# The Timeless Primordial Universe & Its Combinatorial Origin





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Instituto Galego de Física de Altas Enerxías 27 May 2025 - Planck 2025

## The Timeless Priomordial Universe



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## The Timeless Primordial Universe

Can we put constraints on which state can propagate during inflation in a completely model independent way?

• What is the imprint of the inflationary physics in the analytic structure of the relevant observables?

3 What are the rules governing physical processes at energies as large as  $H|_{infl} \sim 10^{14} \, GeV?$ 

What does it even mean?



A theory is defined by..



What does it even mean?



A theory is defined by.. a Lagrangian



What does it even mean?



**Computables** Enjoy *sufficiently* physical features (e.g.: physical dofs, gauge inv.,...)

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What does it even mean?



**Computables** Enjoy *sufficiently* physical features (e.g.: physical dofs, gauge inv.,...)

2 Basic guiding principles that *any reasonable theory* ought to satisfy

What does it even mean?



**Computables** Enjoy *sufficiently* physical features (e.g.: physical dofs, gauge inv.,...)

2 Basic guiding principles that *any reasonable theory* ought to satisfy



 $\label{eq:Constraints} \ensuremath{\mathbb{E}} \ensuremath{\mathsf{Constraints}} \ensuremath{\mathbb{E}} \ensuremath{\mathsf{Constraints}} \ensuremath{\mathsf{simultaneous}} \ensuremath{\mathsf{validity}} \ensuremath{\mathsf{of}} \ensuremath{\mathsf{these}} \ensuremath{\mathsf{principles}} \ensuremath{\mathsf{simultaneous}} \ensuremath{\mathsf{validity}} \ensuremath{\mathsf{of}} \ensuremath{\mathsf{these}} \ensuremath{\mathsf{of}} \ensuremath{\mathsf{these}} \ensuremath{\mathsf{validity}} \ensuremath{\mathsf{of}} \ensuremath{\mathsf{constraints}} \ensuremath{\mathsf{validity}} \ensuremath{\mathsf{of}} \ensuremath{\mathsf{these}} \ensuremath{\mathsf{of}} \ensuremath{\mathsf{validity}} \ensuremath{\mathsf{of}} \ensuremath{\mathsf{validity}} \ensuremath{\mathsf{of}} \ensuremath{\mathsf{validity}} \ensuremath{\mathsf{of}} \ensuremath{\mathsf{of}} \ensuremath{\mathsf{validity}} \ensuremath{\mathsf{of}} \ensuremath{\mathsf{validity}} \ensuremath{\mathsf{of}} \ensuremath{\mathsf{constraints}} \ensuremath{\mathsf{validity}} \ensuremath{\mathsf{of}} \ensuremath{\mathsf{validity}} \ensuremath{\mathsf{of}} \ensuremath{\mathsf{of}}$ 

#### **Guiding Principles**



Cluster decomposition



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#### **Guiding** Principles



Cluster decomposition



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#### **Guiding Principles**



Cluster decomposition



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#### **Guiding** Principles





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#### **Guiding Principles**



Cluster decomposition



Theorems

Spin 1 Self-interaction just for different species

Spin 2 Graviton uniqueness  $\mathcal{N} = 1$  Sugra Spin >2 No self-interactions No interactions with  $s \leq 2$ No elementary massive particles

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#### **Guiding Principles**



Cluster decomposition





Spin 1: Charge conservation

Spin 2: Equivalence principle

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Poincaré invariance

Unitarity









Poincaré invariant operator:  $\hat{O}$  $\langle 3, 4|\hat{O}|1, 2\rangle := \delta (p_1 + p_2 - p_3 - p_4)$  Unitary S-matrix:  $\hat{S} \mid \hat{S}\hat{S}^{\dagger} = \hat{\mathbb{I}} = \hat{S}^{\dagger}\hat{S}$  $\hat{S} := e^{i\lambda\hat{O}}$ 

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Poincaré invariant operator: 
$$\hat{O}$$
  
