



Cosmological Chameleons,
String Theory

the Swampland

based on 2406.07614 [hep-th] with Miguel Montero and Gonzalo F. Casas

Ignacio Ruiz,

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1. Introduction: Desperately Seeking de Sitter

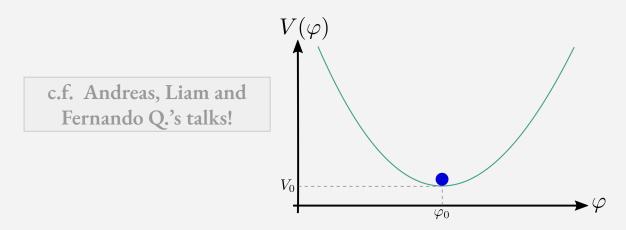
Why do we care about de Sitter?

Experimental measurements seem to indicate our Universe is accelerating, featuring a positive vacuum energy. [Supernova Search Team,'99]

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Several top-down constructions have been proposed in String Theory for a **positive** minimum $V(\vec{\varphi}_0) > 0$, but complete control remains elusive.



[de Alwis, Andriot, Balasubramanian, Bena, Bento, Berglund, Blåbäck, Blumenhagen, Cicoli, Chakraborty, Conlon, Flauger, Gao, Gligovic, Gorbenko, Graña, Gupta, Hebecker, Joyce, Junghans, Kachru, Kaddachi, Kallosh, Kovenski, Linde, Lüst, Maharana, McAllister, Obied, Ooguri, Polchinski, Parameswaran, Quevedo, Randall, Ruehle Sethi, Shiu, Silverstein, Spodyneiko, Toulikas, Tribedi, Vafa, Valandro, Wiesner, Xu, Zavala,...]

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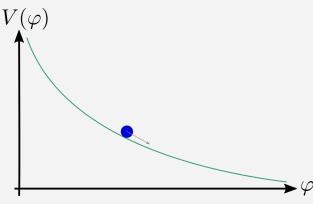
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Other option is a *flat enough* runaway potential $V(\vec{\varphi}) = V_0 e^{-\vec{\lambda} \cdot \vec{\varphi}}$ (*i.e.* quintessence)

with $|\vec{\lambda}| < \frac{2}{\sqrt{d-2}}$.

c.f. Susha and Francisco's talks!

c.f. Flavio's parallel!



[Andriot, Bedroya, Calderón-Infante, Cremonini, Cribiori, Erkinger, Gonzalo, Hebecker, Hertzberg, Kachru, Maldacena, Nuñez, Obied, Ooguri, Rajaguru, van Riet, **I.R.**, Schreyer, Seo, Shiu, Spodyneiko, Tang, Taylor, Tegmark, Tonioni, Tran, Vafa, Valenzuela, Venken, Wrase ...]

Technical difficulties in realizing this in String Theory remain!

Swampland Constraints on de Sitter and Quintessence

Difficulties in dS/Quintessence constructions and bottom-up arguments [Obied, Ooguri, tobservables, etc) motivate the **Swampland asymptotic** dS conjecture: Spodyneiko, Vafa, '18]

$$\lambda = \frac{\|\nabla V(\vec{\varphi})\|}{V(\vec{\varphi})} \ge c_d = \mathcal{O}(1) \text{ in asymptotic regions of } \mathcal{M}$$

For
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On the other hand the **Transplanckian Censorship Conjecture** prevents too long-lived epochs of accelerated expansion:

$$t_{\text{accel.}} \le \frac{1}{H_0} \log \left(\frac{M_{\text{Pl,d}}}{H_0} \right)$$

[Bedroya, Vafa, '19]

What about transient dS?

The current accelerated era has only lasted around $N \sim 0.5 e$ -folds!

$$t_{\rm accel.} \approx 4 \cdot 10^9 \text{ years} \ll 4.5 \cdot 10^{12} \text{ years}$$

While we *possibly* live in an asymptotic region of \mathcal{M} , we *might not* be in the *asymptotically accelerated regime!*

What if we are currently experiencing a transient dS era? [DESI collaboration, '24]

[Gomes, Hardy, Parameswaran,'23]; [Andriot, Tsimpis, Wrase, '23] [Andriot, Parameswaran, Tsimpis, Wrase, Zabala, '24], [Bhattacharya, Borghetto, Malhotra, Parameswaran, Tasinato, '24], [Alestas, Akrami, Delgado, Montero, Nesseris, I.R.,'24]; [Emparan, Garriga, Gutperle, Kallosh, Linde, Ohta, Roy, Russo, Strominger, Townsend, Wohlfarth ...] Review: [Cicoli, Conlon, Mahrana, Parameswaran, Quevedo, '24]

c.f. Joaquim's parallel!

