Based on **2401.04064** with Joe Conlon, Ed Copeland, Martin Mosny and Filippo Revello

String theory and the first half of the Universe, Part I

Fien Apers

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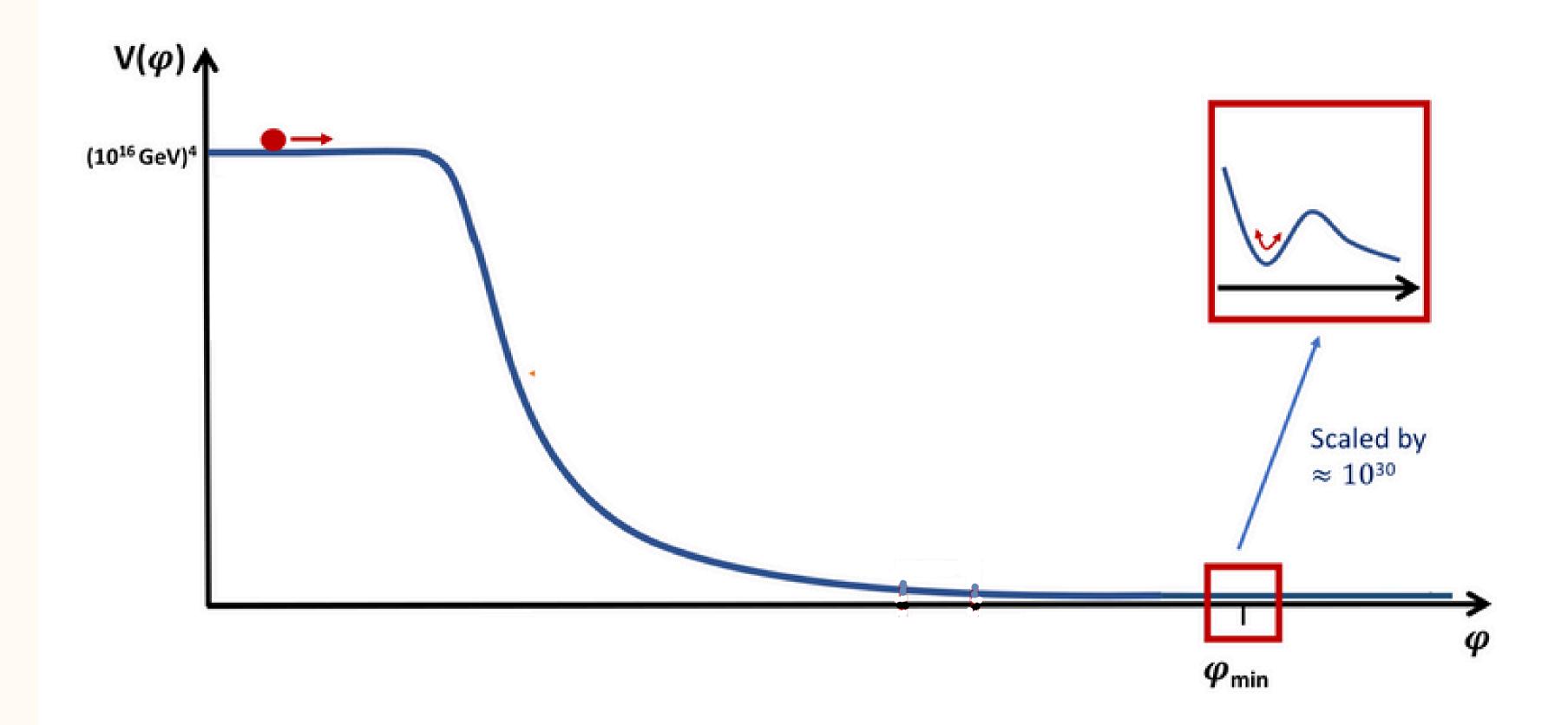
See Filippo Revellos' talk for Part II!

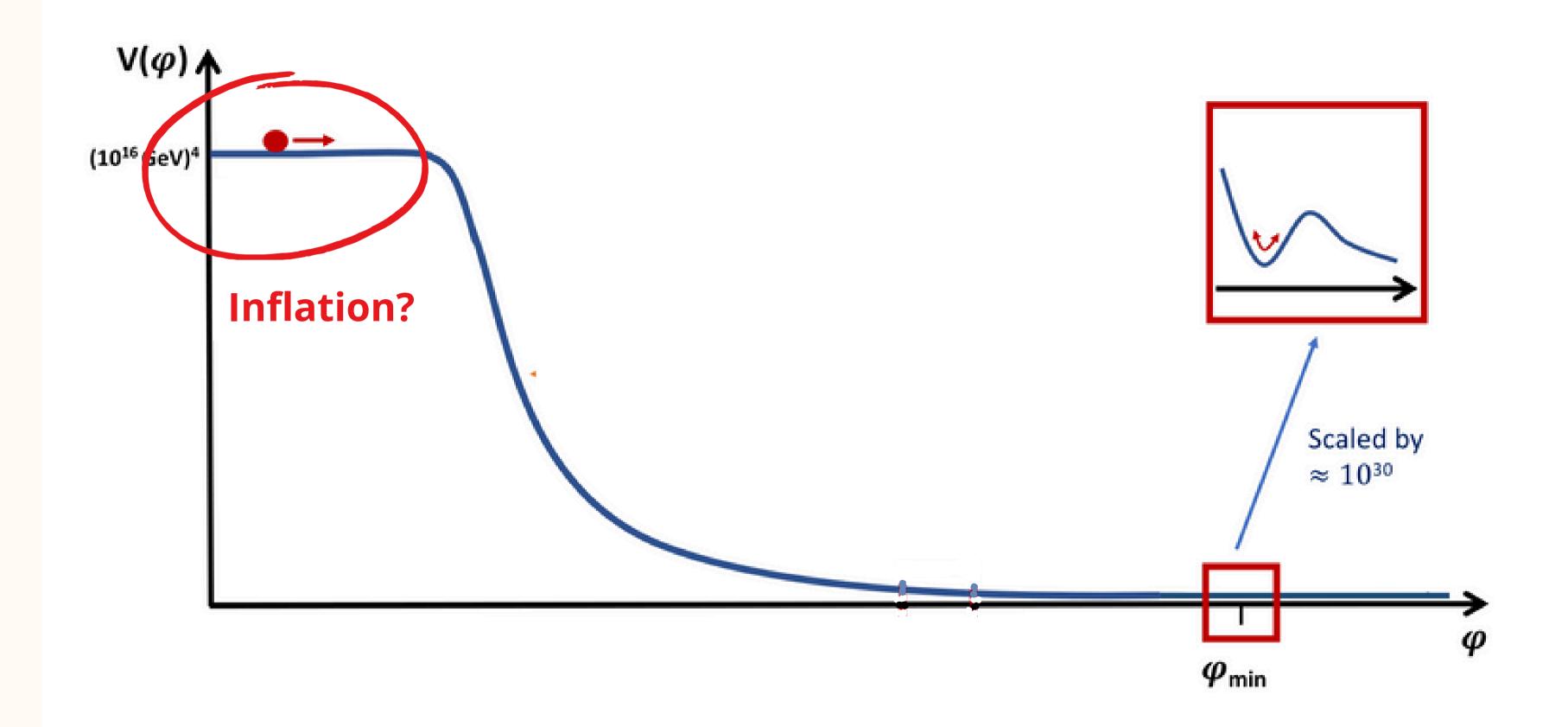
Goal: String Theory predictions for Cosmology