 $\langle 3, 4|\hat{O}|1, 2\rangle := \delta (p_1 + p_2 - p_3 - p_4)$ 
Unitary S-matrix:  $\hat{S} \mid \hat{S}\hat{S}^{\dagger} = \hat{\mathbb{I}} = \hat{S}^{\dagger}\hat{S}$   
 $\hat{S} := e^{i\lambda\hat{O}}$ 
 $\langle 3, 4|\hat{S}|1, 2\rangle = \delta (p_1 + p_2 - p_3 - p_4) \left\{ 1 + i\lambda + \lambda^2 \vartheta(s - 4m^2) + \ldots \right\}$ 

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Poincaré invariant operator:  $\hat{O}$   $\langle 3, 4|\hat{O}|1, 2\rangle := \delta (p_1 + p_2 - p_3 - p_4)$   $\langle 3, 4|\hat{S}|1, 2\rangle = \delta (p_1 + p_2 - p_3 - p_4) \left\{ 1 + i\lambda + \lambda^2 \vartheta (s - 4m^2) + \ldots \right\}$  $1 - \frac{3}{2} - \frac{3}{4}$ 

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Poincaré invariant operator: $\hat{\mathcal{O}}$ Unitary S-matrix: $\hat{S} \mid \hat{S}\hat{S}^{\dagger} = \hat{\mathbb{I}} = \hat{S}^{\dagger}\hat{S}$  $\langle 3, 4 \mid \hat{\mathcal{O}} \mid 1, 2 \rangle := \delta \left( p_1 + p_2 - p_3 - p_4 \right)$  $\hat{S} := e^{i\lambda\hat{\mathcal{O}}}$ 

$$\langle 3,4|\hat{S}|1,2\rangle = \delta\left(p_1+p_2-p_3-p_4\right)\left\{1+i\lambda+\lambda^2\vartheta(s-4m^2)+\ldots\right\}$$



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# The particle physics lesson



#### Simplest processes + constraints



Language:

#### Fields:

redundant organisation of redundant degrees of freedom

**Riemannian geometry:** mathematical tool for the explotation of the equivalence principle

[S. Weinberg, Gravitation and Cosmology]

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Language:

**Fields:** redundant organisation of redundant degrees of freedom **Riemannian geometry:** mathematical tool for the explotation of the equivalence principle

[S. Weinberg, Gravitation and Cosmology]

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Problems as our inability to extrapolate high energy knowledge from our low energy one:

(Quantum) gravity;

Higgs: it is necessary to extrapolate high energy physics from the low energy one

# Deeper understanding of the physics encoded into cosmological observables

Novel rules which can allow to go beyond the regime in which the combinatorial description has been formulated

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$$\langle \Phi(\vec{p}_1) \cdots \Phi(\vec{p}_n) \rangle = \int \mathcal{D}\Phi \mathfrak{P}[\Phi] \Phi(\vec{p}_1) \cdots \Phi(\vec{p}_n)$$

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$$\langle \Phi(\vec{p}_1) \cdots \Phi(\vec{p}_n) \rangle = \int \mathcal{D} \Phi \mathfrak{P}[\Phi] \Phi(\vec{p}_1) \cdots \Phi(\vec{p}_n)$$
  
Probability  
distribution

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$$\langle \Phi(\vec{p}_1) \cdots \Phi(\vec{p}_n) \rangle = \frac{\int \mathcal{D}\Phi \,\Psi^{\dagger}[\Phi] \Phi(\vec{p}_1) \cdots \Phi(\vec{p}_n) \Psi[\Phi]}{\int \mathcal{D}\Phi \,|\Psi[\Phi]|^2}$$
$$\Psi[\Phi] := \langle \Phi | \hat{\mathcal{T}} \exp\left\{-i \int_{-\infty}^0 d\eta \,H(\eta)\right\} |0\rangle$$

Wavefunction of the universe (transition amplitude from  $|0\rangle$  to  $\langle\Phi|)$ 

