



Finanziato
dall'Unione europea
NextGenerationEU



Ministero
dell'Università
e della Ricerca



Italiadomani
PIANO NAZIONALE
DI RIPRESA E RESILIENZA



Istituto Nazionale di Fisica Nucleare
Laboratori Nazionali di Frascati

Interplay of vertical and horizontal gauge symmetry for a high-quality axion

Vasja Susić

Laboratori Nazionali di Frascati (LNF), INFN

PRIN 2022K4B58X

2025-05-29

Based on 2503.16648 with L. Di Luzio, G. Landini, F. Mescia

PLANCK 2025 — Padova, May 26-30 2025

Motivation — Peccei-Quinn and axions

■ Strong CP problem:

- Standard Model (SM) allows the term

$$\mathcal{L} \supset \frac{\alpha_s}{8\pi} \bar{\theta} G_{\mu\nu} \tilde{G}^{\mu\nu}; \quad \bar{\theta} = \theta + \arg \det M_q \quad (1)$$

- Bound from neutron EDM: $\boxed{\bar{\theta} \lesssim 10^{-10}}$. Why is $\bar{\theta}$ so small?

Motivation — Peccei-Quinn and axions

■ Strong CP problem:

→ Standard Model (SM) allows the term

$$\mathcal{L} \supset \frac{\alpha_s}{8\pi} \bar{\theta} G_{\mu\nu} \tilde{G}^{\mu\nu}; \quad \bar{\theta} = \theta + \arg \det M_q \quad (1)$$

→ Bound from neutron EDM: $\boxed{\bar{\theta} \lesssim 10^{-10}}$. Why is $\bar{\theta}$ so small?

■ Axion solution: via Peccei-Quinn (PQ) [1, 2, 3, 4]

- (1) Implement a global $U(1)_{PQ}$ anomalous under QCD
- (2) Break $U(1)_{PQ}$: low energy theory has a Goldstone boson a — the **axion**

$$\mathcal{L} \supset \frac{\alpha_s}{8\pi} (\bar{\theta} + N a/f_a) G_{\mu\nu} \tilde{G}^{\mu\nu} \quad (2)$$

- (3) Axion potential \mathcal{V}_a and mass m_a generated by QCD instantons.

$$\mathcal{V}_a(a) \text{ minimum: } \bar{\theta} + N \langle a \rangle / f_a = 0, \quad m_a f_a \simeq m_\pi f_\pi \quad (3)$$

- (4) Hence $\boxed{\bar{\theta}_{\text{eff}} = 0}$ and strong CP is conserved.

Motivation — PQ quality and model building

■ PQ quality:

- $U(1)_{PQ}$ is a **global symmetry**, thus **not fundamental**
- **PQ**: non-renormalizable \mathcal{O} from e.g. gravity (M_{Planck} suppressed)
- contribution to axion potential \mathcal{V}_a : **shift in vacuum** $\langle \textcolor{violet}{a} \rangle \Rightarrow$ shift in $\bar{\theta}_{\text{eff}}$

Motivation — PQ quality and model building

■ PQ quality:

- $U(1)_{PQ}$ is a **global symmetry**, thus **not fundamental**
- ~~PQ~~: non-renormalizable \mathcal{O} from e.g. gravity (M_{Planck} suppressed)
- contribution to axion potential \mathcal{V}_a : **shift in vacuum** $\langle \alpha \rangle \Rightarrow$ shift in $\bar{\theta}_{\text{eff}}$

■ Goal for model building:

- $U(1)_{PQ}$: accidental at renormalizable level (not imposed)
- **high quality** PQ: $\bar{\theta}_{\text{eff}} < 10^{-10}$ from ~~PQ~~ operators

Motivation — PQ quality and model building

■ PQ quality:

- $U(1)_{PQ}$ is a **global symmetry**, thus **not fundamental**
- PQ : non-renormalizable \mathcal{O} from e.g. gravity (M_{Planck} suppressed)
- contribution to axion potential \mathcal{V}_a : **shift in vacuum** $\langle \textcolor{violet}{a} \rangle \Rightarrow$ shift in $\bar{\theta}_{\text{eff}}$

■ Goal for model building:

- $U(1)_{PQ}$: accidental at renormalizable level (not imposed)
- **high quality** PQ: $\bar{\theta}_{\text{eff}} < 10^{-10}$ from PQ operators

■ This talk: PQ quality from **vertical** and **horizontal** symmetry

- More strongly connected to SM than adding unrelated symmetry
- Vertical and horizontal symmetry for Yukawa sector: Berezhiani [7, 8, 9]
- Vertical for quality: Vecchi [10], Babu+ [11]
- Horizontal for quality: Darmé+ [12, 13]

Motivation — PQ quality and model building

■ PQ quality:

- $U(1)_{PQ}$ is a **global symmetry**, thus **not fundamental**
- ~~PQ~~: non-renormalizable \mathcal{O} from e.g. gravity (M_{Planck} suppressed)
- contribution to axion potential \mathcal{V}_a : **shift in vacuum** $\langle \textcolor{violet}{a} \rangle \Rightarrow$ shift in $\bar{\theta}_{\text{eff}}$

■ Goal for model building:

- $U(1)_{PQ}$: accidental at renormalizable level (not imposed)
- **high quality** PQ: $\bar{\theta}_{\text{eff}} < 10^{-10}$ from ~~PQ~~ operators

■ This talk: PQ quality from **vertical** and **horizontal** symmetry

- Our model [5]: $SU(4)_{PS} \times SU(2)_L \times SU(2)_R \times SU(3)_{f_R}$

Motivation — PQ quality and model building

■ PQ quality:

- $U(1)_{PQ}$ is a **global symmetry**, thus **not fundamental**
- ~~PQ~~: non-renormalizable \mathcal{O} from e.g. gravity (M_{Planck} suppressed)
- contribution to axion potential \mathcal{V}_a : **shift in vacuum** $\langle \textcolor{violet}{a} \rangle \Rightarrow$ shift in $\bar{\theta}_{\text{eff}}$

■ Goal for model building:

- $U(1)_{PQ}$: accidental at renormalizable level (not imposed)
- **high quality** PQ: $\bar{\theta}_{\text{eff}} < 10^{-10}$ from ~~PQ~~ operators

■ This talk: PQ quality from **vertical** and **horizontal** symmetry

- Our model [5]: $SU(4)_{PS} \times SU(2)_L \times SU(2)_R \times SU(3)_{f_R}$
- Inspired by $SO(10) \times SU(3)_f$ from Di Luzio [6]

Motivation — PQ quality and model building

■ PQ quality:

- $U(1)_{PQ}$ is a **global symmetry**, thus **not fundamental**
- ~~PQ~~: non-renormalizable \mathcal{O} from e.g. gravity (M_{Planck} suppressed)
- contribution to axion potential \mathcal{V}_a : **shift in vacuum** $\langle \textcolor{violet}{a} \rangle \Rightarrow$ shift in $\bar{\theta}_{\text{eff}}$

