

Advanced tensor network methods for many body physics
Quantum dynamics in two dimensions etc.

Luka Pavešić, University of Padova

Quantum many body problems from the viewpoint of tensor network methods

1D models

nearest-neighbour interactions

2+ D models

long-range interactions

long time dynamics (thermalization)

generic quantum chaotic systems

volume-law entangled systems



easy

unreasonable

Quantum many body problems from the viewpoint of tensor network methods



1D models

nearest-neighbour interactions

MPS

2+ D models

long-range interactions

long time dynamics (thermalization)

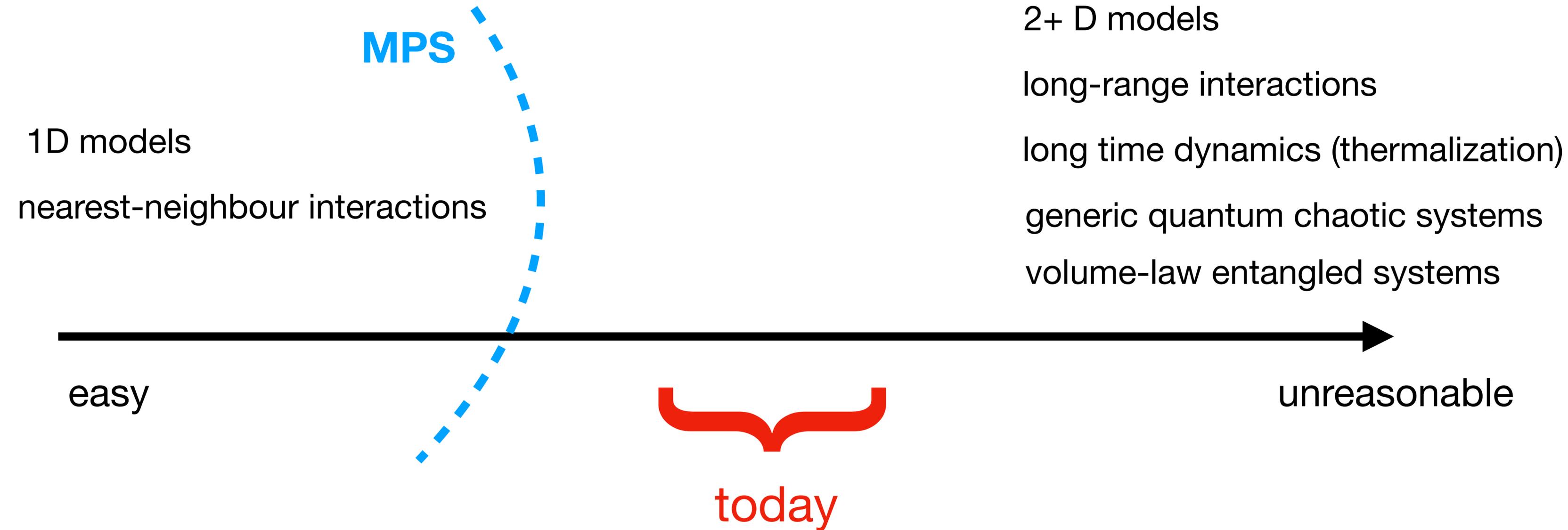
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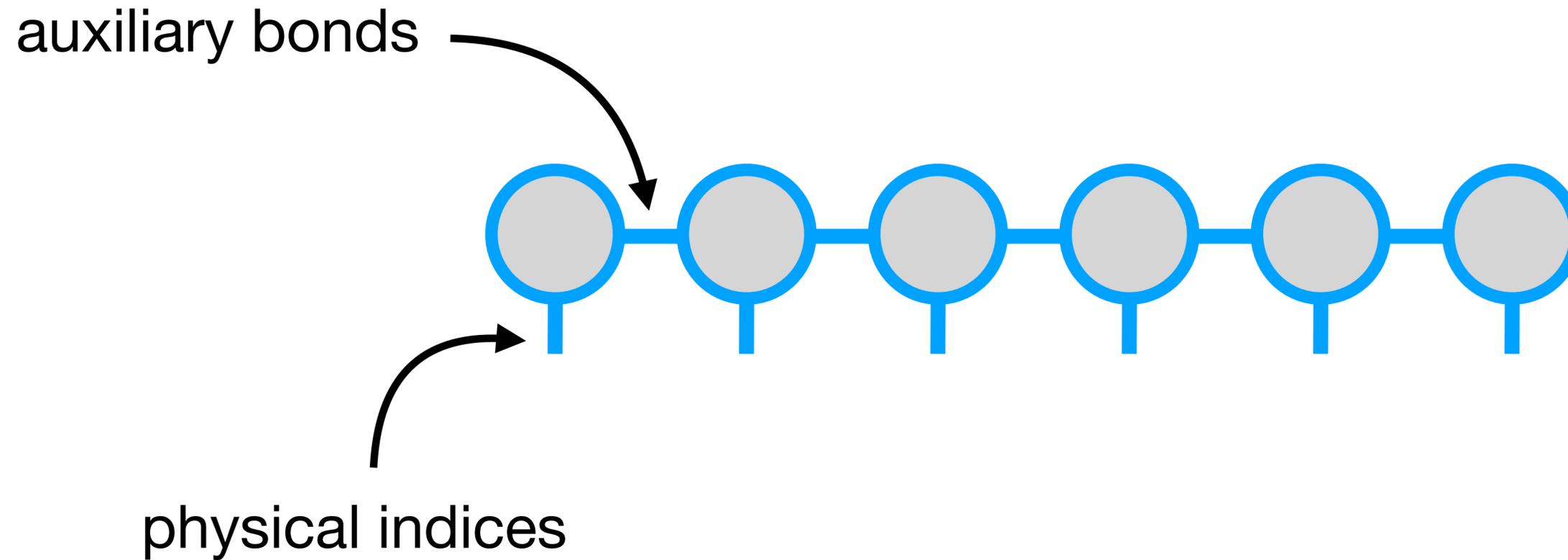
unreasonable

Quantum many body problems from the viewpoint of tensor network methods



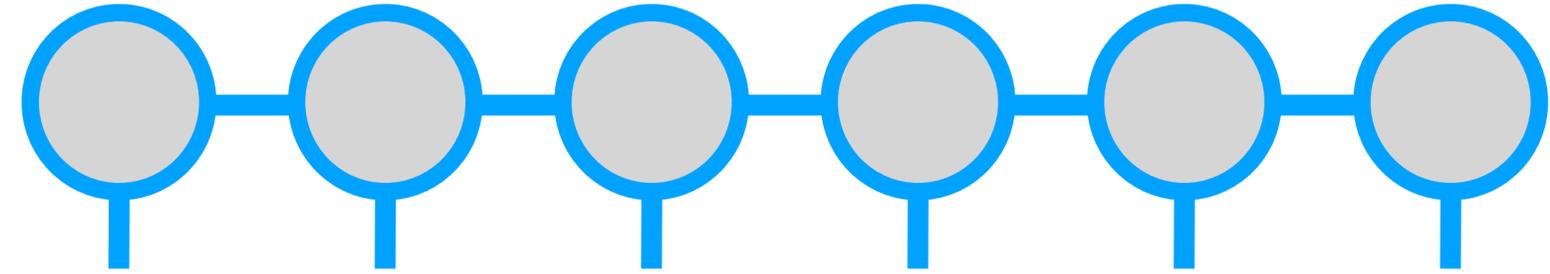
Tensor networks

are an efficient way to compress the wavefunction.



There is a fundamental relation between the ***size of the auxiliary bonds (compression)*** and the **amount of correlations captured**.

Why does it work?

