

Cogenesis of DM & BAU by Majoron (=QCD axion)

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The banner for the Planck 2025 conference. It features a colorful, stylized illustration of various European architectural landmarks, including domes, towers, and classical buildings, rendered in a watercolor-like style with a vibrant color palette. The text "PLANCK 2025" is prominently displayed in large, bold, red capital letters at the bottom of the banner.

Outline

- Introduction to pNGB as DM and BAU generator.
QCD axion vs. Majoron
- Cogenesis by Majoron EJC, Jung, 2311.09005
- Cogenesis by QCD axion=Majoron In preparation
- Large kinetic misalignment from AD mechanism
or symmetry non-restoration EJC, Das, He, Jung, Sun, 2406.04180
- Conclusion

pNGB from U(1) Breaking

- A massless phase field from spontaneous breaking
- QCD axion (KSVZ) vs. Majoron
- Potential from (tiny) explicit breaking
- Derivative and anomaly coupling
- Cosmic evolution

$$\Phi = \frac{f_a}{\sqrt{2}} e^{i\theta} ; \quad \theta \equiv \frac{a}{f_a}$$

$$y_Q \Phi Q Q^c \quad \text{vs.} \quad \frac{y_N}{2} \Phi N N$$

$$V = m_a^2 f_a^2 (1 - \cos\theta)$$

$$\mathcal{L}_\partial = x_\psi \partial_\mu \theta \bar{\psi} \bar{\sigma}^\mu \psi$$

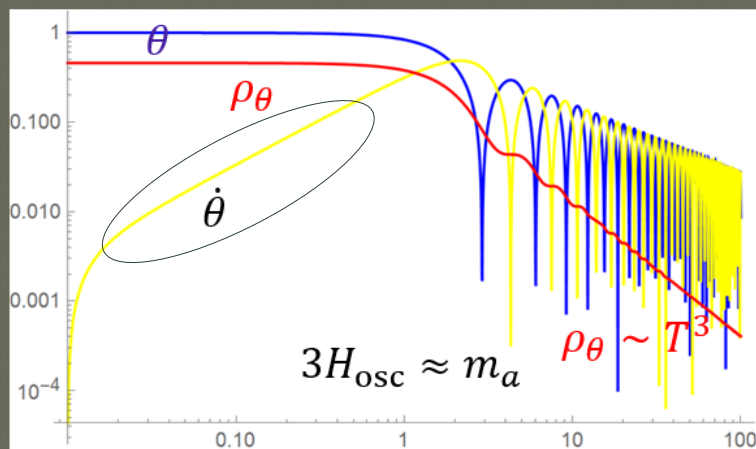
$$\mathcal{L}_A = \frac{c_A}{32\pi^2} \theta F_A^{\mu\nu} F_{A\mu\nu}$$

$$\ddot{\theta} + \left(3H + 2 \frac{\dot{f}_a}{f_a} \right) \dot{\theta} + m_a^2 \sin\theta = 0$$

pNGB as CDM

Conventional misalignment

Preskill, Wise, Wilczek; Abbott, Sikivie;
Dine, Fischler, 1983

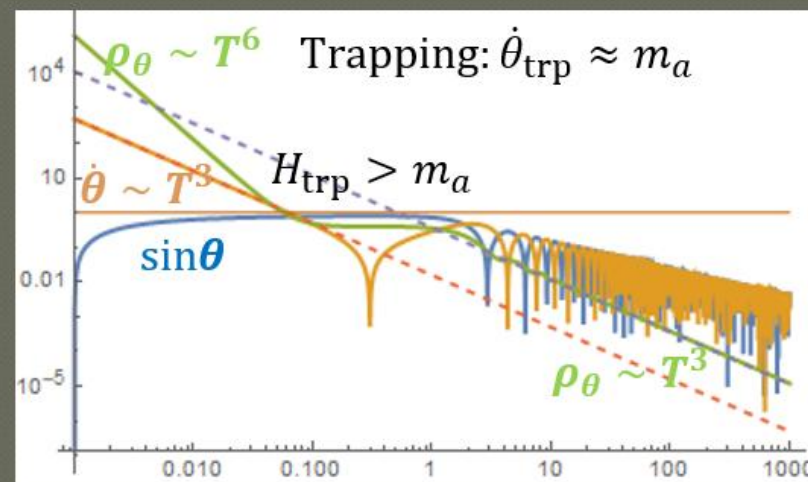


$$\rho_\theta = f_a^2 \left(\frac{1}{2} \dot{\theta}^2 + m_a^2 (1 - \cos\theta) \right)$$

