ASPECTS OF DYNAMICAL Cobordism in Ads/CFT

Jesús Huertas Instituto de Física Teórica (IFT), Madrid

String Pheno 2024, Padova, 25 June 2024



Based on [JH, Uranga, 2306.07335]

INTRODUCTION

COBORDISM CONJECTURE

[McNamara, Vafa, '19]

The cobordism classes of any solution of Quantum Gravity have to be trivial



[Angius, Buratti, Calderón-Infante, Delgado, **JH**, Uranga, '21-22] [Blumenhagen, Cribiori, Kneissl, Makridou, Wang, '22-23]

• "Dynamical": The cobordism is a spacetime solution

[Angius, Buratti, Calderón-Infante, Delgado, **JH**, Uranga, '21-22] [Blumenhagen, Cribiori, Kneissl, Makridou, Wang, '22-23]

• "Dynamical": The cobordism is a spacetime solution



[Angius, Buratti, Calderón-Infante, Delgado, **JH**, Uranga, '21-22] [Blumenhagen, Cribiori, Kneissl, Makridou, Wang, '22-23]

• "Dynamical": The cobordism is a spacetime solution



See Andriana, Bjoern, Christian and Roberta's talks!!

[Angius, Buratti, Calderón-Infante, Delgado, **JH**, Uranga, '21-22] [Blumenhagen, Cribiori, Kneissl, Makridou, Wang, '22-23]

• "Dynamical": The cobordism is a spacetime solution



See Andriana, Bjoern, Christian and Roberta's talks!!

These solutions have some issues...

 $V \to \infty$

Sometimes the EFT breaks down near the ETW brane

THIS TALK: COBORDISM IN ADS/CFT

• We will present and explain the solution that interpolates between $AdS_5 \times S^5$ and nothing.

THIS TALK: COBORDISM IN ADS/CFT

• We will present and explain the solution that interpolates between $AdS_5 \times S^5$ and nothing.

• We will study this setup in holography, the ETW brane (the cobordism to nothing), will have a holographic dual, so it is defined in the full quantum gravity, not only in the EFT.

A NOTE OF CAUTION



A NOTE OF CAUTION



WHAT IS THE STRING THEORY INTERPRETATION OF THE ETW?



WHAT IS THE STRING THEORY INTERPRETATION OF THE ETW?



WHAT IS THE STRING THEORY INTERPRETATION OF THE ETW?





[Karch, Randall, '00]

 AdS_4



THE SOLUTION

[D'Hoker, Estes, Gutperle, '07]

BAGPIPE SOLUTIONS



BAGPIPE SOLUTIONS

The theory has the symmetries

 $SO(2,3) \times SO(3) \times SO(3)$

 $\begin{array}{c} A_{dS_5} + S^5 \\ A_{dS_5} \times S^5 \\ Bag \\ Bag \\ Bag \\ A_{dS_5} \times S^5 \\ Bag \\ B$

[D'Hoker, Estes, Gutperle, '07]

BAGPIPE SOLUTIONS

The theory has the symmetries

 $SO(2,3) \times SO(3) \times SO(3)$

which in spacetimes translate to the manifold

 $\operatorname{AdS}_4 \times \mathbf{S}_1^2 \times \mathbf{S}_2^2 \times \Sigma$

where everything is warped over the oriented Riemann surface Σ . The metric is:



[D'Hoker, Estes, Gutperle, '07]

BAGPIPE SOLUTIONS

The theory has the symmetries

 $SO(2,3) \times SO(3) \times SO(3)$

which in spacetimes translate to the manifold

 $\operatorname{AdS}_4 \times \mathbf{S}_1^2 \times \mathbf{S}_2^2 \times \Sigma$

where everything is warped over the oriented Riemann surface Σ . The metric is:

$$ds^{2} = f_{4}^{2} ds_{AdS_{4}}^{2} + f_{1}^{2} ds_{\mathbf{S}_{1}^{2}}^{2} + f_{2}^{2} ds_{\mathbf{S}_{2}^{2}}^{2} + ds_{\Sigma}^{2}$$



[D'Hoker, Estes, Gutperle, '07]

[D'Hoker, Estes, Gutperle, '07]

