

# One-loop running in bosonic theories

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# Outline

- 1 Introduction
- 2 Lagrangian
- 3 RGEs
- 4 Applications
- 5 Summary

based on: 2502.14030 in collaboration with  
Luigi Bresciani and Nudzeim Selimovic

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# Series of papers

## Renormalization of general Effective Field Theories

Diagrammatic, also fermion couplings

Fonseca/Olgoso/Santiago: 2501.13185

## One-loop Renormalization Group Equations in Generic Effective Field Theories

Diagrammatic, even-dimensional

Misiak/Nalecz: 2501.17134

## Anomalous Dimension of a General Effective Gauge Theory

OS-methods/geometric, up to dim-6

JA/Bresciani/Selimovic: 2502.14030

# Renormalization Group Equations (RGEs)

## RGEs

Scale-dependence of parameters

## Log resummation

$$\sum_N (\alpha \text{Log}(\frac{m^2}{\Lambda^2}))^N$$

## Renormalization scale dependence

cancelled in matching + running

# Idea: General one-loop RGEs

## Fields

scalars, vectors

## Symmetries

Gauge, Poincaré

## Computation

once and for all

# Renormalizable Lagrangian

## Gauge

2-loop

Luo/Wang/Xiao: hep-ph/0211440  
Schienbein/Staub/Steudtner/Svirina: 1809.06797

3-loop

Pickering/Gracey/Jones: hep-ph/0104247  
Poole/Thomsen: 1906.04625

4-loop

Bednyakov/Pikelner: 2105.09918

## Yukawa

3-loop

Davies/Herren/Thomsen: 2110.05496  
Poole/Thomsen: 1901.02749

## Scalar theory

6-loop

Bednyakov/Pikelner: 2102.12832  
Bednyakov: 2501.14087

# On-shell methods

## Helicity spinors

$$\not{p} = \begin{pmatrix} 0 & p_\mu (\sigma^\mu)_{\alpha\dot{\alpha}} \\ p_\mu (\bar{\sigma}^\mu)^{\dot{\alpha}\alpha} & 0 \end{pmatrix} = \begin{pmatrix} 0 & p_{\alpha\dot{\alpha}} \\ p^{\dot{\alpha}\alpha} & 0 \end{pmatrix}$$

$$p_{\alpha\dot{\alpha}} = \lambda_\alpha \tilde{\lambda}_{\dot{\alpha}}, \quad p^{\dot{\alpha}\alpha} = \tilde{\lambda}^{\dot{\alpha}} \lambda^\alpha$$

## Lorentz-invariant products

$$\lambda_i^\alpha \lambda_{j\alpha} = \langle ij \rangle$$

$$\tilde{\lambda}_{i\dot{\alpha}} \tilde{\lambda}_j^{\dot{\alpha}} = [ij]$$

## Mandelstam invariants

$$s_{ij} = 2p_i \cdot p_j = \langle ij \rangle [ji]$$

# Helicity formalism: RGEs

## Form factor

$$F_i(\vec{n}; q) = \langle \vec{n} | \mathcal{O}_i(q) | 0 \rangle$$

## Fundamental relations

Miró/Ingoldby/Riembau: 2005.06983

Analyticity:  $F_i^*(\{s_{ij} - i\epsilon\}) = F_i(\{s_{ij} + i\epsilon\})$

Unitarity:  $\sum_{\vec{n}} \int d\text{LIPS}_n |\vec{n}\rangle \langle \vec{n}| = \mathbb{1}$

CPT Theorem:  $\langle \vec{n}; \text{out} | \mathcal{O}_i(x) | 0 \rangle = \langle 0 | \mathcal{O}_i^\dagger(-x) | \vec{n}; \text{in} \rangle$

## RGEs

Caron-Huot/Wilhelm: 1607.06448

$$\left( \delta_{ij} \mu \frac{\partial}{\partial \mu} + \frac{\partial \beta_i}{\partial C_j} - \delta_{ij} \gamma_{i, \text{IR}} + \delta_{ij} \beta_g \frac{\partial}{\partial g} \right) F_i = 0$$

