

STRING THEORY IN THE FIRST HALF OF THE UNIVERSE

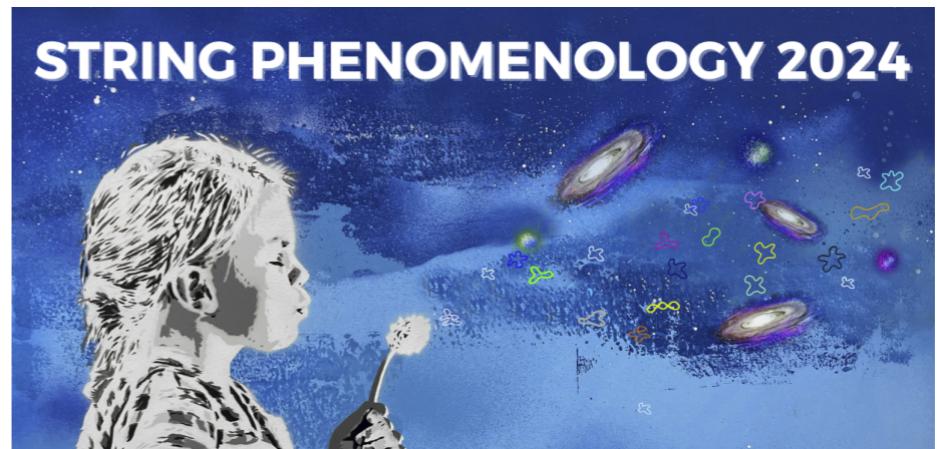
PART II

*See Part I,
by Fien Apers*

Filippo Revello, Utrecht University



**Utrecht
University**



Based on:

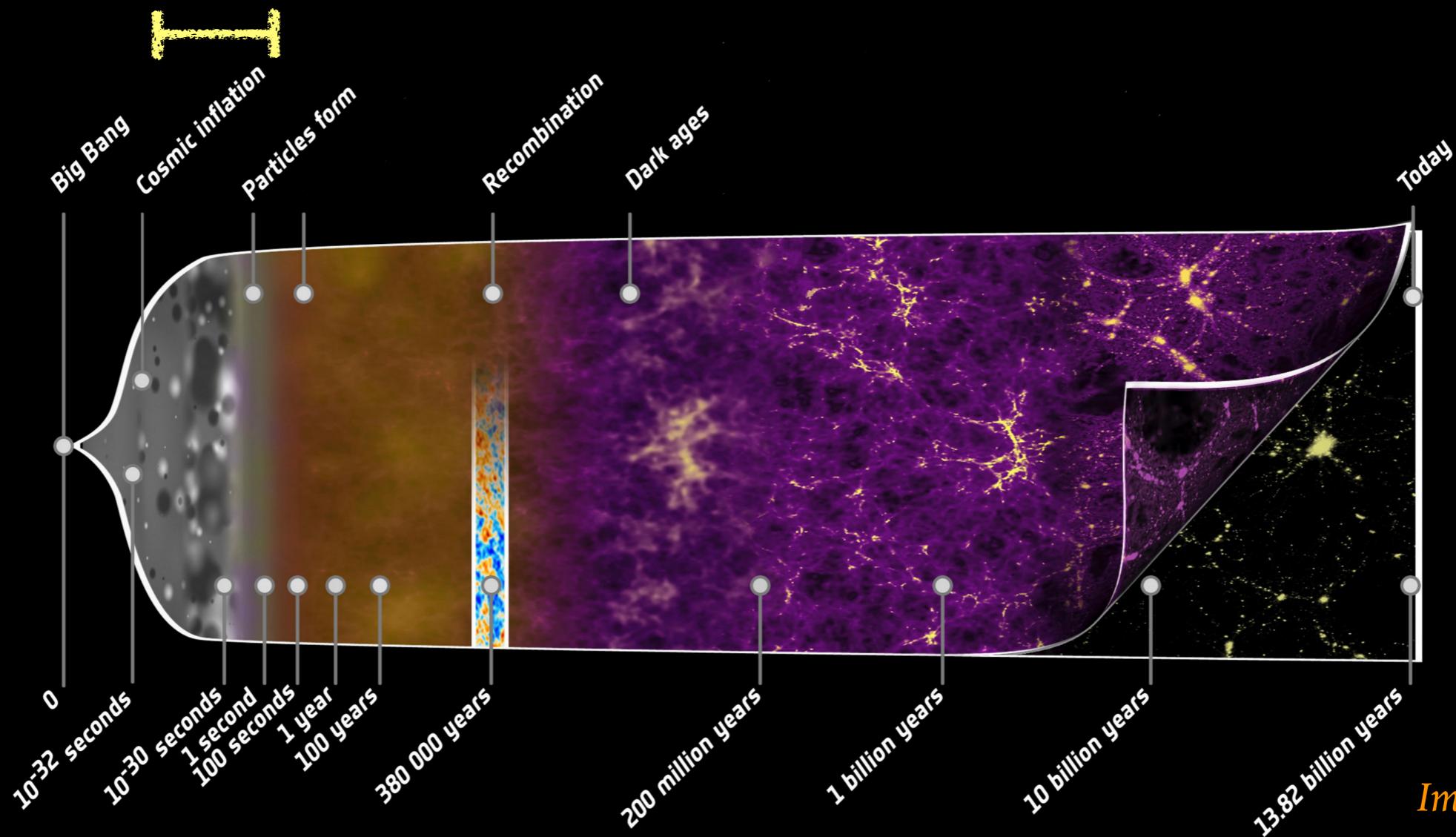
2401.04064 [Apers, Conlon, Copeland, Mosny, FR]
24xx.xxxxxx [Ghoshal, FR, Villa]