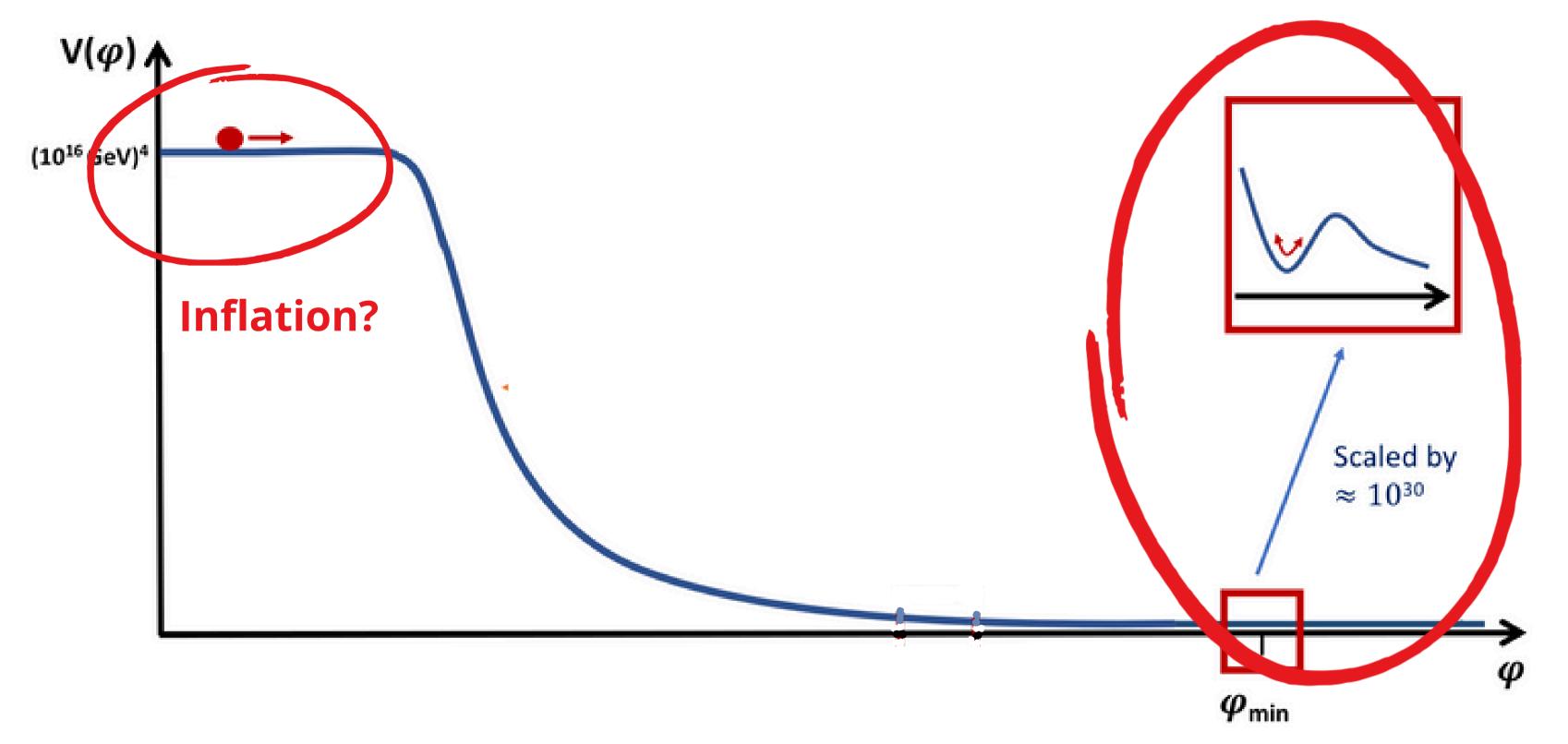
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 - fast rolling moduli
 - steep exponential potentials

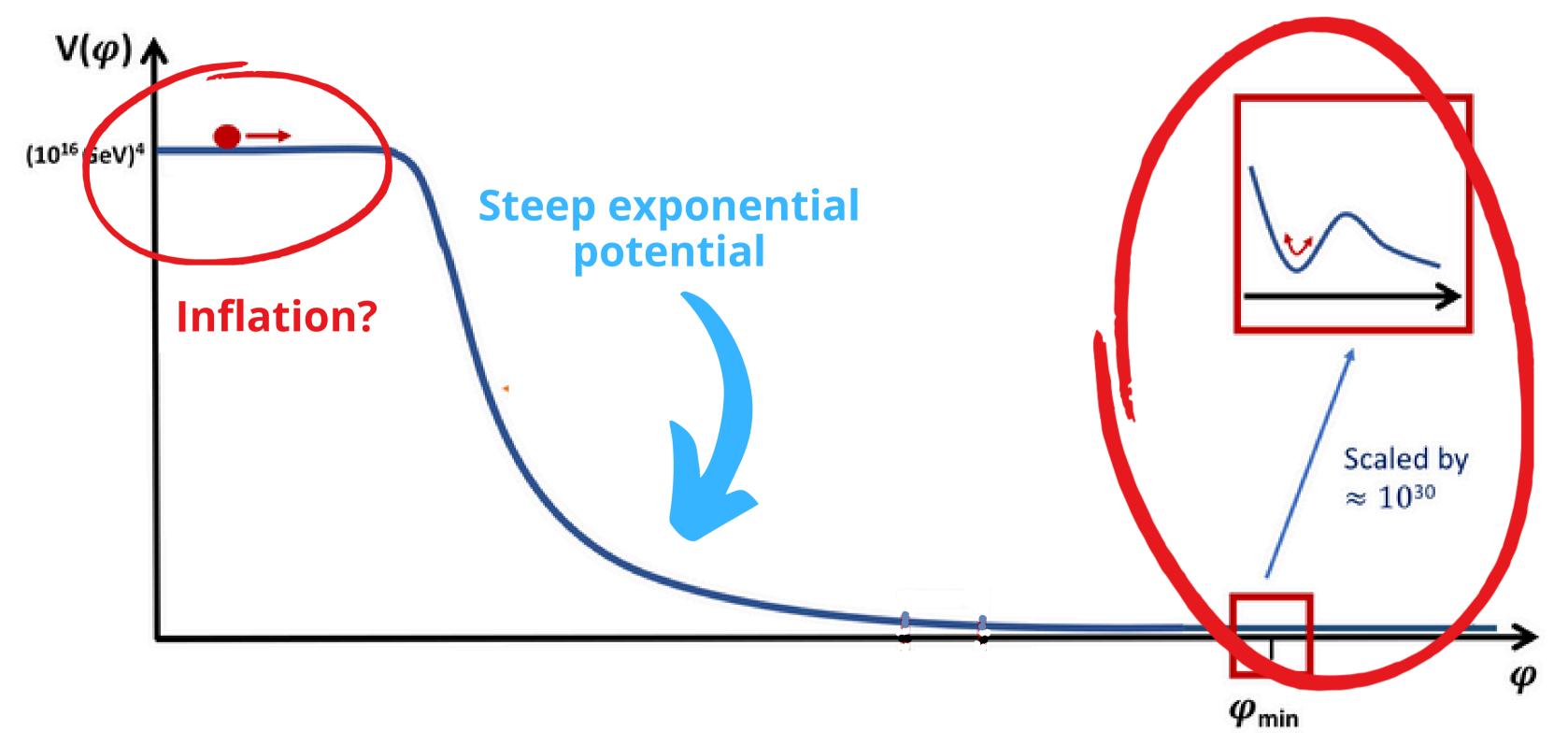
- Goal: String Theory predictions for Cosmology
- Focus on typical dynamics that can easily be obtained from string theory, such as
 - fast rolling moduli
 - steep exponential potentials
- We will not focus on things that hard to get in String Theory such as
 - late-time accelerated expansion
 - inflation







Late time accelerated expansion?



Late time accelerated expansion?



Before radiation domination?

- Inflation ending instantaneously in reheating?
- Less observationally constrained



Radiation domination the

the Hot Big Bang



Matter domination

Structure formation



Dark energy domination Current epoch

t



'Stringy Transition'

Ending in reheating

w = ?



Radiation domination

the Hot Big Bang

w = 1/3

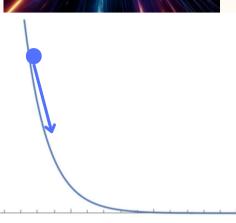
Matter domination

Structure formation w=0



Dark energy domination Current epoch $w \le -0.85$





'Stringy Transition' Energy density: kinetic energy of string moduli



Radiation domination the Hot Big Bang

w = 1/3

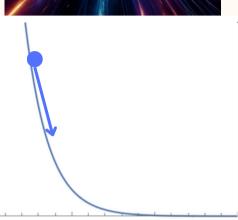
Matter domination

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Dark energy domination Current epoch $w \le -0.85$







'Stringy Transition' Ending in reheating $t_i \approx 10^{-35} s - t_f \approx 10^{-5} s$



Radiation domination the Hot Big Bang $t_i \approx 10^{-5} s - t_f \approx 10^{12} s$

Matter domination

Structure formation $t_i \approx 10^{12} s - t_f \approx 10^{18} s$

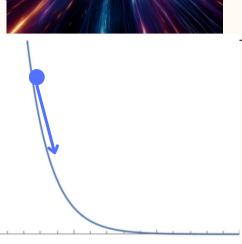


Dark energy domination Current epoch

 $t_i \approx 10^{18} s$



$$t_f \approx 10^{-35} s$$



'Stringy Transition'

Ending in reheating
$$t_i \approx 10^{-35} s - t_f \approx 10^{-5} s$$

`The first half of the Universe' $\frac{t_f}{t_i} \approx 10^{30}$



Radiation domination

the Hot Big Bang
$$t_i pprox 10^{-5} s - t_f pprox 10^{12} s$$

Matter domination

Structure formation $t_i \approx 10^{12} s - t_f \approx 10^{18} s$



Dark energy domination Current epoch

 $t_i \approx 10^{18} s$

Rolling towards the end of the world

Our universe: weak couplings, hierarchies etc.

e.g.
$$\Lambda_{\rm EW} \sim 10^{-16} M_p$$
, $\Lambda_{\rm cc} \sim 10^{-120} M_p^4$, $\theta_{QCD} \sim 10^{-10}$

In string theory, couplings are determined by the values of moduli fields.