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$$\langle \Phi(\vec{p}_1)\cdots\Phi(\vec{p}_n)\rangle = \frac{\int \mathcal{D}\Phi \,\Psi^{\dagger}[\Phi]\Phi(\vec{p}_1)\cdots\Phi(\vec{p}_n)\Psi[\Phi]}{\int \mathcal{D}\Phi \,|\Psi[\Phi]|^2}$$

$$\Psi[\Phi] := \langle \Phi | \hat{\mathcal{T}} \exp \left\{ -i \int_{-\infty}^{0} d\eta \, H(\eta) 
ight\} | 0 
angle$$

Wavefunction of the universe (transition amplitude from  $|0\rangle$  to  $\langle\Phi|)$ 

#### Perturbation theory





 $\eta = -\infty$ 






#### Observables & Their Analytic Structure



External kinematics: 
$$X_s := \sum_{j \in s} |\vec{p}^{(j)}|, y_e := \left| \sum_{j \in s_e} \vec{p}^{(j)} \right| (e \in \mathcal{E} \setminus \{\mathcal{E}^{(L)}\})$$
  
Loop momenta:  $y_{e_1} := |\vec{l}|, y_{e_2} := |\vec{l} + \vec{p}^{(2)}|, \dots (e \in \mathcal{E}^{(L)})$ 

$$q_{\mathfrak{g}}(x,y) := \sum_{s \in \mathcal{V}_{\mathfrak{g}}} x_s + \sum_{e \in \mathcal{E}_{\mathfrak{g}}^{ext}} y_e$$



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O Cosmological integrands: Singularities, combinatorics & computation

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#### 2 The IR/UV structure of cosmological integrals



One loop corrections without integration

Cosmological integrands: Singularities, combinatorics & computation

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Cosmological Integrals



External kinematics: 
$$X_s := \sum_{j \in s} |\vec{p}^{(j)}|, y_e := \left| \sum_{j \in s_e} \vec{p}^{(j)} \right| (e \in \mathcal{E} \setminus \{\mathcal{E}^{(L)}\})$$
  
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Cosmological Integrals



Cosmological Integrals













$$------ x_1 + x_2 = 0$$





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A flavour of cosmological polytopes



The singularities form a bounded region to which a function  $\Omega$  is naturally associated

$$\Omega = \frac{1}{(x_1 + x_2)(x_1 + y)(y + x_2)} \equiv \frac{\mathfrak{n}_{\delta}}{q_{\mathcal{G}}q_{\mathfrak{g}_1}q_{\mathfrak{g}}}$$

A flavour of cosmological polytopes



A flavour of cosmological polytopes





[P.B., W. Torres Bobadilla; '21]

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$$\Omega = \prod_{\mathfrak{g} \in \mathfrak{G}_{o}} \frac{1}{q_{\mathfrak{g}}(x, y)} \sum_{\{\mathfrak{G}_{c}\}} \prod_{\mathfrak{g}' \in \mathfrak{G}_{c}} \frac{1}{q_{\mathfrak{g}'}(x, y)}$$

[**P.B.**, W. Torres Bobadilla; '21]

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$$\Omega = \prod_{\mathfrak{g} \in \mathfrak{G}_{\circ}} \frac{1}{q_{\mathfrak{g}}(x, y)} \sum_{\{\mathfrak{G}_{c}\}} \prod_{\mathfrak{g}' \in \mathfrak{G}_{c}} \frac{1}{q_{\mathfrak{g}'}(x, y)}$$

Example:  $\{\mathfrak{G}_{\circ}\}$  for a triangle graph



[P.B., W. Torres Bobadilla; '21]



$$\Omega = \prod_{\mathfrak{g} \in \mathfrak{G}_{\circ}} \frac{1}{q_{\mathfrak{g}}(x, y)} \sum_{\{\mathfrak{G}_{c}\}} \prod_{\mathfrak{g}' \in \mathfrak{G}_{c}} \frac{1}{q_{\mathfrak{g}'}(x, y)}$$

[P.B., W. Torres Bobadilla; '21]

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Compatibility conditions allow to:

- write all the possible representations without spurious singularities
- 2 make manifest the symmetries that maps a simplex into another one
- improve analytical/numerical efficiency of the integration.