THE FIRST HALF OF THE UNIVERSE (ON A LOG SCALE)

String Theory motivates exotic cosmologies in the Early Universe

Unconstrained by current observations

Huge opportunity for ST !



This talk: consequences for perturbations & cosmic strings/GWs

A NON-STANDARD PICTURE

Example of full cosmological history compatible with ST

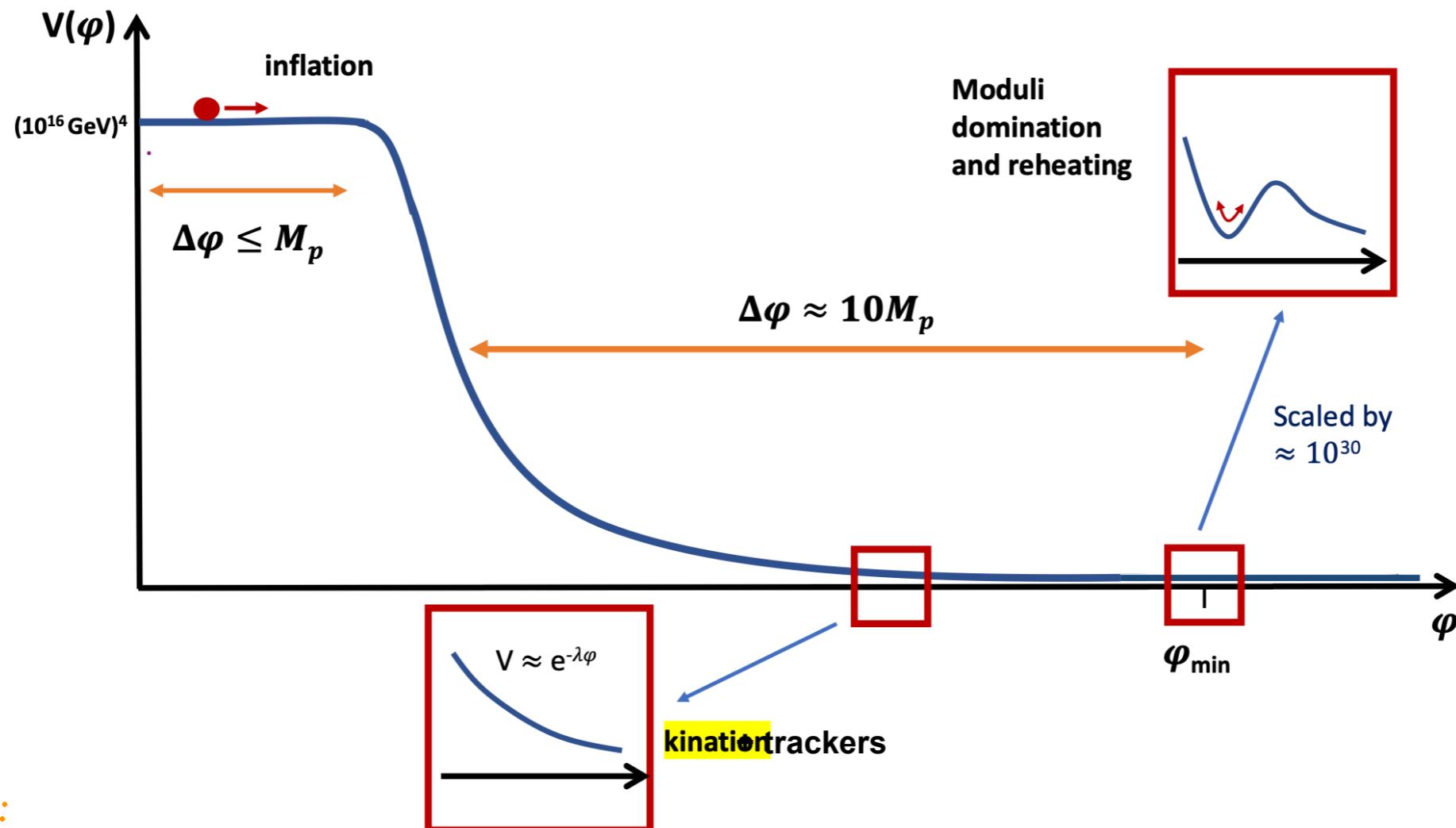


Figure from:

[Cicoli, Conlon, Maharana, Quevedo, Parameswaran, Zavala '23] [Apers, Conlon, Copeland, Mosny, FR '24]

Candidate/example: volume modulus in LVS

(In general \neq inflaton!) E.g. waterfall field [Burgess, Quevedo '22]

COSMOLOGICAL PERTURBATIONS 1

Most general metric and matter pert. + gauge invariance

Newtonian gauge:*

“Gravitational potential”

$$ds^2 = a^2(\eta) [-(1 + 2\Phi)d\eta^2 + (1 - 2\Phi)\delta_{ij}dx^i x^j]$$

$$T_0^0 \equiv -(\bar{\rho} + \delta\rho) \quad T_i^0 \equiv (\bar{\rho} + \bar{P})\partial_i v \quad T_j^i \equiv (\bar{P} + \delta P)\delta_j^i$$

Figure taken from [Baumann '22]

Table 6.1 Summary of the evolution of cosmological perturbations.