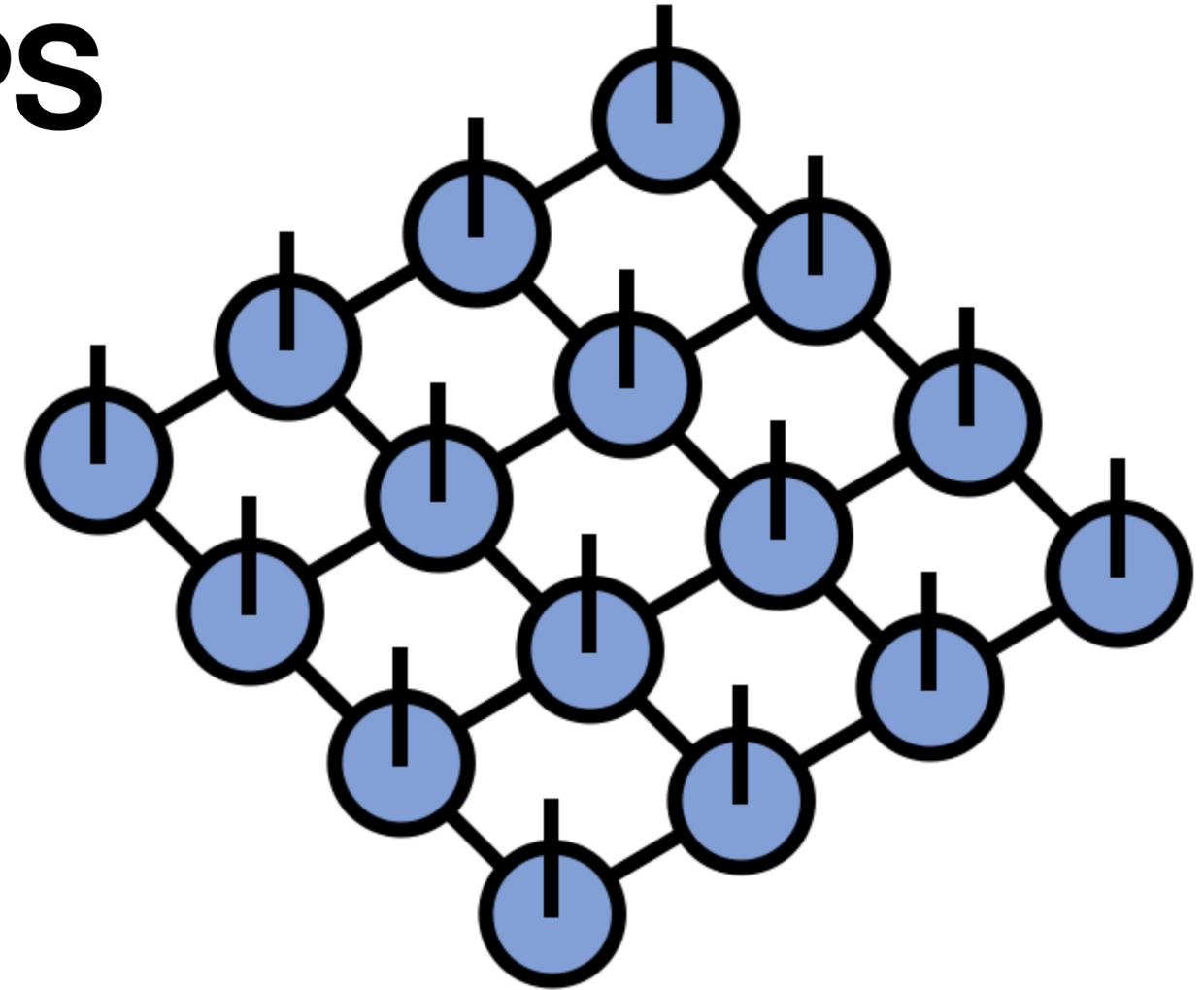


- (1) Ground states of systems with short range interactions have dominantly short-range correlations.
- (2) Efficient algorithms for ground state search and time evolution.
- (3) Mimics the connectivity of the system being represented.

Tensor networks in 2D: PEPS

(projected entangled pair state)

Can represent area-law entangled states.



Tensor networks in 2D: PEPS

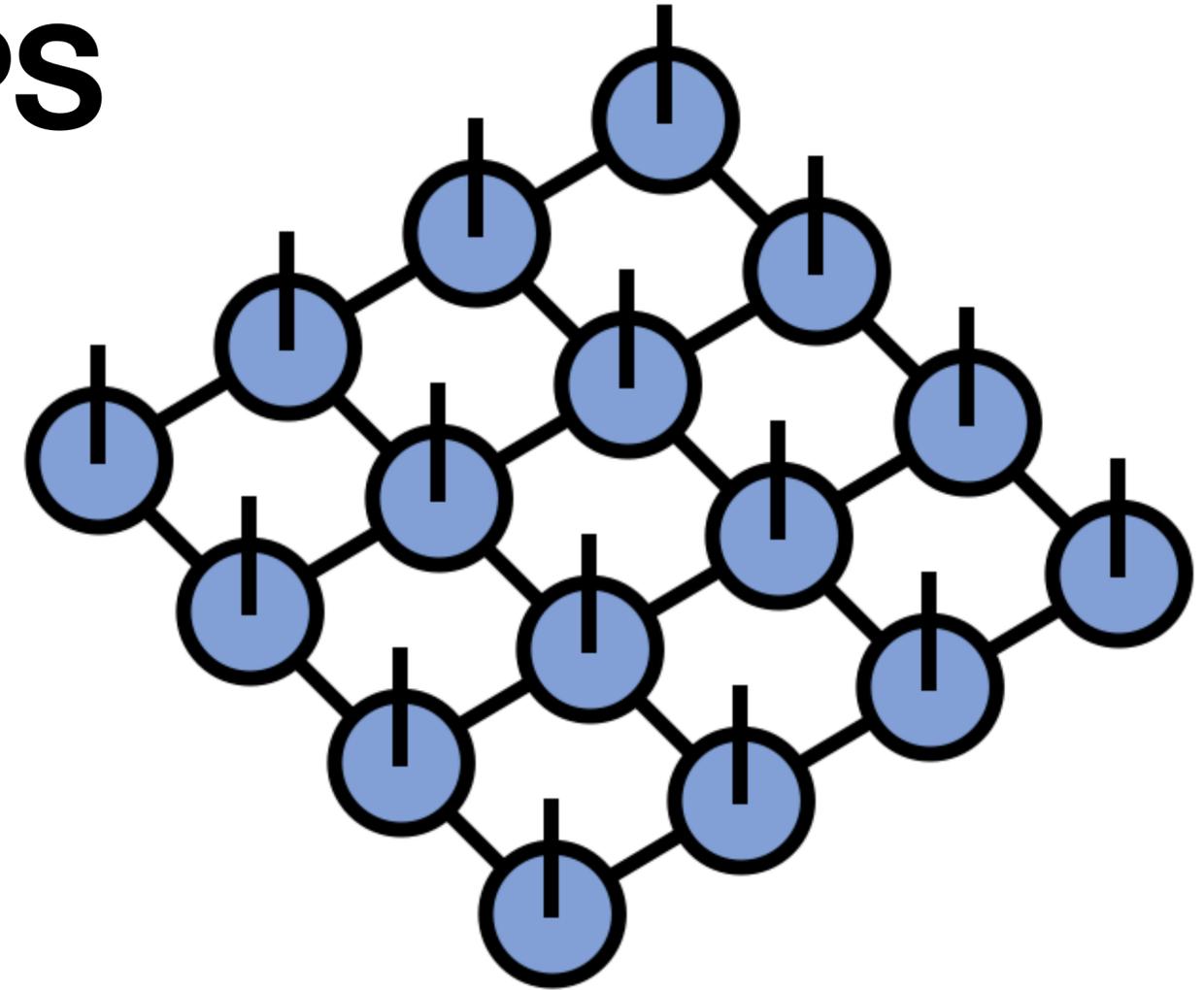
(projected entangled pair state)

Can represent area-law entangled states.

BUT:

Cannot be efficiently contracted. $O(\chi^L)$

Approximate optimization is unfeasible: $O(\chi^{10})$

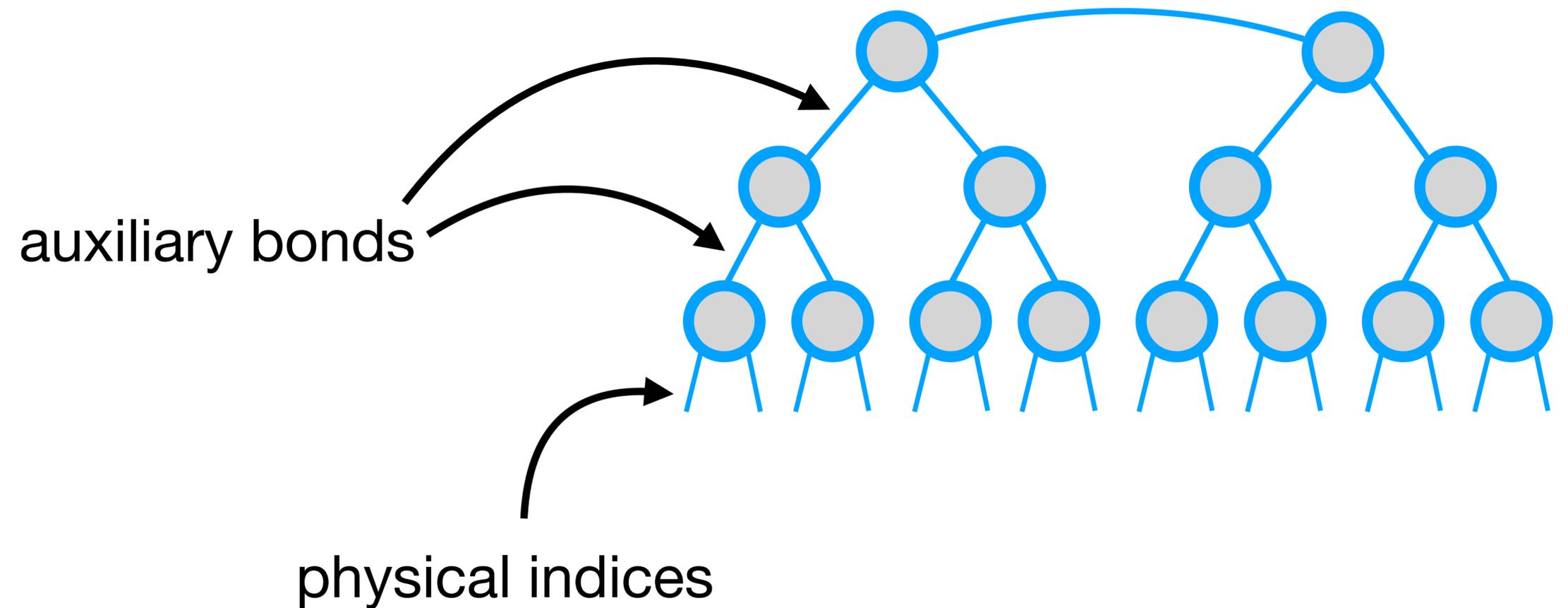


*This is all because PEPS has **loops!***

A compromise: tree tensor networks

A network with *as many links* as possible,
without any *loops*.

