

## STRING THEORY AND THE FIRST HALF OF THE UNIVERSE

Joseph Conlon Planck 2025 Padua

Sanchez-Gonzales, JC, Copeland, Hardy, 2505.14187, JC, Copeland, Hardy, Sanchez-Gonzales 2406.12637, also see Apers, JC, Copeland, Mosny, Revello 2401.04064)









#### Cf Brunelli parallel talk

## I.THE FIRST HALF OF THE UNIVERSE

# THE KNOWN HISTORY OF THE UNIVERSE



# THE KNOWN HISTORY OF THE UNIVERSE



# THE KNOWN HISTORY OF THE UNIVERSE

- State of universe during inflation (10^-30s? Depending on scale of inflation) is constrained by structure of the CMB
- State of universe during nucleosynthesis (10<sup>-2</sup>s to O(100)s ) is constrained by primordial element abundances
- Minimal constraints exist between these two epochs; up to thirty orders of magnitude in time and half the lifetime of the universe on a logarithmic scale

## II. STRING THEORY, PHENOMENOLOGY AND COSMOLOGY

#### STRING THEORY: THE CASE FOR THE PROSECUTION

String theory claims to be a theory of quantum gravity, relevant for physics near the Planck scale  $M_P \simeq 2.4 \times 10^{18} \,\text{GeV}.$ 

Inflation occurs at  $\Lambda_{inf} \lesssim 10^{16} \, {\rm GeV}$  and all other subsequent physics scales are even smaller.

For physicists interested in `ordinary' sub-Planckian physics, why care about string theory?

#### STRING THEORY: THE CASE FOR THE PROSECUTION

If  $\Lambda_{\rm everything\,else} < M_P$  , what can string theory offer that is not already provided by

- Low-energy effective quantum field theory
- The Standard Model and Beyond-the-Standard-Model particle physics (susy, axions, WISPs, etc etc)
- The Standard Cosmology and extensions (dark radiation, quintessence etc)

## WHAT STRING THEORY OFFERS

Theories of quantum gravity allow Planck-scale computation and so give control over Planck-suppressed operators (cf Fermi theory and electroweak theory)

• Control of Planck-suppressed operators implies that in an expansion  $\mathscr{L}(\phi) \to \mathscr{L}(\phi) \left(1 + \frac{\alpha \phi}{M_P} + \frac{\beta \phi^2}{M_P^2} + \frac{\gamma \phi^3}{M_P^3} + \dots\right)$ 

we are able to determine the coefficients  $\alpha$ ,  $\beta$  and  $\gamma$ .

- Such operators need Planck-scale theories; cannot determine them just from low-energy EFT.
- When do they matter?

## WHEN DOES STRING THEORY MATTER?

When does control over Planck-suppressed operators matter?

• Long-lived particles (e.g. moduli) with only gravitational-strength interactions

Their decay rate is  $\Gamma_{\Phi} \sim \frac{1}{8\pi} \frac{m_{\Phi}^3}{M_P^2}$  and they can dominate the energy density of the early universe (moduli problem).

- $\eta$  problem of inflation: how to control contributions to potential of form.  $\delta V(\phi) = \frac{\phi^2}{M_P^2} V(\phi)$  which can destroy flatness of inflationary potential
- Transplanckian field excursions  $\Delta \phi \geq M_P$  in field space.

TRANSPLANCKIAN FIELD EXCURSIONS AND COSMOLOGY Much of cosmology involves scalar fields  $\Phi_i$  coupled to general relativity (inflation / quintessence / dynamical dark energy). Which epochs of the universe involve transPlanckian field excursions?

- Large-field inflation models (e.g. chaotic inflation) resulting in large CMB tensor B modes)
- Extended epochs where scalar field kinetic energy  $\dot{\Phi}^2$  is a large contribution to universe energy density: occurs through either kination or tracker epochs

#### KINATION

• During roll, with universe in kination epoch, field evolves as

$$\Phi(t) = \Phi_0 + \sqrt{\frac{2}{3}} M_P \ln\left(\frac{t}{t_0}\right)$$

• Field moves through  $\sim M_P$  in field space each Hubble time

Extended kination epoch implies large transPlanckian field excursions

- String theorists should **care!** any extended kination epoch requires trans-Planckian field excursions  $\Delta \Phi \gg M_P$ .
- Cosmologists should **care!** any extended kination epoch requires theory of the Planck scale to control it.
- Relatively little work on understanding kination epochs in string theory