■ Goal for model building:

- $U(1)_{PQ}$: accidental at renormalizable level (not imposed)
- **high quality** PQ: $\bar{\theta}_{\text{eff}} < 10^{-10}$ from ~~PQ~~ operators

■ This talk: PQ quality from **vertical** and **horizontal** symmetry

- Our model [5]: $SU(4)_{PS} \times SU(2)_L \times SU(2)_R \times SU(3)_{f_R}$
- Inspired by $SO(10) \times SU(3)_f$ from Di Luzio [6]
- Pati-Salam (PS) easier to **work out details**; only R -flavor gauged

Motivation — PQ quality and model building

■ PQ quality:

- $U(1)_{PQ}$ is a **global symmetry**, thus **not fundamental**
- ~~PQ~~: non-renormalizable \mathcal{O} from e.g. gravity (M_{Planck} suppressed)
- contribution to axion potential \mathcal{V}_a : **shift in vacuum** $\langle \alpha \rangle \Rightarrow$ shift in $\bar{\theta}_{\text{eff}}$

■ Goal for model building:

- $U(1)_{PQ}$: accidental at renormalizable level (not imposed)
- **high quality** PQ: $\bar{\theta}_{\text{eff}} < 10^{-10}$ from ~~PQ~~ operators

■ This talk: PQ quality from **vertical** and **horizontal** symmetry

- Our model [5]: $SU(4)_{PS} \times SU(2)_L \times SU(2)_R \times SU(3)_{f_R}$
- Inspired by $SO(10) \times SU(3)_f$ from Di Luzio [6]
- Pati-Salam (PS) easier to **work out details**; only R -flavor gauged
- **Interesting features**: anomalies, quality mechanism can be probed ...

Constructing the model

- Field content: $G_{PS} \equiv \text{SU}(4)_{PS} \times \text{SU}(2)_L \times \text{SU}(2)_R$

Field	Lorentz	$G_{PS} \times \text{SU}(3)_{f_R}$	copies	$\text{U}(1)_{\text{PQ}}$
\bar{Q}_L	(0, 1/2)	($\bar{\mathbf{4}}, \mathbf{2}, \mathbf{1}, \mathbf{1}$)	3	+3
Q_R	(0, 1/2)	($\mathbf{4}, \mathbf{1}, \mathbf{2}, \mathbf{3}$)	1	+1
Ψ	(0, 1/2)	($\mathbf{1}, \mathbf{1}, \mathbf{1}, \bar{\mathbf{3}}$)	8	+2
Φ	(0, 0)	($\mathbf{1}, \mathbf{2}, \mathbf{2}, \bar{\mathbf{3}}$)	$N_\Phi \geq 1$	+2
Σ	(0, 0)	($\mathbf{15}, \mathbf{2}, \mathbf{2}, \bar{\mathbf{3}}$)	$N_\Sigma \geq 2$	+2
Δ	(0, 0)	($\mathbf{10}, \mathbf{1}, \mathbf{3}, \mathbf{6}$)	1	+2
χ	(0, 0)	($\mathbf{4}, \mathbf{1}, \mathbf{2}, \bar{\mathbf{3}}$)	1	-1
ξ	(0, 0)	($\mathbf{15}, \mathbf{1}, \mathbf{3}, \mathbf{1}$)	1	0

- Not much freedom in choice of irreps for a realistic model...

Constructing the model

- Field content: $G_{\text{PS}} \equiv \text{SU}(4)_{\text{PS}} \times \text{SU}(2)_L \times \text{SU}(2)_R$

Field	Lorentz	$G_{\text{PS}} \times \text{SU}(3)_{f_R}$	copies	$\text{U}(1)_{\text{PQ}}$
$\rightarrow \bar{Q}_L$	(0, 1/2)	($\bar{\mathbf{4}}, \mathbf{2}, \mathbf{1}, \mathbf{1}$)	3	+3
$\rightarrow Q_R$	(0, 1/2)	($\mathbf{4}, \mathbf{1}, \mathbf{2}, \mathbf{3}$)	1	+1
Ψ	(0, 1/2)	($\mathbf{1}, \mathbf{1}, \mathbf{1}, \bar{\mathbf{3}}$)	8	+2
Φ	(0, 0)	($\mathbf{1}, \mathbf{2}, \mathbf{2}, \bar{\mathbf{3}}$)	$N_\Phi \geq 1$	+2
Σ	(0, 0)	($\mathbf{15}, \mathbf{2}, \mathbf{2}, \bar{\mathbf{3}}$)	$N_\Sigma \geq 2$	+2
Δ	(0, 0)	($\mathbf{10}, \mathbf{1}, \mathbf{3}, \mathbf{6}$)	1	+2
χ	(0, 0)	($\mathbf{4}, \mathbf{1}, \mathbf{2}, \bar{\mathbf{3}}$)	1	-1
ξ	(0, 0)	($\mathbf{15}, \mathbf{1}, \mathbf{3}, \mathbf{1}$)	1	0

- Fermions:**
 - $\rightarrow \bar{Q}_L$ and Q_R contain SM fermions and ν_R (standard in PS)
 - \rightarrow only Q_R transforms under gauged flavor, \bar{Q}_L in 3 copies

Constructing the model

- Field content: $G_{\text{PS}} \equiv \text{SU}(4)_{\text{PS}} \times \text{SU}(2)_L \times \text{SU}(2)_R$

Field	Lorentz	$G_{\text{PS}} \times \text{SU}(3)_{f_R}$	copies	$\text{U}(1)_{\text{PQ}}$
\bar{Q}_L	(0, 1/2)	($\bar{\mathbf{4}}, \mathbf{2}, \mathbf{1}, \mathbf{1}$)	3	+3
Q_R	(0, 1/2)	($\mathbf{4}, \mathbf{1}, \mathbf{2}, \mathbf{3}$)	1	+1
$\rightarrow \Psi$	(0, 1/2)	($\mathbf{1}, \mathbf{1}, \mathbf{1}, \bar{\mathbf{3}}$)	8	+2
Φ	(0, 0)	($\mathbf{1}, \mathbf{2}, \mathbf{2}, \bar{\mathbf{3}}$)	$N_\Phi \geq 1$	+2
Σ	(0, 0)	($\mathbf{15}, \mathbf{2}, \mathbf{2}, \bar{\mathbf{3}}$)	$N_\Sigma \geq 2$	+2
Δ	(0, 0)	($\mathbf{10}, \mathbf{1}, \mathbf{3}, \mathbf{6}$)	1	+2
χ	(0, 0)	($\mathbf{4}, \mathbf{1}, \mathbf{2}, \bar{\mathbf{3}}$)	1	-1
ξ	(0, 0)	($\mathbf{15}, \mathbf{1}, \mathbf{3}, \mathbf{1}$)	1	0

- Anomalons Ψ : \rightarrow Needed to cancel anomalies due to $\text{SU}(3)_{f_R}$ factor
 \rightarrow They are SM singlets