$$\frac{\rho_{\text{DM}}}{s} = m_a Y_\theta \approx 0.44 \text{ eV}$$

Kinetic misalignment

Co, Hall, Harigaya, 1910.14152



$$Y_\theta \equiv \frac{n_\theta}{s} = \frac{\dot{\theta} f_a^2}{s} = \text{conserved}$$

pNGB generating BAU

- ◉ In the background of $\dot{\theta} \neq 0$ violating C & CPT, a fermion ψ gets an “external chemical potential” $x_\psi \dot{\theta}$ ($E_{\psi/\bar{\psi}} = E_0 \mp x_\psi \dot{\theta}$).

- ◉ When B(L)NV interactions are in equilibrium, it can be fed into the internal chemical potential $\mu_\psi = c_\psi \dot{\theta}$ generating baryon asymmetry $\mu_B = c_B \dot{\theta}$.

Cohen-Kaplan, 1987, 88

- ◉ QCD axion:
(B+L)NV by **weak spharelon in equilibrium**

$$\sum_{\psi \in A} \mu_\psi = c_A \dot{\theta}$$

Co-Harygaya, 1910.02080
Domcke et.al., 2006.04138

- ◉ Majoron:
(B-L)NV by neutrino Yukawa coupling $N \leftrightarrow lH$ **in equilibrium**

$$\mu_l + \mu_H = x_N \dot{\theta}$$

EJC, Jung 2311.09005

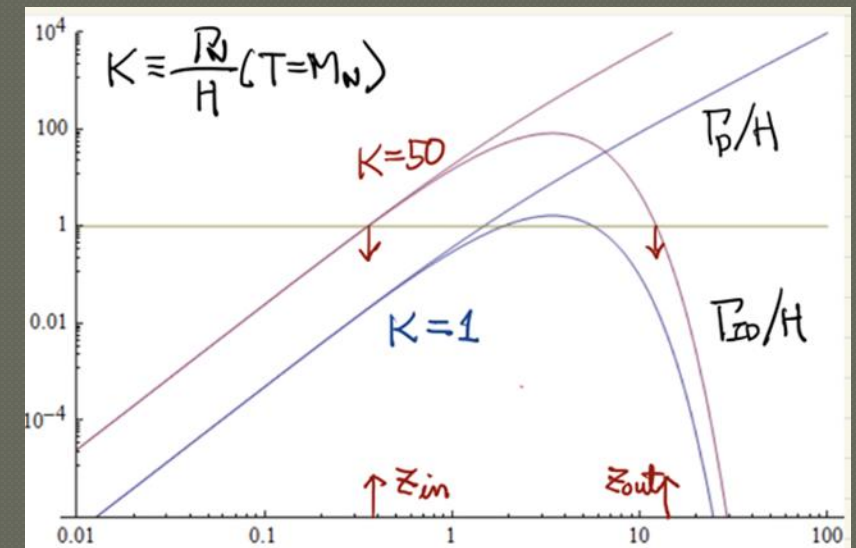
Seesaw & B-L violation

- Seesaw mechanism explaining tiny Majorana neutrinos mass:

$$\mathcal{L} = y_\nu l N H + \frac{1}{2} M_N N N + h.c. \Rightarrow m_\nu = y_\nu \frac{v_H^2}{M_N} y_\nu^T$$

- Define $K \equiv \left(\frac{\Gamma_N}{H} \right)_{T=M_N} \equiv \frac{\tilde{m}_\nu}{\text{meV}}$
- B-L violation by decay/inverse-decay
 $N \leftrightarrow l H$ is in equilibrium around $T \sim M_N$ for $K > 1$.