# (Weighted) Cosmological Polytopes & $I_G$ : A dictionary

Cosmological Polytope  $\mathcal{P}_{\mathcal{G}}$ 

Canonical form  $\boldsymbol{\omega}$ 

Triangulations

Boundaries (Faces)

Canonical form preserving transformations

> Paths along contiguous vertices

Cosmological Integral  $\mathcal{I}_\mathcal{G}$ 

Integrand of  $\mathcal{I}_{\mathcal{G}}$ 

Representations for the integrand

Residues of the integrands

Symmetries of the integrand

Symbols for  $\mathcal{I}_{\mathcal{G}}$ 

# (Weighted) Cosmological Polytopes & $I_G$ : A dictionary

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Residues of the integrands

Symmetries of the integrand

Symbols for  $\mathcal{I}_{\mathcal{G}}$ 

Trascendental<br/>function $f_k = \int_a^b d \log R_1 \circ \ldots \circ d \log R_k$ Iterated<br/>integralSymbols $\mathcal{S}(f_k) := R_1 \otimes \ldots \otimes R_k$ 

# Towards a combinatorial RG: The JR/UV structure of cosmological integrals

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[P.B., F. Vazão; 24]

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$$\underbrace{\begin{array}{c} \begin{array}{c} y_{12} \\ \mathbf{x}_1 \end{array}}_{X_1} x_2 = \int_{\mathbb{R}_+} \prod_{j=1}^2 \left[ \frac{dx_j}{x_j} x_j^{\alpha} \right] \frac{1}{(x_1 + x_2 + \mathcal{X}_{\mathcal{G}})(x_1 + \mathcal{X}_{g_1})(x_2 + \mathcal{X}_{22})} \end{array}$$

[P.B., F. Vazão; 24]

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$$\underbrace{ \sum_{x_1}^{y_{12}} }_{x_2} = \int_{\mathbb{R}_+} \prod_{j=1}^2 \left[ \frac{dx_j}{x_j} x_j^{\alpha} \right] \frac{1}{(x_1^1 x_2^0 + x_1^0 x_2^1 + \mathcal{X}_{\mathcal{G}} x_1^0 x_2^0)(x_1^1 x_2^0 + \mathcal{X}_{\mathfrak{g}_1} x_1^0 x_2^0)(x_1^0 x_2^1 + \mathcal{X}_{\mathfrak{g}_2} x_1^0 x_2^0)}$$



The integral converges for values of  $\alpha$  that identifies points inside the Newton polytope

[P.B., F. Vazão; 24]

$$\underbrace{ \begin{array}{c} & y_{12} \\ \bullet \\ x_1 \end{array} }_{X_1} \underbrace{ \begin{array}{c} y_{12} \\ x_2 \end{array} }_{X_2} = \int_{\mathbb{R}_+} \prod_{j=1}^2 \left[ \frac{dx_j}{x_j} x_j^{\alpha} \right] \frac{1}{(x_1^1 x_2^0 + x_1^0 x_2^1 + \mathcal{X}_{\mathcal{G}} x_1^0 x_2^0)(x_1^1 x_2^0 + \mathcal{X}_{\mathfrak{g}_1} x_1^0 x_2^0)(x_1^0 x_2^1 + \mathcal{X}_{\mathfrak{g}_2} x_1^0 x_2^0)}$$





The integral converges for values of  $\alpha$  that identifies points inside the Newton polytope

$$\mathfrak{W}^{(12)} = \begin{pmatrix} 2\alpha - 3\\ 1\\ 1 \end{pmatrix}, \ \mathfrak{W}^{(1)} = \begin{pmatrix} \alpha - 2\\ 1\\ 0 \end{pmatrix}, \ \mathfrak{W}^{(2)} = \begin{pmatrix} \alpha - 2\\ 0\\ 1 \end{pmatrix}, \ \mathfrak{W}^{'(1)} = \begin{pmatrix} -\alpha\\ -1\\ 0 \end{pmatrix}, \ \mathfrak{W}^{'(2)} = \begin{pmatrix} -\alpha\\ 0\\ -1 \end{pmatrix},$$