(DMRG and TDVP map directly from MPS)

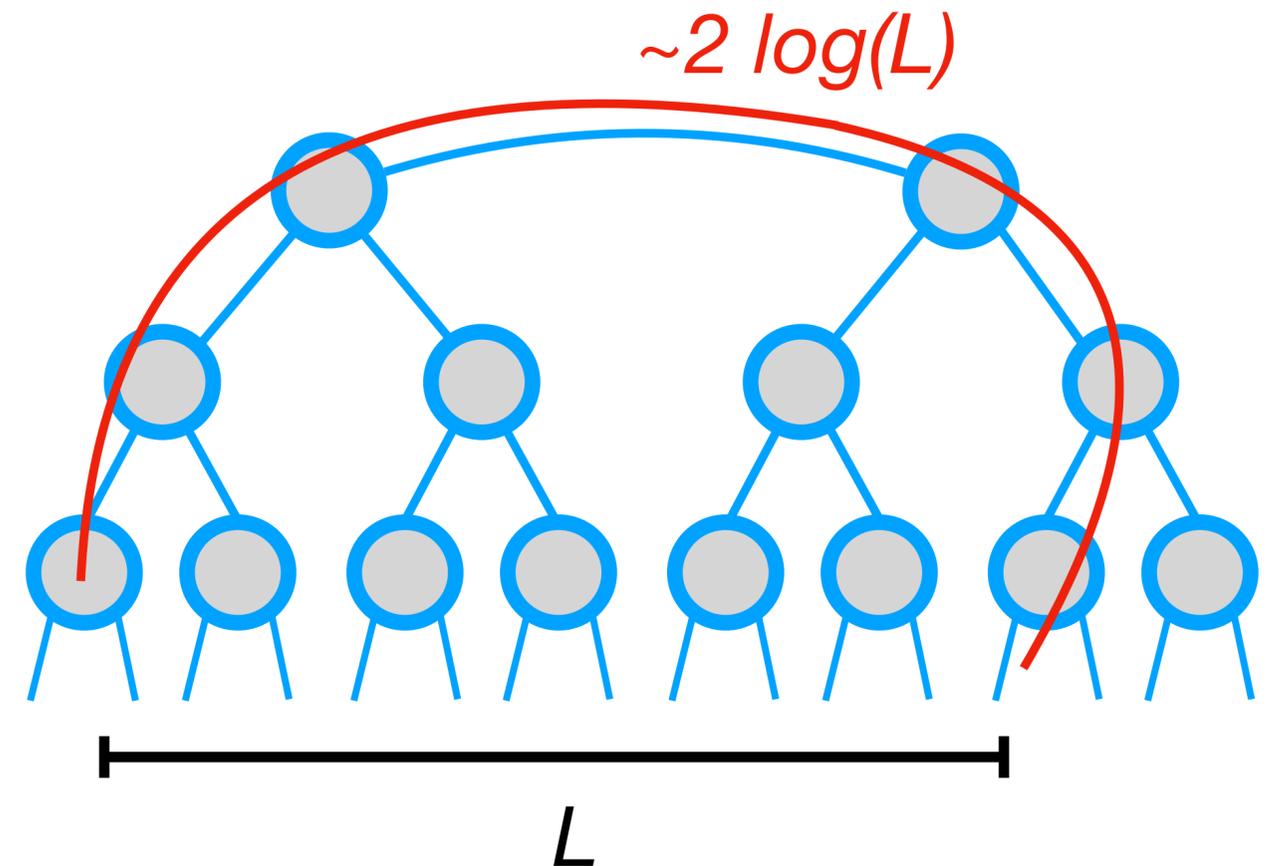


A compromise: tree tensor networks

A network with *as many links* as possible, *without* any *loops*.

(DMRG and TDVP map directly from MPS)

It is better than MPS at encoding long-range interactions:



We expect $\sim \log(N)$ better results than with MPS.

physics

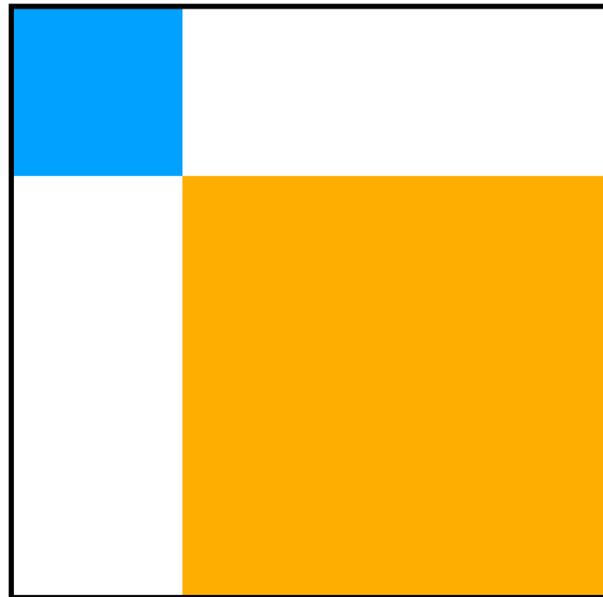
Systems with dynamical constraints

where the dynamics is (*approximately*) constrained to a subspace of the Hilbert space

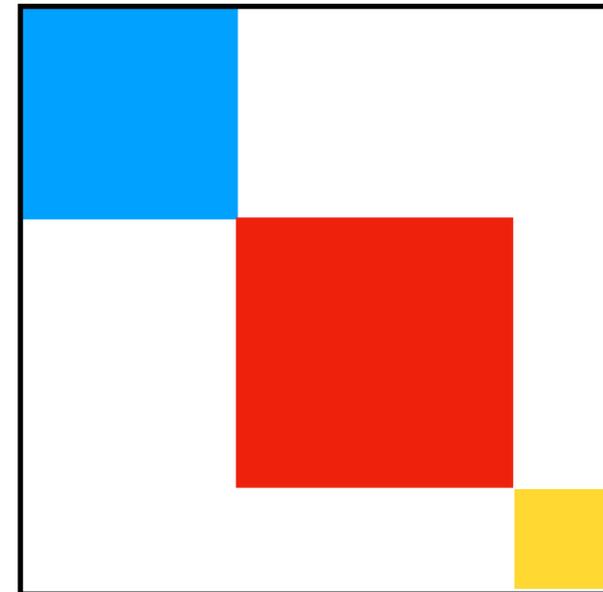
Systems with dynamical constraints

where the dynamics is (*approximately*) constrained to a subspace of the Hilbert space

quantum scars



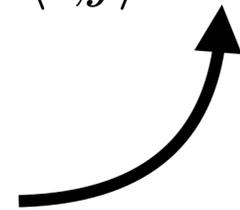
emergent dynamical constraints



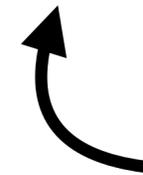
2D quantum Ising on a square lattice

$$H = -J \sum_{\langle i,j \rangle} Z_i Z_j - g \sum_i X_i \quad \text{with } g < g_c = 3J$$

energy cost of
domain walls



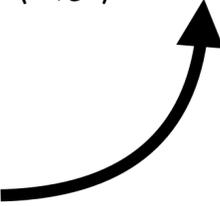
spin flipping



2D Ising on a square lattice

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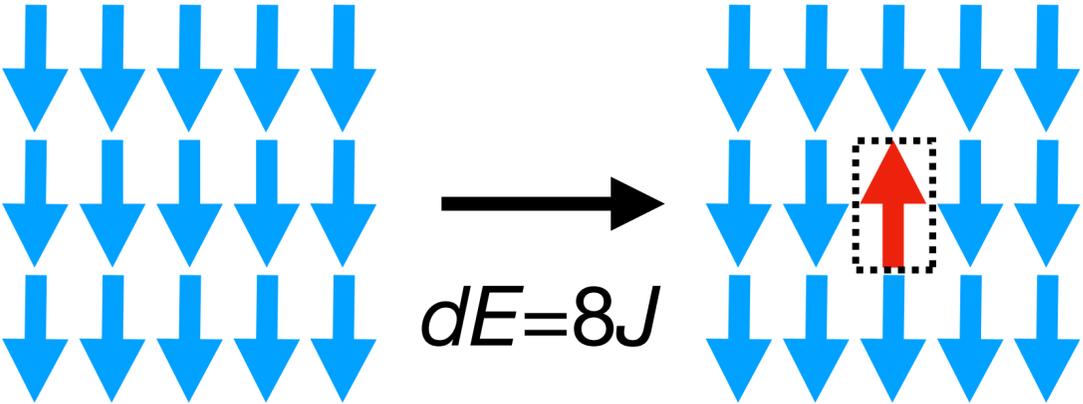
energy cost of domain walls



spin flipping



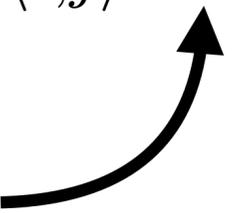
Creating domains is energetically expensive:



2D Ising on a square lattice

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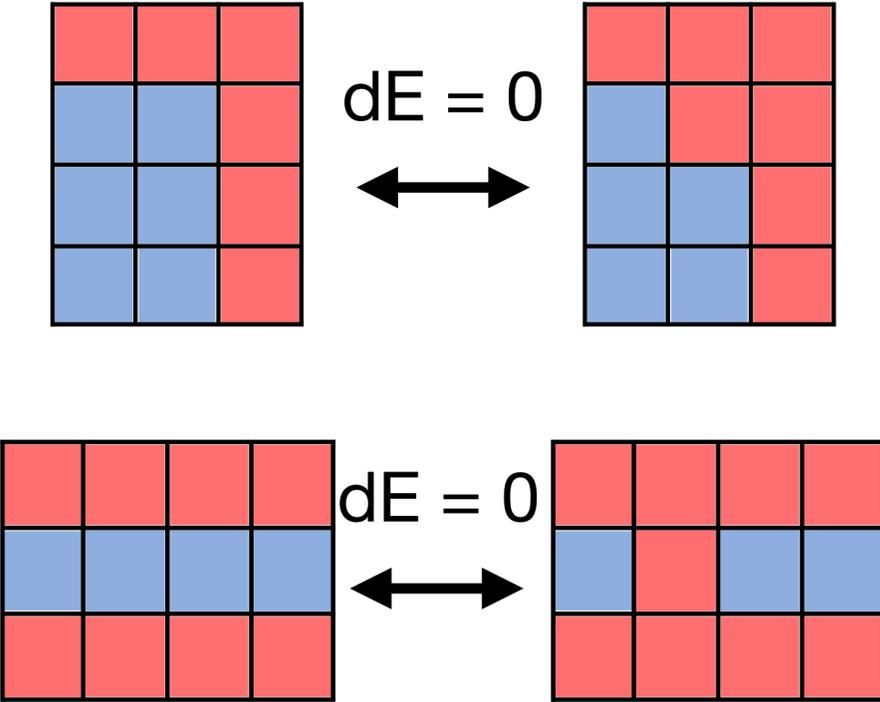
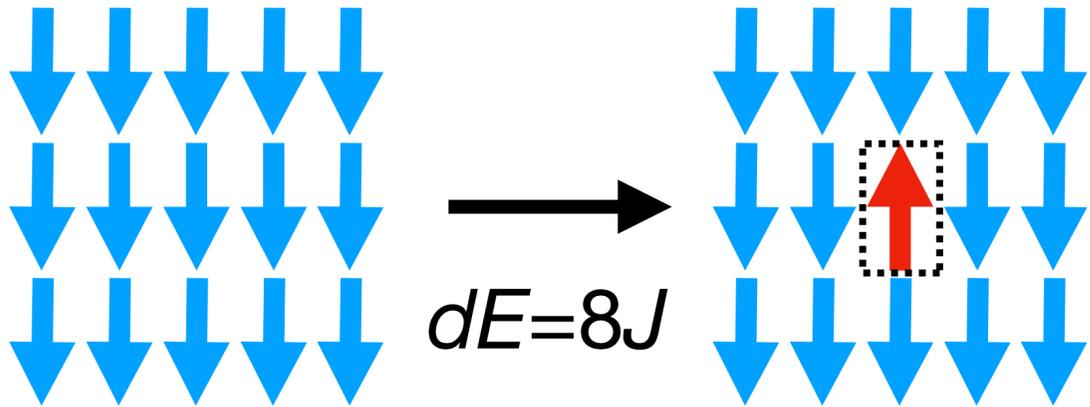