• So let's assume we live near an asymptotic boundary of moduli space $\frac{\phi}{M_p} o {
m large}$

• and that this stringy transition epoch brings us there

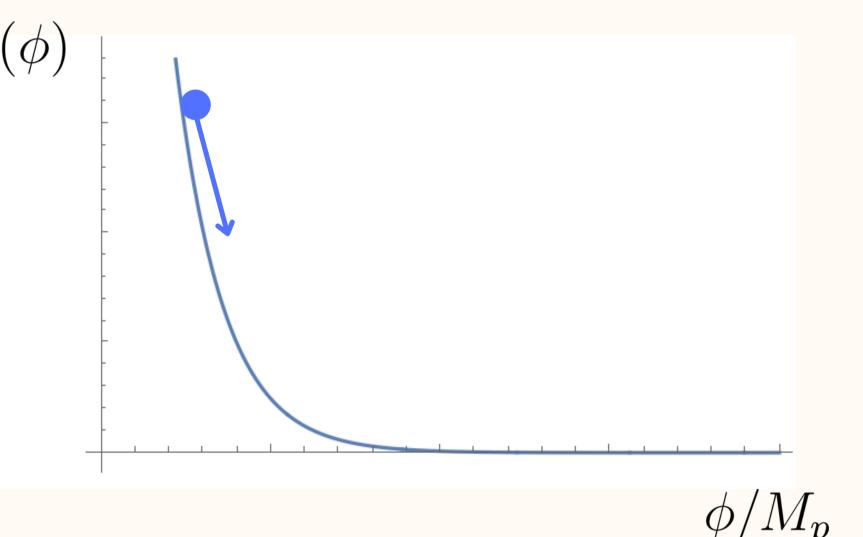
[Copeland, Liddle, Wands 1998]

scalar field

$$V(\phi) = e^{-\lambda \phi} \qquad \lambda \geq \sqrt{6}$$

a small amount of radiation

$$w = P/\rho = 1/3$$



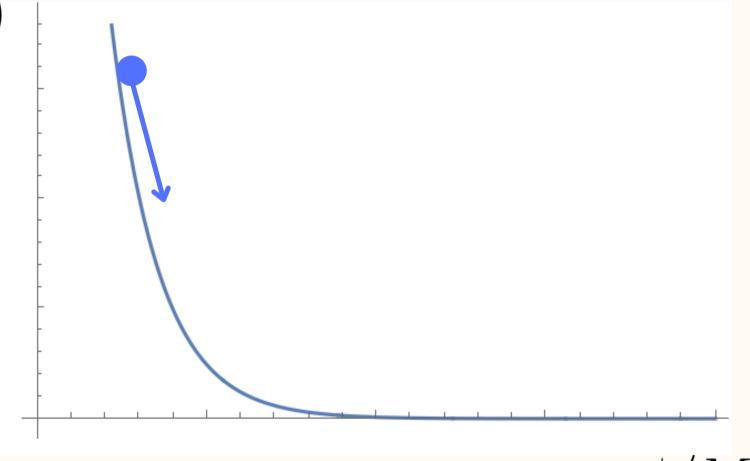
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 ϕ/M_p

Possible sources of radiation: axions (the volume axion), radiation from cosmic strings, SM degrees of freedom

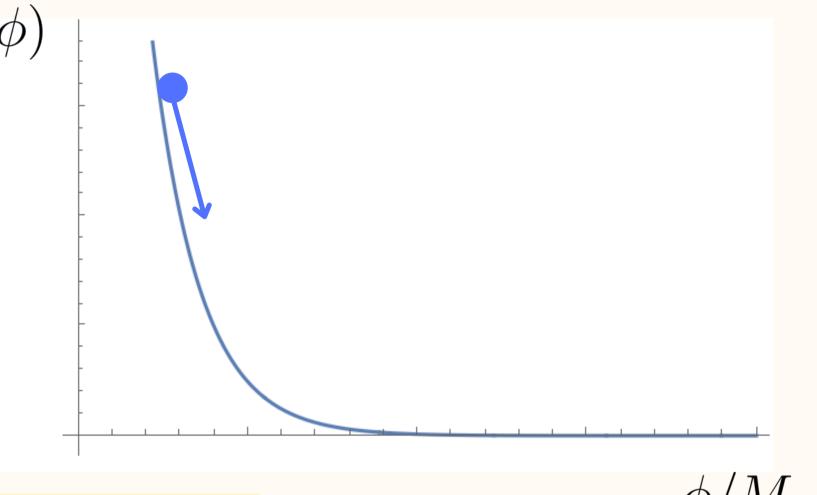
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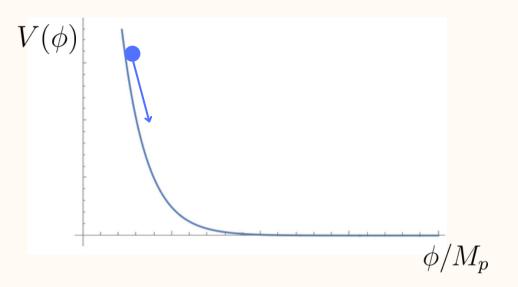
a small amount of radiation

$$w = P/\rho = 1/3$$



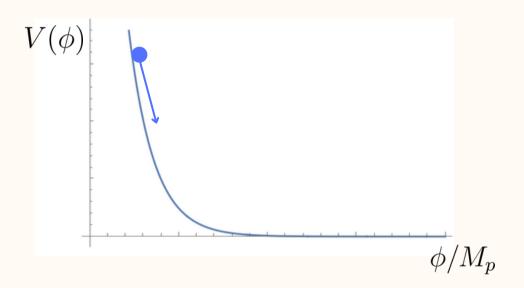
What happens?

Possible sources of radiation: axions (the volume axion), radiation from cosmic strings, SM degrees of freedom



1. Kination

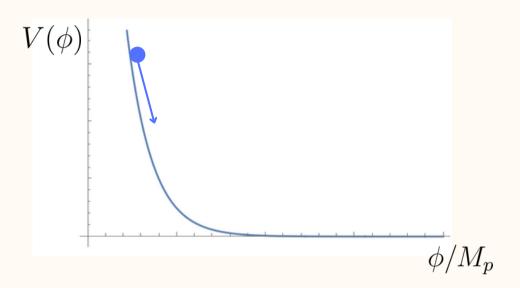
When rolling down a steep exponential potential, the scalar field will gain speed and we go into a regime that is **dominated by the kinetic energy** of the scalar field.



1. Kination

When rolling down a steep exponential potential, the scalar field will gain speed and we go into a regime that is **dominated by the kinetic energy** of the scalar field.

- ullet Energy density dominated by the kinetic energy of a scalar field $\dot{\phi}^2/2\gg V$
- Equation of state $w = \frac{P}{\rho} = \frac{\dot{\phi}^2/2 V}{\dot{\phi}^2/2 + V} = 1$
- Slowest possible expansion $a(t) \sim t^{\frac{2}{3(1+w)}} = t^{\frac{1}{3}}$



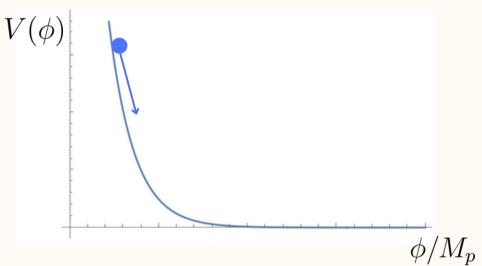
1. Kination

When rolling down a steep exponential potential, the scalar field will gain speed and we go into a regime that is **dominated by the kinetic energy** of the scalar field.