The integral diverges in the direction  $\mathfrak{e}$  if the related  $\lambda$  is  $\geq 0$ 

[P.B., F. Vazão; 24]

$$\underbrace{ \begin{array}{c} {}_{y_{12}}}_{x_1} \quad {}_{x_2} = \int_{\mathbb{R}_+} \prod_{j=1}^2 \left[ \frac{dx_j}{x_j} \, x_j^{\alpha} \right] \frac{1}{(x_1^1 x_2^0 + x_1^0 x_2^1 + \mathcal{X}_{\mathcal{G}} x_1^0 x_2^0)(x_1^1 x_2^0 + \mathcal{X}_{\mathfrak{g}1} x_1^0 x_2^0)(x_1^0 x_2^1 + \mathcal{X}_{\mathfrak{g}2} x_1^0 x_2^0)}$$



E.g.: if  $\lambda^{(12)} \longrightarrow 0$ : sector decomposition

$${\cal I}^{
m div}_{\Delta_{j,12}} \ = \ \int_{0}^{1} {d\zeta_{j}\over \zeta_{j}} \ {(\zeta_{j})^{-\lambda^{(j)}}\over 1+\zeta_{j}} \ imes \ \int_{0}^{1} {d\zeta_{12}\over \zeta_{12}} \ (\zeta_{12})^{-\lambda^{(12)}}$$



[P.B., F. Vazão; 24]

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$$x_{1} \bigoplus x_{2} = \int_{\mathbb{R}_{+}} \prod_{j=1}^{2} \left[ \frac{dx_{j}}{x_{j}} x_{j}^{\alpha} \right] \int_{\Gamma} \prod_{e \in \mathcal{E}} \left[ \frac{dy_{e}}{y_{e}} y_{e}^{\beta} \right] \mu(y) \times$$
$$\times \frac{2(x_{1} + x_{2} + y_{a} + y_{b} + \mathcal{X}_{\mathcal{G}})}{(x_{1} + x_{2} + \mathcal{X}_{\mathcal{G}})(x_{1} + x_{2} + y_{a} + \mathcal{X}_{\mathcal{G}})(x_{1} + x_{2} + y_{b} + \mathcal{X}_{\mathcal{G}})(x_{1} + y_{a} + y_{b} + \mathcal{X}_{1})(x_{2} + y_{a} + y_{b} + \mathcal{X}_{2})}$$

$$\mu(y) \sim \left[\frac{\text{Vol}^2 \Sigma_2(y_e^2, P^2)}{\text{Vol}^2 \Sigma_1(P^2)}\right]^{\frac{d-3}{2}}$$

$$y_a \swarrow y_b$$

$$P$$

$$\Gamma \implies \text{Volume of the triangle, and all its side, are positive}$$

$$(y_a + y_b + P)(y_a + y_b - P)(y_a - y_b + P)(-y_a + y_b + P) \ge 0,$$

$$y_a \ge 0, \quad y_b \ge 0, \quad P \ge 0$$

[P.B., F. Vazão; 24]

$$\begin{aligned} x_1 & \longrightarrow x_2 = \int_{\mathbb{R}_+} \prod_{j=1}^2 \left[ \frac{dx_j}{x_j} x_j^{\alpha} \right] \int_{\Gamma} \prod_{e \in \mathcal{E}} \left[ \frac{dy_e}{y_e} y_e^{\beta} \right] \mu(y) \times \\ & \times \frac{2(x_1 + x_2 + y_a + y_b + \mathcal{X}_{\mathcal{G}})}{(x_1 + x_2 + \mathcal{X}_{\mathcal{G}})(x_1 + x_2 + y_a + \mathcal{X}_{\mathcal{G}})(x_1 + x_2 + y_b + \mathcal{X}_{\mathcal{G}})(x_1 + y_a + y_b + \mathcal{X}_1)(x_2 + y_a + y_b + \mathcal{X}_2)} \end{aligned}$$