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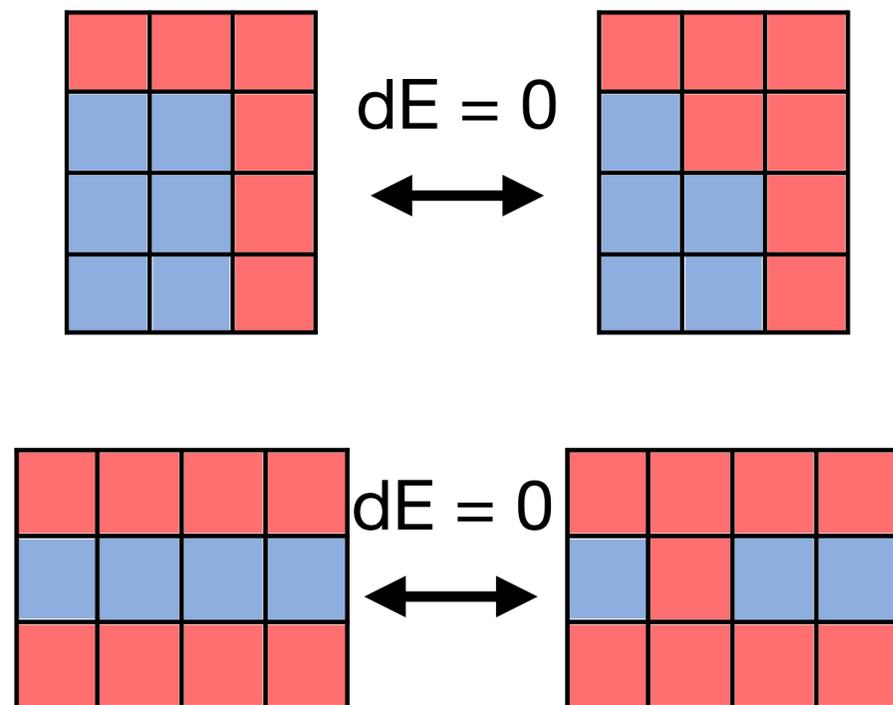
But, there are **resonant** processes:

Creating domains is energetically expensive:

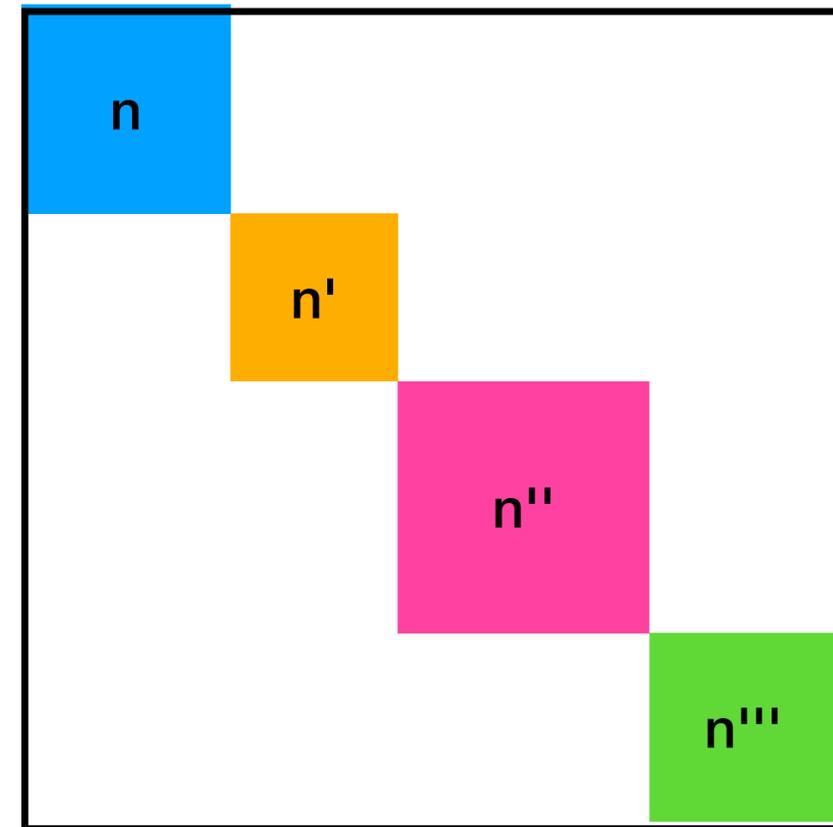


Resonant processes conserve the length of the domain wall

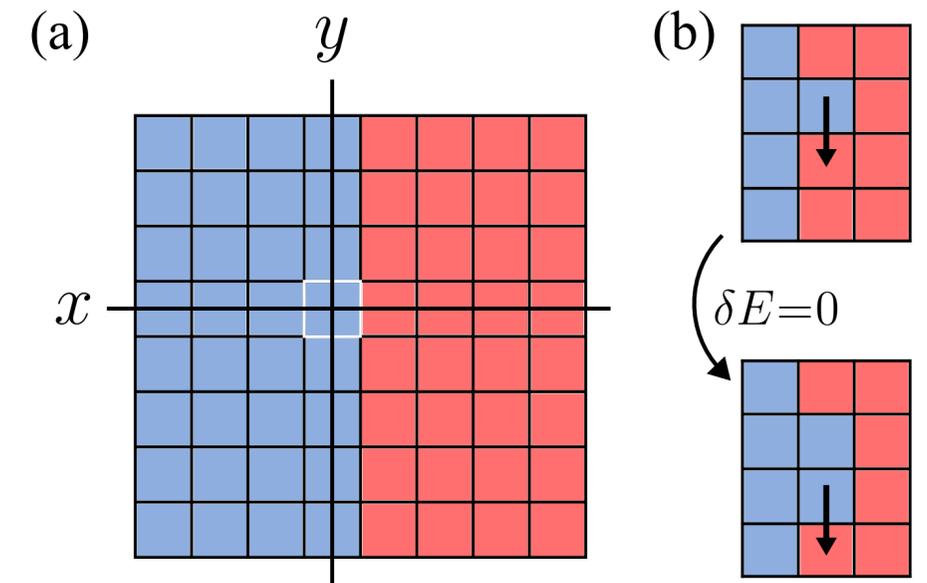
But, there are **resonant** processes:



Fragmentation into sectors with same **total domain wall length** n .



Dynamics from a domain wall?



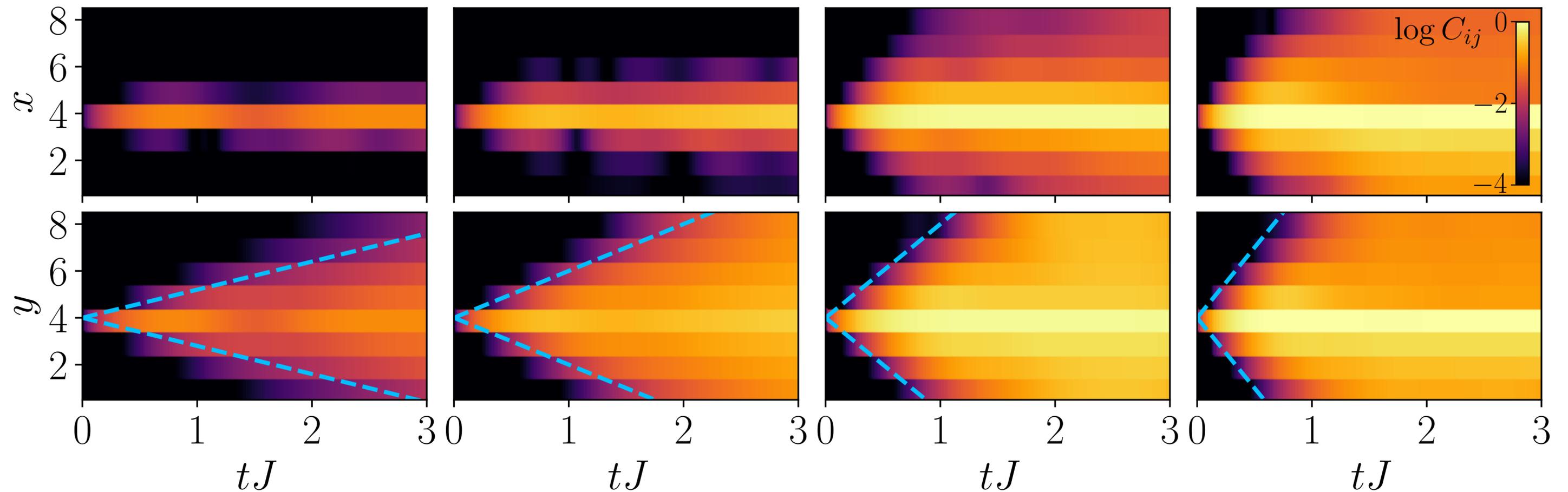
$$C_{ij} = \langle Z_i Z_j \rangle - \langle Z_i \rangle \langle Z_j \rangle$$