BAGPIPE SOLUTIONS

$$ds^{2} = f_{4}^{2} ds_{AdS_{4}}^{2} + f_{1}^{2} ds_{\mathbf{S}_{1}^{2}}^{2} + f_{2}^{2} ds_{\mathbf{S}_{2}^{2}}^{2} + ds_{\Sigma}^{2}$$
$$ds_{\Sigma}^{2} = 4\rho^{2} (dr^{2} + r^{2} d\varphi^{2})$$



[D'Hoker, Estes, Gutperle, '07]

BAGPIPE SOLUTIONS

$$ds^{2} = f_{4}^{2} ds_{AdS_{4}}^{2} + f_{1}^{2} ds_{\mathbf{S}_{1}^{2}}^{2} + f_{2}^{2} ds_{\mathbf{S}_{2}^{2}}^{2} + ds_{\Sigma}^{2}$$
$$ds_{\Sigma}^{2} = 4\rho^{2} (dr^{2} + r^{2} d\varphi^{2})$$
$$f_{4} \equiv f_{4}(r,\varphi)$$
$$f_{1} \equiv f_{1}(r,\varphi)$$
$$f_{2} \equiv f_{2}(r,\varphi)$$
$$\rho \equiv \rho(r,\varphi)$$



D3- BRANES ENDING ON 5-BRANES

[Aharony, Berdichevsky, Berkooz, Shamir, '11]



BRANE CONFIGURATION

$$ds^{2} = f_{4}^{2}ds_{AdS_{4}}^{2} + f_{1}^{2}ds_{\mathbf{S}_{1}^{2}}^{2} + f_{2}^{2}ds_{\mathbf{S}_{2}^{2}}^{2} + 4\rho^{2}(dr^{2} + r^{2}d\varphi^{2})$$











OUR RESULTS

A STRINGY ETW-BRANE FOR ADS/CFT





But this is a complicated solution!

The 5-branes break the SO(6) rotational invariance of the ${\cal S}^5$

We need to smooth the solution to compare with the Dynamical Cobordism results

The SO(6) symmetry is broken down to SO(3)xSO(3) by the 5-branes

The SO(6) symmetry is broken down to SO(3)xSO(3) by the 5-branes

If we decompose into spherical harmonics, we get plenty of modes

$$g_{\mu
u}(\varphi) = \sum_{i} f_{i}(\varphi)\hat{g}^{i}_{\mu
u}$$

The SO(6) symmetry is broken down to SO(3)xSO(3) by the 5-branes

If we decompose into spherical harmonics, we get plenty of modes

$$g_{\mu\nu}(\varphi) = \sum_{i} f_i(\varphi) \hat{g}^i_{\mu\nu}$$

We keep the zero mode, $\hat{g}^0_{\mu\nu}$, which is the one measured at low energies (in the asymptotic AdS)

The SO(6) symmetry is broken down to SO(3)xSO(3) by the 5-branes

If we decompose into spherical harmonics, we get plenty of modes

$$g_{\mu\nu}(\varphi) = \sum_{i} f_{i}(\varphi)\hat{g}^{i}_{\mu\nu}$$

We keep the zero mode, $\hat{g}^0_{\mu\nu}$, which is the one measured at low energies (in the asymptotic AdS)

This mode parametrizes the overall volume of the $S^{
m b}$

$$ds^{2} = f_{4}^{2} ds_{AdS_{4}}^{2} + f_{1}^{2} ds_{\mathbf{S}_{1}^{2}}^{2} + f_{2}^{2} ds_{\mathbf{S}_{2}^{2}}^{2} + 4\rho^{2} (dr^{2} + r^{2} d\varphi^{2})$$

$$ds^{2} = f_{4}^{2} ds_{AdS_{4}}^{2} + f_{1}^{2} ds_{\mathbf{S}_{1}^{2}}^{2} + f_{2}^{2} ds_{\mathbf{S}_{2}^{2}}^{2} + 4\rho^{2} (dr^{2} + r^{2} d\varphi^{2})$$

We reduce the 2-spheres parametrizing their volumes by radions

$$ds^{2} = e^{-\frac{1}{2}(\sigma_{1} + \sigma_{2})} ds_{6}^{2} + e^{\frac{1}{\sqrt{2}}(a\sigma_{1} - b\sigma_{2})} ds_{\mathbf{S}_{1}^{2}}^{2} + e^{\frac{1}{\sqrt{2}}(a\sigma_{2} - b\sigma_{1})} ds_{\mathbf{S}_{2}^{2}}^{2}$$

$$ds^{2} = f_{4}^{2} ds_{AdS_{4}}^{2} + f_{1}^{2} ds_{\mathbf{S}_{1}^{2}}^{2} + f_{2}^{2} ds_{\mathbf{S}_{2}^{2}}^{2} + 4\rho^{2} (dr^{2} + r^{2} d\varphi^{2})$$