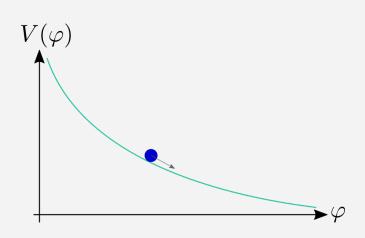
c.f. Yashar's parallel!

c.f. Lilia's parallel!

2. Cosmological Chameleons

We consider d-dimensional EFT with canonically normalized moduli $\{\phi^i\}_{i=1}^K$ and massive states $\{\chi_I\}_I$:

$$S^{(d)} \supseteq \frac{1}{2} \int d^d x \sqrt{-g} \left\{ \kappa_d^{-2} \left[\mathcal{R}_g - \sum_{i=1}^K (\partial \phi^i)^2 \right] - 2V(\vec{\phi}) - \sum_I m_I(\vec{\phi})^2 \chi_I^2 - \sum_I (\partial \chi_I)^2 \right\}$$



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Non-vanishing densities result in effective potential:

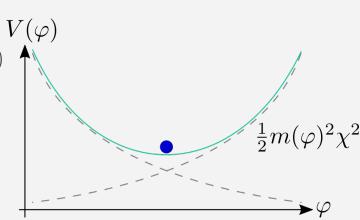
$$V_{\text{eff}}(\vec{\phi}) = V(\vec{\phi}) + \frac{1}{2} \sum_{I} m_{I}(\vec{\phi})^{2} \chi_{I}^{2} = V(\phi) - \frac{1}{d} \sum_{I} T_{\nu}^{(I)\nu}(\vec{\phi})$$

This can result in an effective positive minimum!

[Khouri, Weltman, '03]

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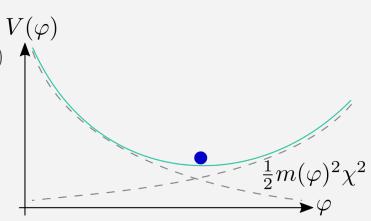
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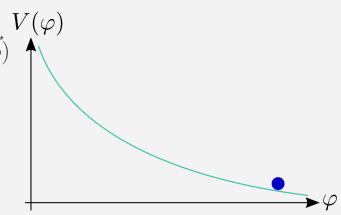
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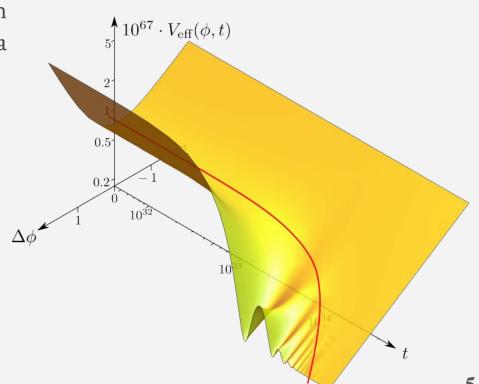
Case 1: Runaway potential + Heavy species

Consider a single scalar ϕ running down an exponential potential and encountering a (stable) scalar state χ becoming heavy:

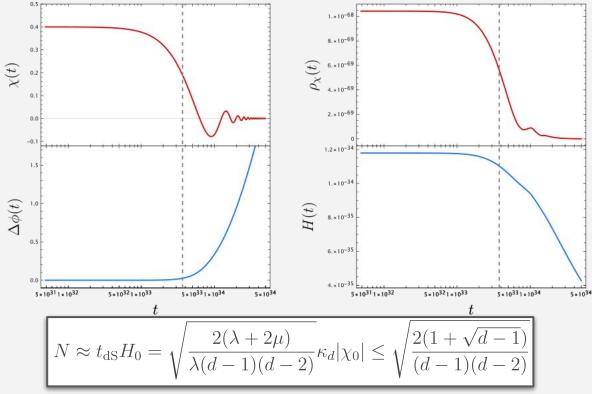
$$V_{\text{eff}}(\phi, t) = V_0 e^{-\lambda \phi} + \frac{1}{2} m_0^2 e^{2\mu \phi} \chi^2$$