(d) $g/J = 0.3$

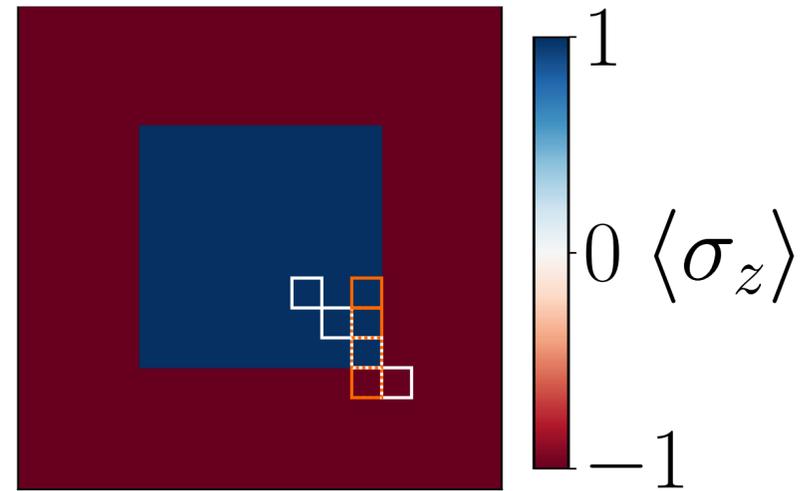
$g/J = 0.5$

$g/J = 1.0$

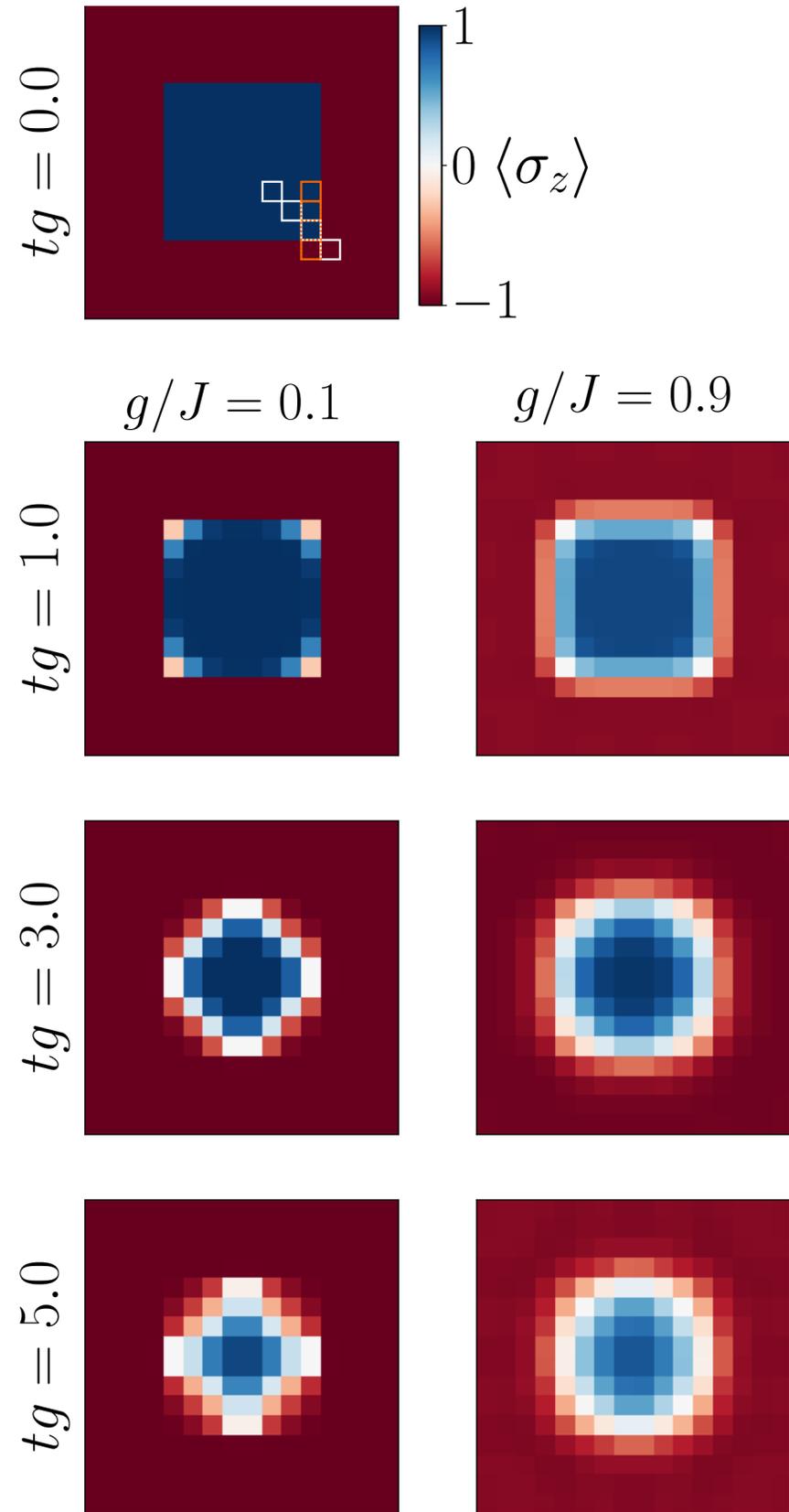
$g/J = 1.5$



Melting of a square



Melting of a square

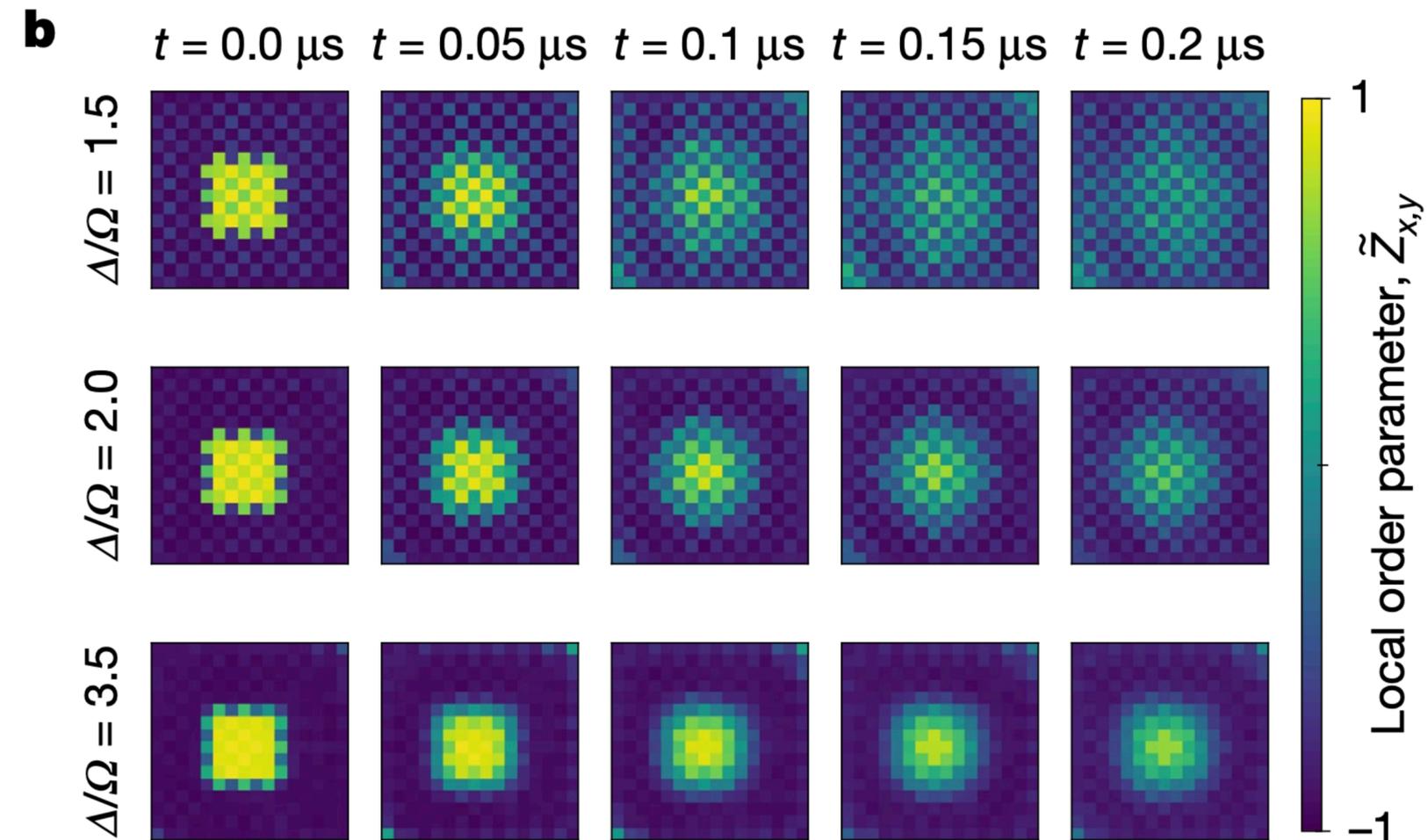


So what?