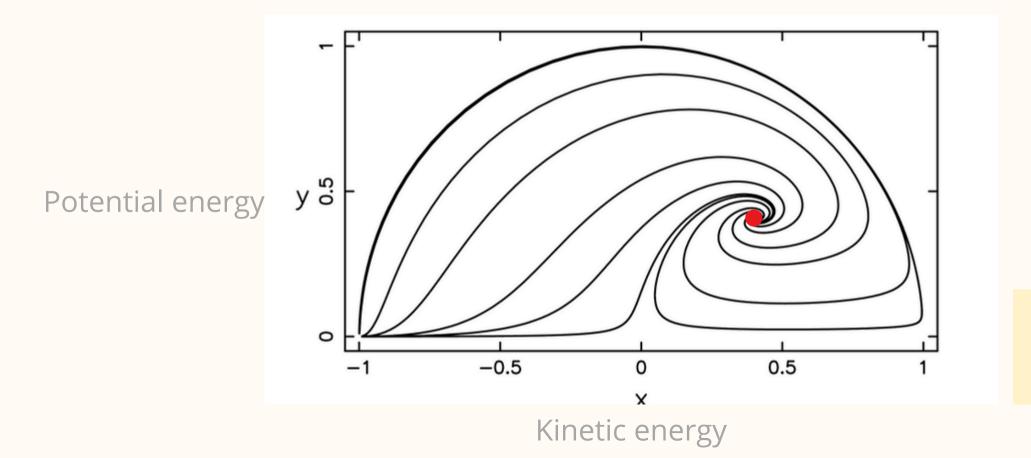
• The scalar equation $\ddot{\phi} + 3H\dot{\phi} = 0$ is solved by

$$\phi = \phi_0 + \sqrt{\frac{2}{3}} \mathbf{M_p} \ln \left(\frac{\mathbf{t}}{\mathbf{t_0}} \right)$$

and the scalar field will likely move multiple Planckian distances.

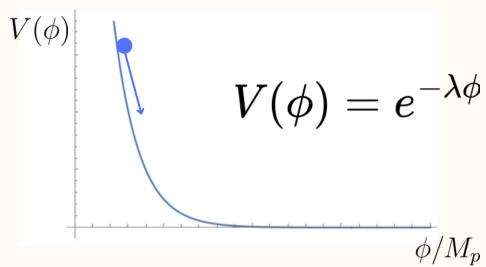


2. Tracker [cosmological scaling solution]

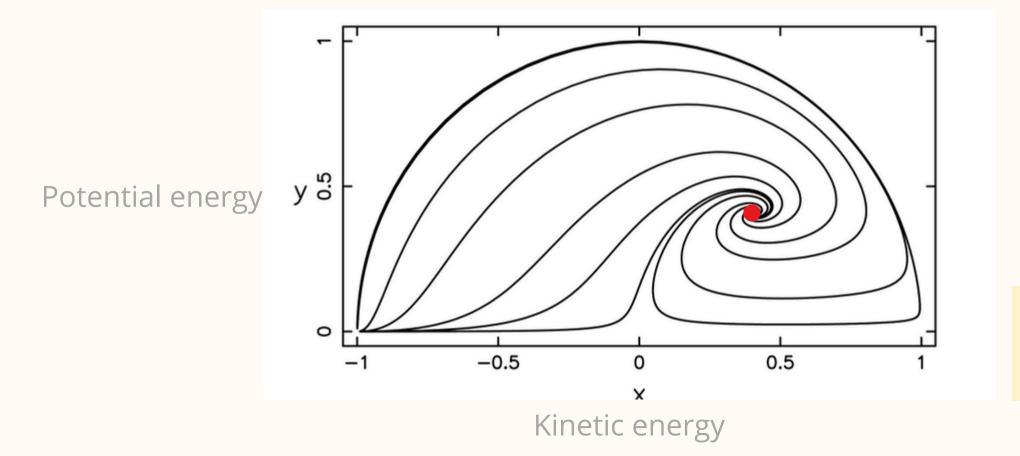


Attractor solution for a scalar field on an exponential potential in the presence of a fluid like radiation.

Scalar field starts tracking/mimicking the radiation

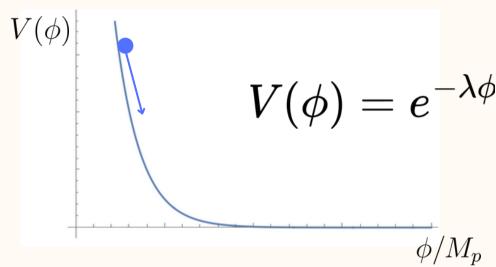


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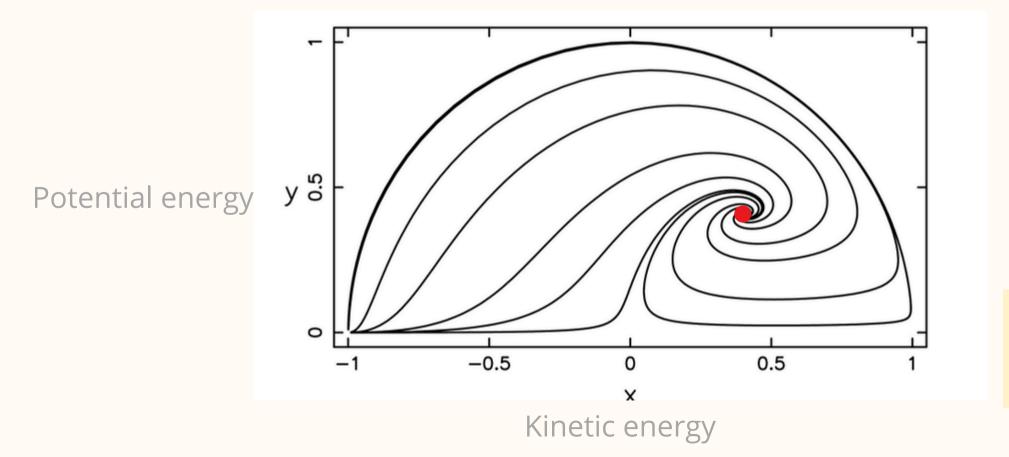


$$\Omega_r = 1 - \frac{4}{\lambda^2}, \quad \Omega_\phi = \frac{4}{\lambda^2}$$

Scalar field starts tracking/**mimicking** the radiation



2. Tracker [cosmological scaling solution]



$$\phi = \phi_0 + \frac{2}{\lambda} M_p \ln\left(\frac{t}{t_0}\right)$$

Again, we can in principle move **multiple Planckian distances**

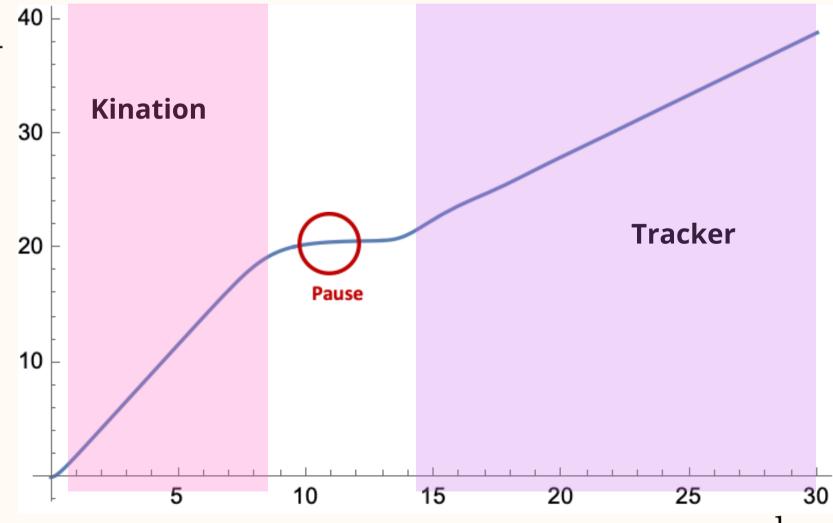
Moving to the asymptotics of moduli space in a cosmological context

1. Kination
$$\phi = \phi_0 + \sqrt{\frac{2}{3}} \mathbf{M_p} \ln \left(\frac{\mathbf{t}}{\mathbf{t_0}} \right)$$

2. Tracker
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 $\frac{\phi}{M_p}$

distances

moving multiple Planckian 20



Moving to the asymptotics of moduli space in a cosmological context

$$\phi = \phi_0 + \sqrt{\frac{2}{3}} \mathbf{M_p} \ln \left(\frac{\mathbf{t}}{\mathbf{t_0}} \right)$$

1. Kination
$$\phi=\phi_0+\sqrt{\frac{2}{3}}\mathrm{M_p\ln\left(\frac{t}{t_0}\right)}$$
 $m_{\mathrm{KK}}(t)\sim t^{-2/3}\sim H(t)^{2/3}$