# Advantages

## **On-shell**

No need for green's basis

## **Gauge invariance**

automatic

## **Zeros in ADM**

helicity, length, angular momentum

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# Lagrangian

## Field content

$\phi, A_\mu$

## Symmetry

gauge, Poincaré

## Mass dimension

$\leq 6$

# Renormalizable Lagrangian

$$\begin{aligned}\mathcal{L}^{(4)} = & -\frac{1}{4}F_{\mu\nu}^A F^{A\mu\nu} + \frac{\theta g^2}{32\pi^2} F_{\mu\nu}^A \tilde{F}^{B\mu\nu} + \frac{1}{2}(D_\mu\phi)_a (D^\mu\phi)^a \\ & - \Lambda - t_a\phi_a - \frac{m_{ab}^2}{2!}\phi_a\phi_b - \frac{h_{abc}}{3!}\phi_a\phi_b\phi_c - \frac{\lambda_{abcd}}{4!}\phi_a\phi_b\phi_c\phi_d\end{aligned}$$

# Operator classification

## Amplitude

$$\mathcal{M} \sim \mathcal{K}_i$$

$$\mathcal{K}_i = \mathcal{K}_i(\{\langle ij \rangle, [ij]\})$$

## Three-point function

Benincasa/Cachazo: 0705.4305

$$\mathcal{M}(1^{h_1}, 2^{h_2}, 3^{h_3}) \sim \begin{cases} \langle 12 \rangle^{a_3} \langle 23 \rangle^{a_1} \langle 31 \rangle^{a_2} & \text{if } h_1 + h_2 + h_3 < 0 \\ [12]^{-a_3} [23]^{-a_1} [31]^{-a_2} & \text{if } h_1 + h_2 + h_3 > 0 \end{cases}$$

$$a_1 = h_1 - h_2 - h_3, \quad a_2 = h_2 - h_3 - h_1, \quad a_3 = h_3 - h_1 - h_2,$$

## Mass dimension

Durieux/Machado: 1912.08827  
Li/Ren/Xiao/Yu/Zheng: 2201.04639

$$[\mathcal{K}_i] = [\mathcal{O}_i] - \ell(\mathcal{O}_i) \geq 0$$

## Example: Dim-5

$$[\mathcal{K}_i] = 5 - \ell(\mathcal{O}_i):$$

$$l_i = 5 \quad \Longrightarrow \quad \mathcal{K}_i = \{1\} \quad \Longrightarrow \quad \mathcal{O}_i = \{\phi_1 \phi_2 \phi_3 \phi_4 \phi_5\}$$

$$l_i = 4 \quad \Longrightarrow \quad \mathcal{K}_i = \{\langle 12 \rangle\} \quad \Longrightarrow \quad \mathcal{O}_i = \{\psi_{1L} \psi_{2L} \phi_3 \phi_4\}$$

$$l_i = 3 \quad \Longrightarrow \quad \mathcal{K}_i = \begin{cases} \langle 12 \rangle^2 \\ \langle 12 \rangle \langle 13 \rangle \end{cases} \quad \Longrightarrow \quad \mathcal{O}_i = \begin{cases} F_{1L} F_{2L} \phi_3 \\ F_{1L} \psi_{2L} \psi_{3L} \end{cases}$$

## Dim-5 Lagrangian

$$\begin{aligned}\mathcal{L}^{(5)} = & \left[ C_{\phi F^2} \right]_a^{A_\alpha B_\beta} \phi_a F_{\mu\nu}^{A_\alpha} F^{B_\beta \mu\nu} + \left[ C_{\phi \tilde{F}^2} \right]_a^{A_\alpha B_\beta} \phi_a F_{\mu\nu}^{A_\alpha} \tilde{F}^{B_\beta \mu\nu} \\ & + \left[ C_{\phi^5} \right]_{abcde} \phi_a \phi_b \phi_c \phi_d \phi_e\end{aligned}$$