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[P.B., F. Vazão; 24]

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$$\mathcal{I}_{\mathcal{G}} = \int_{0}^{+\infty} \prod_{s \in \mathcal{V}} \left[ \frac{dx_s}{x_s} \, x_s^{\alpha_s} \right] \int_{\Gamma} \prod_{e \in \mathcal{E}^{(L)}} \left[ \frac{dy_e}{y_e} \, y_e^{\beta_e} \right] \, \mu_d(y) \, \frac{\mathfrak{n}_{\delta}(x + X, y)}{\prod_{g \subseteq \mathcal{G}} \left[ q_g(x + X, y) \right]^{\tau_g}}$$

[P.B., F. Vazão; 24]

$$\mathcal{I}_{\mathcal{G}} = \int_{0}^{+\infty} \prod_{s \in \mathcal{V}} \left[ \frac{dx_s}{x_s} \, x_s^{\alpha_s} \right] \int_{\Gamma} \prod_{e \in \mathcal{E}^{(L)}} \left[ \frac{dy_e}{y_e} \, y_e^{\beta_e} \right] \, \mu_d(y) \, \frac{\mathfrak{n}_{\delta}(x + X, y)}{\prod_{\mathfrak{g} \subseteq \mathcal{G}} \left[ q_{\mathfrak{g}}(x + X, y) \right]^{\tau_{\mathfrak{g}}}}$$

The asymptotic structure of  $\mathcal{I}_{\mathcal{G}}$  is captured by:

) a nestohedron, which is determined by the underlying cosmological polytope  $\mathcal{P}_{\mathcal{G}}$ , and whose facets are fixed via subgraphs

$$\mathfrak{W}^{(i_1\ldots i_{\boldsymbol{a}_{\mathfrak{s}}^{(\mathfrak{g})}+\boldsymbol{a}_{\mathfrak{s}}^{(\mathfrak{g})})}} = \begin{pmatrix} \lambda^{(i_1\ldots j_{\boldsymbol{a}_{\mathfrak{s}}^{(\mathfrak{g})}+\boldsymbol{a}_{\mathfrak{s}}^{(\mathfrak{g})})} \\ \mathfrak{e}_{(i_1\ldots i_{\boldsymbol{a}_{\mathfrak{s}}^{(\mathfrak{g})}+\boldsymbol{a}_{\mathfrak{s}}^{(\mathfrak{g})})} \end{pmatrix}, \qquad \lambda^{(i_1\ldots i_{\boldsymbol{a}_{\mathfrak{s}}^{(\mathfrak{g})}+\boldsymbol{a}_{\mathfrak{s}}^{(\mathfrak{g})})} = \sum_{\boldsymbol{s}\in\mathcal{V}_{\mathfrak{g}}} \alpha_{\boldsymbol{s}} + \sum_{\boldsymbol{e}\in\mathcal{E}^{(L)}} \beta_{\boldsymbol{e}} - \sum_{\boldsymbol{\mathfrak{g}}'\in(\mathsf{tubings})} \tau_{\boldsymbol{\mathfrak{g}}'}$$

The integral diverges in the direction  $\mathfrak{e}$  if the related  $\lambda$  is  $\geq 0$ 

the contour of the loop integration  $\Gamma$ , which selects the divergent directions among the  $\mathfrak{W}$ 's of the nestohedron

[P.B., F. Vazão; 24]

$$\mathcal{I}_{\mathcal{G}} = \int_{0}^{+\infty} \prod_{s \in \mathcal{V}} \left[ \frac{dx_s}{x_s} \, x_s^{\alpha_s} \right] \int_{\Gamma} \prod_{e \in \mathcal{E}^{(L)}} \left[ \frac{dy_e}{y_e} \, y_e^{\beta_e} \right] \, \mu_d(y) \, \frac{\mathfrak{n}_{\delta}(x + X, y)}{\prod_{\mathfrak{g} \subseteq \mathcal{G}} \left[ q_{\mathfrak{g}}(x + X, y) \right]^{\tau_{\mathfrak{g}}}}$$

This combinatorial picture allows to:

- straightforwardly determine both the directions along which I<sub>G</sub> can diverge and their degree of divergence;
- straightforwardly compute leading and subleading divergences (both in the IR and in the UV) via sector decomposition;
- the leading divergence in the IR are associated to the restriction of the underlying cosmological polytope onto special hyperplanes;
- write a systematic substraction that produces IR-finte quantities.

