2x\eta(1 + \lambda)]dx}{y^2 \eta - x(x - \eta)(x - \eta \lambda)}$$

Analyzing the behavior of $\overline{\mathcal{B}}$ we get

$$D_h = D_v = \emptyset$$

$$D_{log} = D_{\overline{\mathcal{B}}} \cup D_{\infty}$$

With:

$$D_{\overline{\mathcal{B}}} = \overline{\mathcal{E}}_{\lambda} = \{[x : y : \eta] \in \mathbb{P}^2 \mid \overline{\mathcal{B}} = 0\}$$

$$D_{\infty} = \mathbb{P}^1 = \{[x : y : 0] \in \mathbb{P}^2\}$$

Intersecting at $D_{\overline{\mathcal{B}}} \cap D_{\infty} = [0 : 1 : 0]$

The computation of Betti cohomology gives:

$$H_{Betti, glob, \gamma}^{\bullet}(X, \alpha)(\widetilde{X}, S^3, \Pi_*(\mathcal{L}_{\alpha, \gamma})) \cong 0 \oplus \mathbb{C}^2 \oplus \mathbb{C}^2 \oplus 0 \oplus \mathbb{C}^2$$

$$H_{Betti, glob, \gamma}^{\bullet}(X, \alpha)(\widetilde{X}, D_{\overline{\mathcal{B}}}^R, \Pi_*(\mathcal{L}_{\alpha, \gamma})) \cong 0 \oplus 0 \oplus \mathbb{C}^2 \oplus \mathbb{C}^4 \oplus 0$$

- The identification of exponential integrals as periods of twisted de Rham cohomology and Betti homology over complex manifolds, allows to accomodate in the same framework a wide range of physically relevant integrals
- The analysis of the wall crossing structure allows to analytic continue the master integral decomposition in the parameter γ and carefully avoid Stokes' phenomena for a sharp counting of the co-homology dimension

Outlooks

- *Concrete application of this formalism to Feynman integrals in different representation and multiple variables (working progress with Angius, Cacciatori, Mastrolia and Noja)*
- *Analysis of multi parameter dependence (Complex structure moduli of algebraic varieties = kinematic physical variables)*
- *Application to conformal correlators: string amplitudes*

Thank you