Constructing the model

- Field content: $G_{\text{PS}} \equiv \text{SU}(4)_{\text{PS}} \times \text{SU}(2)_L \times \text{SU}(2)_R$

Field	Lorentz	$G_{\text{PS}} \times \text{SU}(3)_{f_R}$	copies	$\text{U}(1)_{\text{PQ}}$
\bar{Q}_L	(0, 1/2)	($\bar{\mathbf{4}}, \mathbf{2}, \mathbf{1}, \mathbf{1}$)	3	+3
Q_R	(0, 1/2)	($\mathbf{4}, \mathbf{1}, \mathbf{2}, \mathbf{3}$)	1	+1
Ψ	(0, 1/2)	($\mathbf{1}, \mathbf{1}, \mathbf{1}, \bar{\mathbf{3}}$)	8	+2
\rightarrow				
Φ	(0, 0)	($\mathbf{1}, \mathbf{2}, \mathbf{2}, \bar{\mathbf{3}}$)	$N_\Phi \geq 1$	+2
\rightarrow	Σ	($\mathbf{15}, \mathbf{2}, \mathbf{2}, \bar{\mathbf{3}}$)	$N_\Sigma \geq 2$	+2
	Δ	($\mathbf{10}, \mathbf{1}, \mathbf{3}, \mathbf{6}$)	1	+2
	χ	($\mathbf{4}, \mathbf{1}, \mathbf{2}, \bar{\mathbf{3}}$)	1	-1
	ξ	($\mathbf{15}, \mathbf{1}, \mathbf{3}, \mathbf{1}$)	1	0

- Scalars with EW VEVs Φ and Σ : contain **SM Higgs**

→ terms $\bar{Q}_L Q_R \Phi$ and $\bar{Q}_L Q_R \Sigma$ for realistic Yukawa sector

Constructing the model

- Field content: $G_{\text{PS}} \equiv \text{SU}(4)_{\text{PS}} \times \text{SU}(2)_L \times \text{SU}(2)_R$

Field	Lorentz	$G_{\text{PS}} \times \text{SU}(3)_{f_R}$	copies	$\text{U}(1)_{\text{PQ}}$
\bar{Q}_L	$(0, 1/2)$	$(\bar{\mathbf{4}}, \mathbf{2}, \mathbf{1}, \mathbf{1})$	3	+3
Q_R	$(0, 1/2)$	$(\mathbf{4}, \mathbf{1}, \mathbf{2}, \mathbf{3})$	1	+1
Ψ	$(0, 1/2)$	$(\mathbf{1}, \mathbf{1}, \mathbf{1}, \bar{\mathbf{3}})$	8	+2
Φ	$(0, 0)$	$(\mathbf{1}, \mathbf{2}, \mathbf{2}, \bar{\mathbf{3}})$	$N_\Phi \geq 1$	+2
Σ	$(0, 0)$	$(\mathbf{15}, \mathbf{2}, \mathbf{2}, \bar{\mathbf{3}})$	$N_\Sigma \geq 2$	+2
\rightarrow				
Δ	$(0, 0)$	$(\mathbf{10}, \mathbf{1}, \mathbf{3}, \mathbf{6})$	1	+2
χ	$(0, 0)$	$(\mathbf{4}, \mathbf{1}, \mathbf{2}, \bar{\mathbf{3}})$	1	-1
ξ	$(0, 0)$	$(\mathbf{15}, \mathbf{1}, \mathbf{3}, \mathbf{1})$	1	0

- Scalar Δ : \rightarrow for ν Majorana mass term $Q_R Q_R \Delta^*$
 \rightarrow involved in PS and flavor breaking

Constructing the model

- Field content: $G_{\text{PS}} \equiv \text{SU}(4)_{\text{PS}} \times \text{SU}(2)_L \times \text{SU}(2)_R$

Field	Lorentz	$G_{\text{PS}} \times \text{SU}(3)_{f_R}$	copies	$\text{U}(1)_{\text{PQ}}$
\bar{Q}_L	(0, 1/2)	($\bar{\mathbf{4}}, \mathbf{2}, \mathbf{1}, \mathbf{1}$)	3	+3
Q_R	(0, 1/2)	($\mathbf{4}, \mathbf{1}, \mathbf{2}, \mathbf{3}$)	1	+1
Ψ	(0, 1/2)	($\mathbf{1}, \mathbf{1}, \mathbf{1}, \bar{\mathbf{3}}$)	8	+2
Φ	(0, 0)	($\mathbf{1}, \mathbf{2}, \mathbf{2}, \bar{\mathbf{3}}$)	$N_\Phi \geq 1$	+2
Σ	(0, 0)	($\mathbf{15}, \mathbf{2}, \mathbf{2}, \bar{\mathbf{3}}$)	$N_\Sigma \geq 2$	+2
Δ	(0, 0)	($\mathbf{10}, \mathbf{1}, \mathbf{3}, \mathbf{6}$)	1	+2
$\rightarrow \chi$	(0, 0)	($\mathbf{4}, \mathbf{1}, \mathbf{2}, \bar{\mathbf{3}}$)	1	-1
ξ	(0, 0)	($\mathbf{15}, \mathbf{1}, \mathbf{3}, \mathbf{1}$)	1	0

- Scalar χ : \rightarrow helps Δ with PS and flavor braking
 \rightarrow to **break PQ and rank of PS**: need Δ and χ both

Constructing the model

- Field content: $G_{\text{PS}} \equiv \text{SU}(4)_{\text{PS}} \times \text{SU}(2)_L \times \text{SU}(2)_R$

Field	Lorentz	$G_{\text{PS}} \times \text{SU}(3)_{f_R}$	copies	$\text{U}(1)_{\text{PQ}}$
\bar{Q}_L	(0, 1/2)	($\bar{\mathbf{4}}, \mathbf{2}, \mathbf{1}, \mathbf{1}$)	3	+3
Q_R	(0, 1/2)	($\mathbf{4}, \mathbf{1}, \mathbf{2}, \mathbf{3}$)	1	+1
Ψ	(0, 1/2)	($\mathbf{1}, \mathbf{1}, \mathbf{1}, \bar{\mathbf{3}}$)	8	+2
Φ	(0, 0)	($\mathbf{1}, \mathbf{2}, \mathbf{2}, \bar{\mathbf{3}}$)	$N_\Phi \geq 1$	+2
Σ	(0, 0)	($\mathbf{15}, \mathbf{2}, \mathbf{2}, \bar{\mathbf{3}}$)	$N_\Sigma \geq 2$	+2
Δ	(0, 0)	($\mathbf{10}, \mathbf{1}, \mathbf{3}, \mathbf{6}$)	1	+2
χ	(0, 0)	($\mathbf{4}, \mathbf{1}, \mathbf{2}, \bar{\mathbf{3}}$)	1	-1
\rightarrow	ξ	($\mathbf{15}, \mathbf{1}, \mathbf{3}, \mathbf{1}$)	1	0