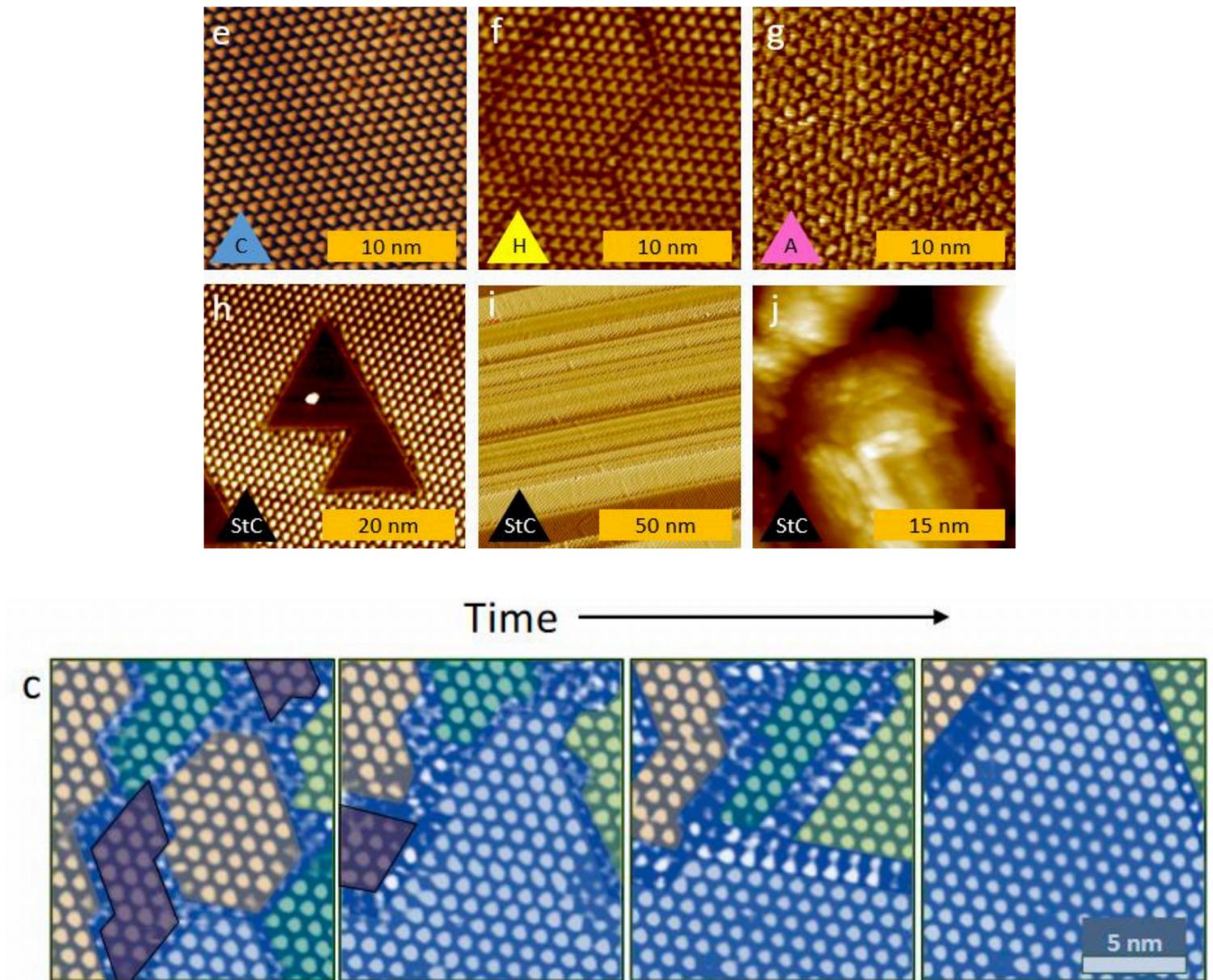
Quantum coarsening and collective dynamics on a programmable simulator

[Tom Manovitz](#), [Sophie H. Li](#), [Sepehr Ebadi](#), [Rhine Samajdar](#), [Alexandra A. Geim](#), [Simon J. Evered](#), [Dolev Bluvstein](#), [Hengyun Zhou](#), [Nazli Ugur Koyluoglu](#), [Johannes Feldmeier](#), [Pavel E. Dolgirev](#), [Nishad Maskara](#), [Marcin Kalinowski](#), [Subir Sachdev](#), [David A. Huse](#), [Markus Greiner](#), [Vladan Vuletić](#) & [Mikhail D. Lukin](#) ✉

Nature **638**, 86–92 (2025) | [Cite this article](#)



Ravnik et al. A time-domain phase diagram of metastable states in a charge ordered quantum material, *Nat. Commun.* 12 2323, 2021



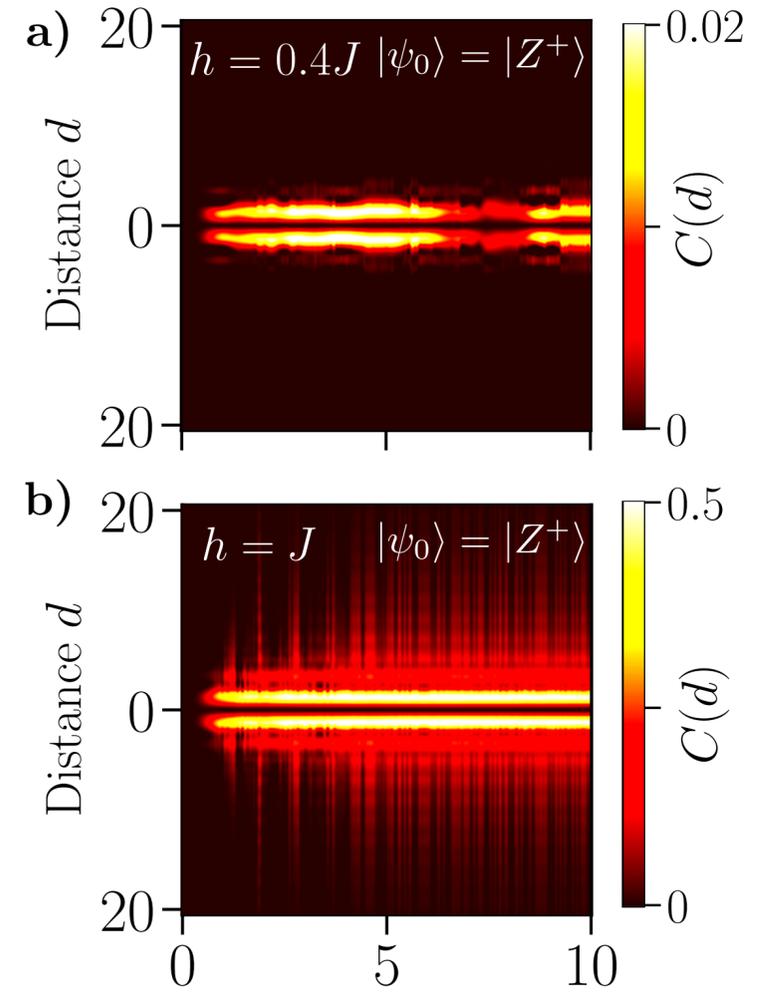
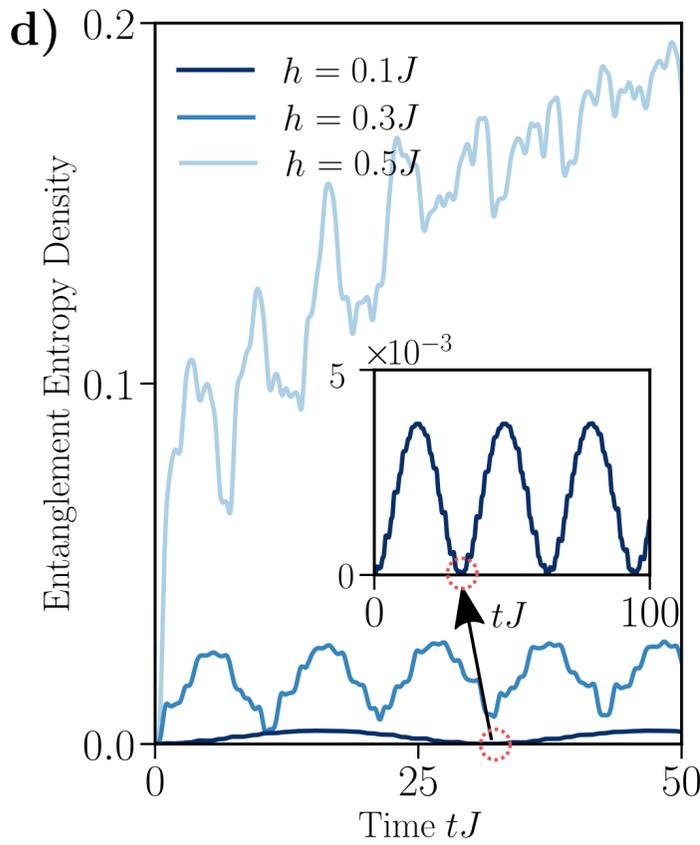
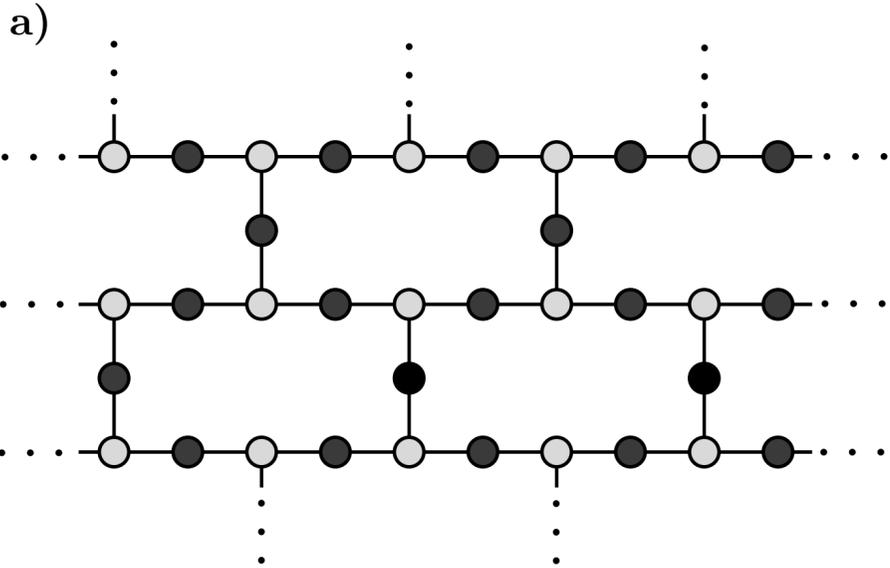
Fun fact