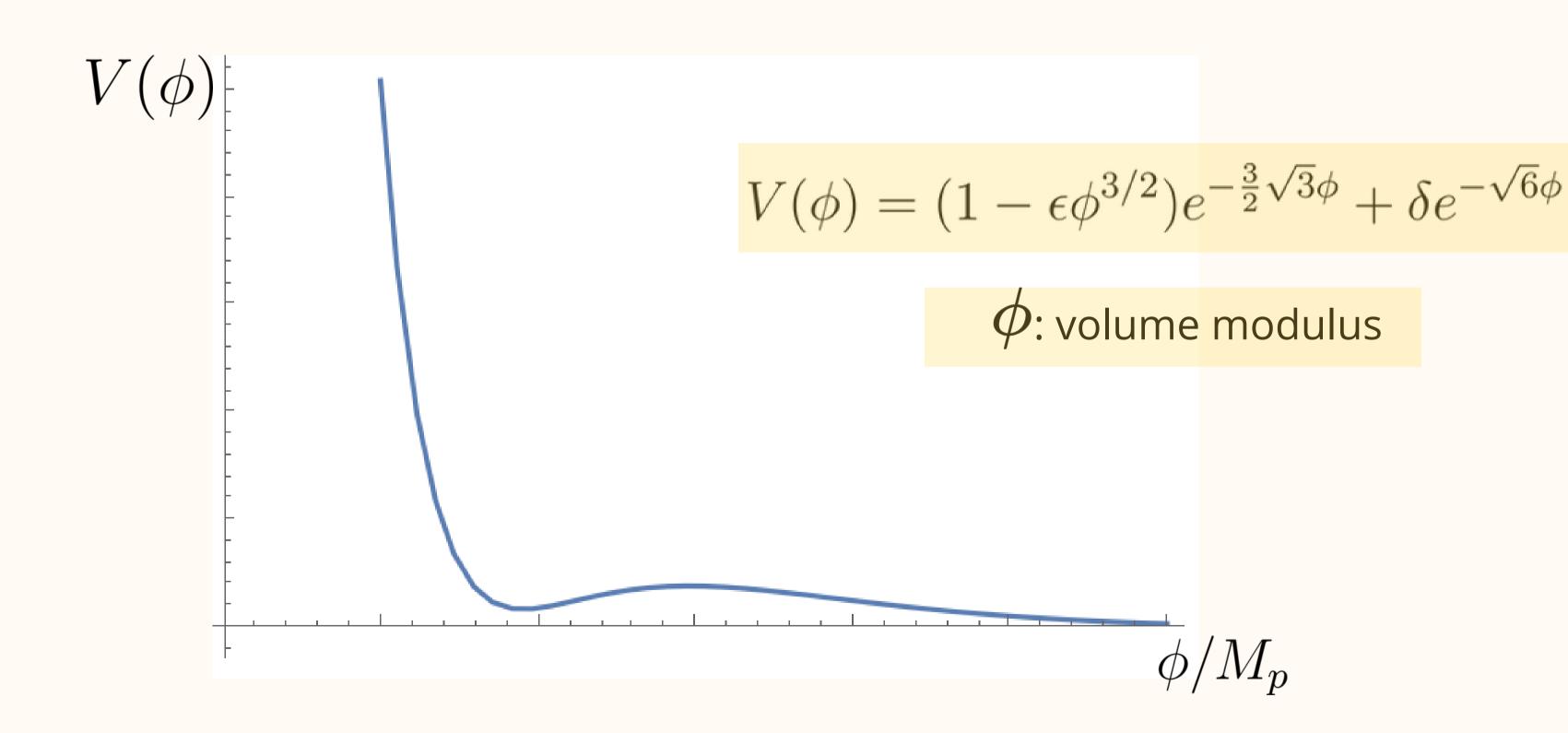
2. Tracker
$$\phi = \phi_0 + \frac{2}{\lambda} M_p \ln \left(\frac{t}{t_0} \right)$$
 $m_{KK}(t) \sim H(t) \sqrt{\frac{2}{3}} \frac{2}{\lambda}$

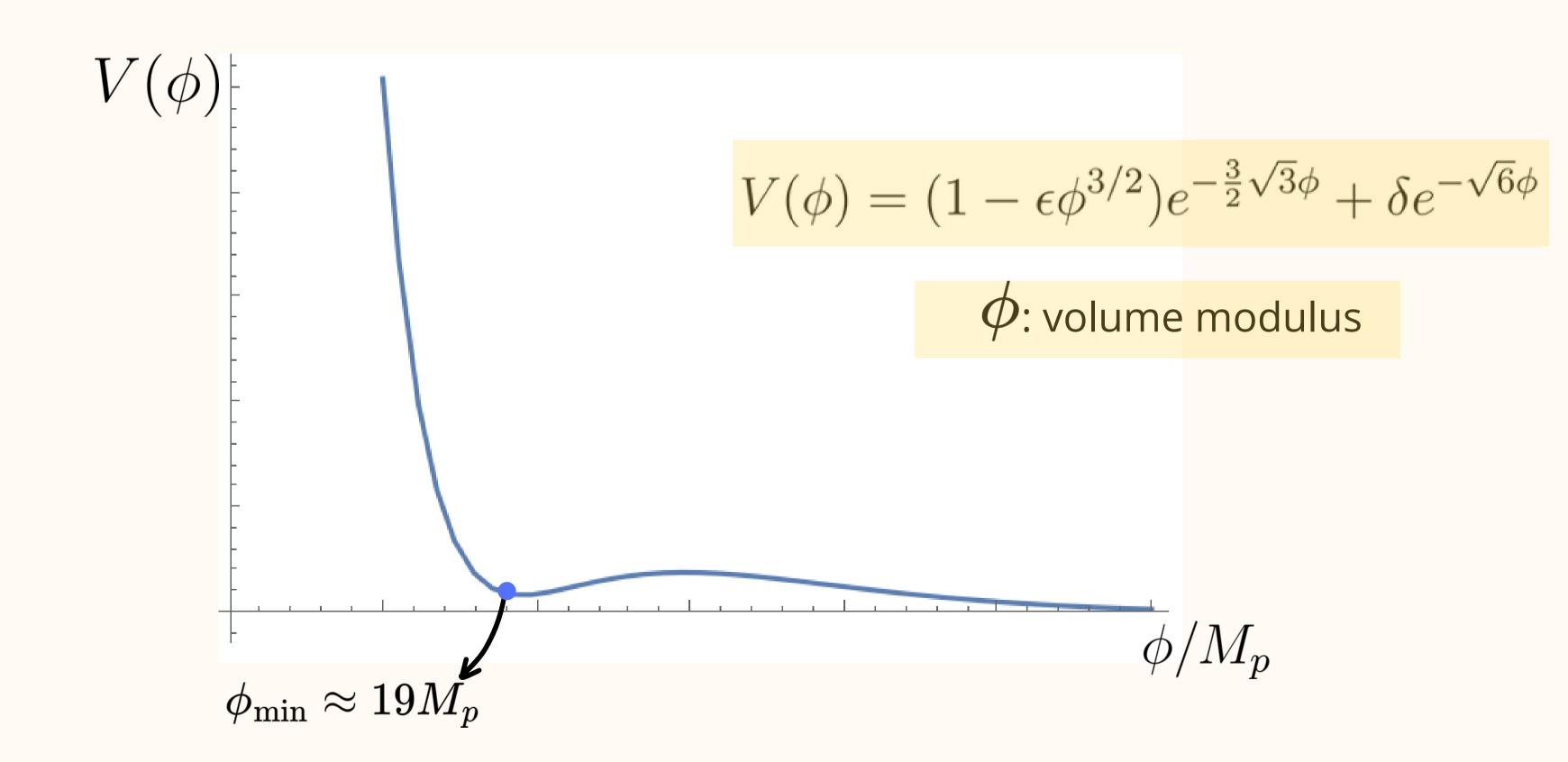
$$m_{KK}(t) \sim H(t)^{\sqrt{\frac{2}{3}}\frac{2}{\lambda}}$$