Gauge invariant variables

$$\Phi$$

$$\Delta = \frac{\delta\rho}{\bar{\rho}} + v \frac{\bar{\rho}'}{\bar{\rho}}$$



		radiation era	matter era
Φ	$k < \mathcal{H}$	const	const
	$k > \mathcal{H}$	$a^{-2} \cos(k\eta/\sqrt{3})$	
Δ_r	$k < \mathcal{H}$	a^2	a
	$k > \mathcal{H}$	$\cos(k\eta/\sqrt{3})$	
Δ_m	$k < \mathcal{H}$	a^2	a
	$k > \mathcal{H}$	$\ln a$	

COSMOLOGICAL PERTURBATIONS 2

Solve perturbed Einstein Eqs

$$\nabla^2 \Phi - 3\mathcal{H}(\Phi' + \mathcal{H}\Phi) = \frac{a^2}{2M_P^2} (\delta\rho_k + \delta\rho_p + \delta\rho_f)$$

$$\Phi'' + 3\mathcal{H}\Phi' + (2\mathcal{H} + \mathcal{H}^2)\Phi = \frac{a^2}{2M_P^2} (\delta\rho_k - \delta\rho_p + (\gamma - 1)\delta\rho_f)$$

+ scalar field EOMs

$$\delta\phi_k'' + 2\mathcal{H}\delta\phi_k' + k^2\delta\phi_k + \lambda^2 a^2 \frac{V(\phi)}{M_P^2} \delta\phi_k = 4\phi'\Phi_k' + 2\lambda a^2 \frac{V(\phi)}{M_P} \Phi_k$$

Sub-horizon

$$k\eta \gg 1$$



*Analytical
solutions*

Super-horizon



$$k\eta \ll 1$$

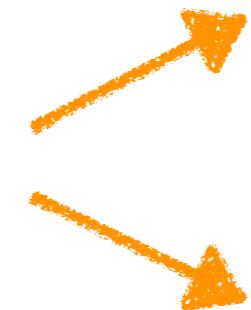
COSMOLOGICAL PERTURBATIONS – RESULTS

Kination:

Very fast!

$$\Phi(\eta) = \frac{C_1}{k\eta} J_1(k\eta) + \frac{C_2}{k\eta} Y_1(k\eta)$$

$$\Delta(\eta) = -\frac{2iH_{\text{inf}}}{M_P \sqrt{\varepsilon_V k^3}} \frac{k\eta J_1(k\eta)}{J_0(k\eta_0)}$$



$a(\eta) \sin(k\eta), \quad k\eta \gg 1$

$a(\eta)^4, \quad k\eta \ll 1$

Radiation Tracker: $k\eta \ll 1$

Very similar to radiation domination, $O(1)$ diff

$$\Phi(\eta) \simeq \frac{C_3 \sin\left(\frac{k\eta}{\sqrt{3}}\right) + C_4 \sin(k\eta + C_5)}{a(\eta)^2} \quad \Delta \simeq C_3 \sin\left(\frac{k\eta}{\sqrt{3}}\right) + C_4 \sin(k\eta + C_5)$$

Matter Tracker: $k\eta \ll 1$

Unlike matter domination, depends on λ

$$\Phi \simeq C_6 a(\eta)^{\frac{1}{4}\left(\frac{\sqrt{25\lambda^2-72}}{\lambda}-5\right)} \sim a(\eta)^{-0.14} \quad \Delta \simeq C_7 a(\eta)^{\frac{1}{4}\left(\frac{\sqrt{25\lambda^2-72}}{\lambda}-1\right)} \sim a(\eta)^{0.86}$$

LVS

LVS

+ various analytical results for transient epochs

TIME-DEPENDENT COUPLINGS

Hallmark of String Theory: couplings depend on moduli

$$\Phi_i \rightarrow \Phi_i(t)$$

Time dependent couplings

Example:

$$\frac{2\pi}{g_{\text{YM}}^2} = \int_{\Sigma_i} e^{-\phi} \sqrt{g}$$

Cycle volumes can vary with time

Relevant for close-to-critical couplings? Higgs mass, quartic...

*THE coupling
in String Theory*

$$m_s \sim \frac{M_P}{\sqrt{\mathcal{V}(t)}}$$

time dependent!

*See talk
by Joe Conlon*

COSMIC SUPERSTRINGS

Cosmic F-strings

[Witten '85] [Polchinski '88] [Copeland, Myers, Polchinski '05]

Conventional wisdom:

cosmic superstrings excluded by $G\mu = \left(\frac{m_s}{M_P}\right)^2 < 10^{-7}$ (CMB)

OK if $\mathcal{V} \gtrsim 10^5 - 10^6$ after CMB



Can be tested with *Gravitational Waves*

[Ghoshal, FR, Villa '24]

Distinctive new pheno from time dependent tension $\mu(t)$

Also relevant for QFT strings (ex. axion strings) ?



Growing loops



Scaling regime?



High freq. GWs

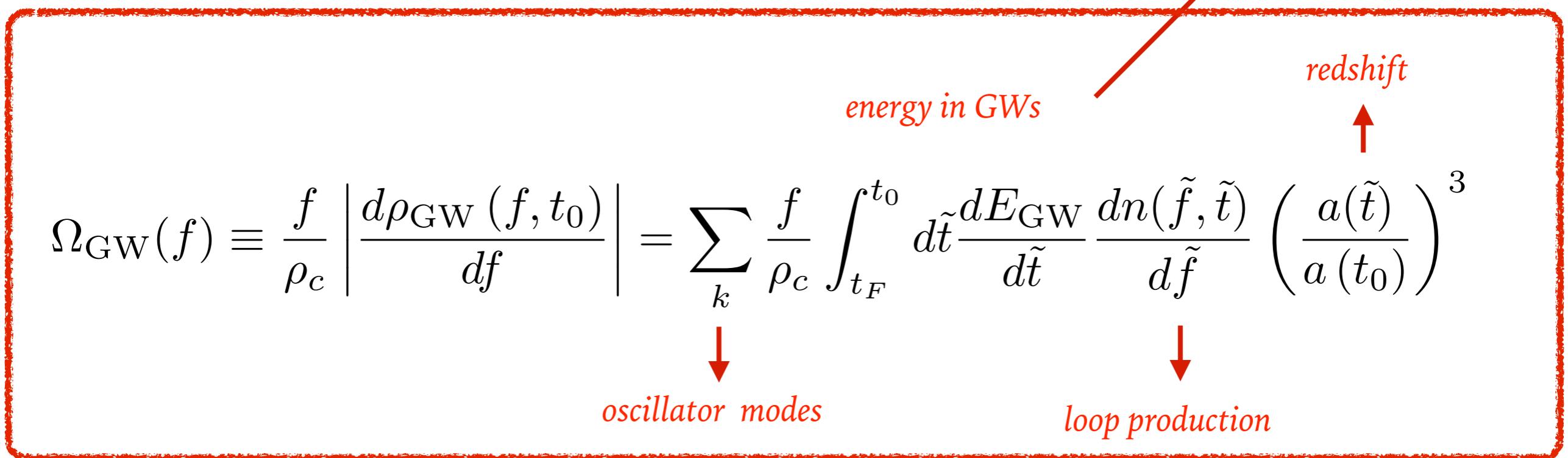
[Conlon, Copeland, Hardy, Sanchez Gonzalez '24] [Ghoshal, FR, Villa '24]