## Dim-6 Lagrangian

$$\begin{aligned}
 \mathcal{L}^{(6)} = & [C_{\phi^2 F^2}]_{ab}^{A_\alpha B_\beta} \phi_a \phi_b F_{\mu\nu}^{A_\alpha} F^{B_\beta \mu\nu} + [C_{\phi^2 \tilde{F}^2}]_{ab}^{A_\alpha B_\beta} \phi_a \phi_b F_{\mu\nu}^{A_\alpha} \tilde{F}^{B_\beta \mu\nu} \\
 & + [C_{\phi^6}]_{abcdef} \phi_a \phi_b \phi_c \phi_d \phi_e \phi_f + [C_{D^2 \phi^4}]_{abcd} (D_\mu \phi)_a (D^\mu \phi)_b \phi_c \phi_d \\
 & + [C_{F^3}]^{A_\alpha B_\alpha C_\alpha} F_\mu^{A_\alpha \nu} F_\nu^{B_\alpha \rho} F_\rho^{C_\alpha \mu} + [C_{\tilde{F}^3}]^{A_\alpha B_\alpha C_\alpha} F_\mu^{A_\alpha \nu} F_\nu^{B_\alpha \rho} \tilde{F}_\rho^{C_\alpha \mu}
 \end{aligned}$$

## Dim-6: Wilson coefficients

Name	Operator	Symmetry
$\mathcal{O}_{\phi^6}$	$\phi_a \phi_b \phi_c \phi_d \phi_e \phi_f$	$[C_{\phi^6}]_{abcdef} = [C_{\phi^6}]_{(abcdef)}$
$\mathcal{O}_{D^2\phi^4}$	$(D_\mu \phi)_a (D^\mu \phi)_b \phi_c \phi_d$	$[C_{D^2\phi^2}]_{abcd} = [C_{D^2\phi^2}]_{(ab)cd} = [C_{D^2\phi^2}]_{ab(cd)}$
$\mathcal{O}_{\phi^2 F^2}$	$\phi_a \phi_b F_{\mu\nu}^{A\alpha} F^{B\beta\ \mu\nu}$	$[C_{\phi^2 F^2}]_{ab}^{A\alpha B\beta} = [C_{\phi^2 F^2}]_{ab}^{(A\alpha B\beta)} = [C_{\phi^2 F^2}]_{(ab)}^{A\alpha B\beta}$
$\mathcal{O}_{\phi^2 \tilde{F}^2}$	$\phi_a \phi_b F_{\mu\nu}^{A\alpha} \tilde{F}^{B\beta\ \mu\nu}$	$[C_{\phi^2 \tilde{F}^2}]_{ab}^{A\alpha B\beta} = [C_{\phi^2 \tilde{F}^2}]_{ab}^{(A\alpha B\beta)} = [C_{\phi^2 \tilde{F}^2}]_{(ab)}^{A\alpha B\beta}$
$\mathcal{O}_{F^3}$	$F_{\mu}^{A\alpha\ \nu} F_{\nu}^{B\alpha\ \rho} F_{\rho}^{C\alpha\ \mu}$	$[C_{F^3}]^{A\alpha B\alpha C\alpha} = [C_{F^3}]^{[A\alpha B\alpha C\alpha]}$
$\mathcal{O}_{\tilde{F}^3}$	$F_{\mu}^{A\alpha\ \nu} F_{\nu}^{B\alpha\ \rho} \tilde{F}_{\rho}^{C\alpha\ \mu}$	$[C_{\tilde{F}^3}]^{A\alpha B\alpha C\alpha} = [C_{\tilde{F}^3}]^{[A\alpha B\alpha C\alpha]}$

