

Higher-curvature Gravities, EFTs and the Swampland

Ángel Jesús Murcia Gil, INFN Padova (Italy)

String Pheno '24, Padova (Italy)



Mainly based on the following works:

- 1) JHEP 11 (2019) 062, w/ P. Bueno, P. Cano & J. Moreno
- 2) JHEP 10 (2020) 125, w/ P. Cano
- 3) JHEP 07 (2022) 010, w/ P. Cano, A. Rivadulla & X. Zhang

String effective actions include **higher-curvature** terms:

$$I_{\text{ST}} = I_{\text{sugra}} + \sum_{k=0}^{\infty} \sum_{n=1}^{\infty} g_s^{2k} (\alpha')^n I_{n,k}.$$

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- 1 **Heterotic String Theory.** Truncating all Yang-Mills fields [[Gross, Sloan '87](#); [Bergshoeff, de Roo '89](#)]:

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- 2 **IIB String Theory.** On $\mathcal{A}_5 \times S^5$ ansatz [Gubser, Klebanov, Tseytlin '98]:

$$I_{\text{IIB}_{\mathcal{A}_5 \times S^5}}[g_{\mathcal{A}_5}] = \frac{1}{16\pi G} \int d^5x \sqrt{|g_{\mathcal{A}_5}|} \left[R + \frac{12}{\ell^2} + \frac{\zeta(3)}{8} \alpha'^3 W^4 \right] + \dots,$$

up to subleading terms in α' . W^4 is certain combination of quartic invariants.

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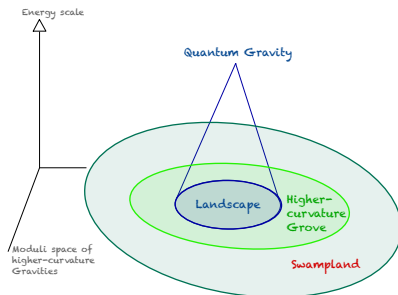
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Within the terrible confines of the Swampland, a magical and mysterious Grove towards the Landscape there is.



Generalized Quasitopological Gravities

We focus on a **special class of theories** with intriguing properties on **static** and **spherically symmetric** ansätze:

$$ds_{N,f}^2 = -N(r)^2 f(r) dt^2 + \frac{1}{f(r)} dr^2 + r^2 d\Omega_{D-2}^2.$$

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Definition (Generalized Quasitopological Gravities)

Let $\mathcal{L}(g^{ab}, R_{abcd})$ be theory of pure gravity. It is a **Generalized Quasitopological Gravity (GQG)** [e.g. Oliva, Ray '10; Myers, Robinson '10; Bueno, Cano '16; Hennigar, Kubizňák, Mann '17; Bueno, Cano '17] iff reduced Lagrangian $L_f = r^{D-2} \mathcal{L}|_{ds_{1,f}^2}$:

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GR is trivially a **GQG**. **Cubic** example in $D = 4$ [Bueno, Cano '16]:

$$\mathcal{P} = 12R_a{}^c{}_b{}^d R_c{}^e{}_d{}^f R_e{}^a{}_f{}^b + R_{ab}{}^{cd} R_{cd}{}^{ef} R_{ef}{}^{ab} - 12R_{abcd} R^{ac} R^{bd} + 8R_a{}^b R_b{}^c R_c{}^a.$$

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Any purely gravitational higher-curvature **theory** can be **mapped** via perturbative **field redefinitions** to a **GQG**
[Bueno, Cano, Moreno, ÁM '19].

Type-IIB action at $\mathcal{O}(\alpha'^3)$ as a GQG

$$I_{\text{IIB}_{\mathcal{A}_5 \times S^5}}[g_{\mathcal{A}_5}] = \frac{1}{16\pi G} \int d^5x \sqrt{|g_{\mathcal{A}_5}|} \left[R + \frac{12}{\ell^2} + \frac{\zeta(3)}{8} \alpha'^3 W^4 \right] + \dots,$$

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We perform the metric redefinition:

$$g_{ab} \rightarrow g_{ab} - \frac{\zeta(3)\alpha'^3}{8} \left(\hat{C}_{ab} - \frac{1}{3} \hat{C} g_{ab} \right),$$

for a specific cubic-curvature symmetric tensor \hat{C}_{ab} and $\hat{C} = \hat{C}_a^a$.