- Scalar ξ : \rightarrow needed for **accidental PQ** to arise

\rightarrow connects Δ and χ in scalar potential: $\Delta \chi^2 \xi$

Constructing the model

- Field content: $G_{\text{PS}} \equiv \text{SU}(4)_{\text{PS}} \times \text{SU}(2)_L \times \text{SU}(2)_R$

Field	Lorentz	$G_{\text{PS}} \times \text{SU}(3)_{f_R}$	copies	$\text{U}(1)_{\text{PQ}}$
\bar{Q}_L	(0, 1/2)	($\bar{\mathbf{4}}, \mathbf{2}, \mathbf{1}, \bar{\mathbf{1}}$)	3	+3
Q_R	(0, 1/2)	($\mathbf{4}, \mathbf{1}, \mathbf{2}, \mathbf{3}$)	1	+1
Ψ	(0, 1/2)	($\mathbf{1}, \mathbf{1}, \mathbf{1}, \bar{\mathbf{3}}$)	8	+2
Φ	(0, 0)	($\mathbf{1}, \mathbf{2}, \mathbf{2}, \bar{\mathbf{3}}$)	$N_\Phi \geq 1$	+2
Σ	(0, 0)	($\mathbf{15}, \mathbf{2}, \mathbf{2}, \bar{\mathbf{3}}$)	$N_\Sigma \geq 2$	+2
Δ	(0, 0)	($\mathbf{10}, \mathbf{1}, \mathbf{3}, \bar{\mathbf{6}}$)	1	+2
χ	(0, 0)	($\mathbf{4}, \mathbf{1}, \mathbf{2}, \bar{\mathbf{3}}$)	1	-1
ξ	(0, 0)	($\mathbf{15}, \mathbf{1}, \mathbf{3}, \mathbf{1}$)	1	0

- PQ accidental in

$$\mathcal{V}_{\mathbb{C}} = \Phi \Sigma^* \xi + \Phi \Sigma^* (|\Sigma|^2 + |\Delta|^2 + |\chi|^2 + \xi^2) + \Sigma^{*2} (\Phi^2 + \Delta^2) + \Delta \chi^2 \xi + \text{h.c.}$$

Constructing the model

- Field content: $G_{\text{PS}} \equiv \text{SU}(4)_{\text{PS}} \times \text{SU}(2)_L \times \text{SU}(2)_R$

Field	Lorentz	$G_{\text{PS}} \times \text{SU}(3)_{f_R}$	copies	$\text{U}(1)_{\text{PQ}}$
\bar{Q}_L	(0, 1/2)	($\bar{\mathbf{4}}, \mathbf{2}, \mathbf{1}, \mathbf{1}$)	3	+3
Q_R	(0, 1/2)	($\mathbf{4}, \mathbf{1}, \mathbf{2}, \mathbf{3}$)	1	+1
Ψ	(0, 1/2)	($\mathbf{1}, \mathbf{1}, \mathbf{1}, \bar{\mathbf{3}}$)	8	+2
Φ	(0, 0)	($\mathbf{1}, \mathbf{2}, \mathbf{2}, \bar{\mathbf{3}}$)	$N_\Phi \geq 1$	+2
Σ	(0, 0)	($\mathbf{15}, \mathbf{2}, \mathbf{2}, \bar{\mathbf{3}}$)	$N_\Sigma \geq 2$	+2
Δ	(0, 0)	($\mathbf{10}, \mathbf{1}, \mathbf{3}, \mathbf{6}$)	1	+2
χ	(0, 0)	($\mathbf{4}, \mathbf{1}, \mathbf{2}, \bar{\mathbf{3}}$)	1	-1
ξ	(0, 0)	($\mathbf{15}, \mathbf{1}, \mathbf{3}, \mathbf{1}$)	1	0

- Breaking: $G_{\text{PS}} \times \text{SU}(3)_{f_R} \xrightarrow{\langle \Delta, \chi, \xi \rangle} G_{\text{SM}} \xrightarrow{\langle \Phi, \Sigma \rangle} \text{SU}(3)_C \times \text{U}(1)_{\text{EM}}$

PQ quality — part 1

- **Quality check:** find dominant \mathcal{PQ} non-renormalizable invariants contributing to the vacuum

PQ quality — part 1

- **Quality check:** find dominant \mathcal{PQ} non-renormalizable invariants contributing to the vacuum
- **Procedure** for finding invariant operators \mathcal{O} :
 - (0) Go through all possible field powers for a given dimension d of \mathcal{O}
 - (1) Necessary condition: trivial under gauge center $\mathbb{Z}_4 \times \mathbb{Z}_2 \times \mathbb{Z}_2 \times \mathbb{Z}_3$
 - (2) Automatized confirmation (GroupMath [14], LieART [15, 16], LiE [17])
 - (3) Explicit index contraction for confirming vacuum contribution $\langle \mathcal{O} \rangle \neq 0$

PQ quality — part 1

- **Quality check:** find dominant \cancel{PQ} non-renormalizable invariants contributing to the vacuum
- **Procedure** for finding invariant operators \mathcal{O} :
 - (0) Go through all possible field powers for a given dimension d of \mathcal{O}
 - (1) Necessary condition: trivial under gauge center $\mathbb{Z}_4 \times \mathbb{Z}_2 \times \mathbb{Z}_2 \times \mathbb{Z}_3$
 - (2) Automatized confirmation (GroupMath [14], LieART [15, 16], LiE [17])
 - (3) Explicit index contraction for confirming vacuum contribution $\langle \mathcal{O} \rangle \neq 0$
- VEVs:

$$\langle \Phi \rangle, \langle \Sigma \rangle \sim v \quad \sim 10^2 \text{ GeV} \quad (\text{EW VEVs}), \quad (4)$$

$$\langle \Delta \rangle, \langle \chi \rangle, \langle \xi \rangle \equiv V_\Delta, V_\chi, V_\xi \quad \sim V \quad (\text{PS VEVs}), \quad (5)$$

$$\Lambda \equiv M_{\text{Planck}} \quad \sim 10^{19} \text{ GeV} \quad (\text{cutoff}) \quad (6)$$

PQ quality — part 2

- **High-quality condition:** small shift in QCD-generated axion potential

$$\langle \mathcal{O} \rangle \lesssim \bar{\theta} \chi_{\text{QCD}}^4 \equiv \bar{\theta} \frac{m_u m_d}{(m_u + m_d)^2} m_\pi^2 f_\pi^2 \simeq 10^{-10} (76 \text{ MeV})^4. \quad (7)$$

PQ quality — part 2

- **High-quality condition:** small shift in QCD-generated axion potential

$$\langle \mathcal{O} \rangle \lesssim \bar{\theta} \chi_{\text{QCD}}^4 \equiv \bar{\theta} \frac{m_u m_d}{(m_u + m_d)^2} m_\pi^2 f_\pi^2 \simeq 10^{-10} (76 \text{ MeV})^4. \quad (7)$$

- Dominant \cancel{PQ} contributions:

\mathcal{O}	$\langle \mathcal{O} \rangle$
$\Delta^4 \Delta^* \chi^{*6}$	V^{11}/Λ^7
$\Phi^{2-k} \Sigma^k \Delta^2 \chi^{*4}$	$v^2 V^6/\Lambda^4$!
$\Phi^{4-k} \Sigma^k \Delta \chi^{*2}$	$v^4 V^3/\Lambda^3$
$\Phi^{4-k} \Sigma^k \Sigma^2$	v^6/Λ^2

PQ quality — part 2

- **High-quality condition:** small shift in QCD-generated axion potential

$$\langle \mathcal{O} \rangle \lesssim \bar{\theta} \chi_{\text{QCD}}^4 \equiv \bar{\theta} \frac{m_u m_d}{(m_u + m_d)^2} m_\pi^2 f_\pi^2 \simeq 10^{-10} (76 \text{ MeV})^4. \quad (7)$$

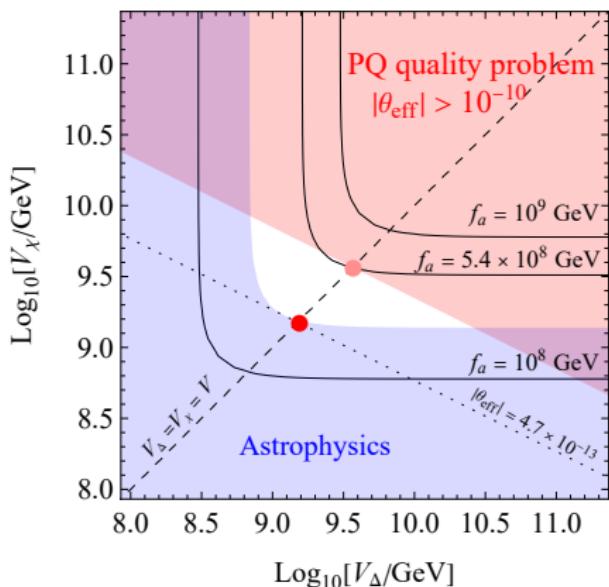
- Dominant PQ contributions:

\mathcal{O}	$\langle \mathcal{O} \rangle$
$\Delta^4 \Delta^* \chi^{*6}$	V^{11}/Λ^7
$\Phi^{2-k} \Sigma^k \Delta^2 \chi^{*4}$	$v^2 V^6/\Lambda^4 !$
$\Phi^{4-k} \Sigma^k \Delta \chi^{*2}$	$v^4 V^3/\Lambda^3$
$\Phi^{4-k} \Sigma^k \Sigma^2$	v^6/Λ^2

- Axion decay constant f_a :

$$f_a = \frac{V_\chi V_\Delta}{3\sqrt{V_\chi^2 + 4V_\Delta^2}} \quad (8)$$

- Constraints: quality and astro



PQ quality — part 2

- **High-quality condition:** small shift in QCD-generated axion potential

$$\langle \mathcal{O} \rangle \lesssim \bar{\theta} \chi_{\text{QCD}}^4 \equiv \bar{\theta} \frac{m_u m_d}{(m_u + m_d)^2} m_\pi^2 f_\pi^2 \simeq 10^{-10} (76 \text{ MeV})^4. \quad (7)$$

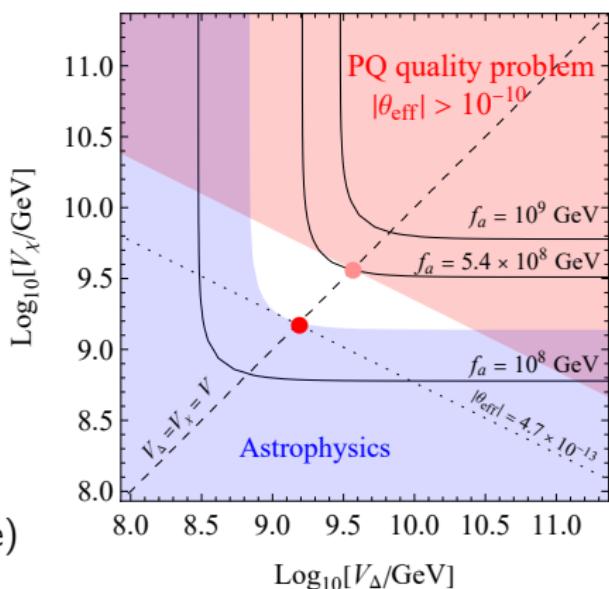
- Dominant PQ contributions:

\mathcal{O}	$\langle \mathcal{O} \rangle$
$\Delta^4 \Delta^* \chi^{*6}$	V^{11}/Λ^7
$\Phi^{2-k} \Sigma^k \Delta^2 \chi^{*4}$	$v^2 V^6/\Lambda^4 !$
$\Phi^{4-k} \Sigma^k \Delta \chi^{*2}$	$v^4 V^3/\Lambda^3$
$\Phi^{4-k} \Sigma^k \Sigma^2$	v^6/Λ^2

- Window for $V_\Delta = V_\chi \equiv V$:

$$V \in [1.5, 3.7] \times 10^9 \text{ GeV} \quad (8)$$

(larger V degrades quality, but possible)



Axion properties

- Axion a : (mostly) a combination of polar modes in Δ and χ
- Couplings to photons and gluons:

$$\mathcal{L} \supset \frac{\alpha_{\text{EM}} E}{4\pi} \frac{a}{f_a} F \tilde{F} + \frac{\alpha_s N}{4\pi} \frac{a}{f_a} G \tilde{G}. \quad (9)$$

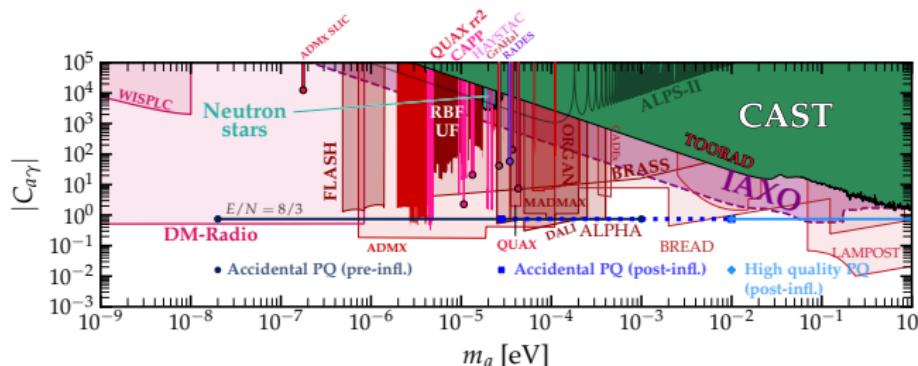
- This model: $E = 16$, $N = 6$, hence E/N=8/3 (DFSZ-like [18, 19])

