Based on a series of works with Ivano Basile, Niccoló Cribiori and Dieter Lüst. [2305.10489], *[2311.12113]*, [2401.06851]

Max-Planck-Institut für Physik (Werner-Heisenberg-Institut)

Carmine Montella, 27/06/2024 Presentation for String Pheno '24

Do black holes know about the emergent string conjecture?

The emergent string conjecture

- From a bottom-up point of view, an infinite distance limit in a space of vacua is a factorisation limit, i.e. a $N−$ point function can be reduced into $N−$ one point functions. *[Stout '21]*
- However, gravity abhors factorisation (due to equivalence principle), thus it must couple to an infinite tower of species. *[Stout '22]* In string theory: tower is asymptotically massless. *(SDC!)*
- The **Emergent String Conjecture** expresses the nature of the tower, stating that any infinite distance limit in the space of vacua is either a decompactification limit or a limit in which there is a weakly coupled (critical) string becoming tensionless. *[Lee, Lerche, Weigand '19]*

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The species scale

- The species scale is (an upper bound to) the cut-off of an effective theory of gravity.
	- 1. The scale at which perturbative gravity breaks down due to the presence of $N_{sp} \gg 1$

2. The cut-off scale Λ_{UV} appears in the higher derivative terms of an effective gravitational action. A modern view in string theory context: *[van de Heisteeg, Vafa, Wiesner, Wu '22-'23]*:

species *[Dvali '09]*:

$$
=\frac{M_{pl,d}}{N_{sp}^{d-2}}
$$

$$
S_{\text{EFT}} \sim \frac{M_{\text{pl},d}^{d-2}}{2} \int d^d x \sqrt{-g} \left(R + \sum_n \frac{c_n}{\Lambda_{\text{UV}}^{2n-2}} O_n(g, \text{Riem}, \nabla) \right).
$$

A depiction of the correspondence or transition between (asymptotically) massless species and minimal black hole

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- Fixing the energy E_{sp^\prime} then the micro-canonical entropy can be calculated as $S_{sp} = \log D(E_{sp})$, i.e. $Z(q) = \sum q^M D(M) = \prod$ *M*

 $S_{sp} \sim N_{sp} + \sum$ *n*≤*N*

$$
\prod_{n\leq N}\left(1-q^{\chi(n)}\right)^{-d_n}
$$

$$
\int_{N} d_{n} \log \frac{E_{sp}}{N_{sp} m_{n}} + \text{corr.}
$$

species energy as $E_{sp} \equiv \langle \mathcal{H} \rangle_0 \sim \sum d_n m_n + \text{th. corr.}$ ̂

• In order to study a possible ToS/black hole transition (or correspondence), we define the

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\sum_{n \leq N} d_n m_n + \text{th. corr.}
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A. The question we asked ourselves is: given leading towers at infinite distance limit, when is it possible to have a correspondence/transition with the smallest black hole in the EFT?

In a given EFT, any consistent tower at infinite distance limit must allow a tower-black hole transition (or

• Due to $S_{sp} \sim \Lambda_{sp}^{2-d}$ + corr. for the leading towers, we impose $E_{sp} = \gamma \Lambda_{sp}^{3-d}$ + corr, for <u>any point</u> at

- **correspondence) and viceversa [Basile, Lüst, Montella '23]**
- infinite distance of the space of vacua.

A bottom-up approach

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- This is a parametrisation of a KK-tower due to a p-torus! [Castellano, Herráez, Ibáñez '21]

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- This correction is lead by higher derivative corrections, and by massless species [Tian, Xiao '21]

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S_{\text{EFT}} = \int_{\mathcal{M}} \frac{R}{16\pi} + \left[c_1(\mu)R^2 + c_2(\mu)R_{\mu\nu}R^{\mu\nu} + c_3(\mu)R_{\mu\nu\rho\sigma}R^{\mu\nu\rho\sigma} \right] - \left[\alpha R \ln(\frac{\Box}{\mu^2})R + \beta R_{\mu\nu} \ln(\frac{\Box}{\mu^2})R^{\mu\nu} + \zeta R_{\mu\nu\alpha\beta} \ln(\frac{\Box}{\mu^2})R^{\mu\nu\alpha\beta} \right],
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$$

• Again, for large $N: \chi(N) = \frac{1}{\Gamma(N+1)} \frac{(N+1)(N+1)}{(N+1)} \sim N^{\frac{1}{p}}$, now with $\Gamma(N + \sigma + 1)\Gamma(\alpha + 1)$ $\Gamma(N+1+\alpha)\Gamma(\sigma+1)$

$$
\sqrt{E_{sp}}: \quad E_{sp} = \gamma \Lambda_{sp}^{3-d} + \omega \Lambda^{\alpha \ge 0}
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$$
\frac{(r+1)}{(r+1)} \sim N^{\frac{1}{p}}, \text{ now with } p \geq 1!
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• In general, <u>every</u> different solution (m_n, d_n) always satisfies the <u>same</u> implicit pattern between $\Lambda_{\rm SD}$ and *m*tow

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 $\hat{p} = + \infty$

• We can state that at infinite distance limit only **two** classes of towers allow a tower-black holes transition (or correspondence) [Basile, Lüst, Montella '23]

$$
\Lambda_{\text{sp}} \sim m_{\text{tow}}^{\frac{\hat{p}}{\hat{p}+d-2}} M_{\text{pl,d}}^{\frac{d-2}{d-2+\hat{p}}} \qquad \Lambda_{\text{sp}} \sim m_{\text{tow}}
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• We can state that at infinite distance limit only **two** classes of towers allow a tower-black holes transition (or correspondence) [Basile, Lüst, Montella '23]

- A. From a bottom-up approach, \hat{p} $(\hat{c}(m_n, d_n)) \ge 1$ represent a generic parameter defined by the black hole thermodynamics, and define different parametrizations of the "microstates".
-
-

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\Lambda_{\text{sp}} \sim m_{\text{tow}}^{\frac{\hat{p}}{\hat{p}+d-2}} M_{\text{pl,d}}^{\frac{d-2}{d-2+\hat{p}}} \qquad \Lambda_{\text{sp}} \sim m_{\text{tow}}
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$$
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 $M_{\rm s}$ *g*2 *s*

$$
T_{sp} \sim M_s \equiv T_{Hag} \qquad M_c \sim M_s^{3-d} = \frac{m_s}{a^2} \quad \text{in Planck units.}
$$

 $\hat{p} = + \infty$

B. From a top-down perspective it corresponds to the number of extra dimensions! $\longrightarrow m_{\rm low} \equiv m_{\rm KK}$, $\Lambda_{\rm sp} \sim M_{pl,d+\hat{p}}$

The limit $\hat{p}\to\infty$ is not defined in terms of mass—degeneracy, but it is always well defined for every thermodynamics quantities. The tower-black hole transition returns (the already well known) string-BH transition —*→* $m_{\rm{LOW}}$ $\equiv M_{\rm{S}}$ *,* Λ_{sp} $\sim M_{\rm{S}}$

Thanks!

Questions?