#### TRACKER EPOCHS

• A rolling scalar field on an exponential potential  $V = V_0 e^{-\lambda \Phi}$  in a radiation background reaches a tracker solution in which field evolves as

$$\Phi(t) = \Phi_0 + \frac{2M_P}{\lambda} \ln\left(\frac{t}{t_0}\right)$$

- Field motion is slightly slower than for kination but field still moves through  $~\sim M_P$  in field space each Hubble time

Extended tracker epoch implies large transPlanckian field excursions

- String theorists should **care!** any extended tracker epoch requires trans-Planckian field excursions  $\Delta \Phi \gg M_P$
- Cosmologists should care! any extended tracker epoch requires theory of the Planck scale to control it

III. THE EDGE OFTHE WORLD

#### STRING THEORY: CENTRE OF THE WORLD?



#### WHERE IS THE CENTRE OF THE WORLD?



#### WHERE IS THE CENTRE OF THE WORLD?



#### STRING THEORY: CENTRE OF THE WORLD?



## OUR HOME, THE UNIVERSE

Our universe is filled with hierarchies and small numbers

$$\frac{\Lambda_{EW}}{M_P} \sim 10^{-16}$$

$$\frac{\delta \rho_{CMB}}{\rho} \sim 10^{-5} \qquad \Lambda_{cc} \sim 10^{-120} M_P^4$$

$$\alpha_{SU(3)} \sim \frac{1}{11}, \alpha_{SU(2)} \sim \frac{1}{30}, \alpha_{U(1)_Y} \sim \frac{1}{60}$$

$$y_e \sim 10^{-5}, y_\mu \sim 10^{-3}, y_\tau \sim 10^{-2}$$

$$m_\nu \sim 10^{-3} \text{eV}$$

$$\theta_{QCD} \lesssim 10^{-10}$$

## OUR HOME, THE UNIVERSE

- The true string vacuum is the vacuum of this universe
- It must contain a method to generate hierarchies, small couplings and small numbers
- This makes the asymptotic boundaries of moduli space appealing

#### FROM CENTRETO END OFTHE WORLD



#### GETTING TO THE END OF THE WORLD



From review 2303.04819 Cicoli, JC, Maharana, Parameswaran, Quevedo, Zavala

## NOVEL COSMOLOGICAL HISTORY



This motivates a distinctive 'stringy' cosmological history quite distinct from the normal assumption of radiation domination after inflation

## NOVEL COSMOLOGICAL HISTORY



We know almost NOTHING observationally about the ~ 30 orders in magnitude in time between end of inflation and the beginning of nucleosynthesis!



## KINATION EPOCHS

• During kination epoch, kinating field evolves as

$$\Phi(t) = \Phi_0 + \sqrt{\frac{2}{3}} M_P \ln\left(\frac{t}{t_0}\right)$$

- Field moves through  $\sim M_P$  in field space each Hubble time

Long kination epoch implies large transPlanckian field excursions

- String theorists should **care!** trans-Planckian field excursions  $\Delta \Phi \gg M_P$  is home territory
- Novel cosmology: real opportunites for string phenomenology

My talk: Cosmic Strings

## IV. COSMIC SUPERSTRINGS IN EARLY UNIVERSE

#### IV. COSMIC SUPERSTRINGS IN EARLY UNIVERSE

#### A. VARYING TENSION AND GROWTH OF STRING LOOPS

#### DYNAMICS OF COSMIC (SUPER)STRINGS

- Cosmic (super)strings long-studied candidate for new stringy cosmologies Brandenberger+Vafa 86 Sarangi+Tye 02 Copeland+Polchinski 05
- Dynamics of closed strings set by Nambu-Goto action in fixed spacetime background,  $\mu$  is the string tension  $S_{NG} = -\int d^2\xi \,\mu \sqrt{-\gamma}$
- What are the dynamics? (assuming stability and FRLW metric)  $ds^{2} = dt^{2} - a(t)^{2}(dx^{2} + dy^{2} + dz^{2})$