We reduce the 2-spheres parametrizing their volumes by radions

$$ds^{2} = e^{-\frac{1}{2}(\sigma_{1} + \sigma_{2})} ds_{6}^{2} + e^{\frac{1}{\sqrt{2}}(a\sigma_{1} - b\sigma_{2})} ds_{\mathbf{S}_{1}^{2}}^{2} + e^{\frac{1}{\sqrt{2}}(a\sigma_{2} - b\sigma_{1})} ds_{\mathbf{S}_{2}^{2}}^{2}$$

where

$$\sigma_1 = \log(f_1^a f_2^b), \quad \sigma_2 = \log(f_1^b f_2^a) \qquad a = 1 + \frac{1}{\sqrt{2}}, \ b = 1 - \frac{1}{\sqrt{2}}$$

$$ds^{2} = f_{4}^{2} ds_{AdS_{4}}^{2} + f_{1}^{2} ds_{\mathbf{S}_{1}^{2}}^{2} + f_{2}^{2} ds_{\mathbf{S}_{2}^{2}}^{2} + 4\rho^{2} (dr^{2} + r^{2} d\varphi^{2})$$

We reduce the 2-spheres parametrizing their volumes by radions

$$ds^{2} = e^{-\frac{1}{2}(\sigma_{1} + \sigma_{2})} ds_{6}^{2} + e^{\frac{1}{\sqrt{2}}(a\sigma_{1} - b\sigma_{2})} ds_{\mathbf{S}_{1}^{2}}^{2} + e^{\frac{1}{\sqrt{2}}(a\sigma_{2} - b\sigma_{1})} ds_{\mathbf{S}_{2}^{2}}^{2}$$

where

$$\sigma_1 = \log(f_1^a f_2^b), \quad \sigma_2 = \log(f_1^b f_2^a) \qquad a = 1 + \frac{1}{\sqrt{2}}, \ b = 1 - \frac{1}{\sqrt{2}}$$

and we get a 6d theory + radions:

$$ds_6^2 = \sqrt{f_1^2 f_2^2} \left[f_4^2 ds_{AdS_4}^2 + 4\rho^2 (dr^2 + r^2 d\varphi^2) \right]$$

We do not yet have cod-1 manifold $ds_6^2 = \sqrt{f_1^2 f_2^2} \left[f_4^2 ds_{AdS_4}^2 + 4\rho^2 (dr^2 + r^2 d\varphi^2) \right]$

We do not yet have cod-1 manifold $ds_6^2 = \sqrt{f_1^2 f_2^2} \left[f_4^2 ds_{AdS_4}^2 + 4\rho^2 (dr^2 + r^2 d\varphi^2) \right]$

$$g_{\mu\nu}(\varphi) = \sum_{i} f_i(\varphi) \hat{g}^i_{\mu\nu}$$

We do not yet have cod-1 manifold $ds_6^2 = \sqrt{f_1^2 f_2^2 [f_4^2 ds_{AdS_4}^2 + 4\rho^2 (dr^2 + r^2 d\varphi^2)]}$

$$g_{\mu\nu}(\varphi) = \sum_{i} f_i(\varphi) \hat{g}^i_{\mu\nu} \qquad f_0(\varphi) = -\frac{2}{\sqrt{\pi}} \sin 2\varphi$$

We do not yet have cod-1 manifold $ds_6^2 = \sqrt{f_1^2 f_2^2 [f_4^2 ds_{AdS_4}^2 + 4\rho^2 (dr^2 + r^2 d\varphi^2)]}$

$$g_{\mu\nu}(\varphi) = \sum_{i} f_{i}(\varphi)\hat{g}_{\mu\nu}^{i} \qquad f_{0}(\varphi) = -\frac{2}{\sqrt{\pi}}\sin 2\varphi \qquad \hat{g}_{\mu\nu} = \int_{\pi/2}^{\pi} f_{0}(\varphi)g_{\mu\nu}(r,\varphi)d\varphi$$