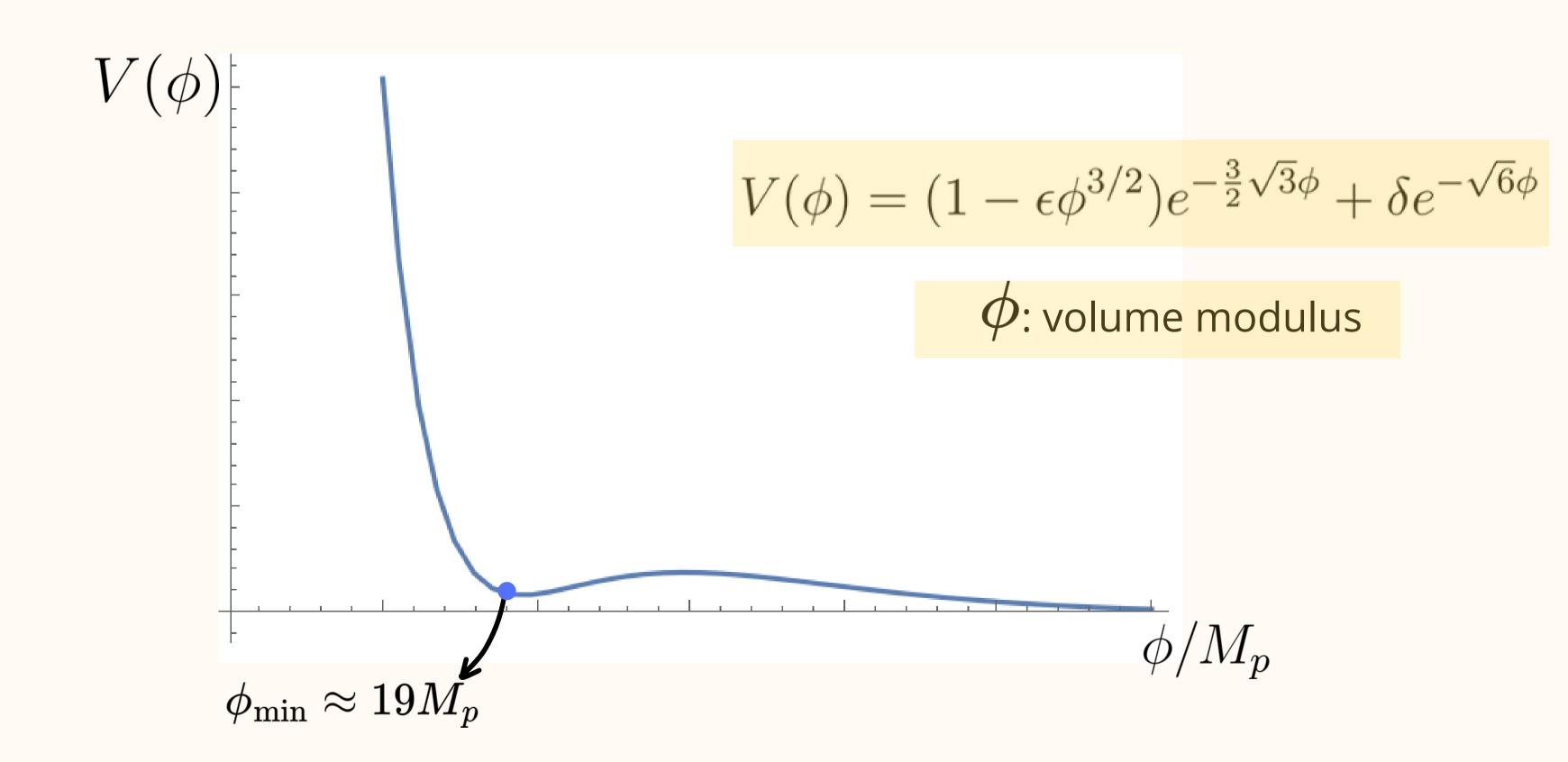
moving multiple Planckian distances

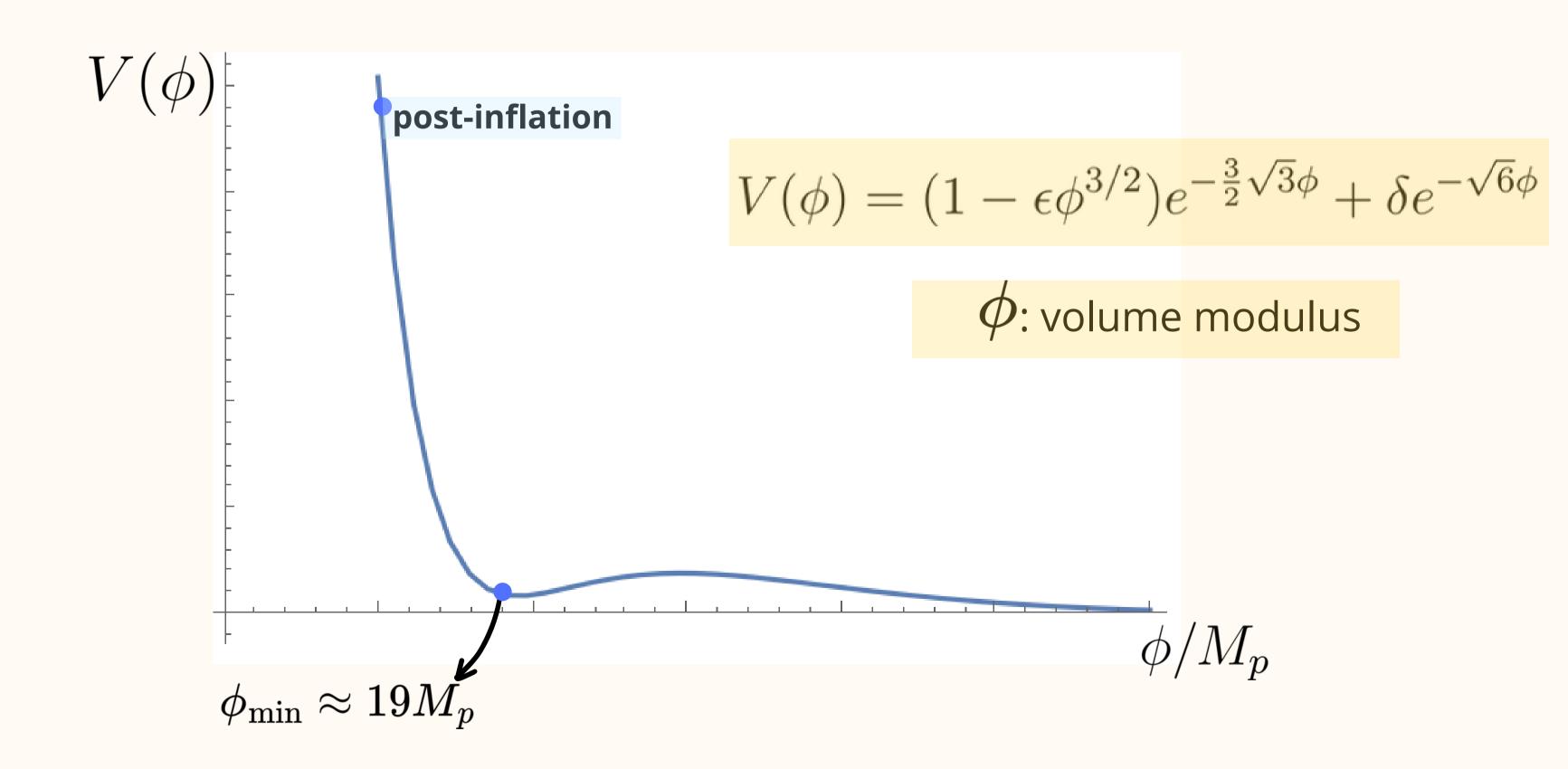
F.e. **decompactification** limit: a Kaluza-Klein tower is becoming light