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# RGEs

## Order

$$\mathcal{O}\left(\frac{1}{\Lambda^2}\right)$$

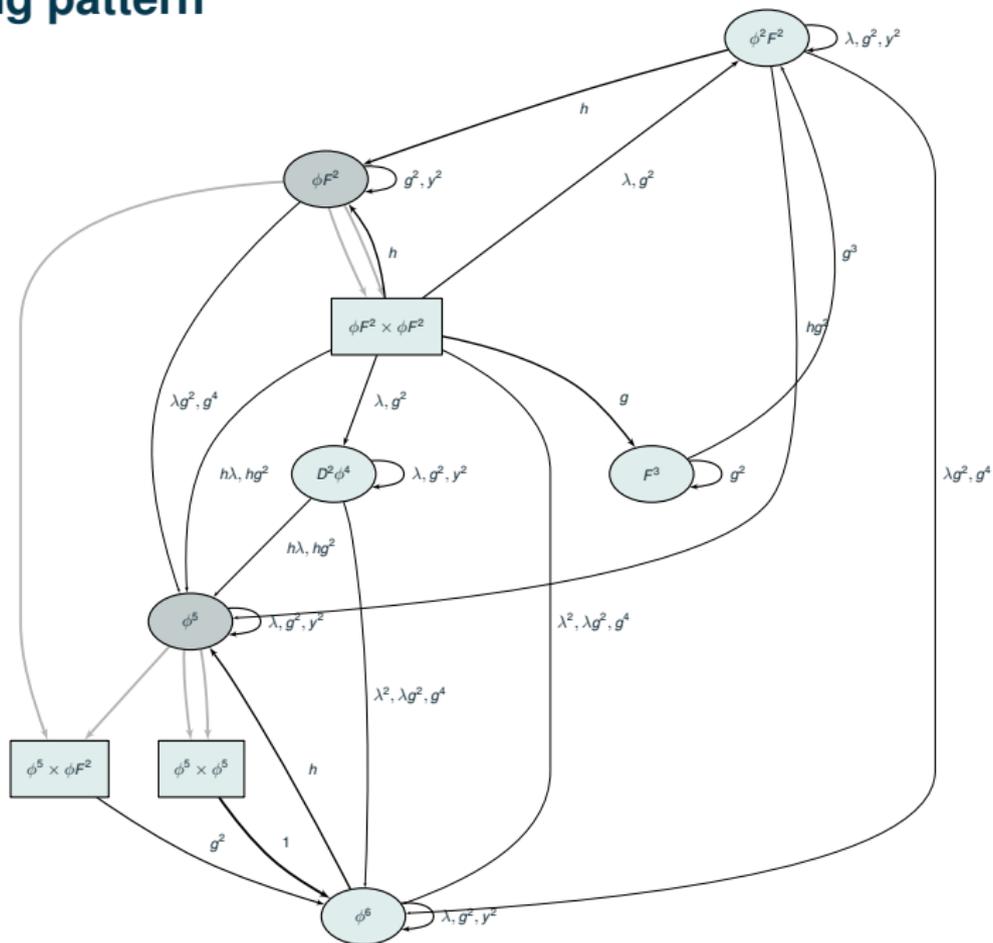
## Renormalizable part

dim-5 & dim-6

## Dim-5, Dim-6

All contributions

# Mixing pattern



**Example:  $F^3 \rightarrow F^3$**

$$\begin{aligned} [\dot{C}_{F^3}]^{ABC} = & 4g^2 \left( F^{BCDF} [C_{F^3}]^{ADF} + F^{CADF} [C_{F^3}]^{BDF} + F^{ABDF} [C_{F^3}]^{CDF} \right) \\ & + \gamma_{C,V}^{AD} [C_{F^3}]^{DBC} + \gamma_{C,V}^{BD} [C_{F^3}]^{ADC} + \gamma_{C,V}^{CD} [C_{F^3}]^{ABD} \end{aligned}$$

$$F^{ABCD} = f^{ABE} f^{CDE}$$

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## SMEFT: $O_G \rightarrow O_G$