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for a specific cubic-curvature symmetric tensor \hat{C}_{ab} and $\hat{C} = \hat{C}^a_a$. The **new theory** is a **GQG** with solution:

$$ds^2 = -f(r)dt^2 + \frac{dr^2}{f(r)} + \frac{r^2}{\ell^2} d\vec{x}_3^2,$$

$$f(r) = \frac{r^2}{\ell^2} \left[1 - \frac{\omega^4}{r^4} + \gamma \left(\frac{25\omega^{12}}{2r^{12}} - \frac{15\omega^{16}}{2r^{16}} \right) \right] + \mathcal{O}(\gamma^2),$$

where $\gamma = \frac{\zeta(3)\alpha'^3}{8\ell^6}$ and ω an integration constant related to the energy.

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The **thermodynamic** properties of the black hole are:

$$T(\omega) = \frac{\omega}{\pi\ell^2} \left(1 - \frac{15}{4}\gamma \right) + \mathcal{O}(\gamma^2),$$

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Higher-curvature gravities belonging to the **swampland** can be extremely useful for **Quantum Gravity!**

Weak Gravity Conjecture for higher-curvature gravities

On the other hand, it is possible to use the **swampland program** to **constrain** the **couplings** of **higher-order terms** in effective actions [e.g. Bellazzini, Lewandowski, Serra '19; Charles '19; Aalsma, Shiu '22].

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Example in $D = 4$ at **quadratic derivative order**:

$$\mathcal{L} = \alpha R F^2 + \beta R_{ab} F^{ac} F^b{}_c - (10\alpha + 2\beta) R_{abcd} F^{ab} F^{cd}.$$

Extremal Black Holes, EQGs and WGC

It is possible to analyze which EQGs are consistent with **Weak Gravity**

Conjecture [Arkani-Hamed, Motl, Nicolis, Vafa '06]: $P/M|_{\text{ext}}$ must not increase as mass increases¹.

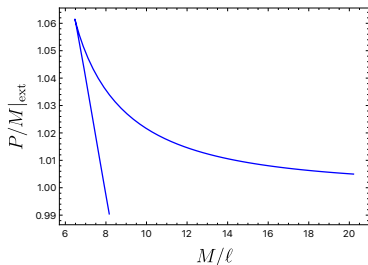
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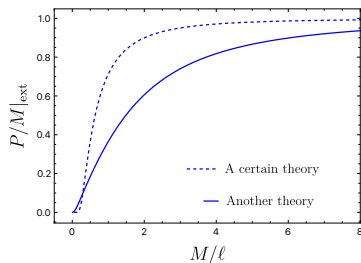
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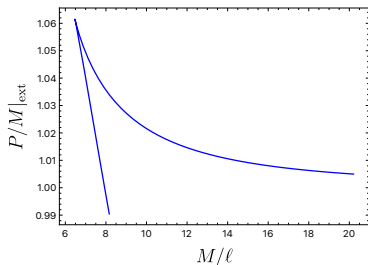
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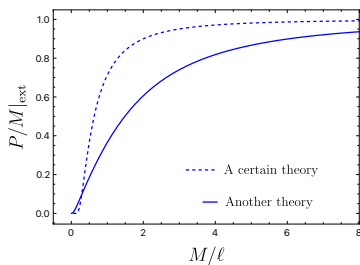
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For EQGs consistent with WGC, **extremal BH** solutions do not exist below a **minimal mass**.

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Holographic EQGs and WGC

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Focus on **Rényi** entropies S_n across an **entangling sphere**. Either for $\mu = 0$ or holographic Einstein-Maxwell, the following inequalities are met [[Hung, Myers, Smolkin, Yale '11](#)]:

$$\frac{\partial}{\partial n} S_n \leq 0, \quad \frac{\partial}{\partial n} \left(\frac{n-1}{n} S_n \right) \geq 0$$
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Grazie mille!