We do not yet have cod-1 manifold $ds_6^2 = \sqrt{f_1^2 f_2^2} \left[f_4^2 ds_{AdS_4}^2 + 4\rho^2 (dr^2 + r^2 d\varphi^2) \right]$

$$g_{\mu\nu}(\varphi) = \sum_{i} f_i(\varphi) \hat{g}^i_{\mu\nu} \qquad f_0(\varphi) = -\frac{2}{\sqrt{\pi}} \sin 2\varphi \qquad \hat{g}_{\mu\nu} = \int_{\pi/2}^{\pi} f_0(\varphi) g_{\mu\nu}(r,\varphi) d\varphi$$

$$ds_5^2 = \hat{f}_4^2 ds_{AdS_4}^2 + 4\hat{\rho}^2 dr^2$$

We do not yet have cod-1 manifold $ds_6^2 = \sqrt{f_1^2 f_2^2} \left[f_4^2 ds_{AdS_4}^2 + 4\rho^2 (dr^2 + r^2 d\varphi^2) \right]$

$$g_{\mu\nu}(\varphi) = \sum_{i} f_i(\varphi) \hat{g}^i_{\mu\nu} \qquad f_0(\varphi) = -\frac{2}{\sqrt{\pi}} \sin 2\varphi \qquad \hat{g}_{\mu\nu} = \int_{\pi/2}^{\pi} f_0(\varphi) g_{\mu\nu}(r,\varphi) d\varphi$$

$$ds_5^2 = e^{2A(x)} ds_{AdS_4}^2 + dx^2$$

EXAMPLE: THE GENUS ONE CASE



EXAMPLE: THE GENUS ONE CASE





$$ds_5^2 = e^{2A(x)} ds_{AdS_4}^2 + dx^2$$



The graviton is bound on the brane!

We get a 4d massive graviton that couples to the 5d massless one

The mass is proportional to the coupling to the 5d space

[Karch, Randall, '00]

$$ds_5^2 = e^{2A(x)} ds_{AdS_4}^2 + dx^2$$



Having a 4d massive graviton, means that holographically, we have a localized energy momentum tensor in our BCFT3 that couples to the CFT4

[Karch, Randall, '00]

$$ds_5^2 = e^{2A(x)} ds_{AdS_4}^2 + dx^2$$



Having a 4d massive graviton, means that holographically, we have a localized energy momentum tensor in our BCFT3 that couples to the CFT4

→ Holographic interpretation of the ETW brane!!

[Karch, Randall, '00]

LOCALIZATION OF GRAVITY IN OUR SETUP



$$ds_5^2 = \hat{f}_4^2(r)ds_{AdS_4}^2 + \hat{\rho}^2(r)dr^2$$

$$ds_5^2 = e^{2A(x)} ds_{AdS_4}^2 + dx^2$$

This is an explicit realization of a KR brane in string theory

LOCALIZATION OF GRAVITY IN OUR SETUP



The double scaling limit makes the 5-form flux disappear and thus closes the AdS throat; so we get a compact geometry!

 $AdS_4 \times \mathbf{X}_6$



Only possible if you have both D5- and NS5-branes



This double scaling limit isolates the dynamics of the ETW-brane!



Which is that of a brane-web



In fact, provides the gravity dual of the 3d BCFT (wedge holography)





[Assel, Bachas, Estes, Gomis '11] [Van Raamsdonk, Waddell, '21]

The fact that we can close the throat in a continuous fashion and make the space compact, is related to continuously making the 4d graviton massless



CONCLUSIONS

• Cobordism defect of $AdS_5 \times S^5 \parallel$

- Cobordism defect of $AdS_5 \times S^5 \parallel$
- New and interesting double scaling limit

- Cobordism defect of $AdS_5 \times S^5 \parallel$
- New and interesting double scaling limit
- Holographic interpretation of a dynamical cobordism

- Cobordism defect of $AdS_5 \times S^5 \parallel$
- New and interesting double scaling limit
- Holographic interpretation of a dynamical cobordism
- The cobordism to nothing of the solution is self-consistent. If we add some consistent stuff on top of it (eg. probe branes) with potentially introduce an anomaly, it has to be miraculously cancelled.

- Cobordism defect of $AdS_5 \times S^5 \parallel$
- New and interesting double scaling limit
- Holographic interpretation of a dynamical cobordism
- The cobordism to nothing of the solution is self-consistent. If we add some consistent stuff on top of it (eg. probe branes) with potentially introduce an anomaly, it has to be miraculously cancelled.