$$O_G = f^{ABC} G_\mu^{A\nu} G_\nu^{B\rho} G_\rho^{C\mu}$$

With

$$[C_{F^3}]^{ABC} = C_G f^{ABC}$$

$$[\gamma_{c,G}]^{AB} = -g_3^2 b_{0,3} \delta^{AB}$$

One finds

$$\begin{aligned} \dot{C}_G f^{ABC} = & 4g_3^2 \left( F^{BCDF} f^{ADF} C_G + F^{CADF} f^{BDF} C_G + F^{ABDF} f^{CDF} C_G \right) \\ & + \gamma_{c,v}^{AD} C_G f^{DBC} + \gamma_{c,v}^{BD} [C_{F^3}]^{ADC} + \gamma_{c,v}^{CD} [C_{F^3}]^{ABD} \end{aligned}$$

$$\Rightarrow \dot{C}_G = (12N_c - 3b_{0,3}) g_3^2 C_G$$

# SMEFT: Higgs sector

## Higgs

$$H = \frac{1}{\sqrt{2}} \begin{pmatrix} \phi_1 + i\phi_3 \\ \phi_2 + i\phi_4 \end{pmatrix}$$

## Generators

$$\theta^A = i \begin{pmatrix} \text{Im}(T^A) & \text{Re}(T^A) \\ -\text{Re}(T^A) & \text{Im}(T^A) \end{pmatrix}$$

## Higgs Lagrangian

$$\mathcal{L}_{\text{SM}} \supset \frac{1}{2} (D_\mu \phi)_a (D^\mu \phi)_a + \frac{1}{4} m_H^2 \delta_{ab} \phi_a \phi_b - \frac{1}{4} \lambda \delta_{(ab} \delta_{cd)} \phi_a \phi_b \phi_c \phi_d$$

# Reproduced results

## SMEFT

Complete bosonic sector

Alonso/Jenkins/Manohar/Trott: 1308.2627, 1312.2014

## ALP-SMEFT

$aX, a\tilde{X} \rightarrow H^6, H^4 D^2, X^2 H^2, X^3$

Galda/Neubert/Renner: 2105.01078

Brescian/Brunello/Levati/Mastrolia/Paradisi: 2412.04160

## $O(n)$ scalar EFT

$\lambda$ -dependence

Cao/Herzog/Melia/Nepveu: 2105.12742

Jenkins/Manohar/Naterop/Pagès: 2310.19883

# New results: ALP-SMEFT

## Lagrangian

$$\begin{aligned}\mathcal{L}_{\text{ALP}}^{(5)} = & C_{aG} a G_{\mu\nu}^A G^{A\mu\nu} + C_{a\tilde{G}} a G_{\mu\nu}^A \tilde{G}^{A\mu\nu} + C_{aW} a W_{\mu\nu}^I W^{I\mu\nu} \\ & + C_{a\tilde{W}} a W_{\mu\nu}^I \tilde{W}^{I\mu\nu} + C_{aB} a B_{\mu\nu} B^{\mu\nu} + C_{a\tilde{B}} a B_{\mu\nu} \tilde{B}^{\mu\nu}\end{aligned}$$

## mass

$$\dot{m}^2 = 8 m^4 \left( C_{aB}^2 + C_{a\tilde{B}}^2 + 3 C_{aW}^2 + 3 C_{a\tilde{W}}^2 + 8 C_{aG}^2 + 8 C_{a\tilde{G}}^2 \right)$$

## dim-6: $X^2 H^2$

$$\begin{aligned}\dot{C}_{H\tilde{B}} &= 4 g_1^2 C_{aB} C_{a\tilde{B}}, \quad \dot{C}_{H\tilde{W}} = 4 g_2^2 C_{aW} C_{a\tilde{W}}, \\ \dot{C}_{H\tilde{W}B} &= 4 g_1 g_2 (C_{a\tilde{B}} C_{aW} + C_{aB} C_{a\tilde{W}})\end{aligned}$$

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# Summary

## General bosonic EFT

Operators up to dim-6

## Complete RGEs

$\lambda, h, g$

## New results

RGEs for ALP-SMEFT