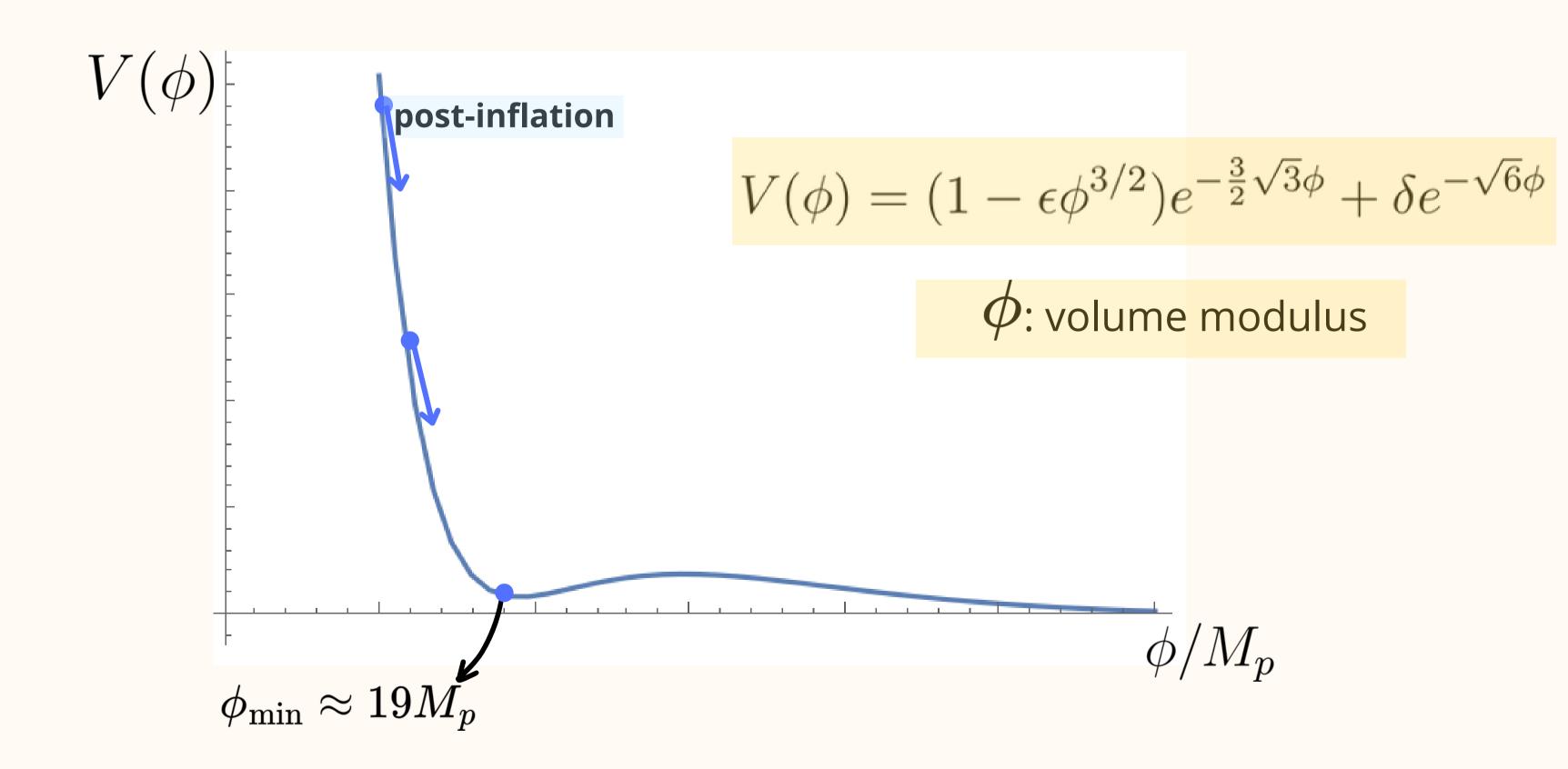
$$m_{
m KK}(t)\gg H(t)$$

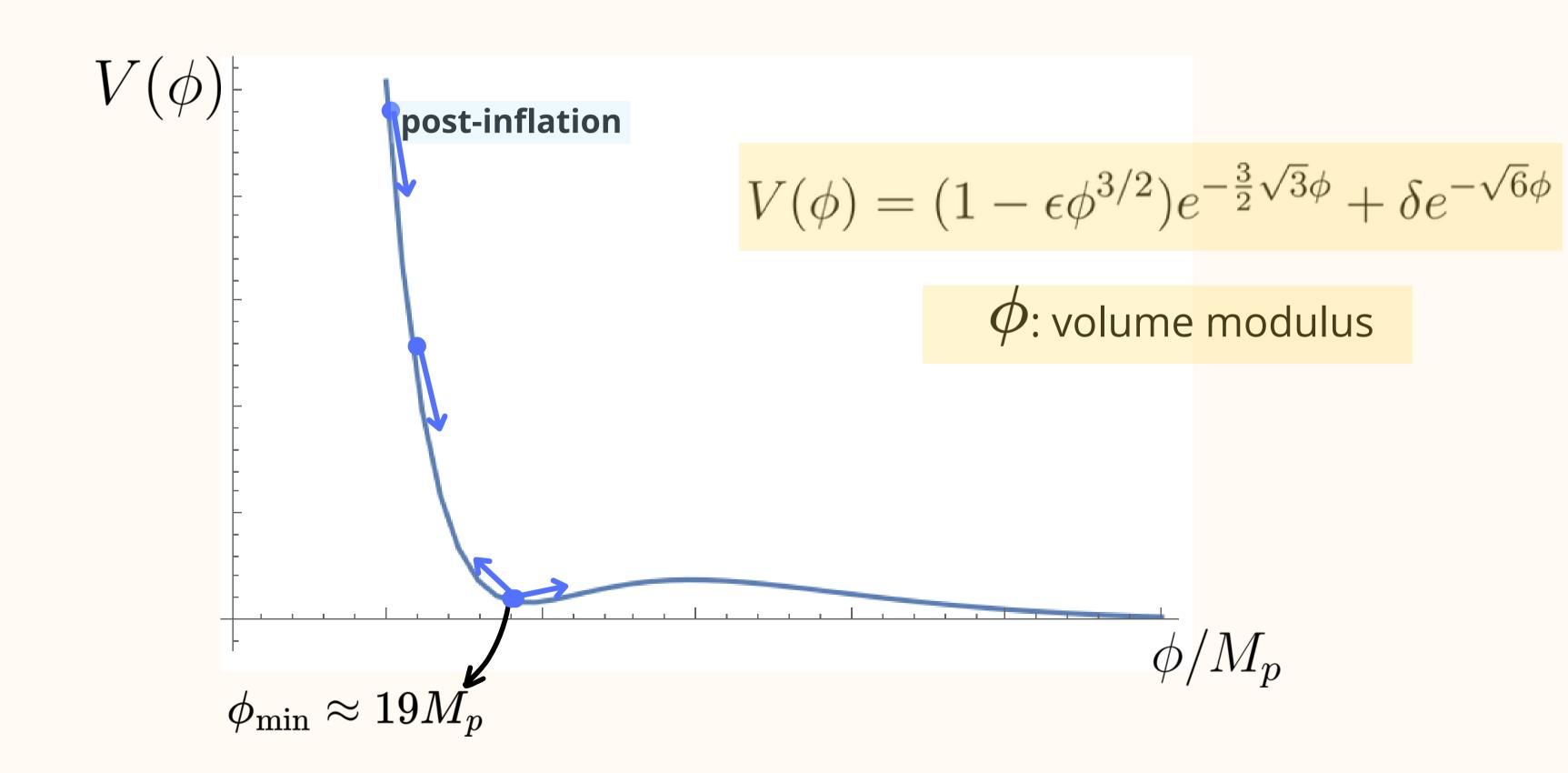


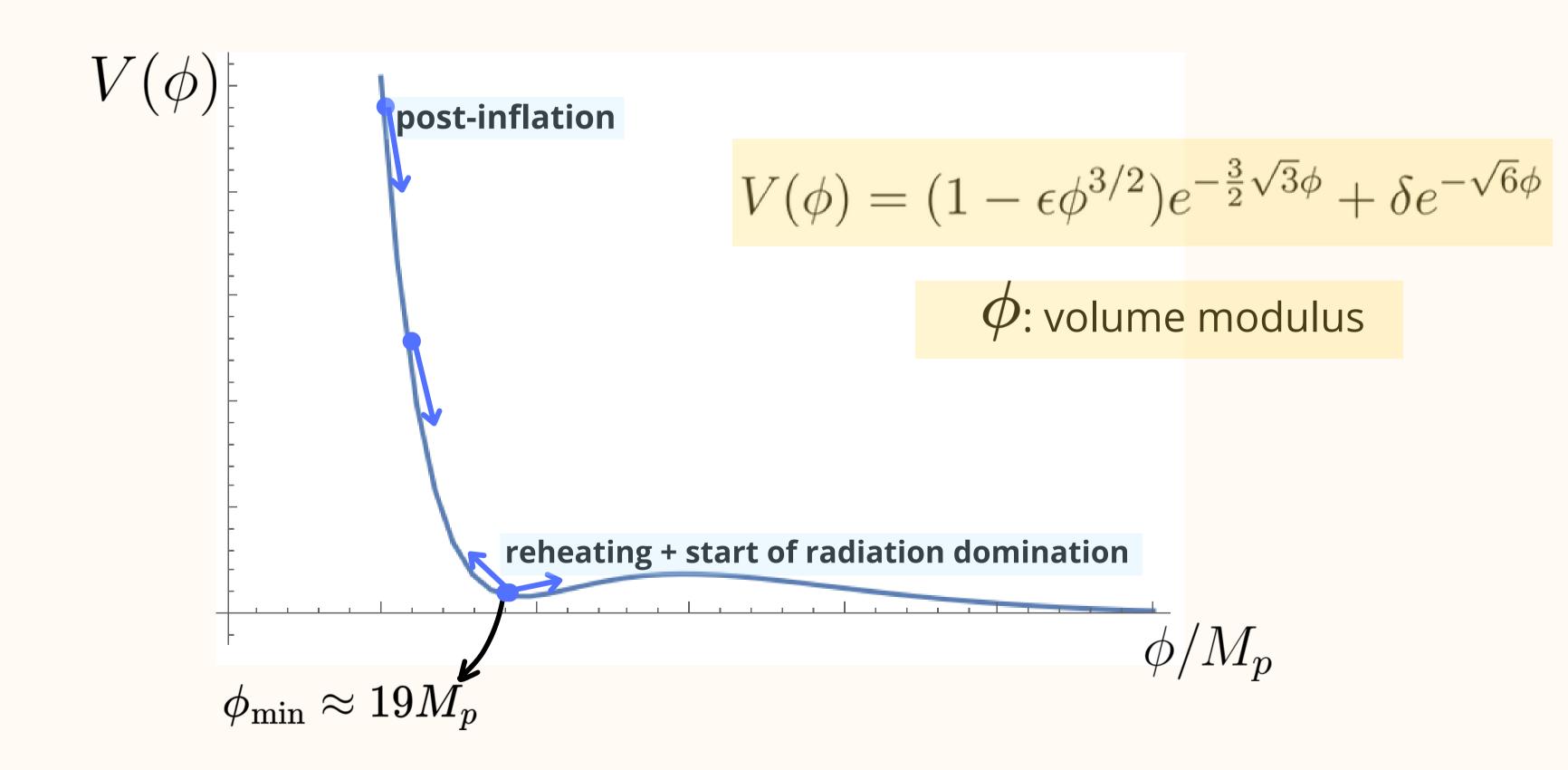


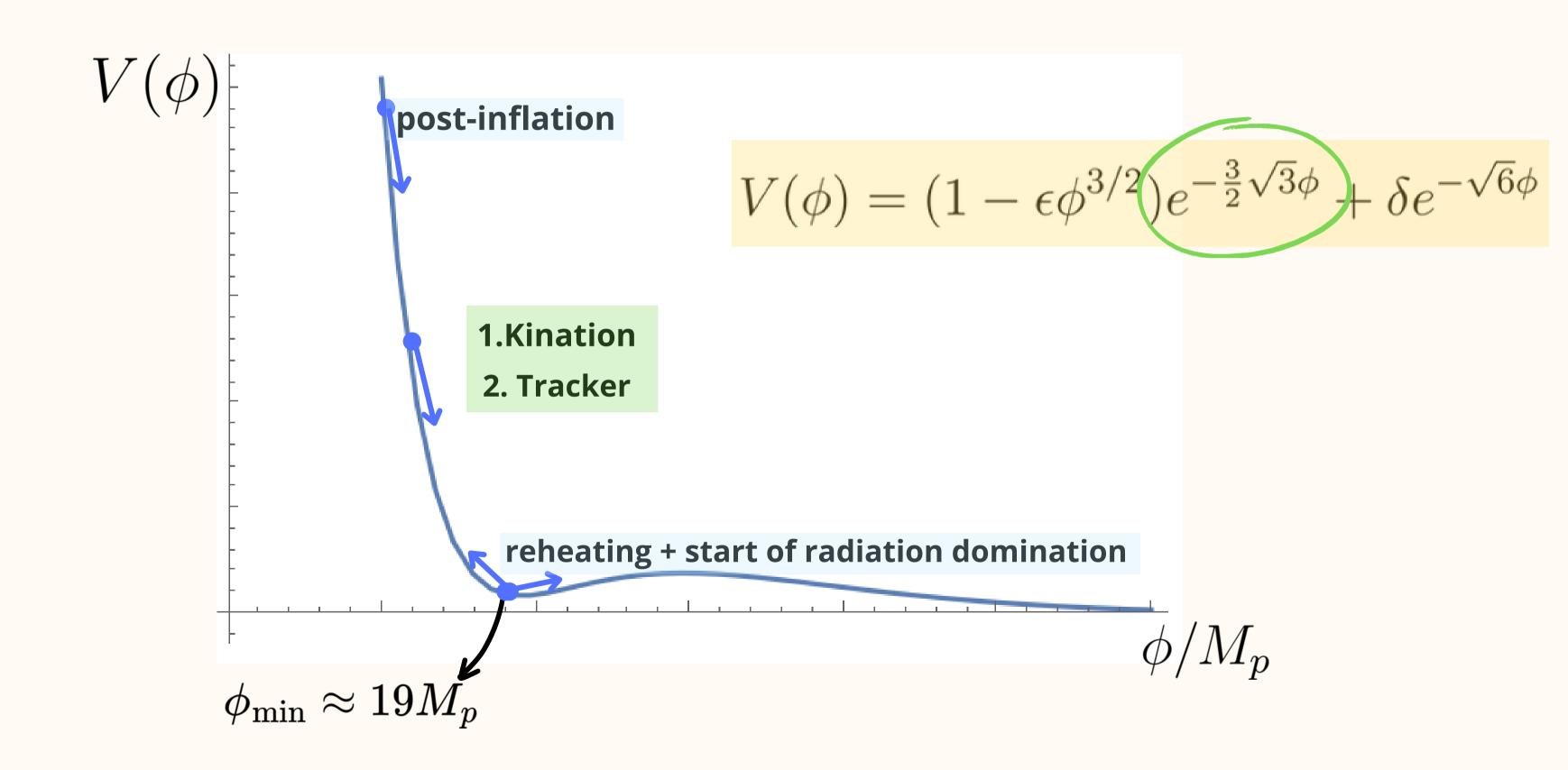












For details... background evolution and cosmological perturbations

String Theory and the First Half of the Universe

Fien Apers,^a Joseph P. Conlon,^a Edmund J. Copeland,^b Martin Mosny^a and Filippo Revello^c

Epoch	a(t)	η	Range of η	$\mathcal{H} = \frac{a'(\eta)}{a(\eta)}$	PE:KE:Rad
Inflation	$e^{H_{inf}t}$	$\sim -e^{-Ht}$	$-\infty < \eta \lesssim 0 \sim \eta_0$	H_{inf}	$\frac{1}{2}:\frac{1}{2}:\varepsilon$ (at end)
Kination	$t^{1/3}$	$\eta \sim t^{2/3}$	$\eta_0 \lesssim \eta \lesssim \frac{\eta_0}{\varepsilon}$	$\frac{1}{2\eta}$	$\begin{array}{c} \varepsilon^{3/2} : \frac{1}{2} : \frac{1}{2} \\ \text{(at end)} \end{array}$
Radiation domination: $PE \le KE$	$t^{1/2}$	$\eta \propto t^{1/2}$	$\frac{\eta_0}{arepsilon} \lesssim \eta \lesssim \frac{\eta_0}{arepsilon^{5/4}}$	$\frac{1}{\eta}$	$\varepsilon^{1/2}$: $\varepsilon^{1/2}$:1 (at end)
Radiation domination: $PE \ge KE$	$t^{1/2}$	$\eta \propto t^{1/2}$	$\frac{\eta_0}{\varepsilon^{5/4}} \lesssim \eta \lesssim \frac{\eta_0}{\varepsilon^{11/8}}$	$\frac{1}{\eta}$	$\frac{1}{2}$: $\varepsilon^{3/4}$: $\frac{1}{2}$ (at end)