Axion properties

- Axion a : (mostly) a combination of polar modes in Δ and χ
- Couplings to photons and gluons:

$$\mathcal{L} \supset \frac{\alpha_{\text{EM}} E}{4\pi} \frac{a}{f_a} \mathbf{F} \tilde{\mathbf{F}} + \frac{\alpha_s N}{4\pi} \frac{a}{f_a} \mathbf{G} \tilde{\mathbf{G}}. \quad (9)$$

- This model: $E = 16$, $N = 6$, hence $E/N=8/3$ (DFSZ-like [18, 19])
- Mass range: **high-quality** PQ vs **accidental** PQ scenario



high-quality:

$$m_a \gtrsim 10 \text{ meV}$$

Yukawa sector — up, down, charged leptons

- Renormalizable operators for UDE sectors:

$$\mathcal{L}_Y \supset \sum_{\alpha=1}^{N_\Phi} \sum_{I=1}^3 Y_I^{\Phi^\alpha} \bar{Q}_L^I Q_R \Phi^\alpha + \sum_{\alpha=1}^{N_\Sigma} \sum_{I=1}^3 2\sqrt{3} Y_I^{\Sigma^\alpha} \bar{Q}_L^I Q_R \Sigma^\alpha \quad (10)$$

→ Indices: I [family], A [gauged R-flavor], α [multiplicity]

Yukawa sector — up, down, charged leptons

- Renormalizable operators for UDE sectors:

$$\mathcal{L}_Y \supset \sum_{\alpha=1}^{N_\Phi} \sum_{I=1}^3 Y_I^{\Phi^\alpha} \bar{Q}_L^I Q_R \Phi^\alpha + \sum_{\alpha=1}^{N_\Sigma} \sum_{I=1}^3 2\sqrt{3} Y_I^{\Sigma^\alpha} \bar{Q}_L^I Q_R \Sigma^\alpha \quad (10)$$

→ Indices: I [family], A [gauged R-flavor], α [multiplicity]

- $N_\Phi \geq 1, N_\Sigma \geq 2$ is **realistic** (based on DOF count, not fit)

[considering masses, CKM, PMNS]

$$(M_U)_{IA} = Y_I^\Phi v_A^{u\Phi} + Y_I^\Sigma v_A^{u\Sigma} + Y_I^{\Sigma'} v_A^{u\Sigma'}, \quad (11)$$

$$(M_D)_{IA} = Y_I^\Phi v_A^{d\Phi} + Y_I^\Sigma v_A^{d\Sigma} + Y_I^{\Sigma'} v_A^{d\Sigma'}, \quad (12)$$

$$(M_E)_{IA} = Y_I^\Phi v_A^{d\Phi} - 3 Y_I^\Sigma v_A^{d\Sigma} - 3 Y_I^{\Sigma'} v_A^{d\Sigma'}, \quad (13)$$

Yukawa sector — up, down, charged leptons

- Renormalizable operators for UDE sectors:

$$\mathcal{L}_Y \supset \sum_{\alpha=1}^{N_\Phi} \sum_{I=1}^3 Y_I^{\Phi^\alpha} \bar{Q}_L^I Q_R \Phi^\alpha + \sum_{\alpha=1}^{N_\Sigma} \sum_{I=1}^3 2\sqrt{3} Y_I^{\Sigma^\alpha} \bar{Q}_L^I Q_R \Sigma^\alpha \quad (10)$$

→ Indices: I [family], A [gauged R-flavor], α [multiplicity]

- $N_\Phi \geq 1$, $N_\Sigma \geq 2$ is **realistic** (based on DOF count, not fit)

[considering masses, CKM, PMNS]

$$(M_U)_{IA} = Y_I^\Phi v_A^{u\Phi} + Y_I^\Sigma v_A^{u\Sigma} + Y_I^{\Sigma'} v_A^{u\Sigma'}, \quad (11)$$

$$(M_D)_{IA} = Y_I^\Phi v_A^{d\Phi} + Y_I^\Sigma v_A^{d\Sigma} + Y_I^{\Sigma'} v_A^{d\Sigma'}, \quad (12)$$

$$(M_E)_{IA} = Y_I^\Phi v_A^{d\Phi} - 3 Y_I^\Sigma v_A^{d\Sigma} - 3 Y_I^{\Sigma'} v_A^{d\Sigma'}, \quad (13)$$

- Unusual** setup: → Y_I are **vectors** in family space (not matrices)

→ EW VEVs v_A are R-flavor **vectors** (not numbers)

Yukawa sector — neutrinos and anomalons

- Anomalons Ψ are EW (and SM) singlets $\Rightarrow \nu \leftrightarrow \Psi$ can mix
- Spectrum:
 - (a) terms with Ψ : only from **non-renormalizable** \mathcal{O}
 - (b) number of Ψ : $8 \times [\bar{\mathbf{3}}] = 24$ in total

Yukawa sector — neutrinos and anomalons

- Anomalons Ψ are EW (and SM) singlets $\Rightarrow \nu \leftrightarrow \Psi$ can mix
- Spectrum:
 - (a) terms with Ψ : only from **non-renormalizable** \mathcal{O}
 - (b) number of Ψ : $8 \times [\bar{\mathbf{3}}] = 24$ in total
 - (c) turns out: “**heavy-light split**” $\Psi = \Psi_{\perp} \oplus \Psi_0$

Yukawa sector — neutrinos and anomalons

- Anomalons Ψ are EW (and SM) singlets $\Rightarrow \nu \leftrightarrow \Psi$ can mix
- Spectrum:
 - (a) terms with Ψ : only from **non-renormalizable** \mathcal{O}
 - (b) number of Ψ : $8 \times [\bar{\mathbf{3}}] = 24$ in total
 - (c) turns out: “**heavy-light split**” $\Psi = \Psi_{\perp} \oplus \Psi_0$
 - (d) **Flavor basis**: $\bar{\nu}_L \oplus \nu_R \oplus \Psi_{\perp} \oplus \Psi_0$ ($3 + 3 + 16 + 8$)

Yukawa sector — neutrinos and anomalons

- Anomalons Ψ are EW (and SM) singlets $\Rightarrow \nu \leftrightarrow \Psi$ can mix
- Spectrum:
 - (a) terms with Ψ : only from **non-renormalizable** \mathcal{O}
 - (b) number of Ψ : $8 \times [\bar{\mathbf{3}}] = 24$ in total
 - (c) turns out: “**heavy-light split**” $\Psi = \Psi_{\perp} \oplus \Psi_0$
 - (d) **Flavor basis**: $\bar{\nu}_L \oplus \nu_R \oplus \Psi_{\perp} \oplus \Psi_0$ ($3 + 3 + 16 + 8$)
 - (e) If mixing small: mass eigenstates denoted by $L, R, \perp, 0$