GRAVITATIONAL WAVES FROM COSMIC SUPERSTRINGS

Dynamical attractor: scaling regime

$\mu(t)$ “quasi” scaling

Long strings $\rho_{\text{strings}} = \frac{\mu}{L^2}$ $L = \xi t$ Loops decay to GWs $\frac{dE}{dt} = \Gamma G \mu^2$



i) Information on cosmic history

E.g. volume kination

$$\Omega_{\text{GW}} h^2 \sim f^5$$

ii) GW emission extended in time

signal extends to low frequencies*

OUTLOOK & FUTURE DIRECTIONS

Exotic epochs from ST: vast and unexplored territory

Kination

Trackers

*Dynamical system
approach*

GWs

Anomalous perturbation growth

*time dependent
couplings*

Radiation seed / self perturbations

DM, baryogenesis

Primordial Black Holes

Axion dynamics

THANK YOU FOR YOUR ATTENTION!

WHAT ABOUT THE DISTANCE CONJECTURE?

Distance conjecture:

$$\Delta\phi \gg 1$$



[Ooguri, Vafa '06]

[Ooguri, Palti, Shiu, Vafa '19]

towers of
light states



Invalidate EFT

Kinematics

KK modes become light, but so do m_ϕ, Λ_ϕ



Dynamics

KK modes above Hubble

$$m_{KK}(t) \gg H(t)$$



Cutoff is adiabatic

$$\left| \frac{d\Lambda(t)_{KK}}{dt} \right| \ll \Lambda(t)_{KK}^2$$



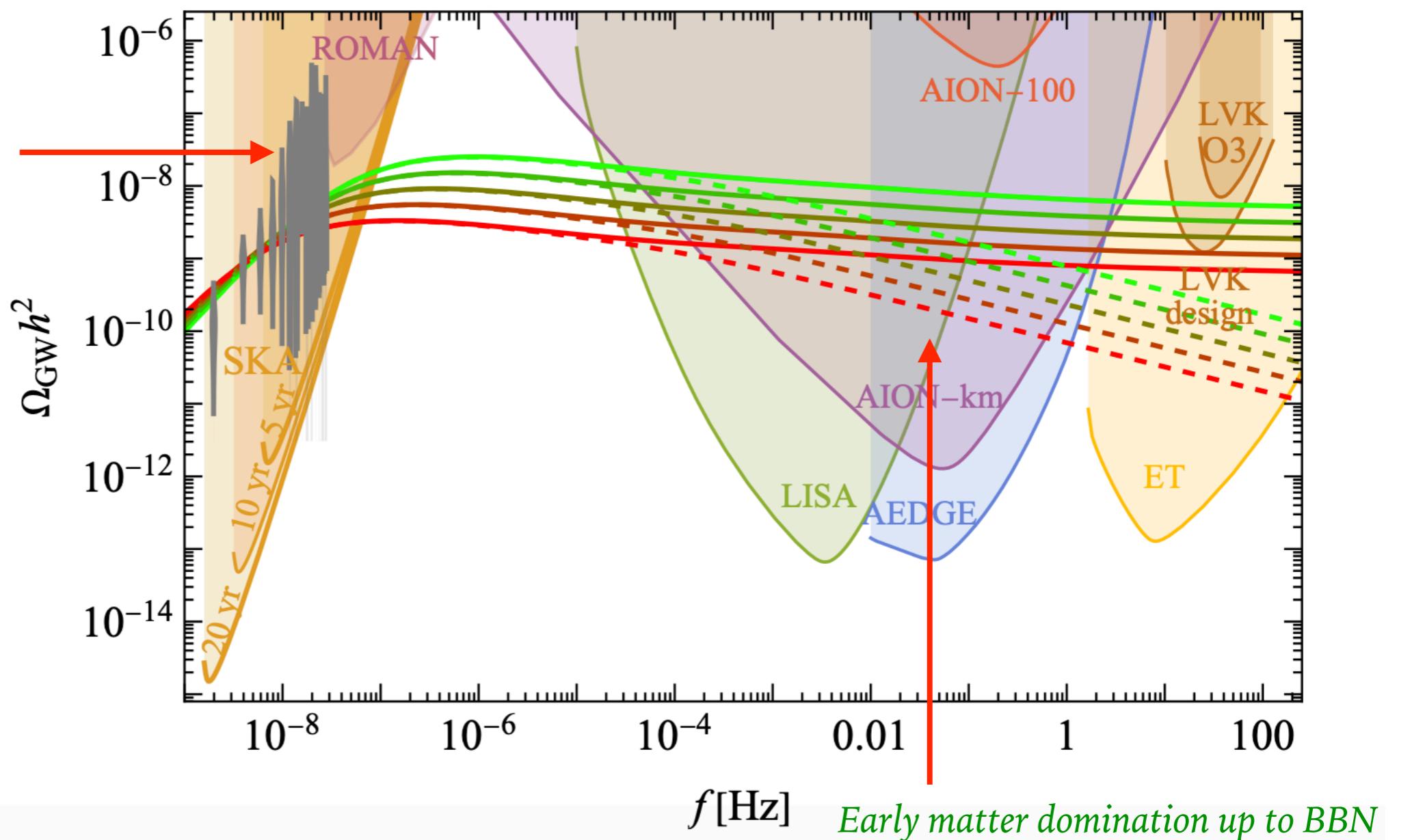
GW FROM COSMIC STRINGS

[Ellis, Lewicki, Lin, Vaskonen '23]

Different fits to Nanograv
for cosmic superstrings

$$G\mu \sim 10^{-11} - 10^{-12}$$

$$P \sim 10^{-1} - 10^{-3}$$



For LVS

$$t_{\text{reh}} \sim \frac{1}{\Gamma_\Phi} = \left(\frac{1}{4\pi} \frac{m_\Phi^3}{M_P^2} \right)^{-1}$$