Radiation Tracker	$t^{1/2}$	$\eta \propto t^{1/2}$	$\frac{\eta_0}{\varepsilon^{11/8}} \lesssim \eta \lesssim m_{\Phi}^{-1/2}$	$\frac{1}{\eta}$	$\begin{vmatrix} \frac{3(2-\gamma)\gamma}{2\lambda^2} : \frac{3\gamma^2}{2\lambda^2} : \\ 1 - \frac{3\gamma}{\lambda^2} \end{vmatrix}$
Matter domina- tion	$t^{2/3}$	$\eta \propto t^{1/3}$	$m_{\Phi}^{-1/2} \lesssim \eta \lesssim \Gamma_{\Phi}^{-1/2}$	$\frac{2}{\eta}$	NA
Reheating to Standard Model	$t^{1/2}$	$\eta \propto t^{1/2}$	$\eta \gtrsim \Gamma_{\Phi}^{-1/2}$	$\frac{1}{\eta}$	0:0:1 (at end)

Summary

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- There is an epoch in the early universe, in between inflation and the start of radiation domination, where large part of the energy density is in string theory moduli
- If we live at the **end of the world** (near boundary of moduli space) there is a **general description** of this epoch with a **scalar field rolling down a steep exponential** $V(\phi) = V_0 e^{-\lambda \phi/M_p}$ and going through periods of
 - kination
 - tracker (cosmological scaling solutions).
- This may be followed by a period of moduli domination.

Kination	$\phi = \sqrt{\frac{2}{3}} \ln t$	$a(t) \sim t^{1/3}$	$\Omega_{\phi} = 1$
Radiation Tracker	$\phi = \frac{2}{\lambda} \ln t$	$a(t) \sim t^{1/2}$	$\Omega_{\phi} = \frac{4}{\lambda^2}$