Yukawa sector — neutrinos and anomalons

- Anomalons Ψ are EW (and SM) singlets $\Rightarrow \nu \leftrightarrow \Psi$ can mix
- Spectrum:
 - (a) terms with Ψ : only from **non-renormalizable** \mathcal{O}
 - (b) number of Ψ : $8 \times [\bar{\mathbf{3}}] = 24$ in total
 - (c) turns out: “**heavy-light split**” $\Psi = \Psi_{\perp} \oplus \Psi_0$
 - (d) **Flavor basis**: $\bar{\nu}_L \oplus \nu_R \oplus \Psi_{\perp} \oplus \Psi_0$ ($3 + 3 + 16 + 8$)
 - (e) If mixing small: mass eigenstates denoted by $L, R, \perp, 0$
- Dominant operators in every entry (schematically):

	\overline{Q}_L	Q_R	Ψ_{\perp}	Ψ_0
\overline{Q}_L	$\Delta(\Phi^2 + \Phi\Sigma + \Sigma^2)$	$\Phi + \Sigma$	$\chi(\Phi + \Sigma)$	$\Delta\chi^*(\Phi^* + \Sigma^*)$
Q_R	$\Phi + \Sigma$	Δ^*	$\Delta^*\chi$	$\Delta^{*2}\Delta\chi$
Ψ_{\perp}	$\chi(\Phi + \Sigma)$	$\Delta^*\chi$	$\Delta^*\chi^2$	$\Phi^{*2} + \Sigma^{*2} + \Delta\Delta^{*2}\chi^2$
Ψ_0	$\Delta\chi^*(\Phi^* + \Sigma^*)$	$\Delta^{*2}\Delta\chi$	$\Phi^{*2} + \Sigma^{*2} + \Delta\Delta^{*2}\chi^2$	$\Phi^{*2} + \Sigma^{*2} + \Delta\Delta^{*2}\chi^2$

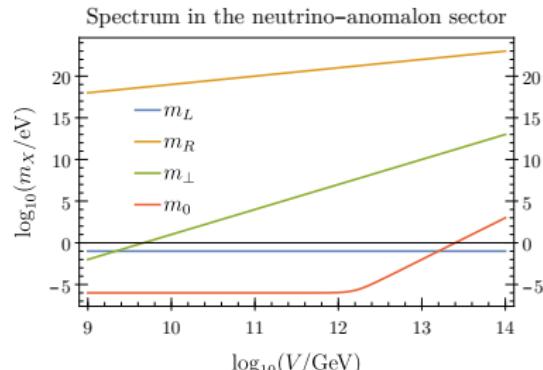
Yukawa sector — neutrinos and anomalons

- Anomalons Ψ are EW (and SM) singlets $\Rightarrow \nu \leftrightarrow \Psi$ can mix
- Spectrum:
 - (a) terms with Ψ : only from **non-renormalizable** \mathcal{O}
 - (b) number of Ψ : $8 \times [\bar{\mathbf{3}}] = 24$ in total
 - (c) turns out: “**heavy-light split**” $\Psi = \Psi_{\perp} \oplus \Psi_0$
 - (d) **Flavor basis**: $\bar{\nu}_L \oplus \nu_R \oplus \Psi_{\perp} \oplus \Psi_0$ ($3 + 3 + 16 + 8$)
 - (e) If mixing small: mass eigenstates denoted by $L, R, \perp, 0$

ν - Ψ mass matrix:

$$M_{\nu\Psi} = \begin{pmatrix} \frac{v^2 V}{\Lambda^2} & y v & l \frac{vV}{\Lambda} & \tilde{l} \frac{vV^2}{\Lambda^2} \\ .. & V & r \frac{V^2}{\Lambda} & \tilde{r} \frac{V^4}{\Lambda^3} \\ .. & .. & \frac{V^3}{\Lambda^2} & \frac{v^2}{\Lambda} + \frac{V^5}{\Lambda^4} \\ .. & .. & .. & \frac{v^2}{\Lambda} + \frac{V^5}{\Lambda^4} \end{pmatrix}$$

→ In upper-left 2×2 : **see-saw type I**



Anomalon cosmology (in high quality scenario)

- Anomalon mass $< 1 \text{ eV}$: $m_\perp \simeq 10 \text{ meV}$, $m_0 \simeq 1 \mu\text{eV}$
- Anomalons act as **dark radiation**: $\Delta N_{\text{eff}} = \frac{8}{7} \left(\frac{11}{4} \right)^{4/3} \frac{\rho_\Psi}{\rho_\gamma}$

Anomalon cosmology (in high quality scenario)

- Anomalon mass $< 1 \text{ eV}$: $m_\perp \simeq 10 \text{ meV}$, $m_0 \simeq 1 \mu\text{eV}$
- Anomalons act as **dark radiation**: $\Delta N_{\text{eff}} = \frac{8}{7} \left(\frac{11}{4} \right)^{4/3} \frac{\rho_\Psi}{\rho_\gamma}$
- **Freeze-out**:
 - If ever in thermal equilibrium: $\Delta N_{\text{eff}}^{\text{TH}} \simeq 1.13 \frac{N_\Psi}{24} \left(\frac{106.75}{g_s(T_{\text{decouple}})} \right)^{4/3}$
 - **Excluded** by Planck'18: $\Delta N_{\text{eff}} \leq 0.285$ (at 95% C.L.) [20]

Anomalon cosmology (in high quality scenario)

- Anomalon mass $< 1 \text{ eV}$: $m_\perp \simeq 10 \text{ meV}$, $m_0 \simeq 1 \mu\text{eV}$
- Anomalons act as **dark radiation**: $\Delta N_{\text{eff}} = \frac{8}{7} \left(\frac{11}{4} \right)^{4/3} \frac{\rho_\Psi}{\rho_\gamma}$
- **Freeze-out**:
 - If ever in thermal equilibrium: $\Delta N_{\text{eff}}^{\text{TH}} \simeq 1.13 \frac{N_\Psi}{24} \left(\frac{106.75}{g_s(T_{\text{decouple}})} \right)^{4/3}$
 - **Excluded** by Planck'18: $\Delta N_{\text{eff}} \leq 0.285$ (at 95% C.L.) [20]
- **Freeze-in**: production of anomalons via...
 - (i) Annihilation of SM fermions mediated by W_{f_R} : $\bar{f} f \xrightarrow{W_{f_R}} \overline{\Psi} \Psi$
 - (ii) Conversion of ν_L via mixing $\theta_{\nu\Psi}$: $e^+ e^- \rightarrow \bar{\nu}_f \nu_f \rightarrow \bar{\nu}_f \Psi_m$
 - (iii) (generalized) Yukawa interaction: $\nu_L \Psi h$, $\Psi \Psi h$, $\nu_L \Psi h h$, $\Psi \Psi \phi$, etc.

processes: $\bar{f} f \xrightarrow{h} \overline{\Psi} \Psi$, $\phi \rightarrow \overline{\Psi} \Psi$, $\phi \phi \rightarrow \overline{\Psi} \Psi$, etc.