Article

Evidence for the utility of quantum computing before fault tolerance

<https://doi.org/10.1038/s41586-023-06096-3> Youngseok Kim^{1,6}, Andrew Eddins^{2,6}, Sajant Anand³, Ken Xuan Wei¹, Ewout van den Berg¹, Sami Rosenblatt¹, Hasan Nayfeh¹, Yantao Wu^{3,4}, Michael Zaletel^{3,5}, Kristan Temme¹ & Abhinav Kandala¹

Received: 24 February 2023
Accepted: 18 April 2023



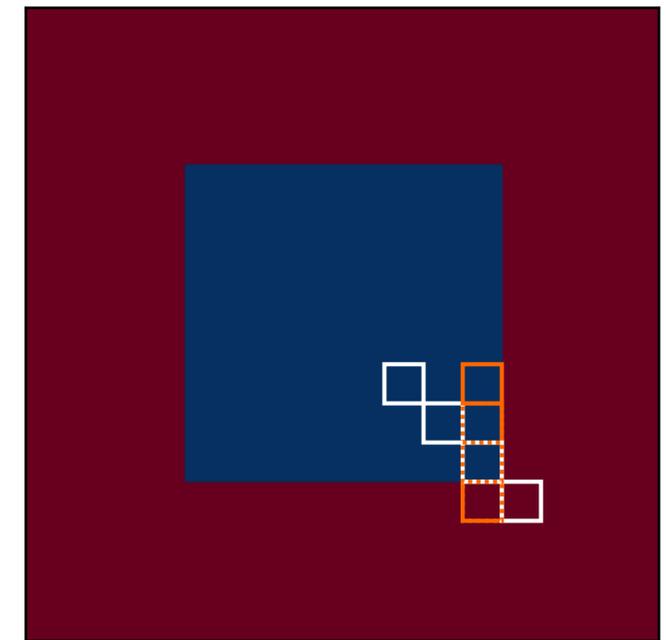
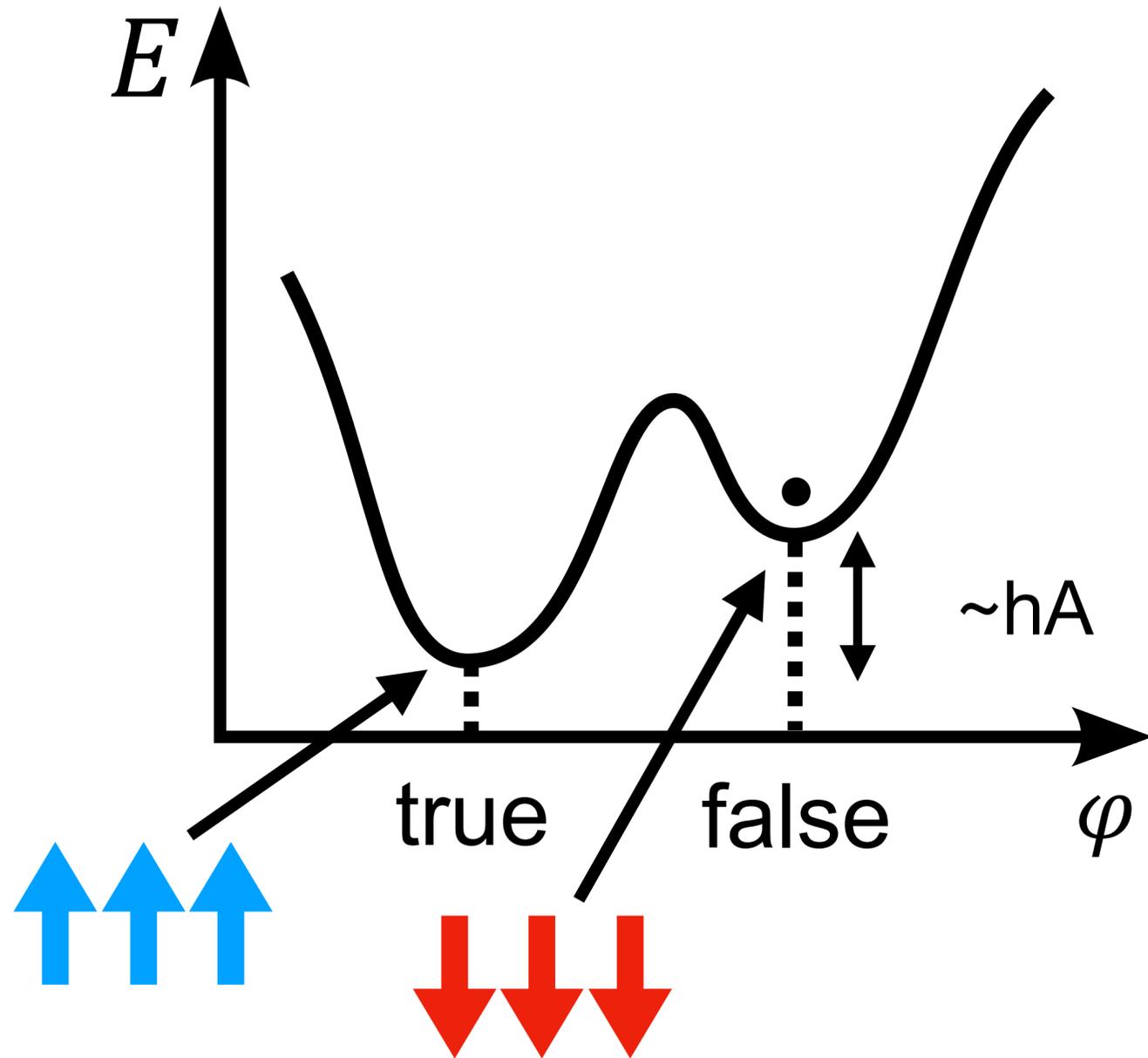
Efficient Tensor Network Simulation of IBM's Eagle Kicked Ising Experiment

Joseph Tindall^{1,*}, Matthew Fishman¹, E. Miles Stoudenmire¹ and Dries Sels^{1,2}
¹Center for Computational Quantum Physics, Flatiron Institute, New York, New York 10010, USA
²Center for Quantum Phenomena, Department of Physics, New York University, 726 Broadway, New York, New York 10003, USA

VS. Confinement in the Transverse Field Ising model on the Heavy Hex lattice

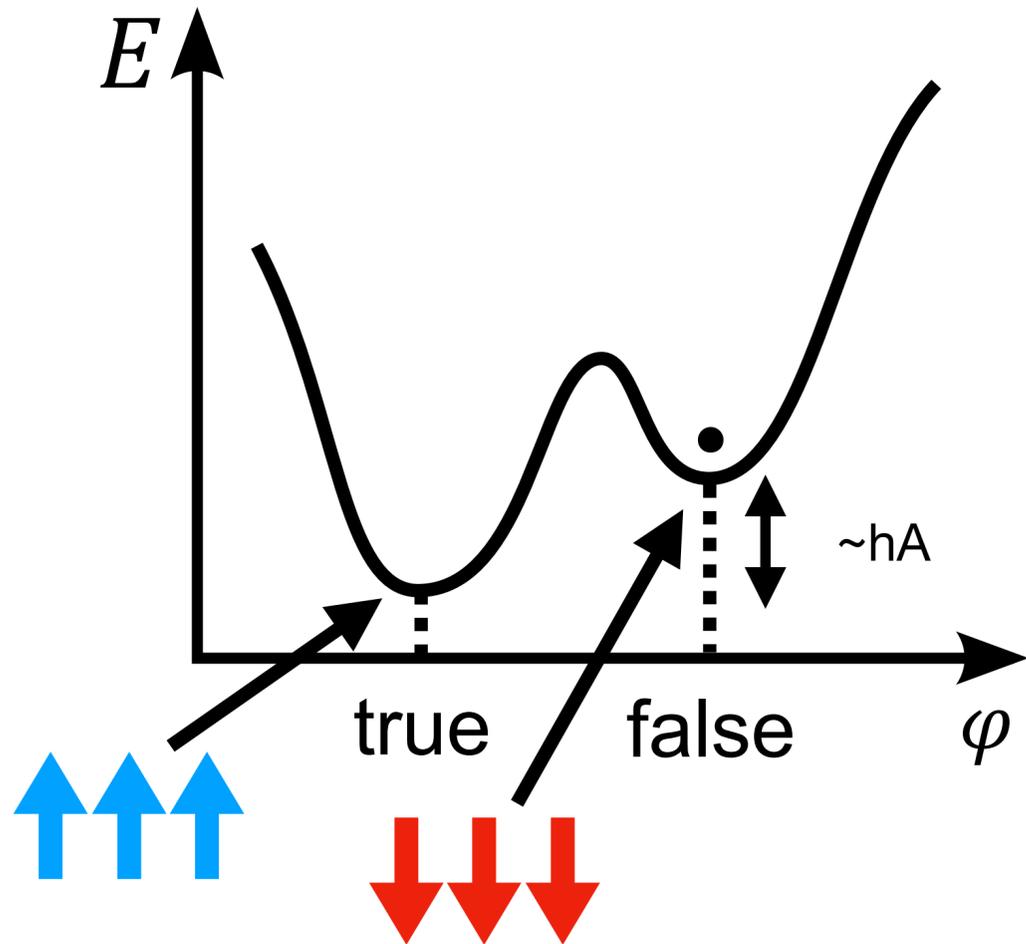
Joseph Tindall¹ and Dries Sels^{1,2}
¹Center for Computational Quantum Physics, Flatiron Institute, New York, New York 10010, USA
²Center for Quantum Phenomena, Department of Physics, New York University, 726 Broadway, New York, NY, 10003, USA
 (Dated: March 6, 2024)