[P.B., G. Brunello, M. K. Mandal, P. Mastrolia, F. Vazão; 24]

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$$\mathcal{I}_{\mathcal{G}} = \int_{0}^{+\infty} \prod_{s \in \mathcal{V}} \left[ \frac{dx_s}{x_s} \, x_s^{\alpha_s} \right] \int_{\Gamma} \prod_{e \in \mathcal{E}^{(L)}} \left[ \frac{dy_e}{y_e} \, y_e^{\beta_e} \right] \, \mu_d(y) \, \frac{\mathfrak{n}_{\delta}(x + X, y)}{\prod_{\mathfrak{g} \subseteq \mathcal{G}} \left[ q_{\mathfrak{g}}(x + X, y) \right]^{\tau_{\mathfrak{g}}}}$$

[P.B., G. Brunello, M. K. Mandal, P. Mastrolia, F. Vazão; 24]

$$\mathcal{I}_{\mathcal{G}} = \int_{0}^{+\infty} \prod_{s \in \mathcal{V}} \left[ \frac{dx_{s}}{x_{s}} \, x_{s}^{\alpha_{s}} \right] \int_{\Gamma} \prod_{e \in \mathcal{E}^{(L)}} \left[ \frac{dy_{e}}{y_{e}} \, y_{e}^{\beta_{e}} \right] \, \mu_{d}(y) \, \frac{\mathfrak{n}_{\delta}(x + X, y)}{\prod_{\mathfrak{g} \subseteq \mathcal{G}} \left[ q_{\mathfrak{g}}(x + X, y) \right]^{\tau_{\mathfrak{g}}}}$$

can be expressed in terms of

$$\begin{array}{ll} \text{(twisted period} \\ \text{integrals)} \end{array} \quad \quad \mathcal{I}_{\{\tau_{\mathfrak{g}}\}}^{(j)} := \int_{\Gamma} \mu_{d} \, \varphi \ \, , \qquad \varphi := \frac{\prod_{e \in \mathcal{E}^{(L)}} dy_{e}}{\prod_{\mathfrak{g} \in \mathfrak{G}^{(j)} \cup \{e\}} [q_{\mathfrak{g}}(y)]^{\tau_{\mathfrak{g}}}},$$

Each of these integrals can be expressed as a *finite* linear combination of master integrals

$$\mathcal{I}_{\{\tau_{\mathfrak{g}}\}}^{(j)} := \sum_{j=1}^{\nu} c_{j} \mathcal{J}_{j}, \qquad d\mathcal{J} = d\mathbb{A} \mathcal{J}$$
  
Canonical form:  $d\mathcal{J} = \varepsilon d\mathbb{A} \mathcal{J} \Rightarrow \mathcal{J} = \mathbb{P} \exp\left\{\varepsilon \int_{\Gamma} d\mathbb{A}\right\} \mathcal{J}_{\circ}$ 

[P.B., G. Brunello, M. K. Mandal, P. Mastrolia, F. Vazão; 24]

The system of differential equations has a block-triangular form, and for each block can be rewritten in terms of a higher order differential equation for a single  $\mathcal{J}_i$ ;

[P.B., G. Brunello, M. K. Mandal, P. Mastrolia, F. Vazão; 24]

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First clues on constraints on cosmological processes: perturbative unitarity, flat-space limit, factorisations, higher-codimensions singularities

General framework to have a direct formulation with IR safe observables & a novel formulation of the RG

We scratched the surface of the one-loop structure: first glimpses of its analytic structure and its space of functions.

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