Predicts deviation!

THE BACKGROUND COSMOLOGY

Epoch	$a(t)$	η	Range of η	$\mathcal{H} = \frac{a'(\eta)}{a(\eta)}$	PE:KE:Rad
Inflation	$e^{H_{inf}t}$	$\sim -e^{-Ht}$	$-\infty < \eta \lesssim 0 \sim \eta_0$	H_{inf}	$\frac{1}{2} : \frac{1}{2} : \epsilon$ (at end)
Kination	$t^{1/3}$	$\eta \sim t^{2/3}$	$\eta_0 \lesssim \eta \lesssim \frac{\eta_0}{\epsilon}$	$\frac{1}{2\eta}$	$\epsilon^{3/2} : \frac{1}{2} : \frac{1}{2}$ (at end)
Radiation domination: PE \leq KE	$t^{1/2}$	$\eta \propto t^{1/2}$	$\frac{\eta_0}{\epsilon} \lesssim \eta \lesssim \frac{\eta_0}{\epsilon^{5/4}}$	$\frac{1}{\eta}$	$\epsilon^{1/2} : \epsilon^{1/2} : 1$ (at end)
Radiation domination: PE \geq KE	$t^{1/2}$	$\eta \propto t^{1/2}$	$\frac{\eta_0}{\epsilon^{5/4}} \lesssim \eta \lesssim \frac{\eta_0}{\epsilon^{11/8}}$	$\frac{1}{\eta}$	$\frac{1}{2} : \epsilon^{3/4} : \frac{1}{2}$ (at end)
Tracker	$t^{1/2}$	$\eta \propto t^{1/2}$	$\frac{\eta_0}{\epsilon^{11/8}} \lesssim \eta \lesssim m_\Phi^{-1/2}$	$\frac{1}{\eta}$	$\frac{3(2-\gamma)\gamma}{2\lambda^2} : \frac{3\gamma^2}{2\lambda^2} : 1 - \frac{3\gamma}{\lambda^2}$
Matter domination	$t^{2/3}$	$\eta \propto t^{1/3}$	$m_\Phi^{-1/2} \lesssim \eta \lesssim \Gamma_\Phi^{-1/2}$	$\frac{2}{\eta}$	NA
Reheating to Standard Model	$t^{1/2}$	$\eta \propto t^{1/2}$	$\eta \gtrsim \Gamma_\Phi^{-1/2}$	$\frac{1}{\eta}$	0:0:1 (at end)