False vacuum decay

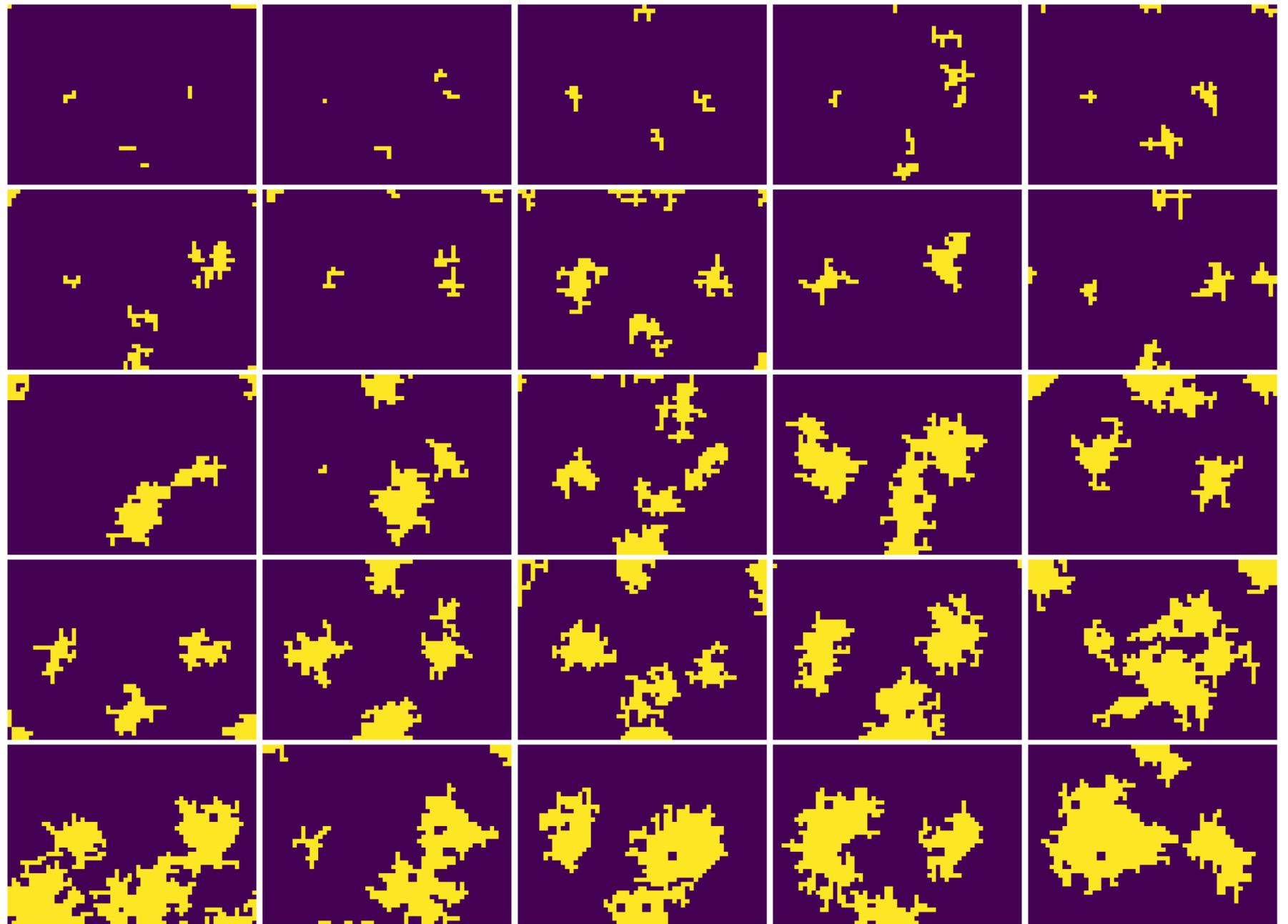


+ longitudinal field

Outlook2: false vacuum decay & metastable states



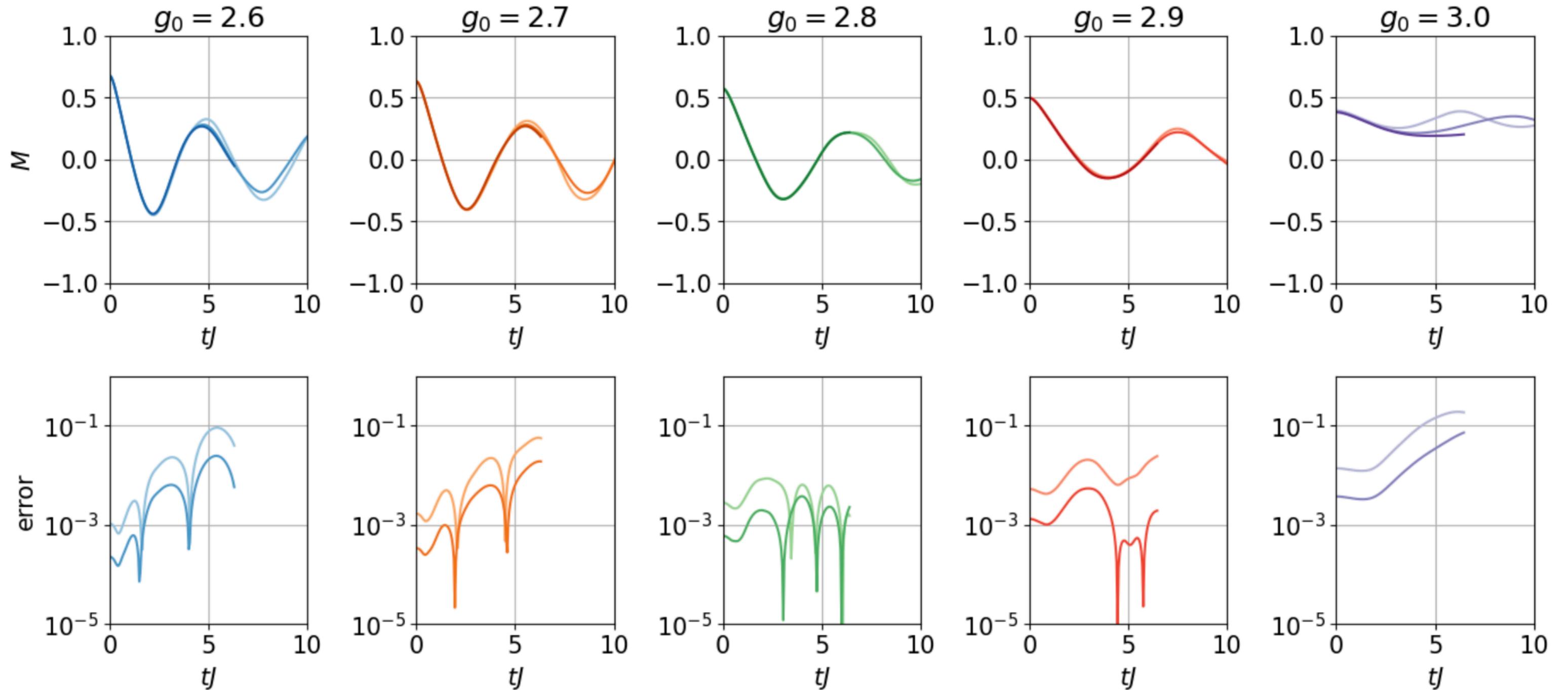
False vacuum decay @dWave? Yes!(!?)



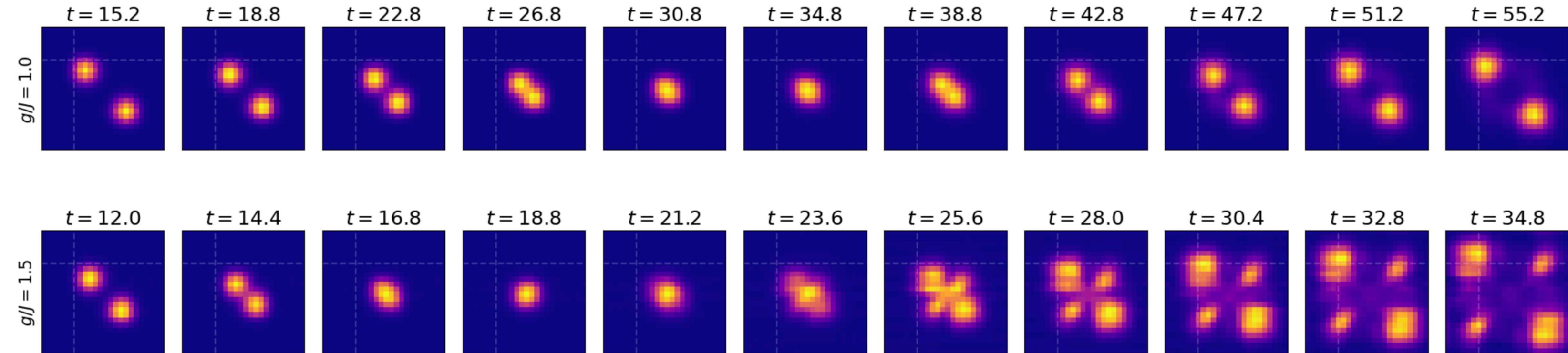
Jaka Vodeb, Gregor Humar, Marko Ljubotina, Jean-Yves Desaulles, Zlatko Papič

2D Ising at the critical point and adS/CFT

$N = 32 \times 32; \quad \chi = 200, 300$



Scattering in 2D



MPS

TTN

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nearest-neighbour interactions

constrained dynamics

metastable states

locally entangled dynamics

2+ D models

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