#### DYNAMICS OF COSMIC (SUPER)STRINGS (FIXED TENSION)

- Equations of motion follow from NG action  $x^{\nu;a}_{,a} + \Gamma^{\nu}_{\beta\rho}(g)\gamma^{ad}x^{\beta}_{,d}x^{\rho}_{,a} = 0$
- Focus on circular string loops  $X^{\mu}(t,\sigma) = R(t)(\cos\sigma,\sin\sigma,0)$
- Study equations of motion in FLRW background (gauge choice identifies worldsheet and spacetime time)

#### DYNAMICS OF COSMIC (SUPER)STRINGS (FIXED TENSION)

• Focus on circular string loops

 $X^{\mu}(t,\sigma) = R(t) \big(\cos\sigma, \sin\sigma, 0\big)$ 

• Equations of motion are

$$\left(\varepsilon = \sqrt{\frac{a^2 R^2}{(1 - a^2 \dot{R}^2)}} \equiv a R_{\max}\right)$$

 $\frac{\dot{\varepsilon}}{\varepsilon} = H - 2a^2 \dot{R}^2 H \qquad \langle a^2 \dot{R}^2 \rangle = 1/2$  $\frac{\dot{\varepsilon}}{\dot{\kappa}} + H\dot{R} + \varepsilon^{-2}R + 2H(1 - a^2 \dot{R}^2)\dot{R} = 0$ 

• Loops oscillate with a fixed maximum (physical) size  $R_{max}$ 

#### EVOLUTION OF A LOOP



Circular string loops oscillate in and out back on themselves at constant physical radius and shrink in comoving coordinates

> More complicated exact solutions also exist (Burden, Kibble+Turok)

Loops are left behind as universe expands (and gradually decay by emission of gravitational waves)

- Equations of motion follow from NG action  $x_{,a}^{\nu;a} + \Gamma^{\nu}_{\beta\rho}(g)\gamma^{ad}x_{,d}^{\beta}x_{,a}^{\rho} + \frac{\mu_{,\rho}}{\mu}\gamma^{ab}x_{,a}^{\rho}x_{,b}^{\nu} - \frac{\mu^{,\nu}}{\mu} = 0,$
- Focus on circular string loops  $X^{\mu}(t,\sigma) = R(t)(\cos\sigma,\sin\sigma,0)$
- Study equations of motion in kinating FLRW background (gauge choice identifies worldsheet and spacetime time)

• Focus on circular string loops

 $X^{\mu}(t,\sigma) = R(t)(\cos\sigma,\sin\sigma,0)$ 

• Equations of motion are

$$\varepsilon = \sqrt{\frac{a^2 R^2}{(1 - a^2 \dot{R}^2)}} \equiv a R_{\text{max}}$$
$$\frac{\dot{\varepsilon}}{\varepsilon} = H - a^2 \dot{R}^2 \left(2H + \frac{\dot{\mu}}{\mu}\right) \qquad \langle a^2 \dot{R}^2 \rangle = 1/2$$
$$\dot{\varepsilon} + H\dot{R} + \varepsilon^{-2}R + \left(2H + \frac{\dot{\mu}}{\mu}\right)(1 - a^2 \dot{R}^2)\dot{R} = 0$$

• High-frequency oscillation at (physical) amplitude  $R_{max}$  but....

• Equations of motion are

$$\varepsilon = \sqrt{\frac{a^2 R^2}{(1 - a^2 \dot{R}^2)}} \equiv a R_{\text{max}}$$
$$\frac{\dot{\varepsilon}}{\varepsilon} = H - a^2 \dot{R}^2 \left(2H + \frac{\dot{\mu}}{\mu}\right) \qquad \langle a^2 \dot{R}^2 \rangle = 1/2$$

- Decreasing tension causes loops to grow with cosmic time
- Right hand side of equation determines precisely how loops compared to scale factor (cf  $\frac{\dot{a}}{a} = H$ )

• Equations of motion are

$$\varepsilon = \sqrt{\frac{a^2 R^2}{(1 - a^2 \dot{R}^2)}} \equiv a R_{\text{max}}$$
$$\frac{\dot{\varepsilon}}{\varepsilon} = H - a^2 \dot{R}^2 \left(2H + \frac{\dot{\mu}}{\mu}\right) \qquad \langle a^2 \dot{R}^2 \rangle = 1/2$$