(Consider $\tilde{m}_\nu = 0.05 \text{eV}$, i.e., $K = 50$.)



Cogenesis of BAU & DM by pNGB

- The final baryon asymmetry is fixed at T_B when BNV interactions go out of

equilibrium: $Y_B \equiv \frac{n_B}{s} = \frac{c_B \dot{\theta} T^2}{s},$

i.e., $Y_B = c_B Y_\theta \left(\frac{T_B}{f_a} \right)^2.$

- Combining with DM condition

$m_a Y_\theta = 0.44 \text{eV},$

$Y_B = c_B \left(\frac{0.44 \text{eV}}{m_a} \right) \left(\frac{T_B}{f_a} \right)^2 \approx 10^{-10}$

- QCD axion:

$T_B = T_{EW} \text{ \& } m_a f_a = m_\pi f_\pi$

$\rightarrow f_a \approx 10^7 c_B \text{GeV (ruled out).}$

- Majoron:

$T_B = M_N / z_{fo}$

$\rightarrow M_N \approx \frac{z_{fo}}{\sqrt{c_B}} \frac{f_a}{10^9 \text{GeV}} \sqrt{\frac{m_a}{6 \text{meV}}}.$

Equilibrium solution with Majoron

- Four Yukawas + EW Sphaleron + charge neutrality (simple case):

$$y_u q u^c H \Rightarrow \mu_q + \mu_{u^c} + \mu_H = 0$$

$$y_d q d^c \tilde{H} \Rightarrow \mu_q + \mu_{d^c} - \mu_H = 0$$

$$y_e l e^c \tilde{H} \Rightarrow \mu_l + \mu_{e^c} - \mu_H = 0$$

$$y_\nu l N H \Rightarrow \mu_l + \mu_H = \dot{\theta} \text{ (LNV)}$$

$$\mathcal{A}_{B+L}(W\tilde{W}) \Rightarrow 3(3\mu_q + \mu_l) = 0$$

$$Y = 0 \Rightarrow 3 \left(\frac{1}{6} 2 \cdot 3\mu_q - \frac{2}{3} 3\mu_{u^c} + \frac{1}{3} 3\mu_{d^c} - \frac{1}{2} 2\mu_l + \mu_{e^c} \right) - \frac{1}{2} 2 \cdot 2 \mu_H = 0$$



$$\mu_B = \frac{1}{3} 3 (2\mu_q - \mu_{u^c} - \mu_{d^c}) = \frac{28}{11} \dot{\theta}$$

$$\mu_L = \frac{1}{3} 3 (2\mu_l - \mu_{e^c}) = -\frac{51}{11} \dot{\theta}$$

$$\mu_{B-L} = \mu_B - \mu_L = \frac{79}{11} \dot{\theta}$$

Cogenesis region

- ❖ Simultaneous generation of BAU & DM:

$$Y_B \approx 0.1 Y_\theta \left(\frac{T_B}{f_a} \right)^2 \approx 0.1 \frac{0.44 \text{eV}}{m_a} \left(\frac{T_B}{f_a} \right)^2$$

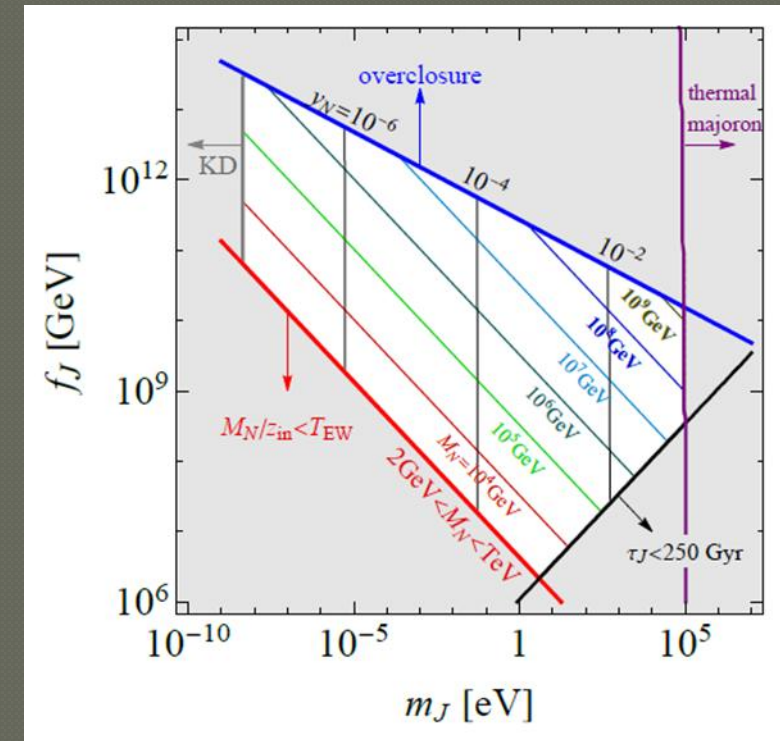
- ❖ $T_B = \frac{M_N}{z_{\text{fo}}} > T_{EW}$:

Trapping condition $m_a \sim 4 \cdot 10^6 \text{ eV } y_N^2$

$$\dot{\theta}_{\text{trp}} \sim m_a > H_{\text{trp}} \quad f_a \lesssim 10^8 y_N^{-1} \text{GeV} \left(\frac{\text{eV}}{m_a} \right)^{1/4}$$

- ❖ $T_B = T_{EW}$ for $M_N < T_{EW}$:

$$f_a \sim 2 \cdot 10^6 \text{GeV} \left(\frac{\text{eV}}{m_a} \right)^{1/2}$$



$$\tau_a^{-1} = \frac{m_\nu^2 m_a}{16\pi f_a^2}$$

Cogenesis by QCD axion=Majoron

- KSVZ+Seesaw

- Chemical equilibrium conditions to be met by separating out the first generation quark Yukawas

- $\mathcal{L}_{PQ} = y_Q \Phi Q Q^c + \frac{1}{2} y_N \Phi N N + h.c.$

PQ charges: $x_Q = x_Q^c = x_N = -x_l = x_e^c = -\frac{1}{2}$.

- $\dot{n}_{q_i} = \dot{n}_{u_i^c} = \dot{n}_{d_i^c} = \dot{n}_{l_i} = \dot{n}_{e_i^c} = 0$
(Yukawas+SS+WS) & charge neutrality.

$$\mu_B = \frac{28}{79} \frac{1}{33} \left(28 c_W - \frac{57 m_u^2 - 15 m_d^2}{m_u^2 + m_d^2} c_S - 153 x_l \right) \dot{\theta} = c_B \dot{\theta}$$

$$c_B = -2.2 \text{ for } c_W = 0 \text{ \& } c_S = N_Q = 1$$

Cogenesis achieved for $Y_B = c_B \left(\frac{0.44 \text{eV}}{m_a} \right) \left(\frac{T_B}{f_a} \right)^2 \Rightarrow M_N \sim 10 \text{TeV} \sqrt{\frac{f_a}{10^9 \text{GeV}}}$

Large $\dot{\theta}$ via AD mechanism

- PQV operator of good quality to generate $Y_\theta = \frac{0.22\text{eV}}{m_a} \approx 400 \left(\frac{f_a}{10^{10}\text{GeV}} \right)$