Anomalon cosmology (in high quality scenario)

- Anomalon mass $< 1 \text{ eV}$: $m_\perp \simeq 10 \text{ meV}$, $m_0 \simeq 1 \mu\text{eV}$
- Anomalons act as **dark radiation**: $\Delta N_{\text{eff}} = \frac{8}{7} \left(\frac{11}{4} \right)^{4/3} \frac{\rho_\Psi}{\rho_\gamma}$
- **Freeze-out**:
 - If ever in thermal equilibrium: $\Delta N_{\text{eff}}^{\text{TH}} \simeq 1.13 \frac{N_\Psi}{24} \left(\frac{106.75}{g_s(T_{\text{decouple}})} \right)^{4/3}$
 - **Excluded** by Planck'18: $\Delta N_{\text{eff}} \leq 0.285$ (at 95% C.L.) [20]
- **Freeze-in**: production of anomalons via...
 - (i) Annihilation of SM fermions mediated by W_{f_R} : $\bar{f} f \xrightarrow{W_{f_R}} \overline{\Psi} \Psi$
 - (ii) Conversion of ν_L via mixing $\theta_{\nu\Psi}$: $e^+ e^- \rightarrow \bar{\nu}_f \nu_f \rightarrow \bar{\nu}_f \Psi_m$
 - (iii) (generalized) Yukawa interaction: $\nu_L \Psi h$, $\Psi \Psi h$, $\nu_L \Psi h h$, $\Psi \Psi \phi$, etc.
 - processes: $ff \xrightarrow{h} \Psi\Psi$, $\phi \rightarrow \Psi\Psi$, $\phi\phi \rightarrow \Psi\Psi$, etc.
- **Challenge**: large enough relic abundance for measurable ΔN_{eff} , but no thermalization → only (ii) may be viable, requires dedicated analysis

Conclusions

- (A) Quality from vertical and horizontal symmetry:
PQ-breaking scale \leftrightarrow flavor and PS breaking scale
- (B) **Technically challenging** to check PQ quality
- (C) Presence of **anomalons** (fermions): they mix with neutrinos
 \rightarrow UV dynamics of PQ quality potentially testable (e.g. ΔN_{eff})
- (D) However, also some **drawbacks** (not discussed):
domain walls, perturbativity

Conclusions

- (A) Quality from vertical and horizontal symmetry:
PQ-breaking scale \leftrightarrow flavor and PS breaking scale
- (B) **Technically challenging** to check PQ quality
- (C) Presence of **anomalons** (fermions): they mix with neutrinos
 \rightarrow UV dynamics of PQ quality potentially testable (e.g. ΔN_{eff})
- (D) However, also some **drawbacks** (not discussed):
domain walls, perturbativity

Thank you for your attention!

References I

- [1] R. D. Peccei and H. R. Quinn, "CP Conservation in the Presence of Instantons," *Phys. Rev. Lett.* **38** (1977) 1440–1443.
- [2] R. D. Peccei and H. R. Quinn, "Constraints Imposed by CP Conservation in the Presence of Instantons," *Phys. Rev. D* **16** (1977) 1791–1797.
- [3] S. Weinberg, "A New Light Boson?," *Phys. Rev. Lett.* **40** (1978) 223–226.
- [4] F. Wilczek, "Problem of Strong P and T Invariance in the Presence of Instantons," *Phys. Rev. Lett.* **40** (1978) 279–282.
- [5] L. Di Luzio, G. Landini, F. Mescia, and V. Susič, "High-quality Peccei-Quinn symmetry from the interplay of vertical and horizontal gauge symmetries," [arXiv:2503.16648 \[hep-ph\]](https://arxiv.org/abs/2503.16648).
- [6] L. Di Luzio, "Accidental SO(10) axion from gauged flavour," *JHEP* **11** (2020) 074, [arXiv:2008.09119 \[hep-ph\]](https://arxiv.org/abs/2008.09119).
- [7] Z. G. Berezhiani, "The Weak Mixing Angles in Gauge Models with Horizontal Symmetry: A New Approach to Quark and Lepton Masses," *Phys. Lett. B* **129** (1983) 99–102.
- [8] Z. G. Berezhiani, "Horizontal Symmetry and Quark - Lepton Mass Spectrum: The SU(5) \times SU(3)-h Model," *Phys. Lett. B* **150** (1985) 177–181.
- [9] Z. G. Berezhiani and M. Y. Khlopov, "The Theory of broken gauge symmetry of families. (In Russian)," *Sov. J. Nucl. Phys.* **51** (1990) 739–746.
- [10] L. Vecchi, "Axion quality straight from the GUT," *Eur. Phys. J. C* **81** no. 10, (2021) 938, [arXiv:2106.15224 \[hep-ph\]](https://arxiv.org/abs/2106.15224).
- [11] K. S. Babu, B. Dutta, and R. N. Mohapatra, "Hybrid SO(10) Axion Model without Quality Problem," *Phys. Rev. Lett.* **134** no. 11, (2025) 111803, [arXiv:2410.07323 \[hep-ph\]](https://arxiv.org/abs/2410.07323).
- [12] L. Darmé and E. Nardi, "Exact accidental U(1) symmetries for the axion," *Phys. Rev. D* **104** no. 5, (2021) 055013, [arXiv:2102.05055 \[hep-ph\]](https://arxiv.org/abs/2102.05055).
- [13] L. Darmé, E. Nardi, and C. Smarra, "The axion flavour connection," *JHEP* **02** (2023) 201, [arXiv:2211.05796 \[hep-ph\]](https://arxiv.org/abs/2211.05796).

References II

- [14] R. M. Fonseca, "GroupMath: A Mathematica package for group theory calculations," *Comput. Phys. Commun.* **267** (2021) 108085, [arXiv:2011.01764 \[hep-th\]](#).
- [15] R. Feger and T. W. Kephart, "LieART—A Mathematica application for Lie algebras and representation theory," *Comput. Phys. Commun.* **192** (2015) 166–195, [arXiv:1206.6379 \[math-ph\]](#).
- [16] R. Feger, T. W. Kephart, and R. J. Saskowski, "LieART 2.0 – A Mathematica application for Lie Algebras and Representation Theory," *Comput. Phys. Commun.* **257** (2020) 107490, [arXiv:1912.10969 \[hep-th\]](#).
- [17] M. van Leeuwen, A. M. Cohen, and B. Lisser, *LiE, A Package for Lie Group Computations*. Computer Algebra Nederland, Amsterdam, ISBN 90-74116-02-7, 1992.
- [18] A. R. Zhitnitsky, "On Possible Suppression of the Axion Hadron Interactions. (In Russian)," *Sov. J. Nucl. Phys.* **31** (1980) 260.
- [19] M. Dine, W. Fischler, and M. Srednicki, "A Simple Solution to the Strong CP Problem with a Harmless Axion," *Phys. Lett. B* **104** (1981) 199–202.
- [20] Planck Collaboration, N. Aghanim *et al.*, "Planck 2018 results. VI. Cosmological parameters," *Astron. Astrophys.* **641** (2020) A6, [arXiv:1807.06209 \[astro-ph.CO\]](#). [Erratum: *Astron. Astrophys.* 652, C4 (2021)].