- Decreasing tension causes loops to grow with cosmic time
- If crossed out term vanishes, oscillating loops grow precisely with the scale factor (cf  $\frac{\dot{a}}{a} = H$ )

• Equations of motion are

$$\varepsilon = \sqrt{\frac{a^2 R^2}{(1 - a^2 \dot{R}^2)}} \equiv a R_{\text{max}}$$
$$\frac{\dot{\varepsilon}}{\varepsilon} = H - a^2 \dot{R}^2 \left(2H + \frac{\dot{\mu}}{\mu}\right) \qquad \langle a^2 \dot{R}^2 \rangle = 1/2$$

. When  $2H + \frac{\mu}{\mu} < 0$ , oscillating loops grow faster than the scale factor **and will percolate given enough** time

#### KINATION AND TIME-VARYING TENSION

- We want to make  $2H + \frac{\mu}{\mu}$  as negative as possible  $\mu$
- This requires

(a) 
$$H \equiv \frac{\dot{a}}{a}$$
 as small as possible  
(b)  $\frac{\dot{\mu}}{\mu}$  as large and negative as possible

Kination epochs are ideal as
(a) a(t) ~ t<sup>1/3</sup> and so growth is as slower than any other fluid
(b) All energy is in kinetic evolution of a modulus and so maximises rate of change of tension vev

#### KINATION AND TIME-VARYING TENSION

During volume modulus kination, volume grows with time

$$\frac{\mathcal{V}}{\mathcal{V}_0} = \frac{t}{t_0}$$

- For superstrings,  $G\mu \sim m_s^2$  and so  $\mu \propto t^{-1}$  using standard relationship  $m_s \sim \frac{M_P}{\sqrt{\mathcal{V}}}$
- It follows that

$$2H + \frac{\mu}{\mu} = -H < 0$$

and so loops of fundamental strings grow faster than the scale factor!

 Loops of fundamental strings grow in comoving coordinates and can percolate!

• During kination, scale factor and loop radius grow as

$$a(t) \sim t^{1/3}$$
$$R_{max}(t) \sim t^{1/2}$$

• In comoving coordinates,

$$R_{max,comoving} \sim t^{1/6} \sim \left(\frac{\mathcal{V}_f}{\mathcal{V}_i}\right)^{1/6}$$

 Long kination epochs (closely tie to vacua in asymptotic region of moduli space) essential to give percolation











## IV. COSMIC SUPERSTRINGS IN EARLY UNIVERSE

#### B. SUPERSTRING PHASES

Sanchez-Gonzales, JC, Copeland, Hardy, 2505.14187

#### SUPERSTRING PHASES

 Consider a scalar field (volume modulus) rolling on an exponential potential,

$$V = V_0 e^{-\lambda \Phi}$$

• For string theory LVS potential,  $\lambda = \sqrt{\frac{27}{2}}$ 

• What is the stable attractor?

#### LOOP ENERGY DENSITY

Energy in an individual loop is set as

$$E_{loop} = l \times \mu$$

As  $l \propto \sqrt{\frac{\mu_0}{\mu(t)}}$ , overall energy density in loops scales as

$$\rho_{loops} \sim \frac{1}{a^3} \sqrt{\frac{\mu(t)}{\mu_0}}$$

#### LOOP ENERGY DENSITY

#### During kination,

 $\rho_{loops} \sim \frac{1}{a^{9/2}}$ 

#### During radiation tracker,

 $\rho_{loops} \sim \frac{1}{a^{11/3}}$ 

## LOOP ENERGY DENSITY

Loop energy density grows relative to background and stable attractor is a loop tracker with

$$\Omega_{loops} = \frac{3}{4}, \Omega_{kin} = \frac{1}{4}$$

Post-inflationary tracker with three-quarters of universe energy density in form of superstrings!

#### LOOPTRACKER



## CONCLUSIONS

- Universe between inflation and nucleosynthesis is poorly constrained — opportunity for string cosmology
- Evolution of volume modulus causes fundamental string loops to grow and they can percolate into a cosmic string network
- Also a novel post-inflationary string loop tracker with 75% of universe energy density in form of superstring loops
- With LVS final vacuum, string network today with  $G\mu \sim 10^{-10}$
- Such a fundamental cosmic string network with  $10^{-7} \leq G\mu \leq 10^{-11}$ - in reach of upcoming experiments (cf NANOGrav)