[Apers,Conlon,Copeland,Mosny,FR '24]

Analytical treatment of transition epochs

Eg: kination to radiation domination:

$$a(\eta) = a_0 \sqrt{\frac{\eta}{\eta_0} + \frac{\varepsilon}{4} \left(\frac{\eta}{\eta_0}\right)^2}$$

$$\mathcal{H}(\eta) = \frac{1}{2\eta} \frac{1 + \frac{\varepsilon\eta}{2\eta_0}}{1 + \frac{\varepsilon\eta}{4\eta_0}}$$

Useful to analyse e.g. perturbations

EQUATIONS OF MOTION

$$\left\{ \begin{array}{l} \ddot{\phi}^i + \Gamma_{jk}^i \dot{\phi}^j \dot{\phi}^k + (d-1)H\dot{\phi}^i + \partial^i V = 0 \\ \frac{(d-1)(d-2)}{2} H^2 = \frac{1}{2} G_{ij} \partial_\mu \phi^i \partial^\mu \phi^j - V(\phi^i) + \rho_\gamma \\ \dot{\rho}_\gamma + 3\gamma H \rho_\gamma = 0 \end{array} \right.$$

*Generic fluid:
matter, radiation...*

Asymptotics of moduli space (large Vol, CS moduli...)

$$G_{ij,I} = C_I \frac{\delta_{ij}}{(s^I)^2} \quad V(s^I, a^I) = V_0 \prod_{I=1}^N \frac{1}{(s^I)^{\lambda_I}}$$

Simpler expressions, analytic treatment

DYNAMICAL SYSTEM REFORMULATION

$$x_i = \sqrt{\frac{C_i}{(d-1)(d-2)}} \frac{\dot{s}_i}{H s_i} \quad y_i = \sqrt{\frac{C_i}{(d-1)(d-2)}} \frac{\dot{a}_i}{H s_i} \quad w = \frac{1}{H} \sqrt{\frac{2\rho_\gamma}{(d-1)(d-2)}}$$

saxions

axions

fluid

$$\left\{ \begin{array}{l} \frac{d\textcolor{brown}{x}_i}{dM} = -\sqrt{\frac{(d-1)(d-2)}{C_i}} \left[\textcolor{green}{y}_i^2 - \frac{\lambda_i}{2} \left(1 - \sum_{j=1}^N (\textcolor{brown}{x}_j^2 + \textcolor{green}{y}_j^2) - \textcolor{red}{w}^2 \right) \right. \\ \quad \left. - \textcolor{brown}{x}_i(d-1) \left(1 - \sum_{j=1}^N (\textcolor{brown}{x}_j^2 + \textcolor{green}{y}_j^2) - \frac{\gamma}{2} \textcolor{red}{w}^2 \right) \right] \\ \frac{dy_i}{dM} = \sqrt{\frac{(d-1)(d-2)}{C_i}} \textcolor{brown}{x}_i y_i - (d-1) \left(1 - \sum_{j=1}^N (\textcolor{brown}{x}_j^2 + \textcolor{green}{y}_j^2) - \frac{\gamma}{2} \textcolor{red}{w}^2 \right) y_i \\ \frac{dw}{dM} = (d-1) \textcolor{red}{w} \left[\sum_{j=1}^N (\textcolor{brown}{x}_j^2 + \textcolor{green}{y}_j^2) + \frac{\gamma}{2} (\textcolor{red}{w}^2 - 1) \right] \end{array} \right.$$

*See talk
by Flavio Tonioni*

[Copeland, Liddle, Wands '98], [Collinucci, Nielsen, Van Riet '04]

[(Brinkmann), Cicoli, Dibitetto, Pedro '20 x2, 22], [Shiu, Tonioni, Tran '23] x 3, [FR '23]...