Scalar is stabilized at effective minimum:

$$\phi(\chi) = \underbrace{\frac{1}{\lambda + 2\mu} \log \left(\frac{V_0}{m_0^2 \chi_0^2 \mu} \right)}_{\phi_0} - \log \left(\frac{\chi}{\chi_0} \right)^{\frac{2}{\lambda + 2\mu}}$$



Case 1: Runaway potential + Heavy species



for $\mu \le \sqrt{\frac{d-1}{d-2}}$, $\lambda > \frac{2}{\sqrt{d-2}}$ and $\kappa_d |\chi_0| \lesssim 1$.

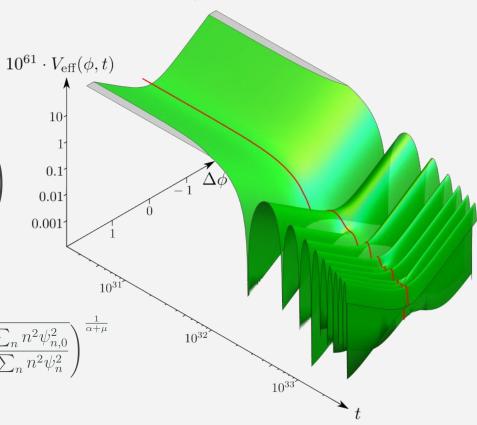
Case 2: Light tower + Heavy species

Now the scalar ϕ controls the masses of a stable state χ becoming heavy, as well as a tower of light scalars $\{\psi_n\}_n$:

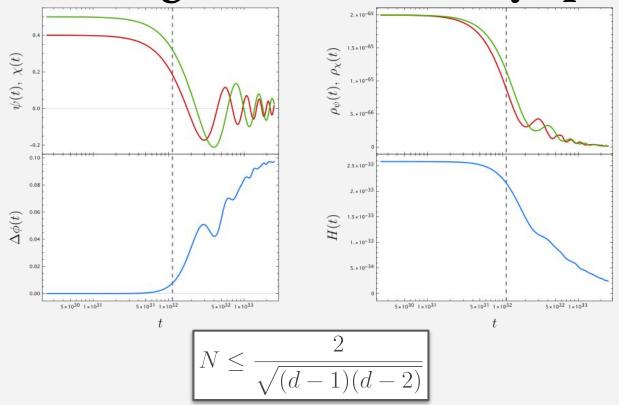
$$V_{\text{eff}}(\phi) = \frac{1}{2} \left(M_0^2 e^{-2\alpha\phi} \sum_n n^2 \psi_n^2 + m_0^2 e^{2\mu\phi} \chi^2 \right)$$

Scalar is stabilized at effective minimum:

$$\phi(\psi_n, \chi) = \underbrace{\frac{1}{\alpha + \mu} \log \left(\sqrt{\frac{\alpha}{\mu}} \frac{M_0}{m_0} \left| \frac{\sqrt{\sum_n n^2 \psi_{n,0}^2}}{\chi_0} \right| \right)}_{\chi_0} - \log \left(\frac{\chi}{\chi_0} \sqrt{\frac{\sum_n n^2 \psi_{n,0}^2}{\sum_n n^2 \psi_n^2}} \right)^{\frac{1}{\alpha + \mu}}$$



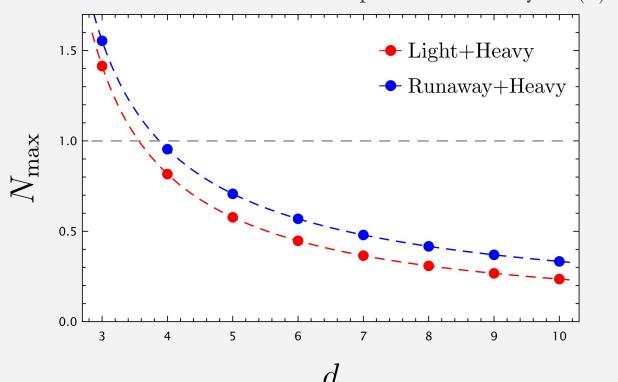
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for $\frac{1}{\sqrt{d-2}} \le \alpha$, $\mu \le \sqrt{\frac{d-1}{d-2}}$ and $\kappa_d |\chi_0| \lesssim 1$.