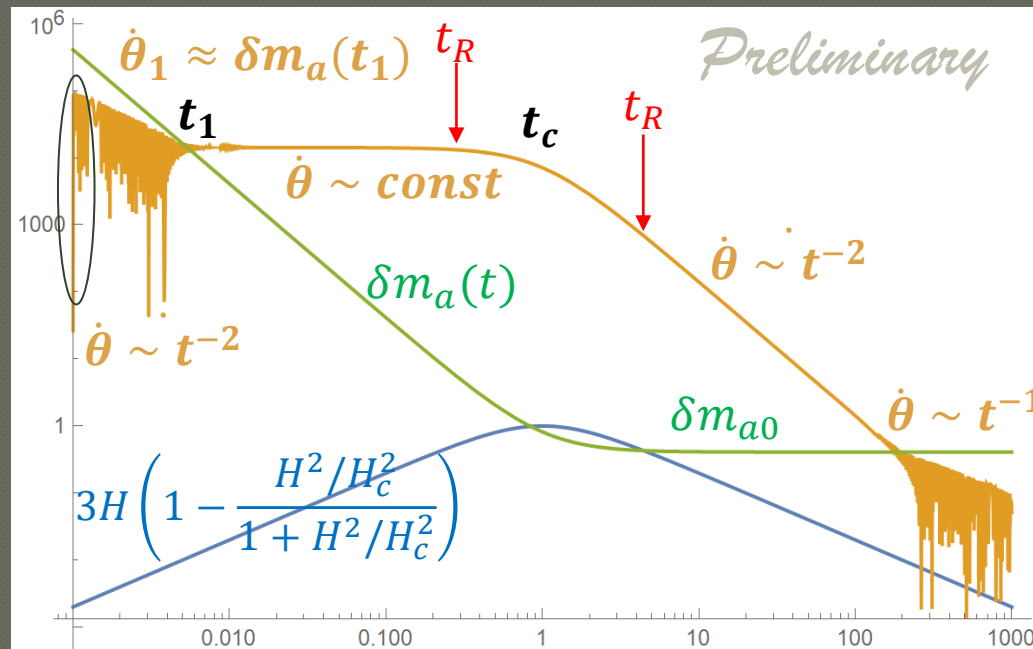
$$V = -\frac{1}{2}(\mu_0^2 + c_H H^2)\phi^2 + \frac{\lambda_0}{4}\phi^4 - \frac{\lambda_n}{\sqrt{2}^n} \frac{\phi^n}{M_P^{n-4}} 2\cos(n\theta) \quad \Phi = \frac{\phi}{\sqrt{2}} e^{i\theta}$$

$$\phi = \sqrt{\frac{\mu_0^2 + c_H H^2}{\lambda_0}} \equiv f_a \sqrt{1 + \frac{H^2}{H_c^2}} \quad \text{with } H_c \equiv \frac{\mu_0}{\sqrt{c_H}} \quad \delta m_{a0}^2 = 2n^2 \frac{\lambda_n}{\sqrt{2}^n} \frac{f_a^n}{M_P^{n-4}} \ll m_a^2$$

During inflation ($H = H_I$) & matter domination ($H = \frac{2}{3t}$)

$$\ddot{\theta} + \frac{2}{t} \left(1 - \frac{t_c^2/t^2}{1 + t_c^2/t^2} \right) \dot{\theta} + \frac{\delta m_{a0}^2}{n} \left(1 + \frac{t_c^2}{t^2} \right)^{\frac{n-2}{2}} \sin(n\theta) = 0$$

Cosmic evolution



$$\odot \quad Y_\theta = \frac{\dot{\theta} f_a^2}{s(T_R)} \approx$$

$$\frac{1}{3} \left(\frac{\delta m_{a0}}{H_c} \right) \left(\frac{H_1}{H_c} \right)^{\frac{n-2}{2}} \left(\frac{f_a}{H_c} \right)^2 \left(\frac{H_c}{M_P} \right)^{\frac{3}{2}} \begin{cases} \left(\frac{H_c}{H_R} \right)^{\frac{3}{2}} & t_R < t_c \\ \left(\frac{H_R}{H_c} \right)^{\frac{1}{2}} & t_R > t_c \end{cases}$$

$$\odot Y_{\text{max}} \approx 30 \left(\frac{\delta m_{a0}}{\mu\text{eV}} \right) \left(\frac{H_1}{10^{13} \text{GeV}} \right)^4 \left(\frac{10^5 \text{GeV}}{H_c} \right)^{\frac{11}{2}}$$

for $n = 10$ & $f_a = 10^{10}$ GeV.

Summary

- Type-I seesaw model with majoron provides an affordable framework for the cogenesis of BAU and DM enjoying the freedom in (m_a, f_a, M_N) .
- Combining PQ and Seesaw, QCD axion can drive cogenesis for the RHN mass around 10~100 TeV.
- Large kinetic misalignment can arise by early dynamics setting $\dot{\theta} = m_a(t) \gg m_a(0)$.