Number of e-folds supported

The total number of *e*-folds that this transient dS phase last is always $\mathcal{O}(1)$:





Attempt 1: M-th on CY3 close to SCFT point

Consider M-theory on $\mathbb{M}^4 \times X_3 \times S^1$, with all moduli stabilized but radion ρ of S^1 .

- Heavy states: M5-brane wrapping $\Sigma_4 \subset X_3$ cycle and S^1 .
- Runaway potential: Casimir energy

$$V_{\text{eff}}(\rho,\chi) = V_0 e^{-4\sqrt{\frac{2}{3}}\rho} + m_0^2 e^{\sqrt{\frac{2}{3}}\rho} \chi^2$$

This would result in $N \approx \frac{1}{2} \sqrt{\frac{5}{3}} \kappa_4 |\chi_0| < 0.6455 e$ -folds.

- ★ How do we stabilize the rest of moduli close to the SCFT point?
- ★ Planckian values of the fields might make corrections important!

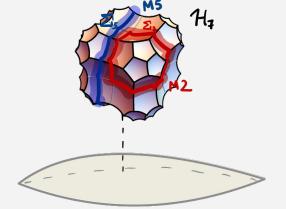
Attempt 2: M-th on hyperbolic manifold

Take M-theory compactified on compact hyperbolic 7-manifold \mathcal{H}_7 with torsion: All moduli are stabilized but volume ρ .

The effective potential comes from negative curvature and M2 and M5 branes wrapped on torsion 2- and 5-cycles:

$$V_{\text{eff}}(\rho,t) = M_{\text{Pl},4}^4 \frac{R_{g_7,0}}{\mathcal{V}_{24,0}^{2/7}} e^{-3\sqrt{\frac{2}{7}}\rho} + \frac{1}{2} \left[m_{\text{M2},0}^2 e^{-\sqrt{\frac{2}{7}}\rho} \psi(t)^2 + m_{\text{M5},0}^2 e^{\sqrt{\frac{2}{7}}\rho} \chi(t)^2 \right]$$

At most
$$N = \sqrt{\frac{2}{3}} \approx 0.82$$
 e-folds are supported.



Attempt 2: M-th on hyperbolic manifold

Possible problems:

- A hierarchy is needed: $Vol(\Sigma_2)^{1/2} \gg Vol(\Sigma_5)^{1/5}$. Interesting study on its own.
- Higher order terms might become important for Planckian values of massive states:
 - Interaction between M2 and M5 particles are suppressed at large volume.
 - Supersymmetry is broken: Not control on self-interactions between M2 and M5 as they intersect their world-volumes. More study is needed!

No full control is achieved!

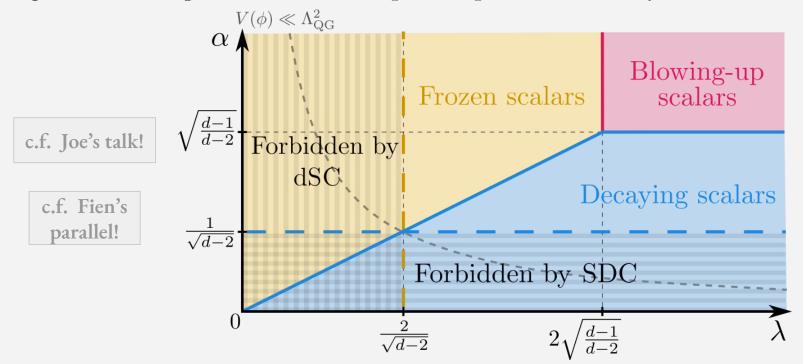


4. Bottom-up constraints from light towers of states.

[2407.XXXXX] with Gonzalo F. Casas

Evolution of light towers with runaway potential

An analogous study can be done of the evolution of light towers in a FLRW background with a potential: Bottom-up Swampland consistency constraints inferred!





Conclusions and Outlook

- Transient dS phases can be consistent with both observational and Swampland principles.
- Cosmological chameleons offer O(1) *e*-folds of accelerated expansion with natural ingredients in asymptotic regions of \mathcal{M} .
- Stringy realizations can be complicated due to higher order corrections and need of "well-behaved" heavy states: More research is needed!
- Other possibilities: monodromic axions?
- Cosmological evolution of towers of states can provide interesting insights!

THINK OUTSIDE THE BOX!

Thanks for listening!

